

Majoranyzacja elektronów w nadprzewodnikach topologicznych

**Tadeusz Domański
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J. Bardeen



E. Majorana

Outline:

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- Majorana fermions:

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⇒ what are they ?

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- Majorana madness ...

Basic notions

– on Majorana fermions

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- P. Dirac (1928)

$$i\dot{\psi} = (\vec{\alpha} \cdot \vec{p} + \beta m) \psi$$

/ relativistic description of fermions /

particles ($E > 0$),

anti-particles ($E < 0$)

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⇒ Basic features:

chargeless

zero-energy

Searching for majoranas

– in particle and nuclear physics

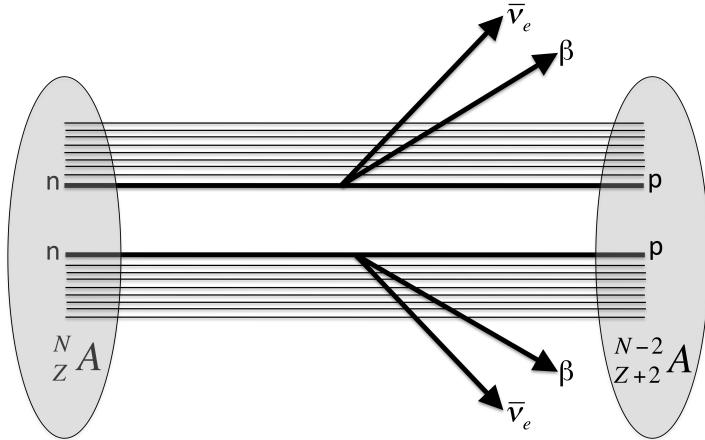
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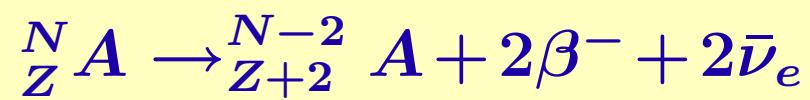
DOUBLE BETA DECAY

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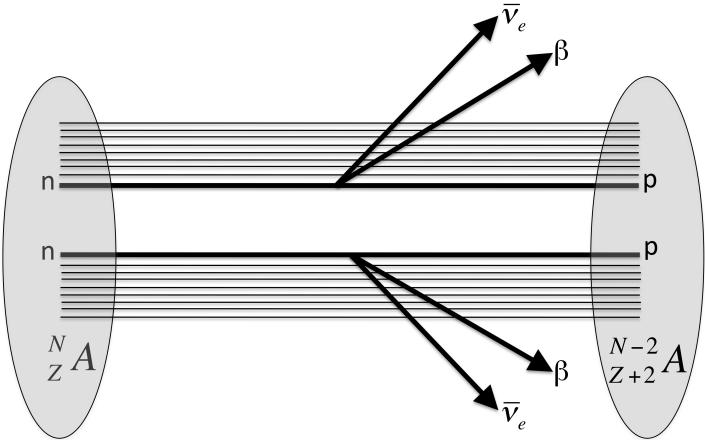


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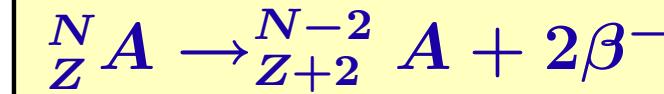
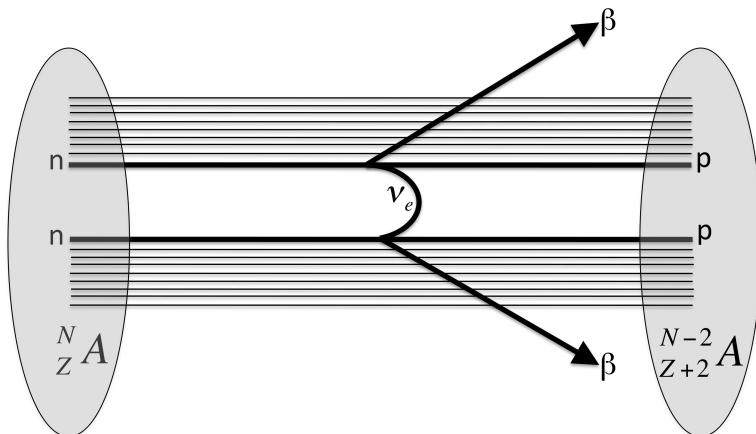


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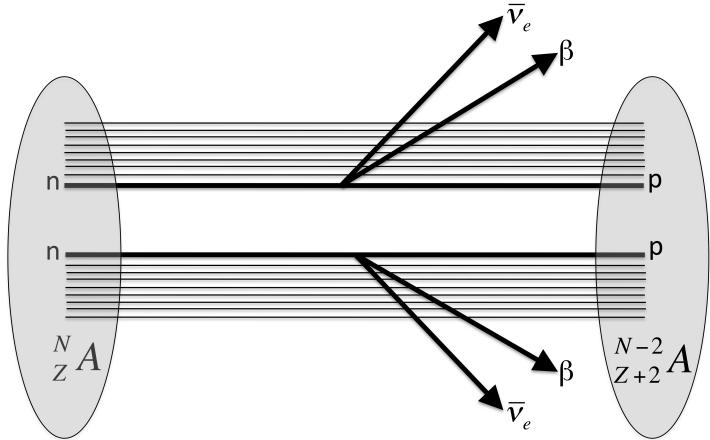
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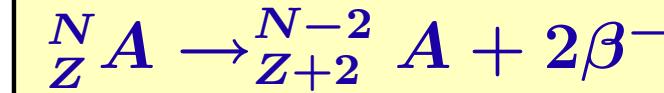
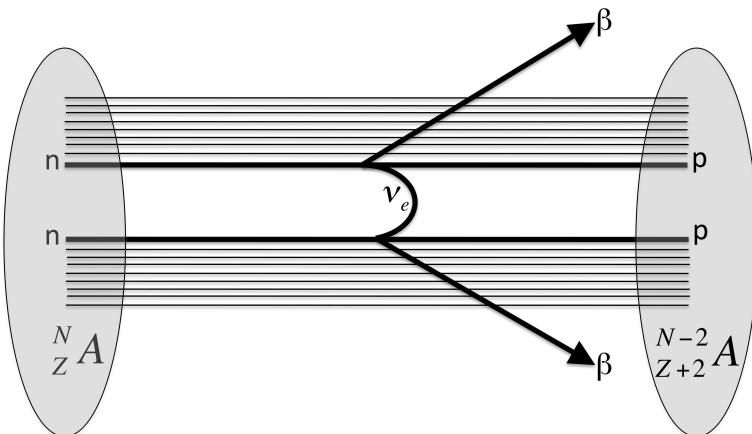
Neutrinoless decay would imply neutrinos to be majoranas.

Searching for majoranas

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DOUBLE BETA DECAY



Neutrinoless decay would imply neutrinos to be majoranas.

This issue is still under dispute.

Quasiparticles

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- ★ So, how about majoranions ?
- ★ Formaly, any usual fermion can be majoranized ...

Majorization – of normal fermions

- Normal fermions (e.g. electrons) obey the anticommutation relations

$$\{\hat{c}_i, \hat{c}_j^\dagger\} = \delta_{i,j}$$

$$\{\hat{c}_i, \hat{c}_j\} = 0 = \{\hat{c}_i^\dagger, \hat{c}_j^\dagger\}$$

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- $c_j^{(\dagger)}$ can be recast in terms of Majorana operators

$$\hat{c}_j \equiv (\hat{\gamma}_{j,1} + i\hat{\gamma}_{j,2}) / \sqrt{2}$$

$$\hat{c}_j^\dagger \equiv (\hat{\gamma}_{j,1} - i\hat{\gamma}_{j,2}) / \sqrt{2}$$

'real' and 'imaginary' parts

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$\hat{\gamma}_{i,n}$ correspond to neutral objects

- Exotic properties

$$\hat{\gamma}_{i,n}^\dagger = \hat{\gamma}_{i,n}$$

$$\{\hat{\gamma}_{i,n}, \hat{\gamma}_{j,m}^\dagger\} = \delta_{i,j}\delta_{n,m}$$

creation = annihilation !

fermionic antisymmetry

Majorization

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- Exotic properties (cd)

$$\hat{\gamma}_{i,n} \hat{\gamma}_{i,n} = 1/2$$

$$\hat{\gamma}_{i,n}^\dagger \hat{\gamma}_{i,n} = 1/2$$

no Pauli principle !

half ‘occupied’ & half ‘empty’

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- Majorana-type quasiparticles

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- resemble the Bogoliubov qps of BCS theory

$$\hat{\beta}_{k\uparrow} \equiv u_k \hat{c}_{k\uparrow} + v_k \hat{c}_{-k\downarrow}^\dagger$$

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quasiparticle charge is $u_k^2 - v_k^2$

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quasiparticle charge is $u_k^2 - v_k^2$

- At the Fermi level $u_{k_F} = v_{k_F} = 1/\sqrt{2}$, and then

$$\hat{\beta}_{k_F\uparrow} \longleftrightarrow \hat{\gamma}_{k_F,1}$$

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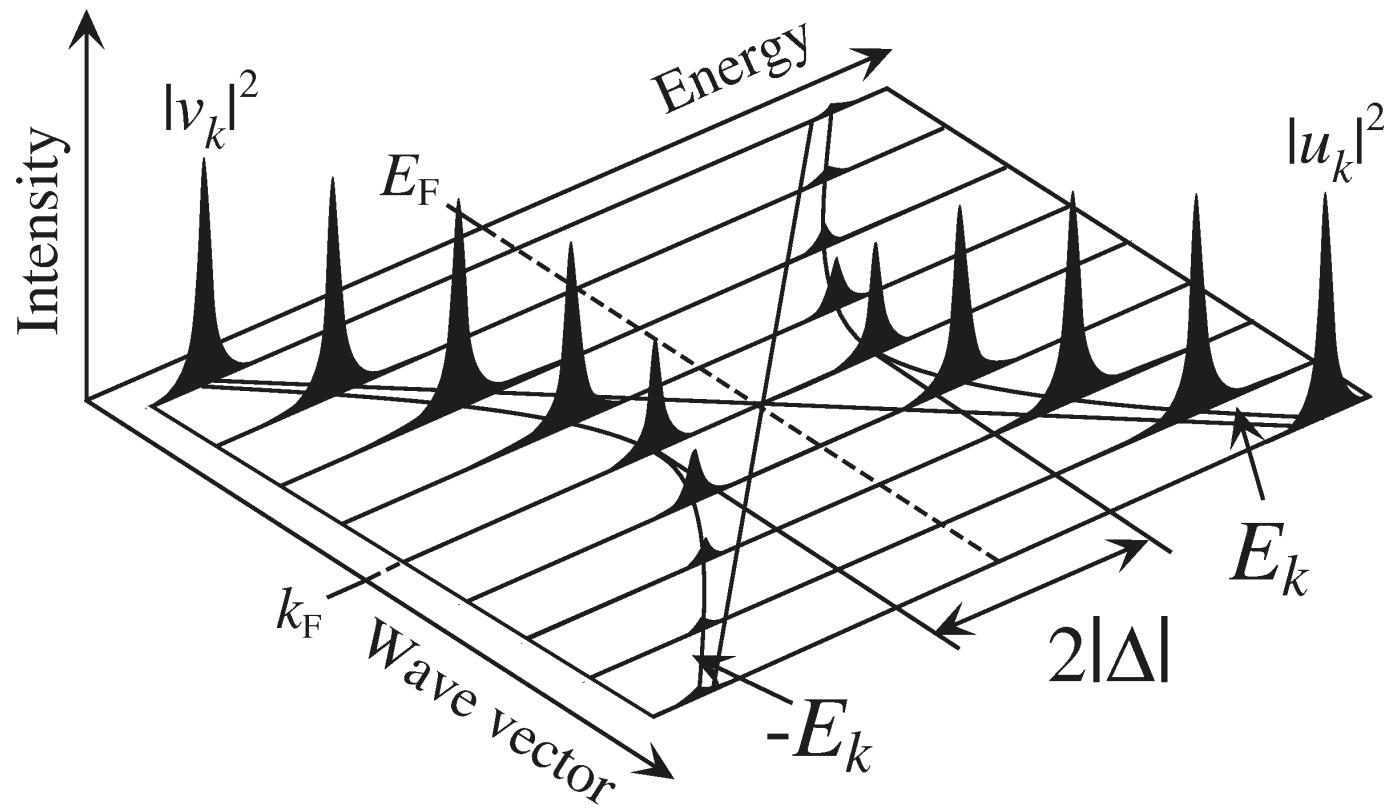
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OK, but what about their energy ?

Bogoliubov quasiparticles

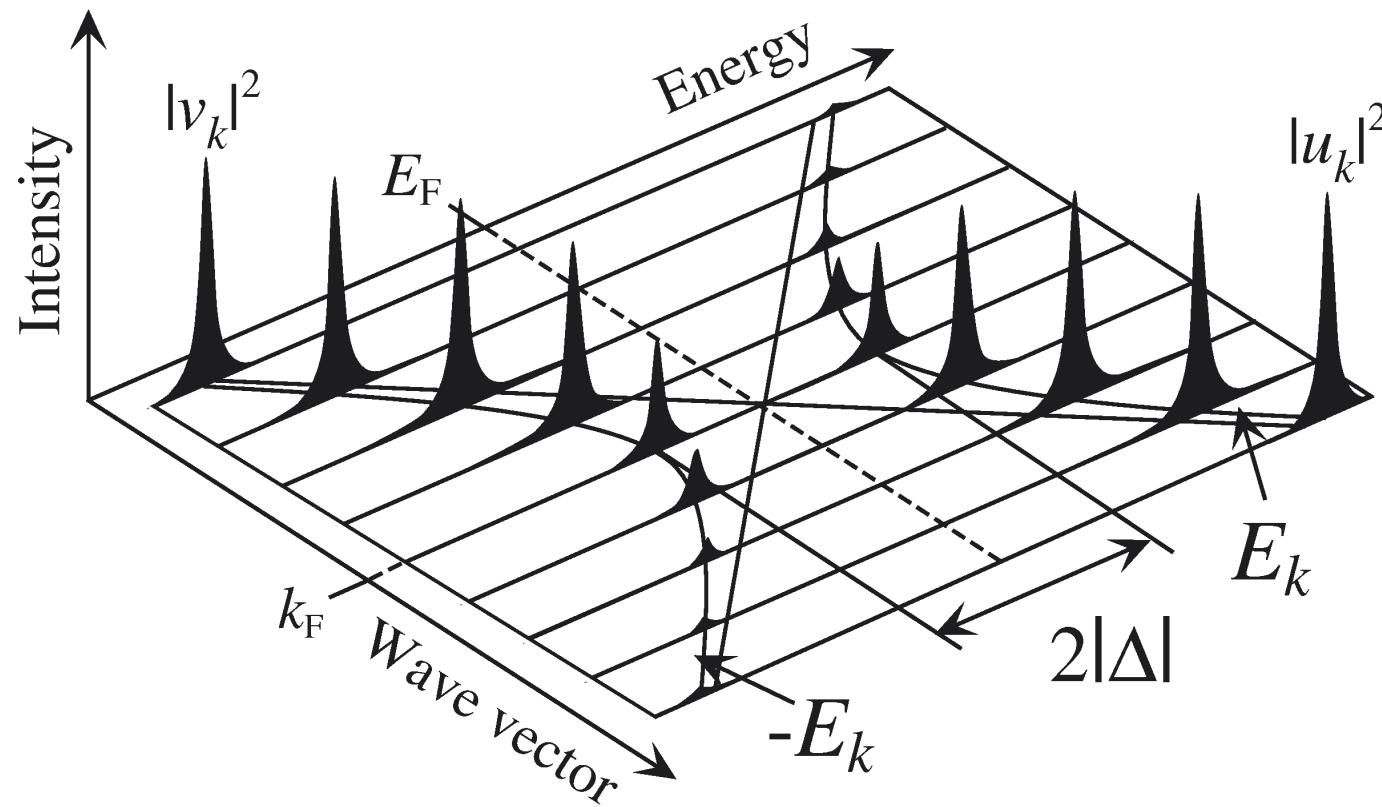
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Bogoliubov qps of s-wave superconductors are gapped $\pm\sqrt{\epsilon_k^2 + \Delta^2}$.

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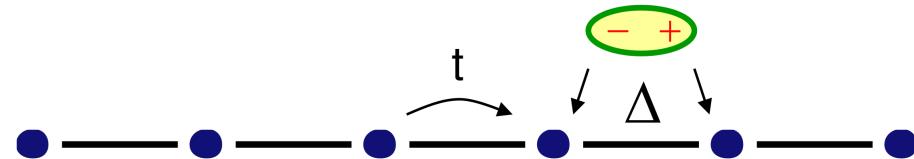


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True majoranas have to be the zero energy ($E_k = 0$) quasiparticles !

Kitaev chain

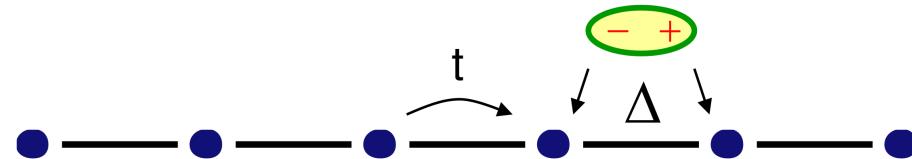
– a paradigm for Majorana modes



p-wave pairing of spinless 1D fermions

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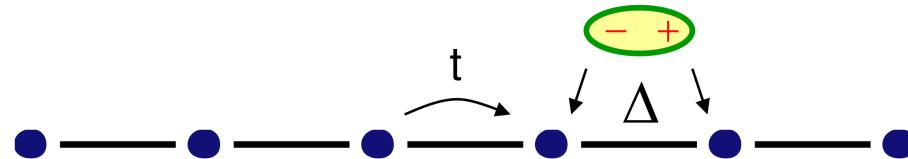


p-wave pairing of spinless 1D fermions

$$\hat{H} = t \sum_i (\hat{c}_i^\dagger \hat{c}_{i+1} + \text{h.c.}) - \mu \sum_i \hat{c}_i^\dagger \hat{c}_i + \Delta \sum_i (\hat{c}_i^\dagger \hat{c}_{i+1}^\dagger + \text{h.c.})$$

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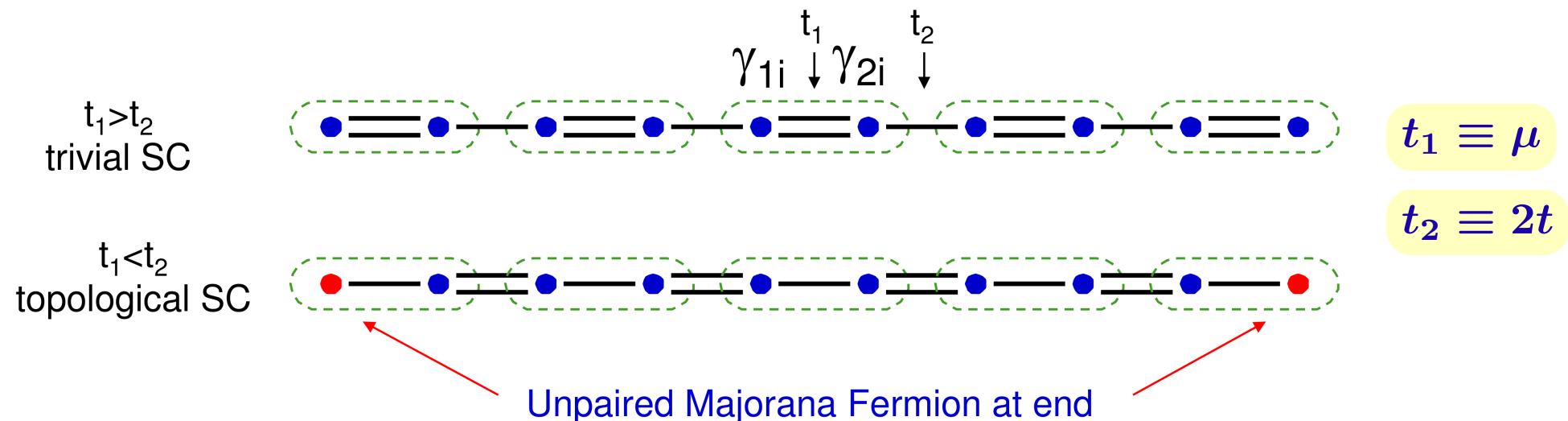
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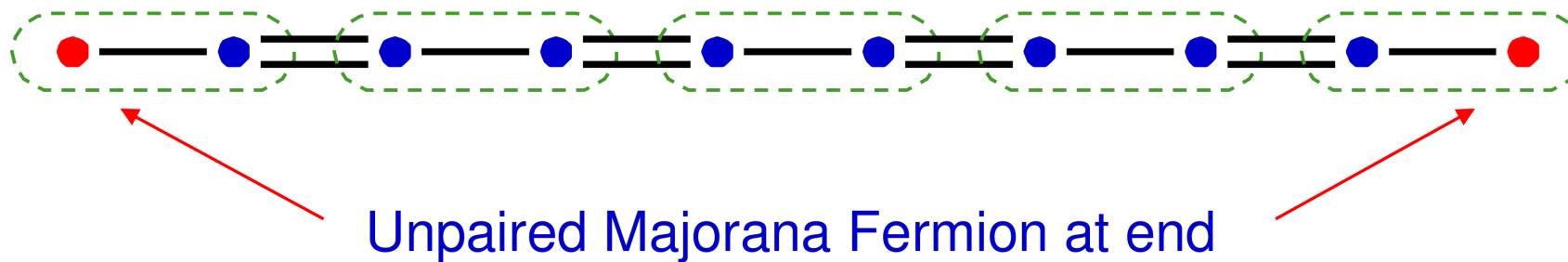
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This toy-model can be exactly solved in Majorana basis. For $\Delta = t$ one obtains:



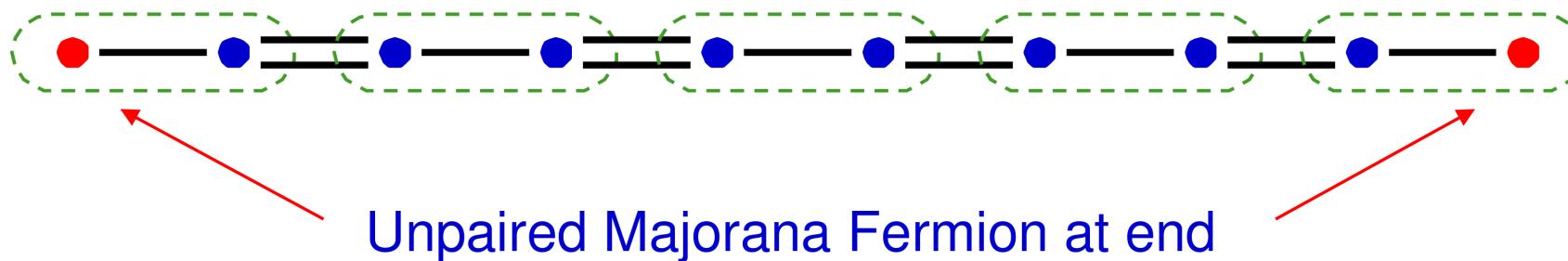
★ In the special case $\Delta = t$ and $|\mu| < 2t$



operators $\hat{\gamma}_{1,1}$ and $\hat{\gamma}_{2,N}$ are decoupled from all the rest. This implies

zero-energy modes appearing at the chain edges

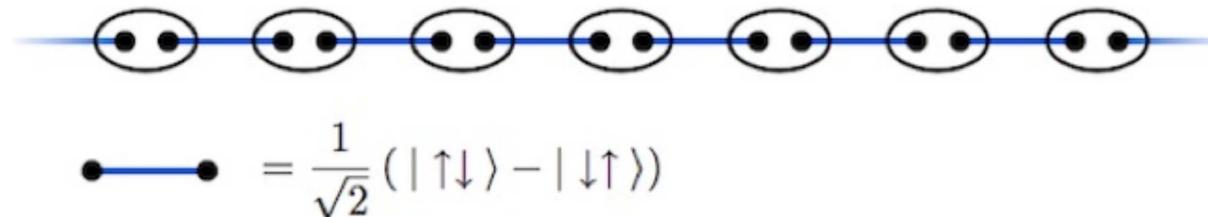
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★ Similar ideas have been considered for 1D Heisenberg chain of 1/2 spins



F.D.M. Haldane, Phys. Rev. Lett. 50, 1153 (1983)

Nobel Prize, 2016

Various scenarios

- for Majorana quasiparticles

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- vortex states in p -wave superconductors

Volovik (1999)

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- **quantum nanowires attached to superconductor**

Alicea (2010); Oreg *et al* (2010); Lutchyn *et al* (2010); Stanescu & Tewari (2013)

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- **edge states of 2D and 3D topological insulators**

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- **magnetic atoms chain on superconducting substrate**

Choy *et al* (2011); Martin & Morpugo (2012); Nadj-Perge *et al* (2013)

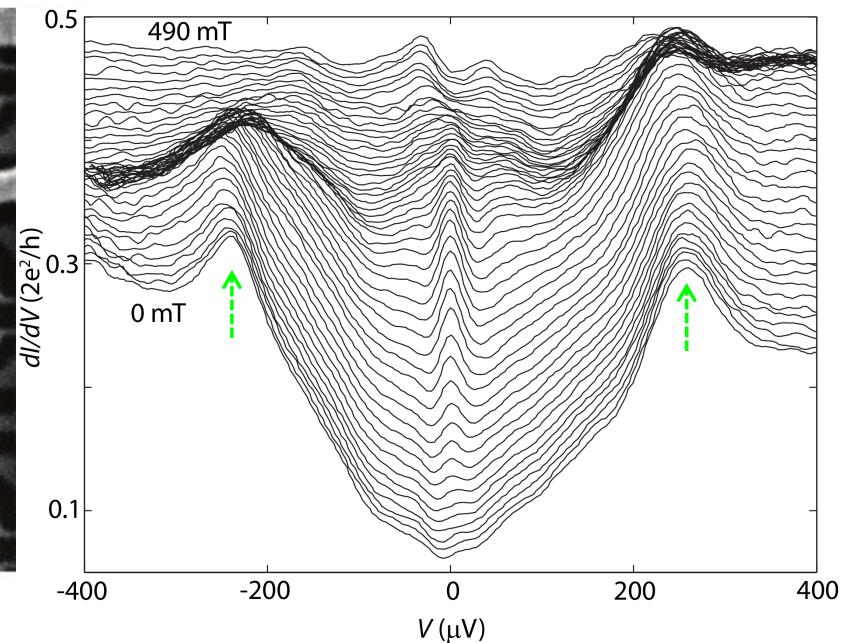
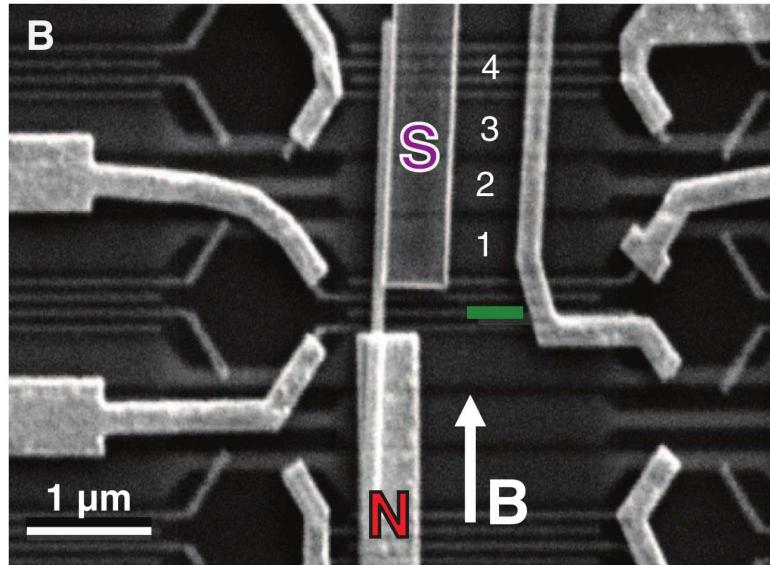
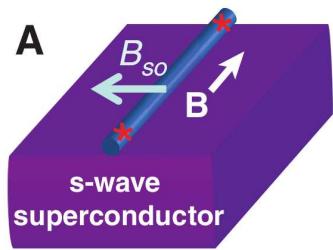
Experimental evidence

– for Majorana quasiparticles

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InSb nanowire between a metal (gold) and a superconductor (Nb-Ti-N)



dI/dV measured at 70 mK for varying magnetic field B indicated:

⇒ a zero-bias enhancement due to Majorana state

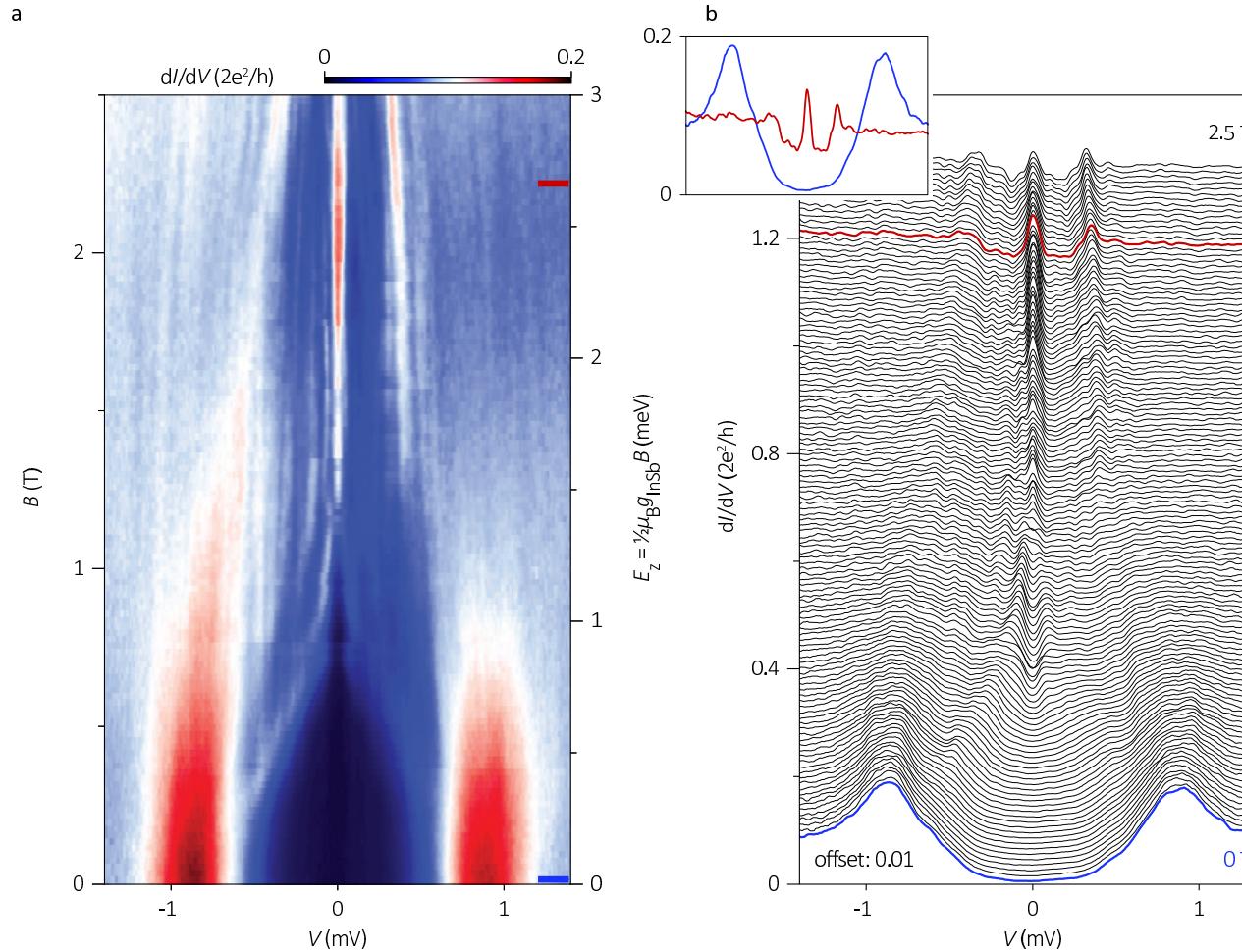
V. Mourik, ..., and L.P. Kouwenhoven, Science **336**, 1003 (2012).

/ Kavli Institute of Nanoscience, Delft Univ., Netherlands /

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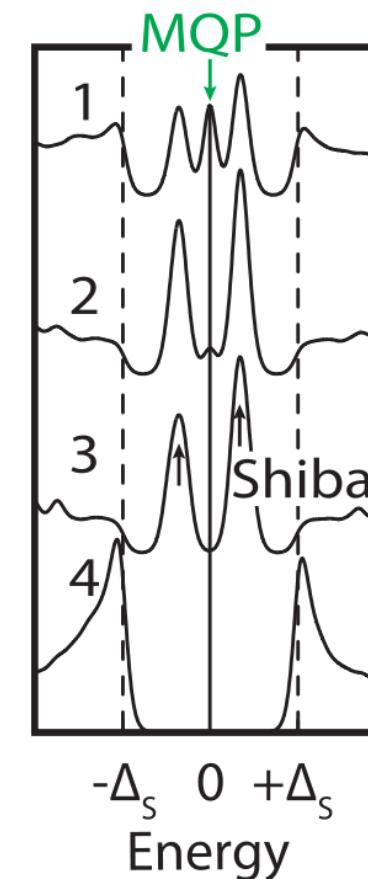
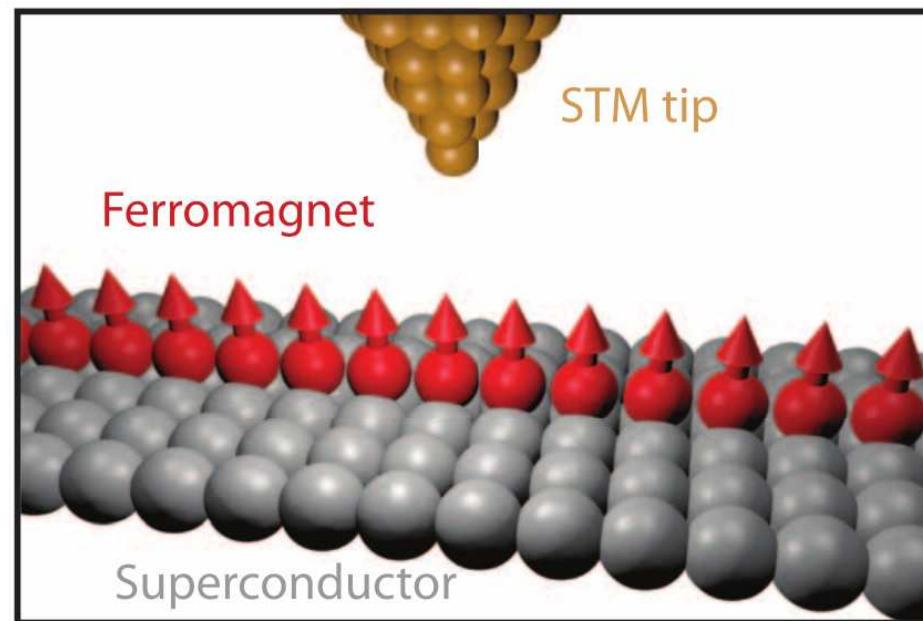
H. Zhang, ..., and L.P. Kouwenhoven, arXiv:1603.04069 (2016).

/ Kavli Institute of Nanoscience, Delft Univ., Netherlands /

Experimental evidence

– for Majorana quasiparticles

A chain of iron atoms deposited on a surface of superconducting lead



STM measurements provided evidence for:

⇒ Majorana bound states at the edges of a chain.

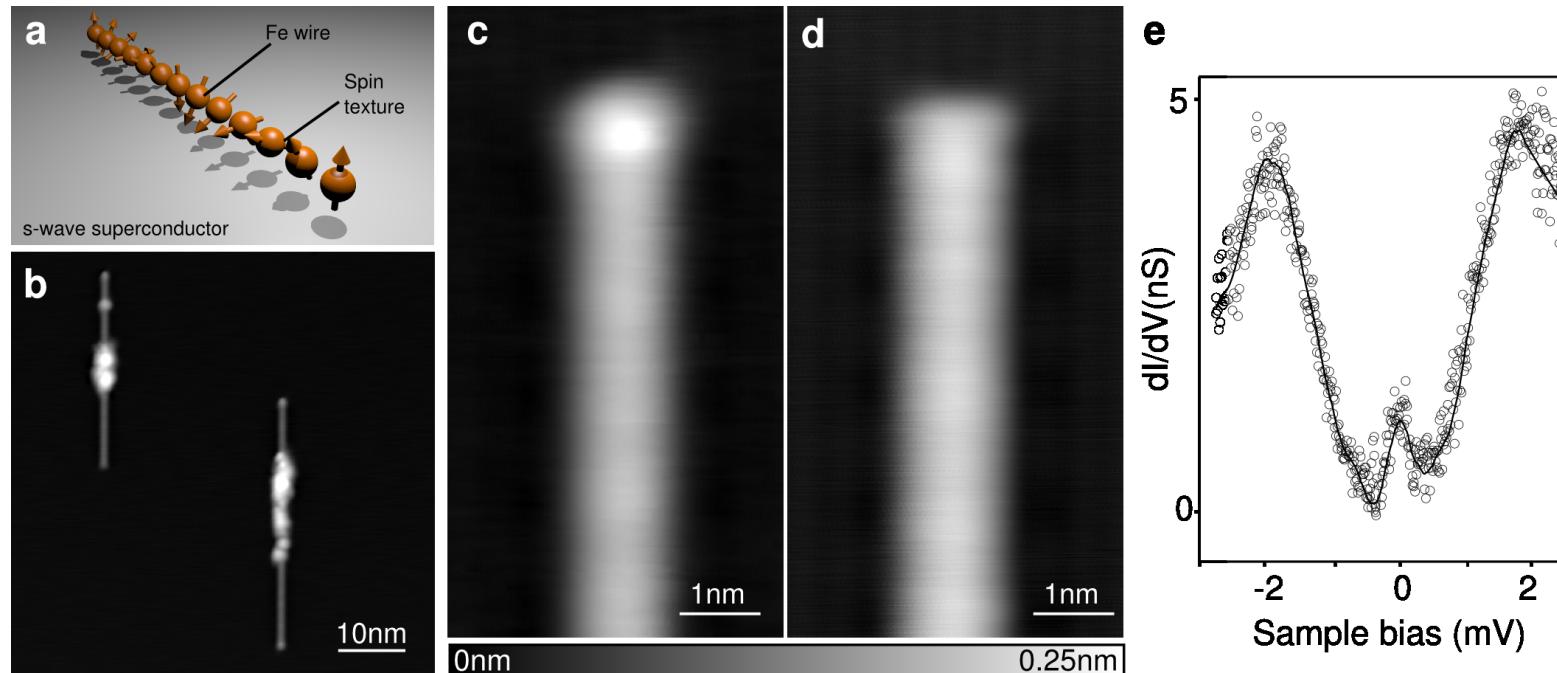
S. Nadj-Perge, ..., and A. Yazdani, Science 346, 602 (2014).

/ Princeton University, Princeton (NJ), USA /

Experimental evidence

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Self-assembled Fe chain on superconducting Pb(110) surface



AFM combined with STM provided evidence for:

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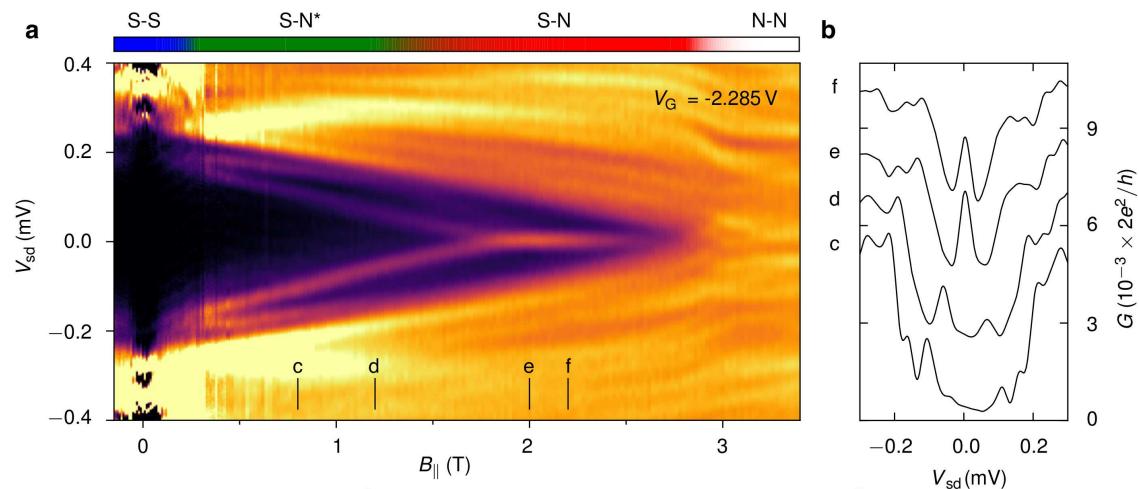
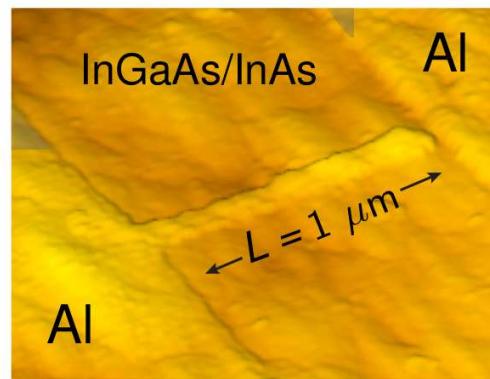
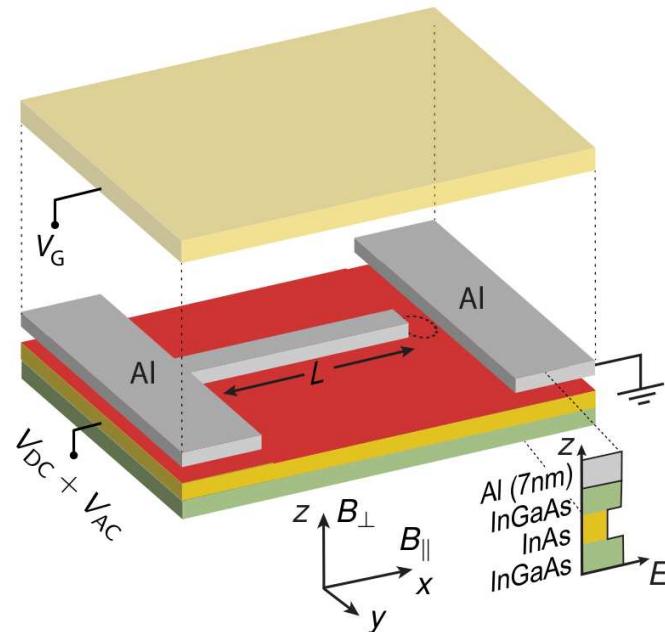
R. Pawlak, M. Kisiel *et al*, npj Quantum Information 2, 16035 (2016).

/ University of Basel, Switzerland /

Experimental evidence

for Majorana quasiparticles

Wire-like device constructed lithographically



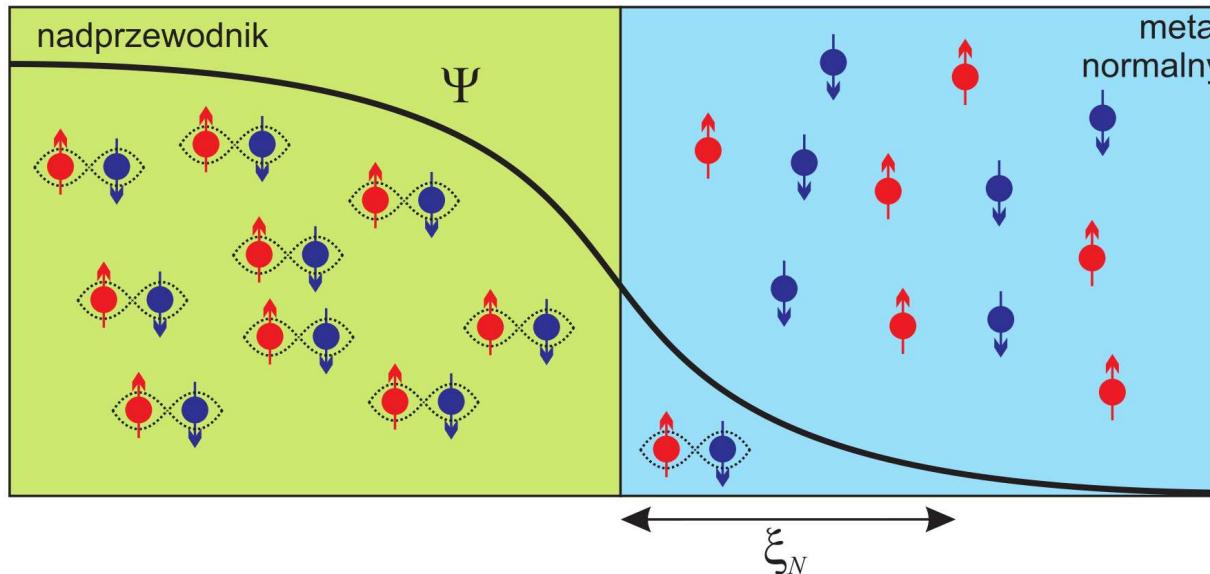
H.J. Suominen *et al*, arXiv:1703.03699 (2017).

/ University of Copenhagen, Denmark /

2. Electron pairing in nanosystems

Superconductivity in nanosystems

Any material contacted with a bulk superconductor absorbs the Cooper pairs

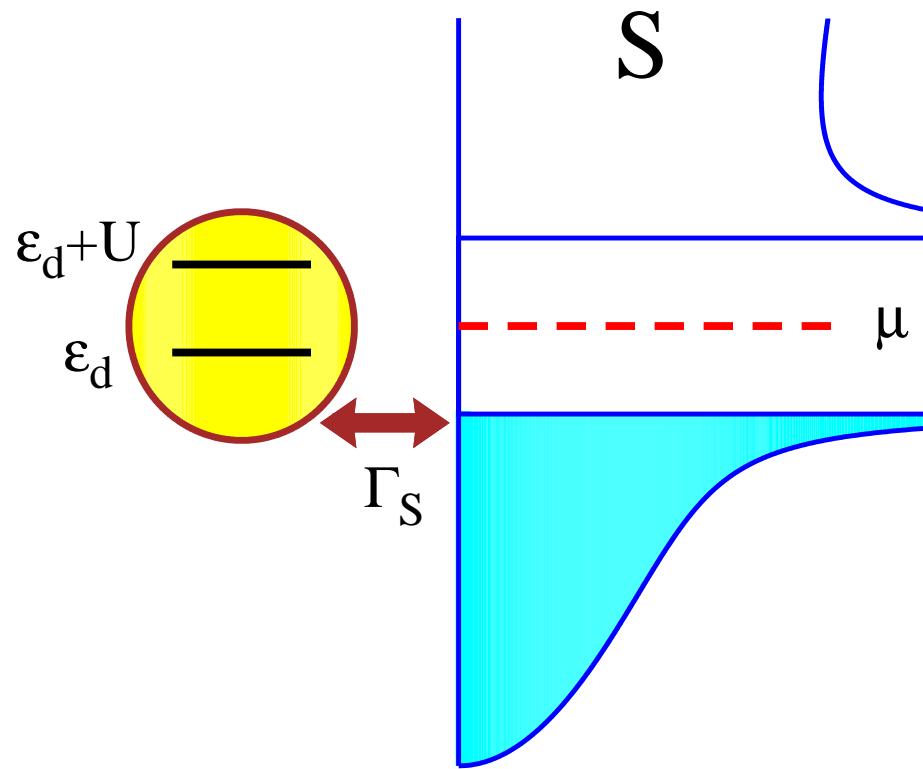


Size of the quantum dots is usually smaller than ξ_N

Prototype model

– single Anderson impurity

The single quantum impurity (dot) coupled to superconducting reservoir



ε_d – energy level, U – Coulomb potential, Γ_S – hybridization

Microscopic model

Anderson-type Hamiltonian

Quantum impurity (dot)

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

coupled with a superconductor

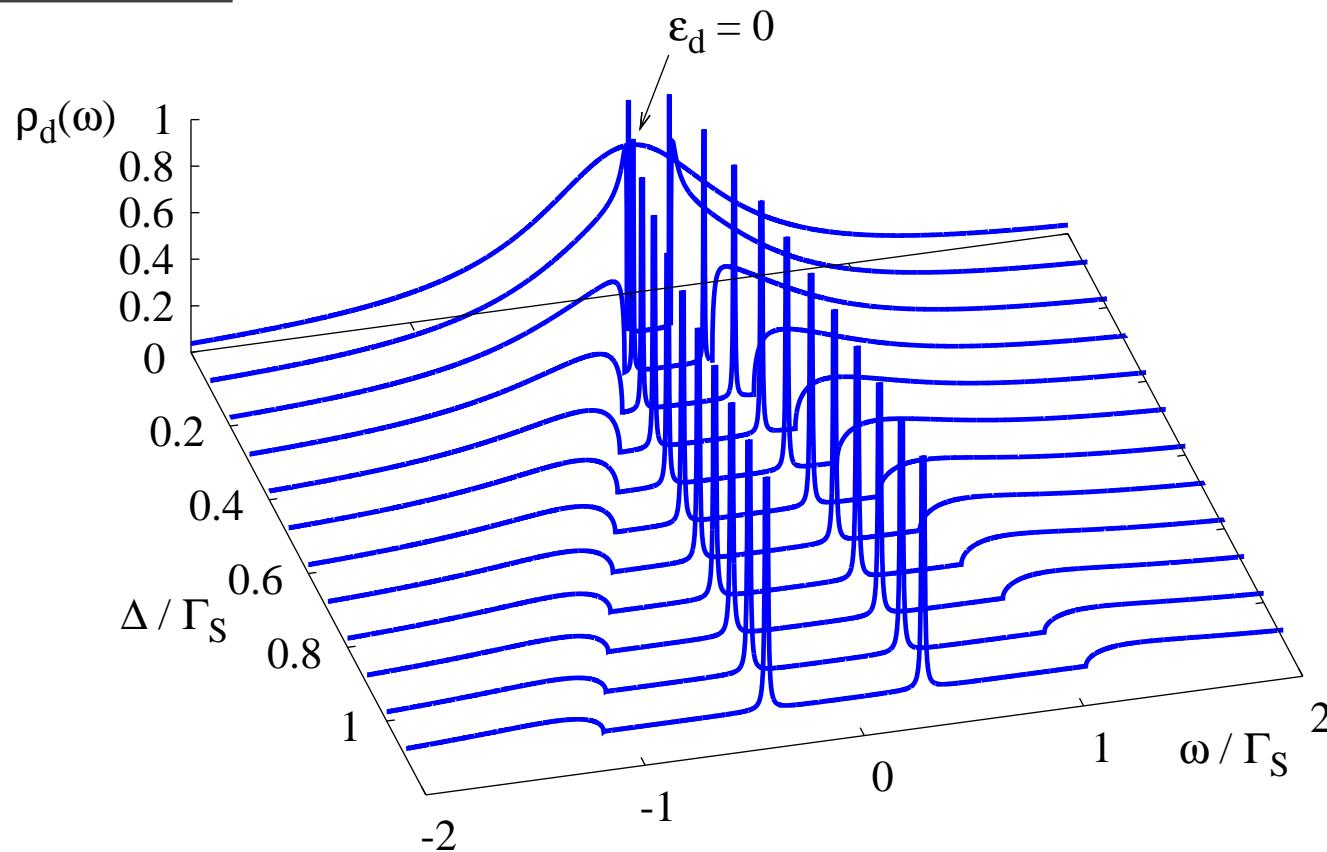
$$\begin{aligned} \hat{H} &= \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} + \hat{H}_S \\ &+ \sum_{\mathbf{k}, \sigma} \left(V_{\mathbf{k}} \hat{d}_{\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + V_{\mathbf{k}}^* \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{d}_{\sigma} \right) \end{aligned}$$

where

$$\hat{H}_S = \sum_{\mathbf{k}, \sigma} (\epsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} - \sum_{\mathbf{k}} \left(\Delta \hat{c}_{\mathbf{k}\uparrow}^{\dagger} \hat{c}_{\mathbf{k}\downarrow}^{\dagger} + \text{h.c.} \right)$$

Uncorrelated QD

$U_d = 0$ (exactly solvable case)

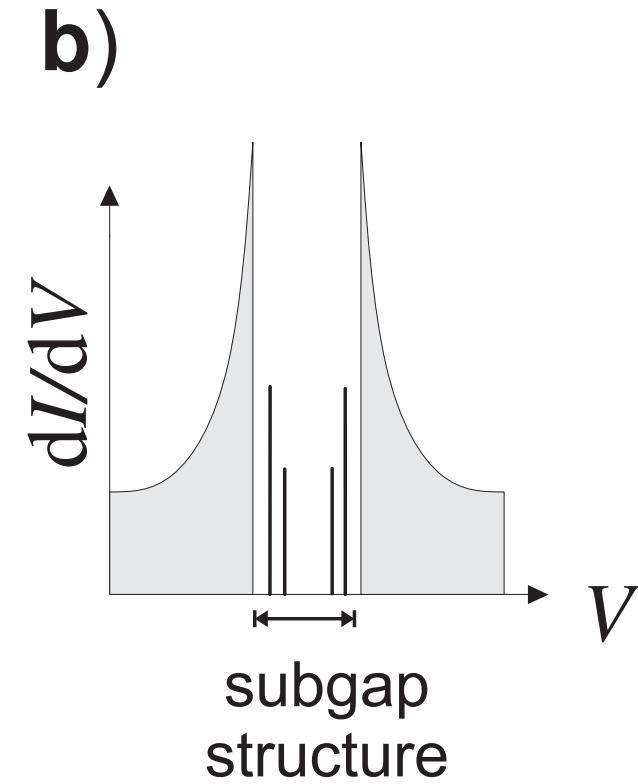
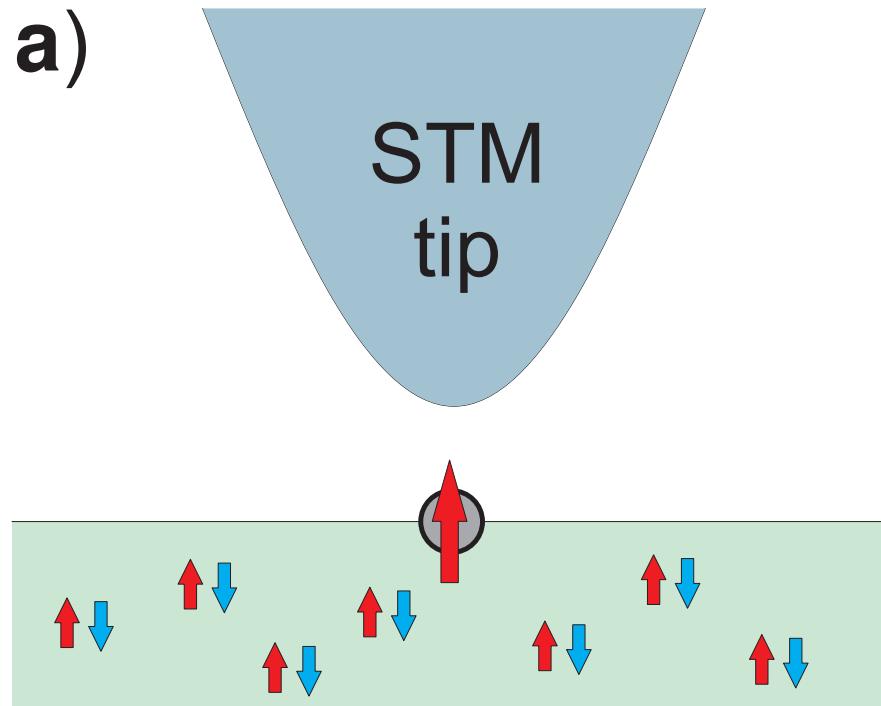


In-gap (Andreev/Shiba) bound states :

- ⇒ always appear in pairs,
- ⇒ appear symmetrically at finite energies.

Subgap states

of multilevel quantum impurities



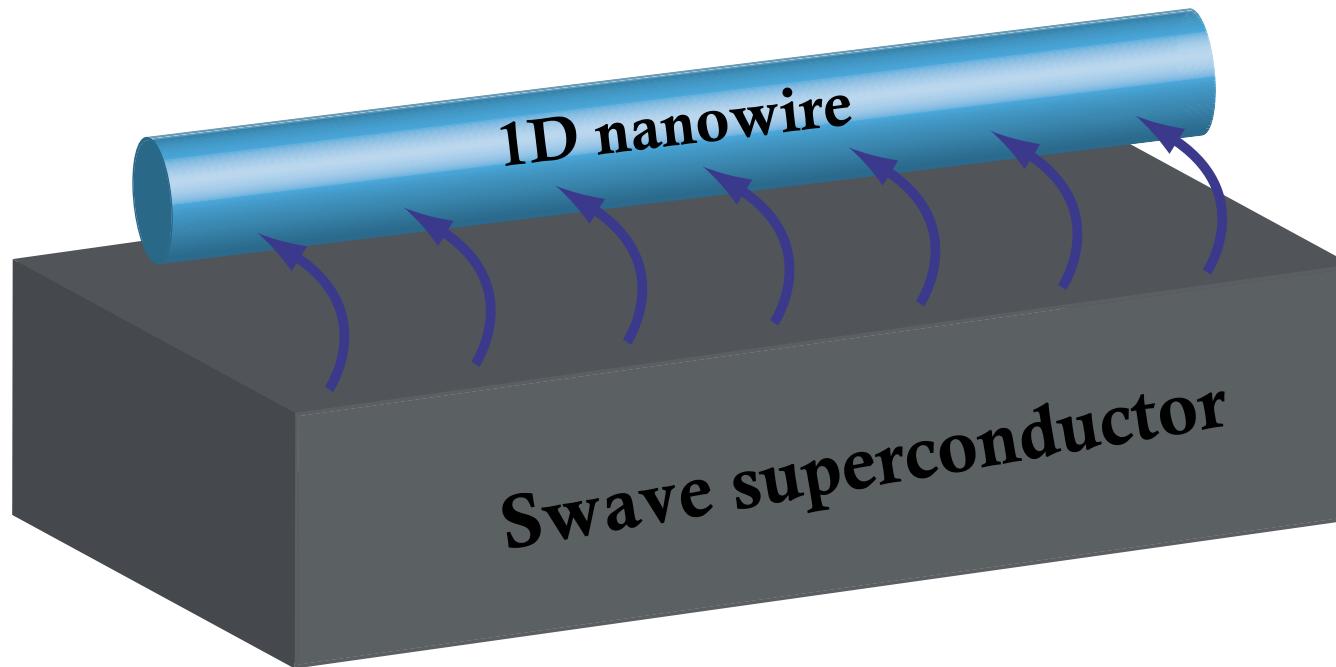
a) STM scheme and b) differential conductance for multilevel quantum impurity adsorbed on surface of bulk superconductor.

R. Žitko, O. Bodensiek, and T. Pruschke, Phys. Rev. B **83**, 054512 (2011).

Andreev vs Majorana states

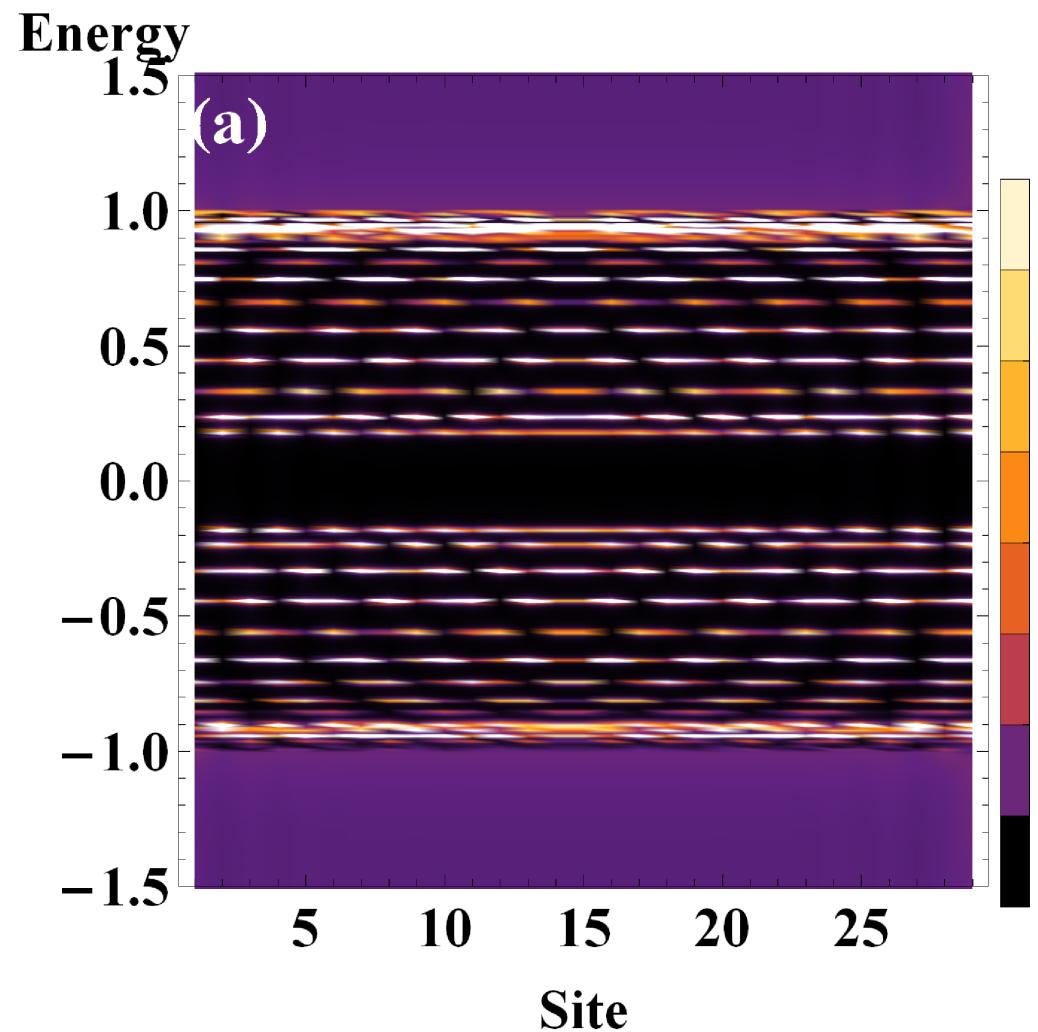
– **'A story of mutation'**

Let us consider:



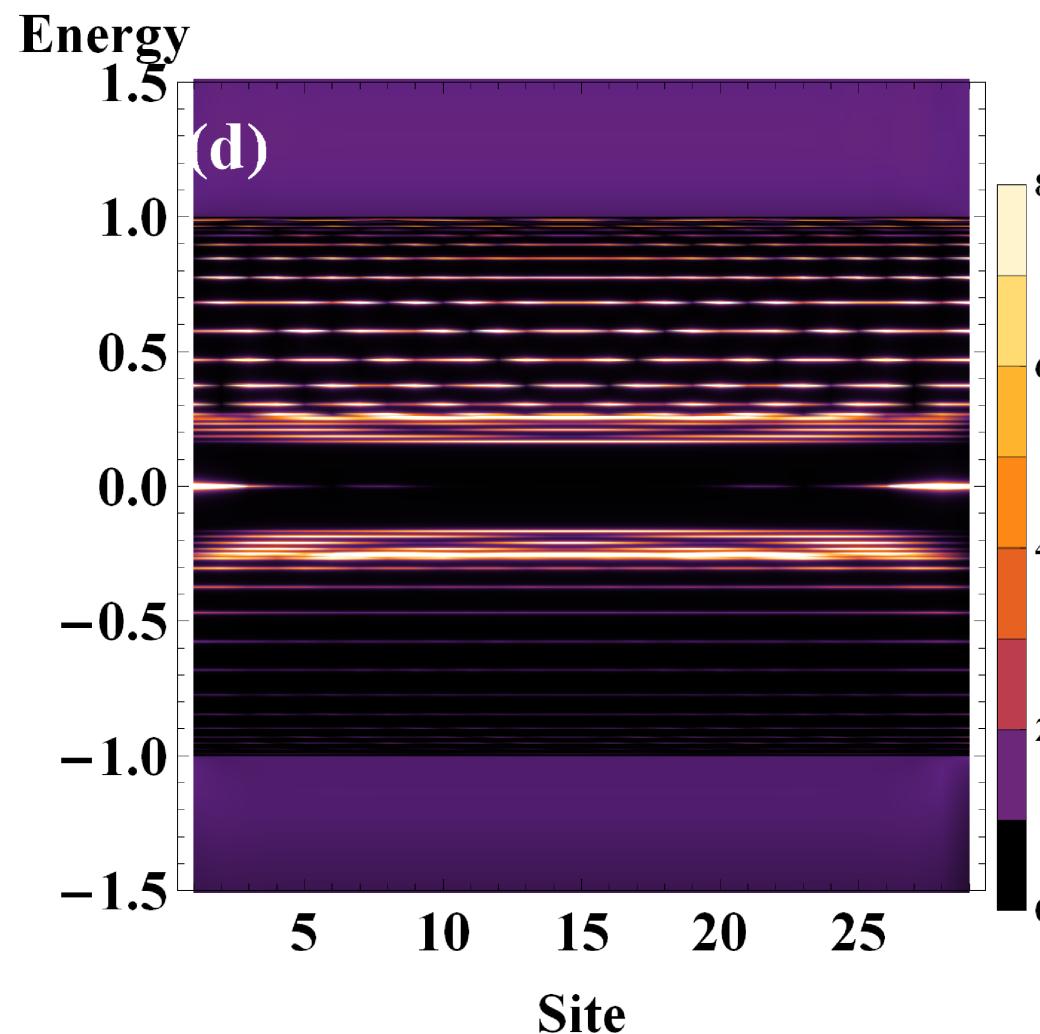
1D quantum wire deposited on s-wave superconductor

*D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B **88**, 165401 (2013).*



Electronic spectrum comprises a series of Andreev states.

D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B 88, 165401 (2013).



Spin-orbit + Zeeman interactions induce the Majorana edge modes.

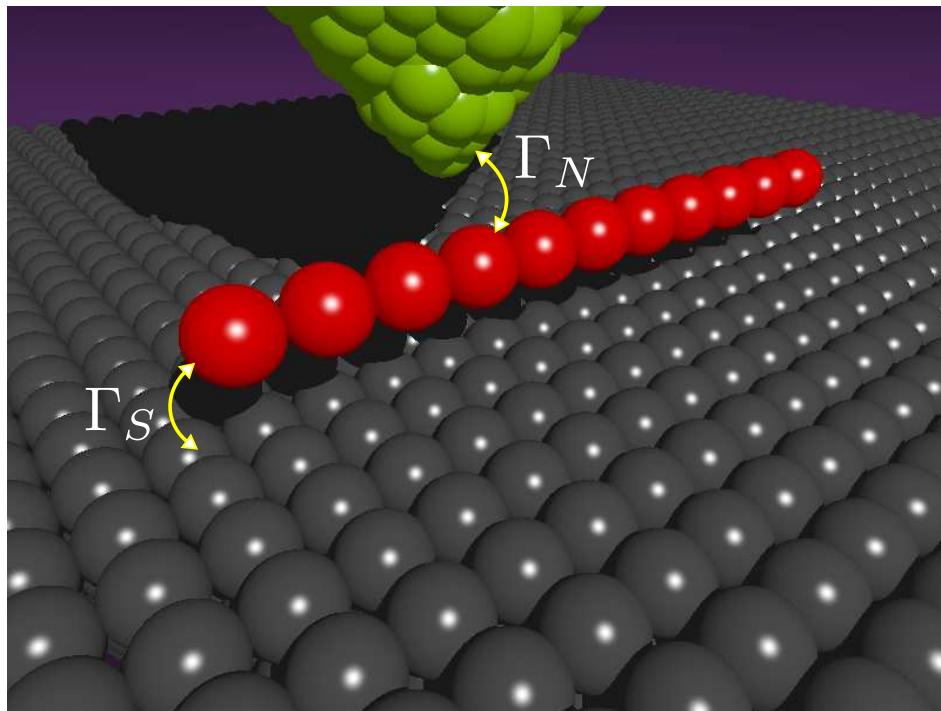
D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B **88**, 165401 (2013).

Towards more realistic situation

/ Rashba chain + pairing /

Towards more realistic situation

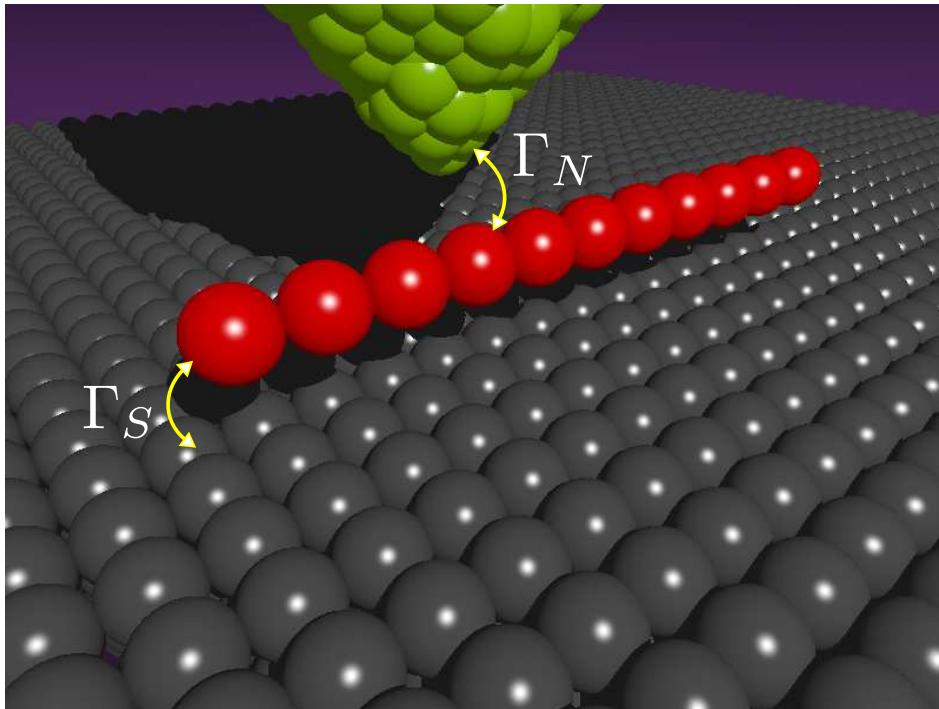
/ Rashba chain + pairing /



Scheme of STM configuration

Towards more realistic situation

/ Rashba chain + pairing /



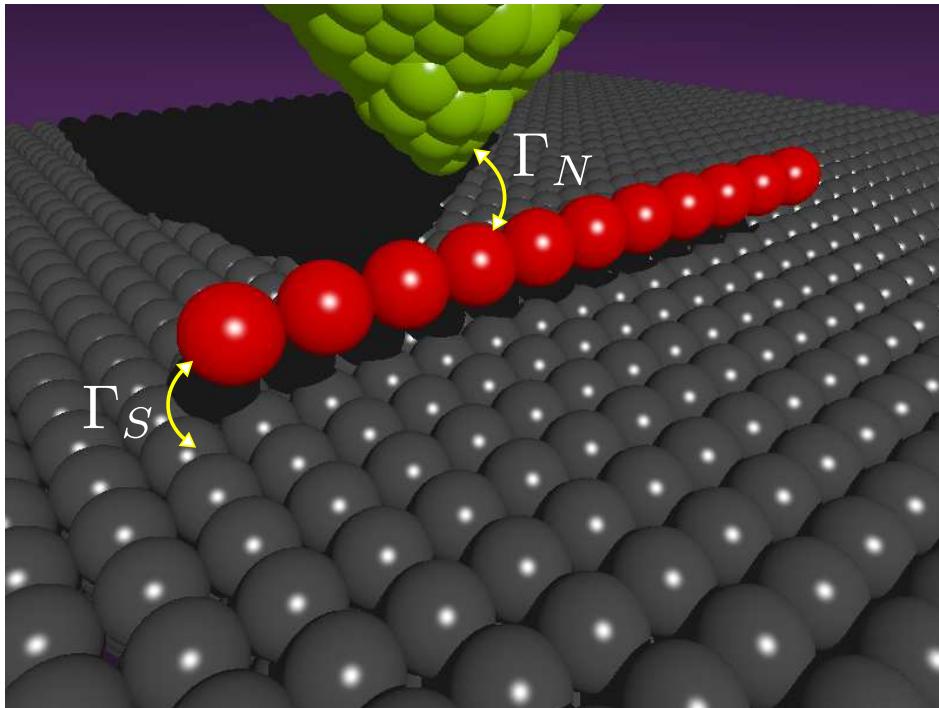
Scheme of STM configuration

$$\hat{H} = \hat{H}_{tip} + \hat{H}_{chain} + \hat{H}_S + \hat{V}_{hybr}$$

We studied this model, focusing on the deep subgap regime $|E| \ll \Delta_{sc}$.

Towards more realistic situation

/ Rashba chain + pairing /



Scheme of STM configuration

where

$$\hat{H}_{chain} = \sum_{i,j,\sigma} (t_{ij} - \delta_{ij}\mu) \hat{d}_{i,\sigma}^\dagger \hat{d}_{j,\sigma} + \hat{H}_{Rashba} + \hat{H}_{Zeeman}$$

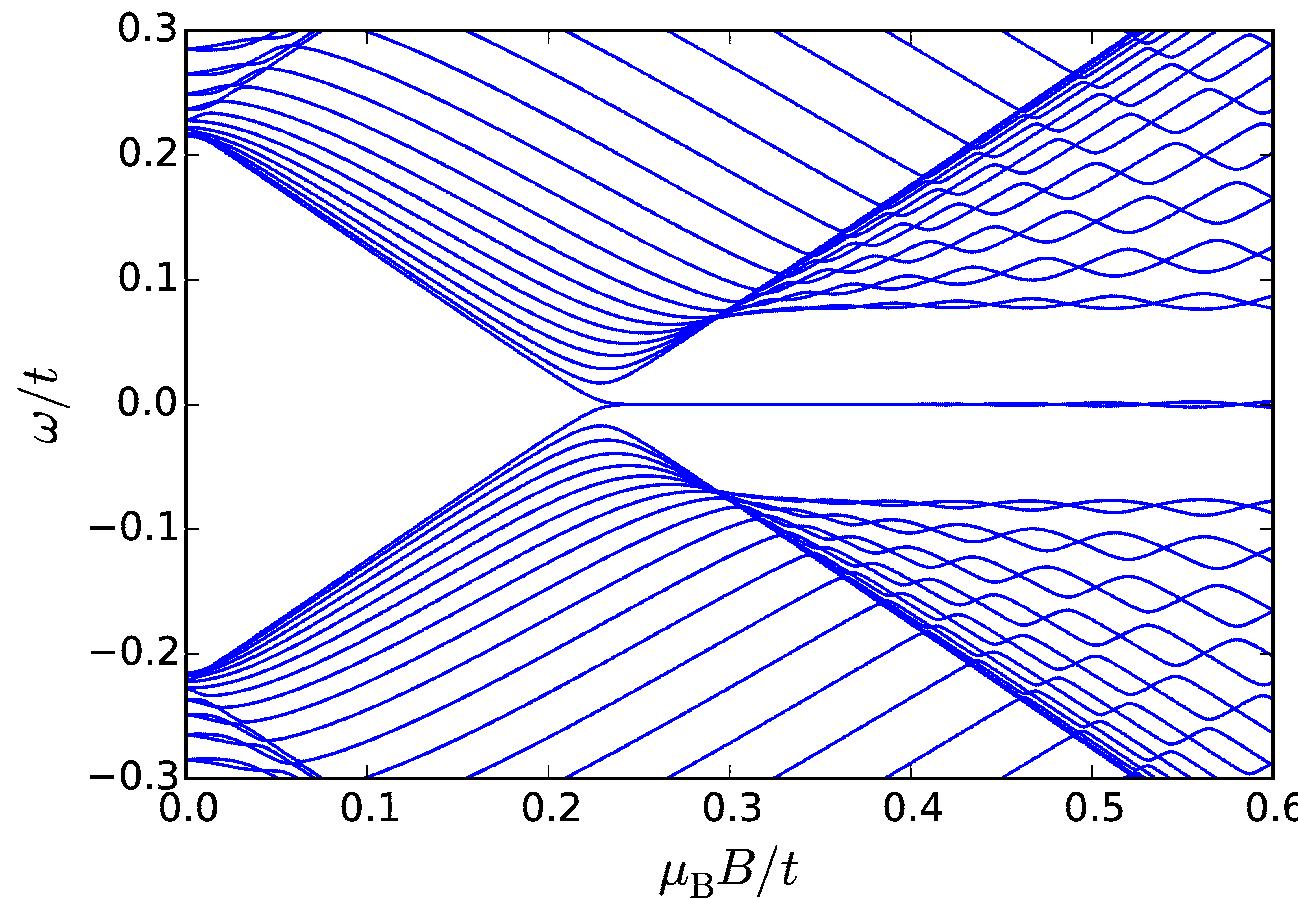
M. Maśka, A. Gorczyca-Goraj, J. Tworzydło and T. Domański, Phys. Rev. B 95, 045429 (2017).

Majorana states – of the Rashba chain

Majorana states

– of the Rashba chain

Mutation of Andreev states into zero-energy (Majorana) mode

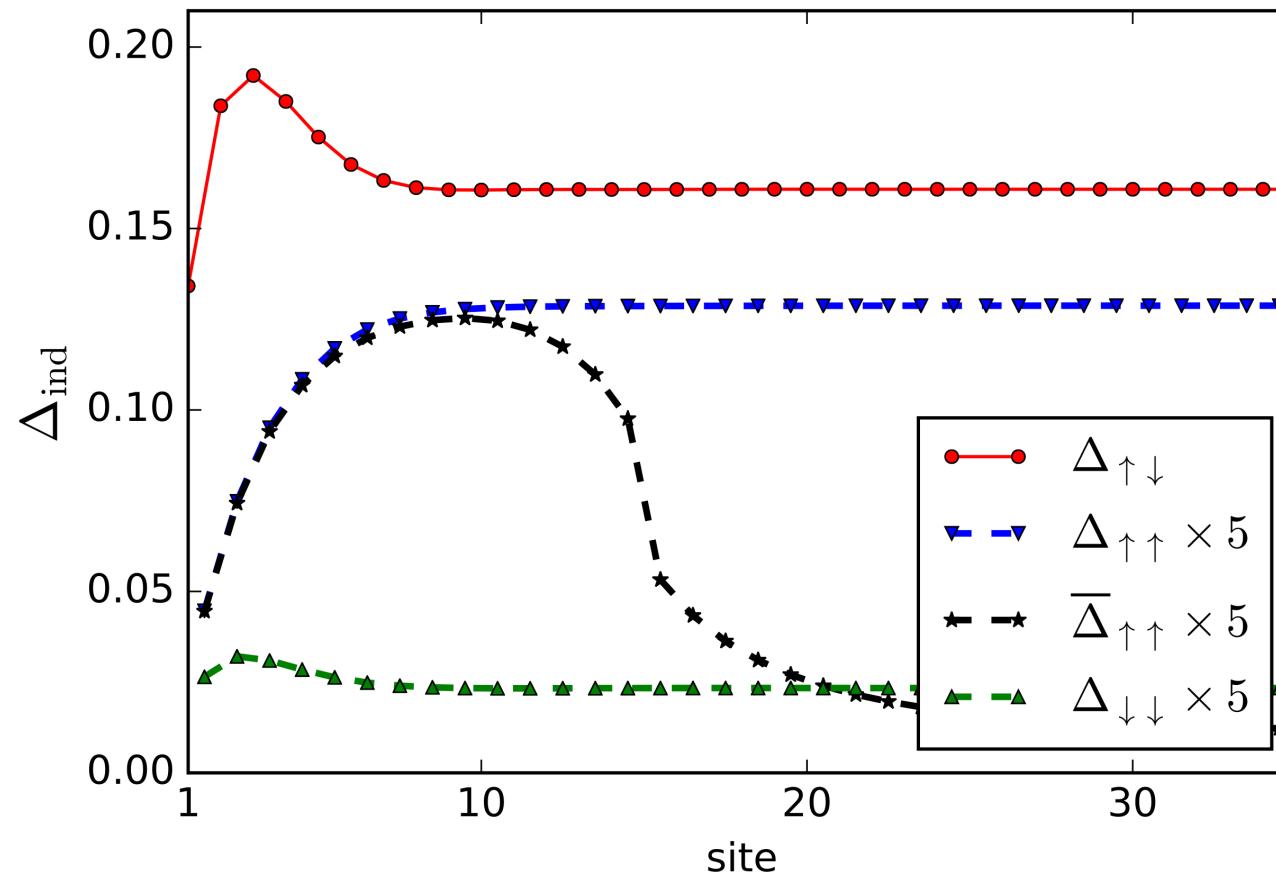


M. Maśka, A. Gorczyca-Goraj, J. Tworzydło and T. Domański, Phys. Rev. B 95, 045429 (2017).

Majorana states

– of the Rashba chain

Spatial variation of the induced pairings $\Delta_{\sigma,\sigma'} = \langle \hat{d}_{i,\sigma} \hat{d}_{i+1,\sigma'} \rangle$

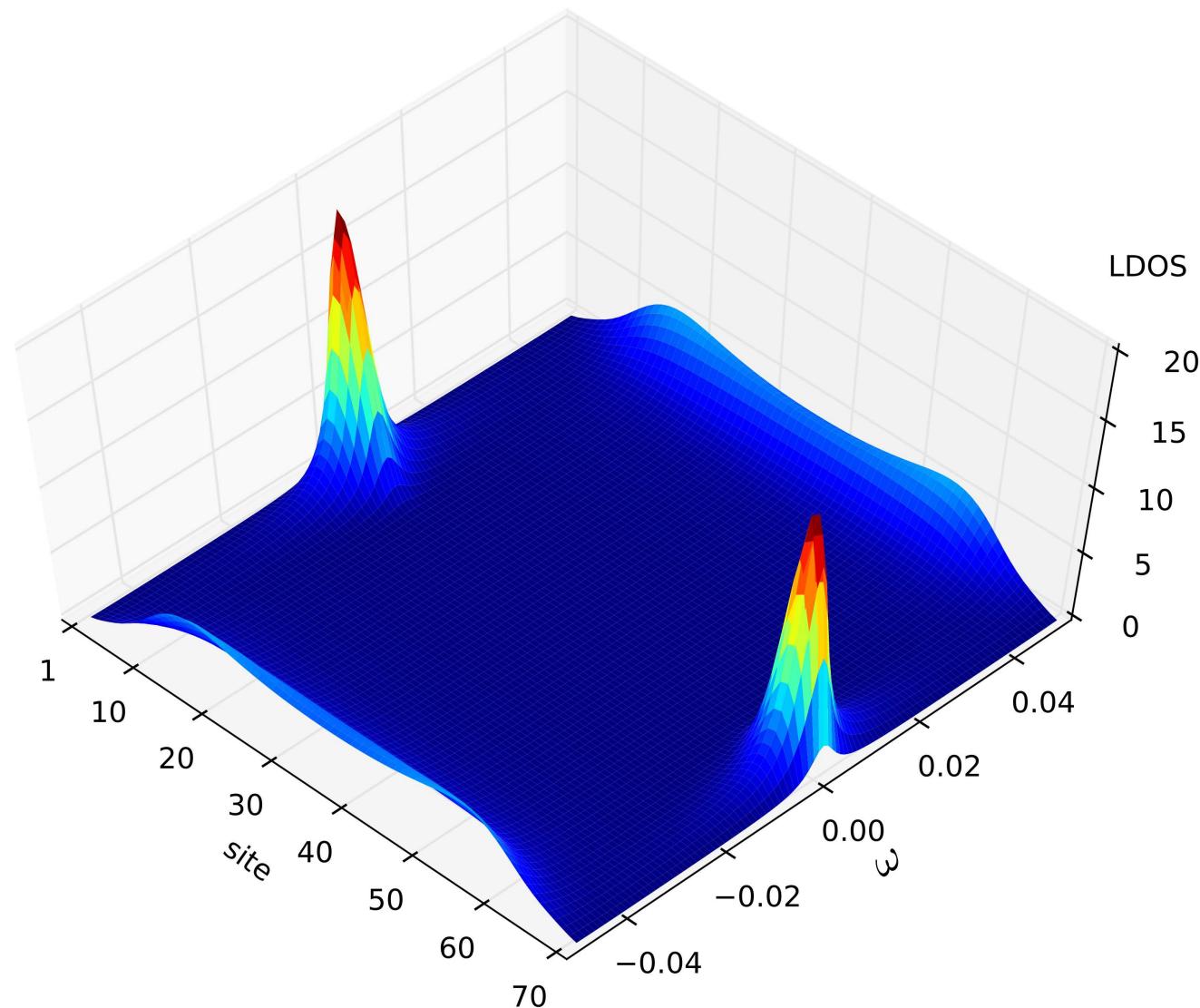


M. Maśka, A. Gorczyca-Goraj, J. Tworzydło and T. Domański, Phys. Rev. B 95, 045429 (2017).

Majorana states

– of the Rashba chain

Spectrum with the edge Majorana quasiparticles

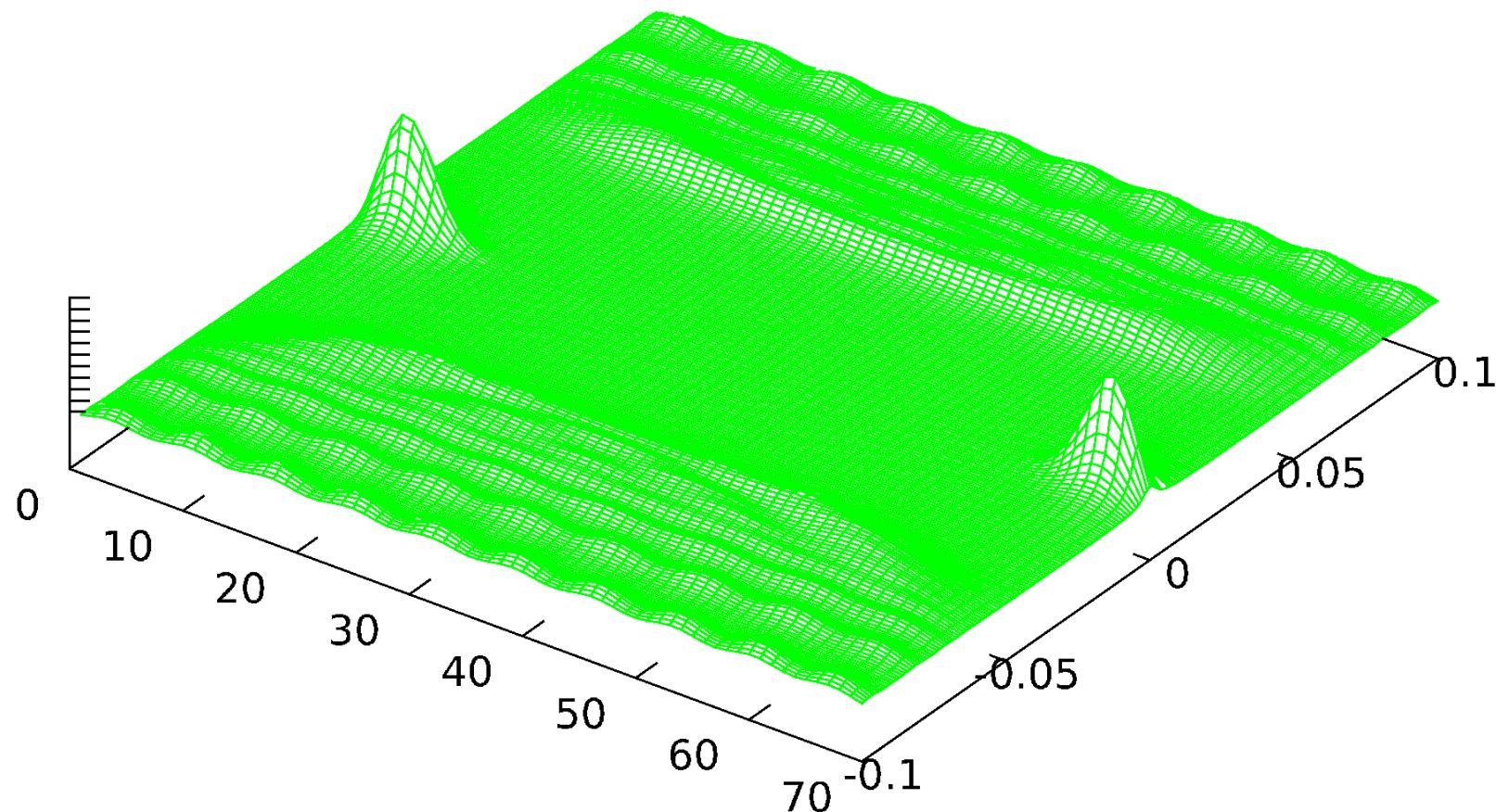


Breaking the Rashba chain

by a reduced hopping at site $i = N/2$

Breaking the Rashba chain

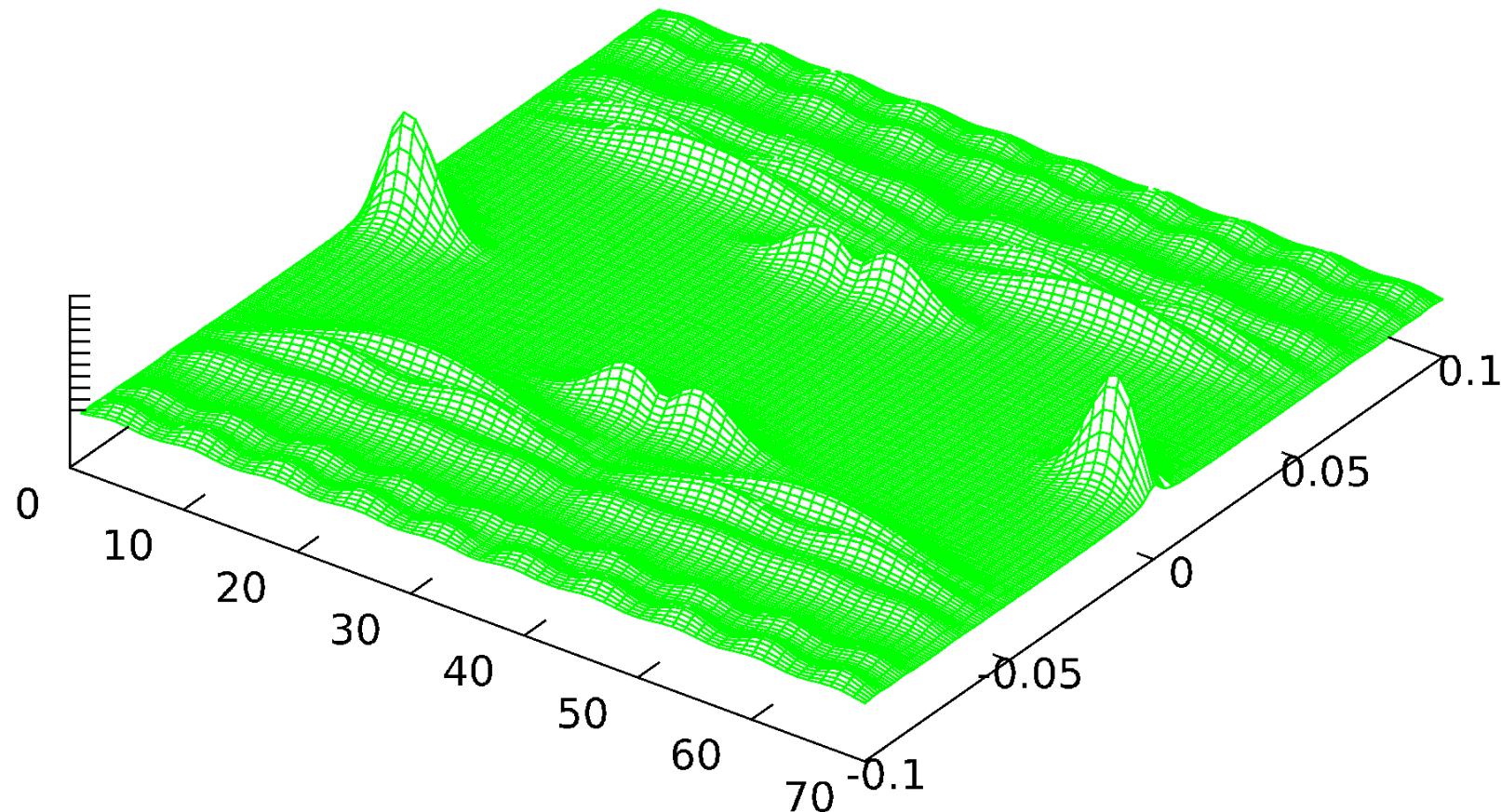
by a reduced hopping at site $i = N/2$



$$t_i/t = 1$$

Breaking the Rashba chain

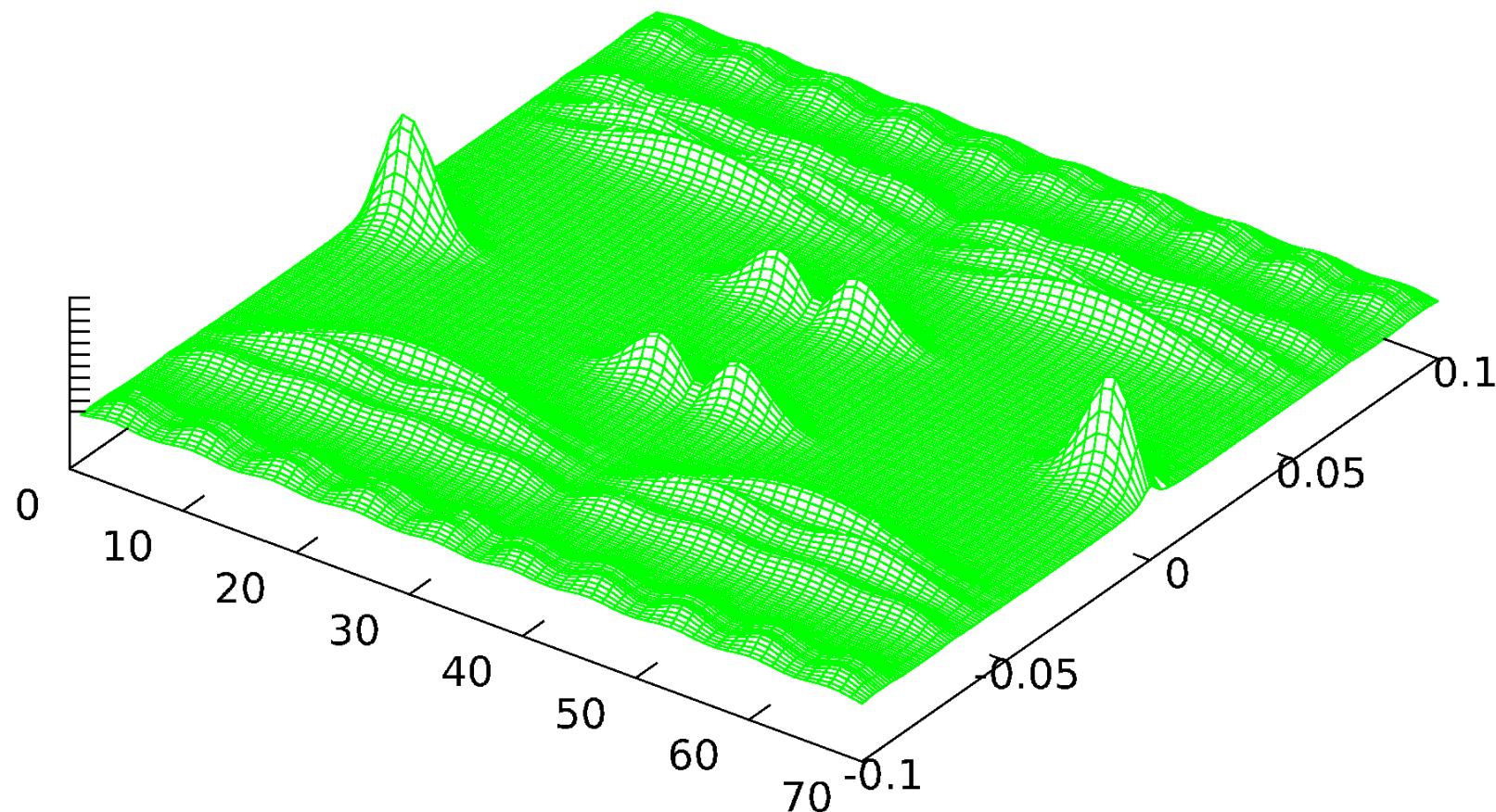
by a reduced hopping at site $i = N/2$



$$t_i/t = 0.8$$

Breaking the Rashba chain

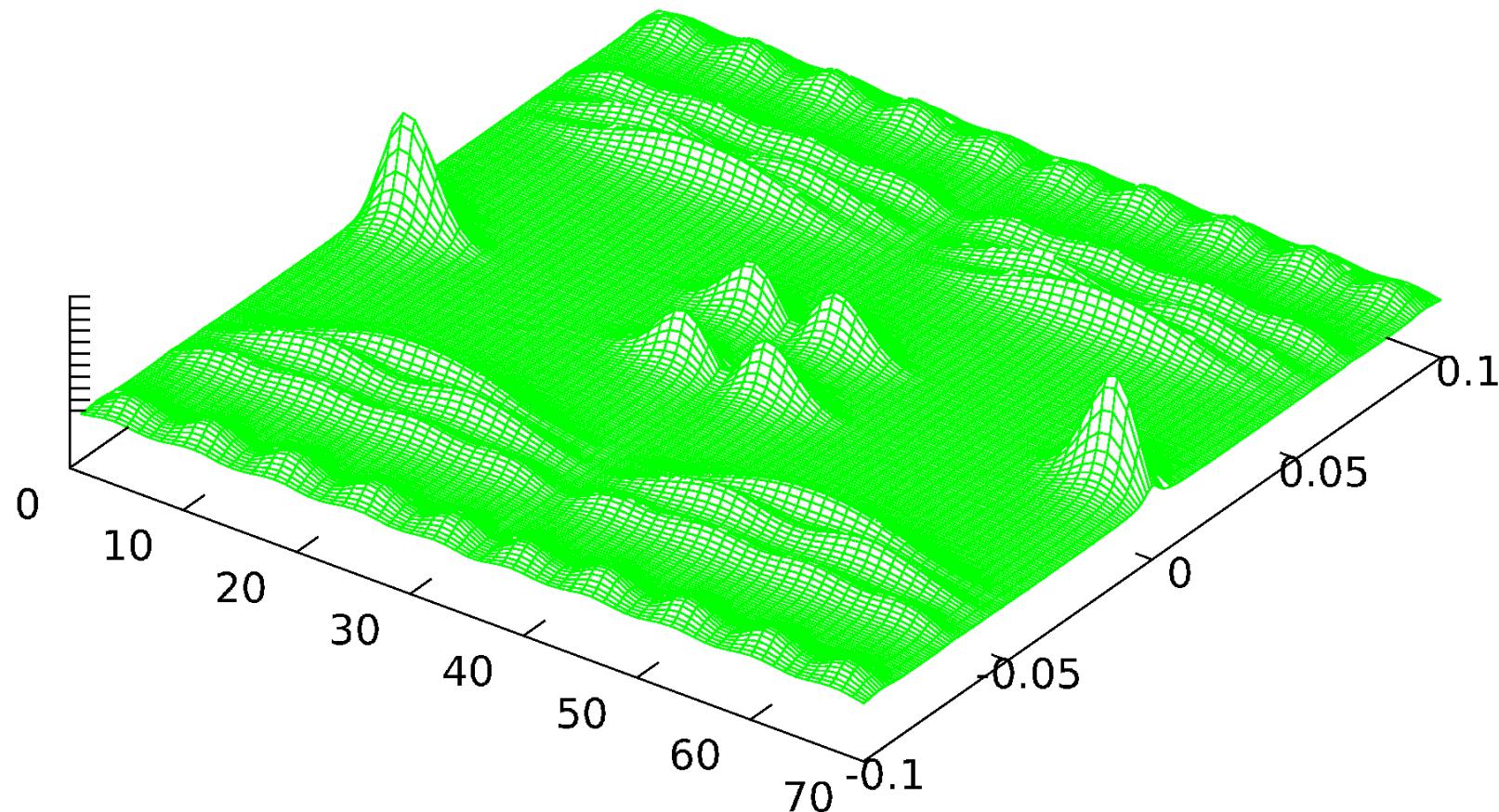
by a reduced hopping at site $i = N/2$



$$t_i/t = 0.6$$

Breaking the Rashba chain

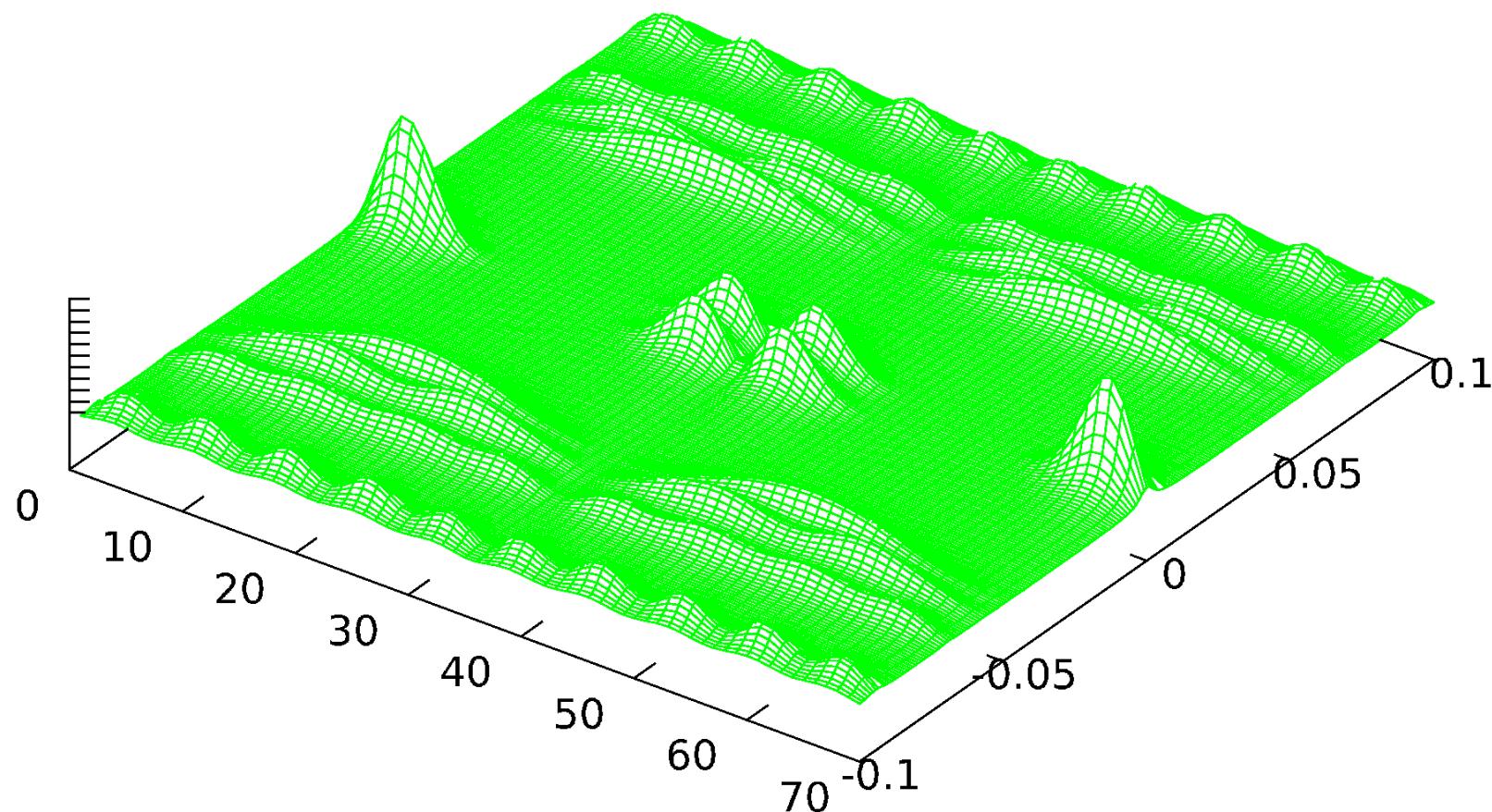
by a reduced hopping at site $i = N/2$



$$t_i/t = 0.4$$

Breaking the Rashba chain

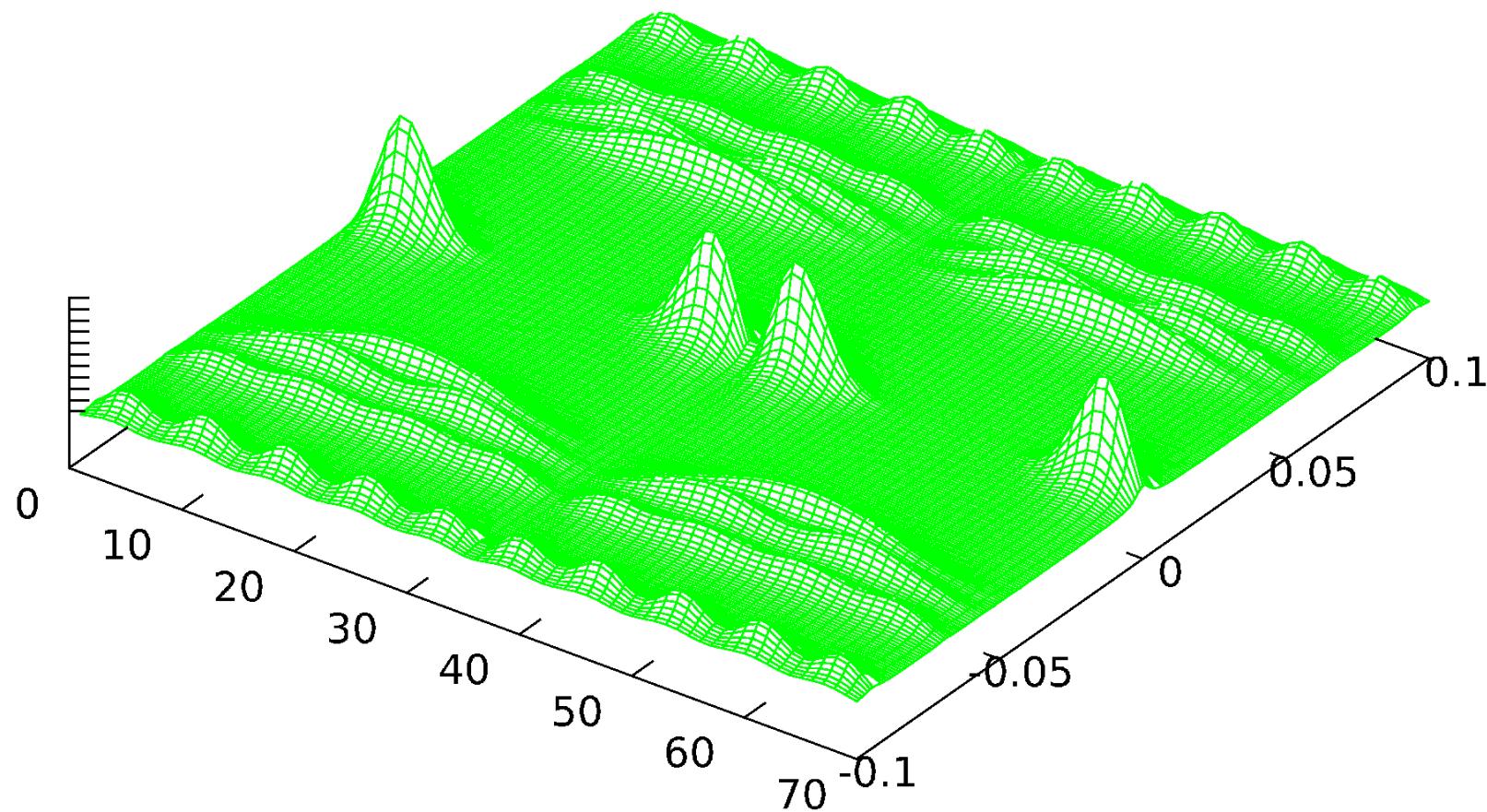
by a reduced hopping at site $i = N/2$



$$t_i/t = 0.2$$

Breaking the Rashba chain

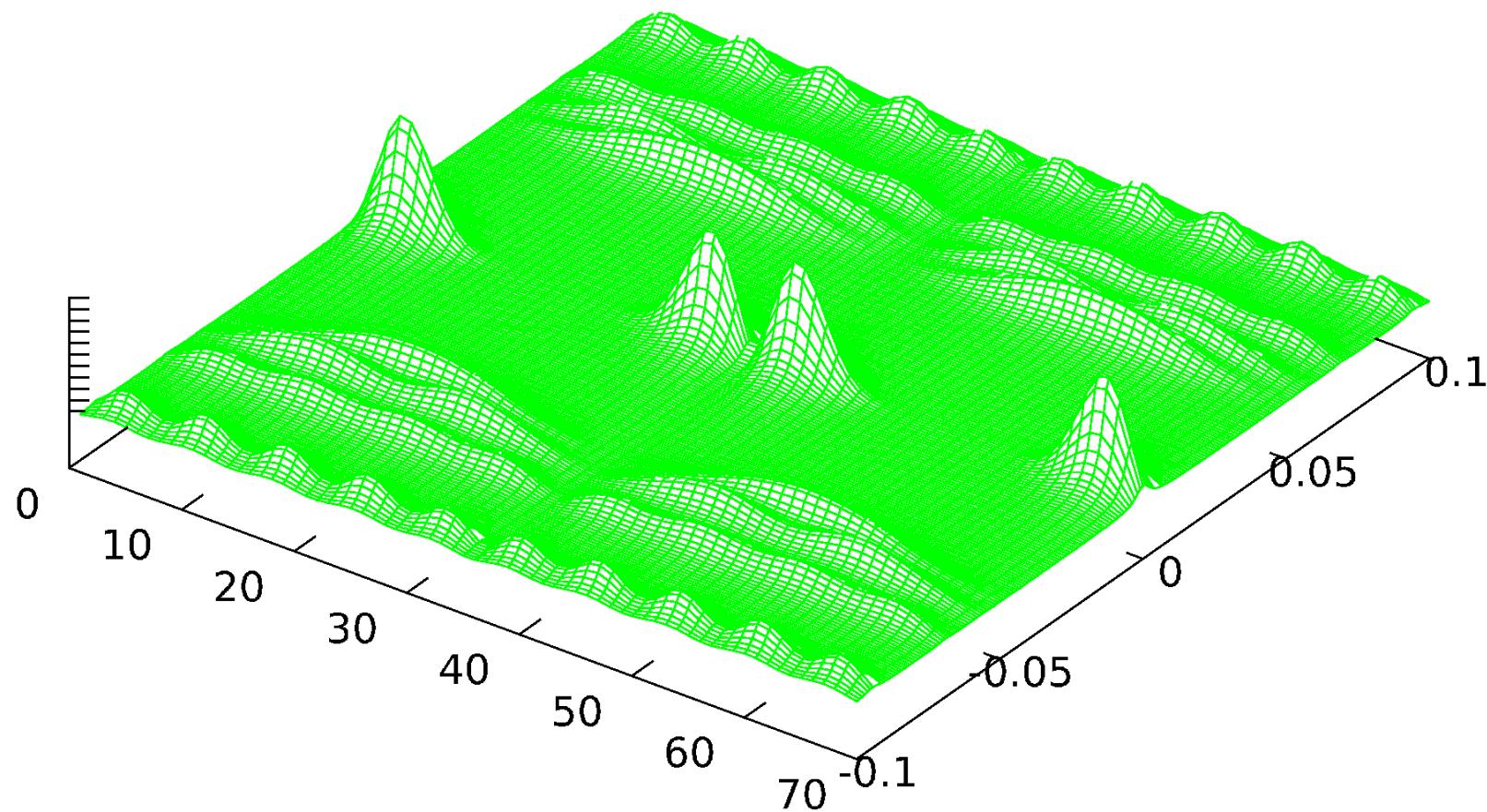
by a reduced hopping at site $i = N/2$



$$t_i/t = 0$$

Breaking the Rashba chain

by a reduced hopping at site $i = N/2$

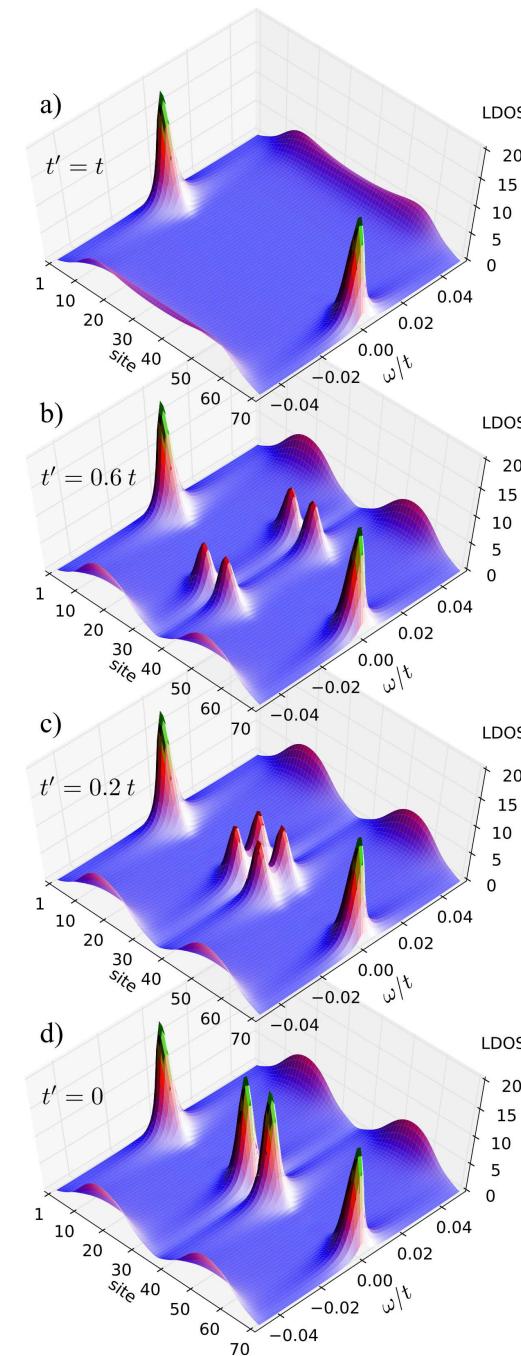


Number of Majorana quasiparticles doubled !

Fussion/splitting of Majorana states

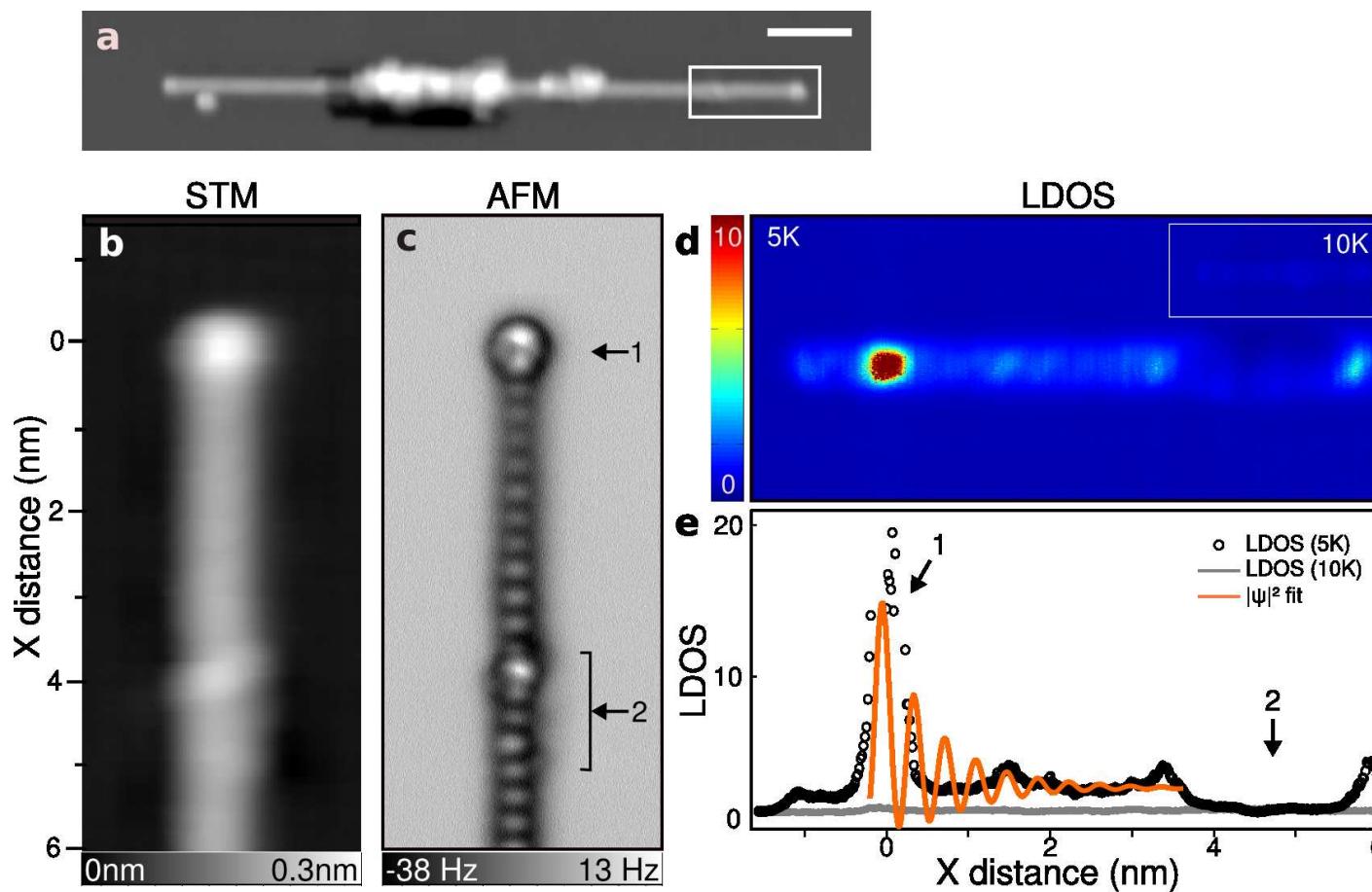
Partitioning
the Rashba chain
into two pieces
via reduced hopping
 t_j at site $j = N/2$

M. Maśka et al, Phys. Rev. B 95, 045429 (2017).



Majoranas at quantum defects

- experimental relevance

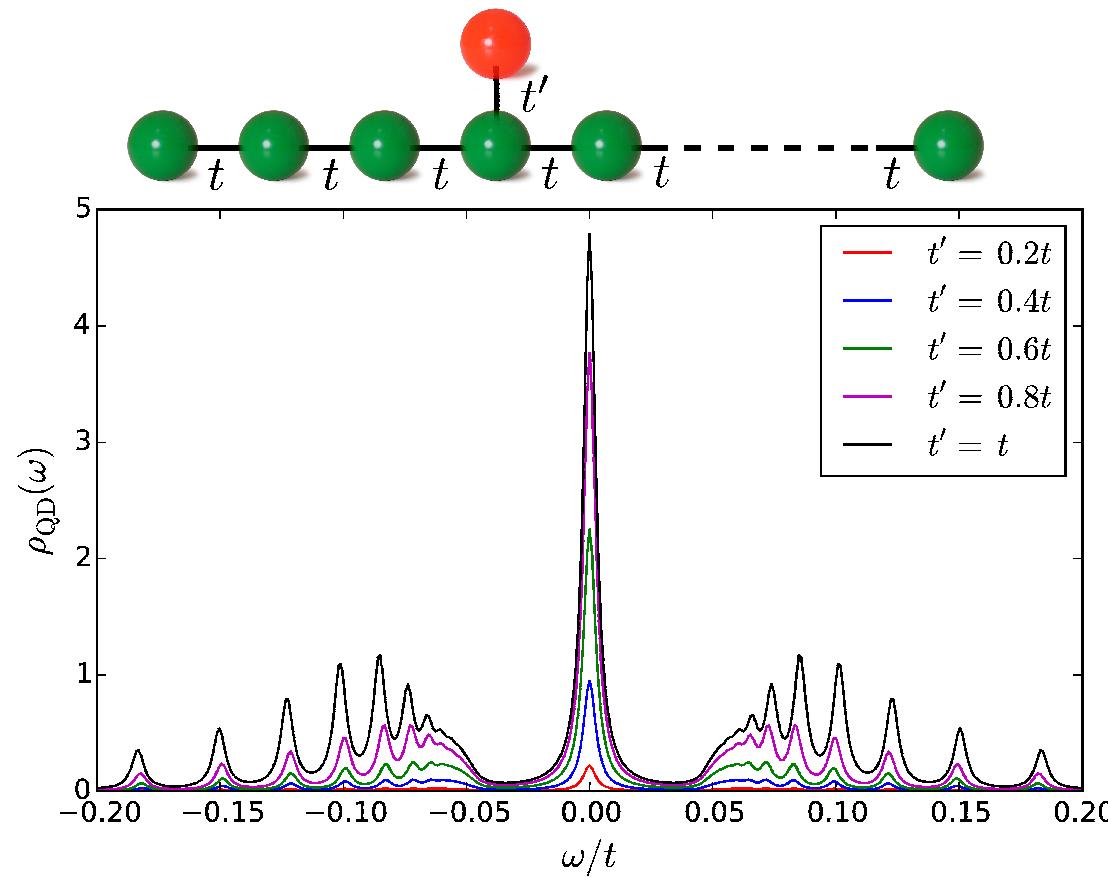


R. Pawlak, M. Kisiel, ..., and E. Meyer, npj Quantum Information **2**, 16035 (2016).

Majorana states

- proximity effect

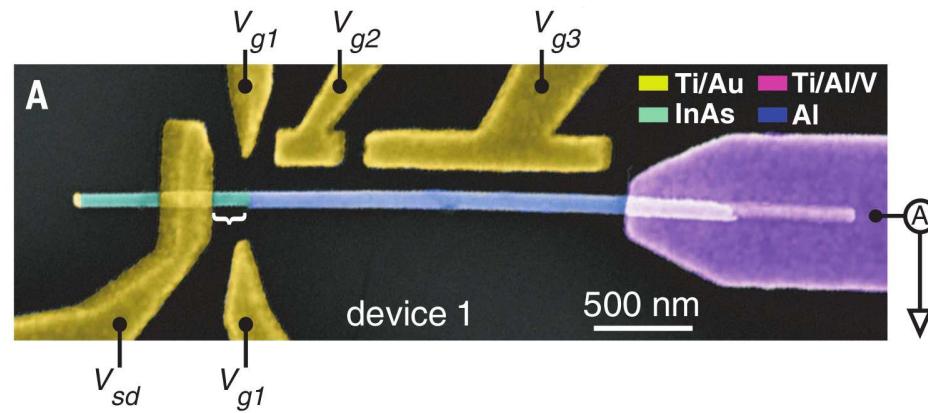
Majorana state can 'leak' into a normal quantum impurity



M. Maśka, A. Gorczyca-Goraj, J. Tworzydło and T. Domański, Phys. Rev. B 95, 045429 (2017).

Proximity effect

– experimental realization (27 Dec 2016)

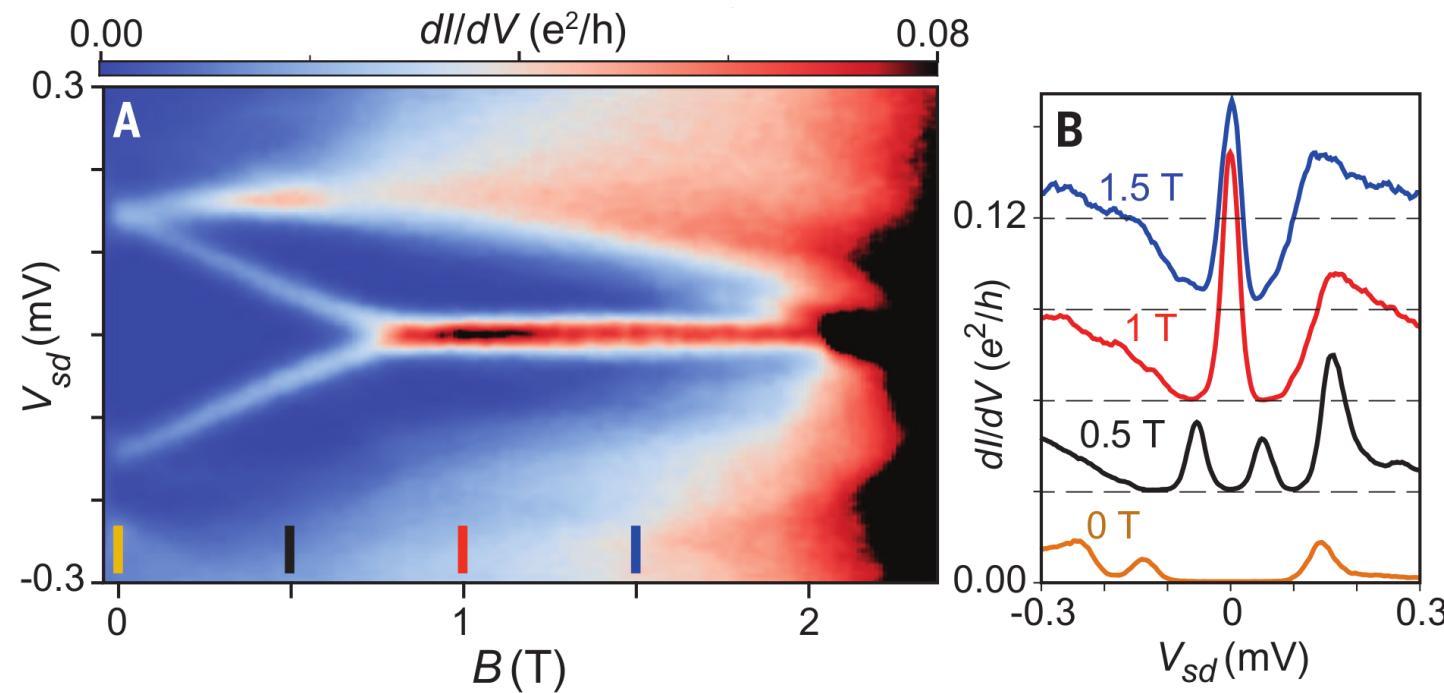


TOPOLOGICAL MATTER

Majorana bound state in a coupled quantum-dot hybrid-nanowire system

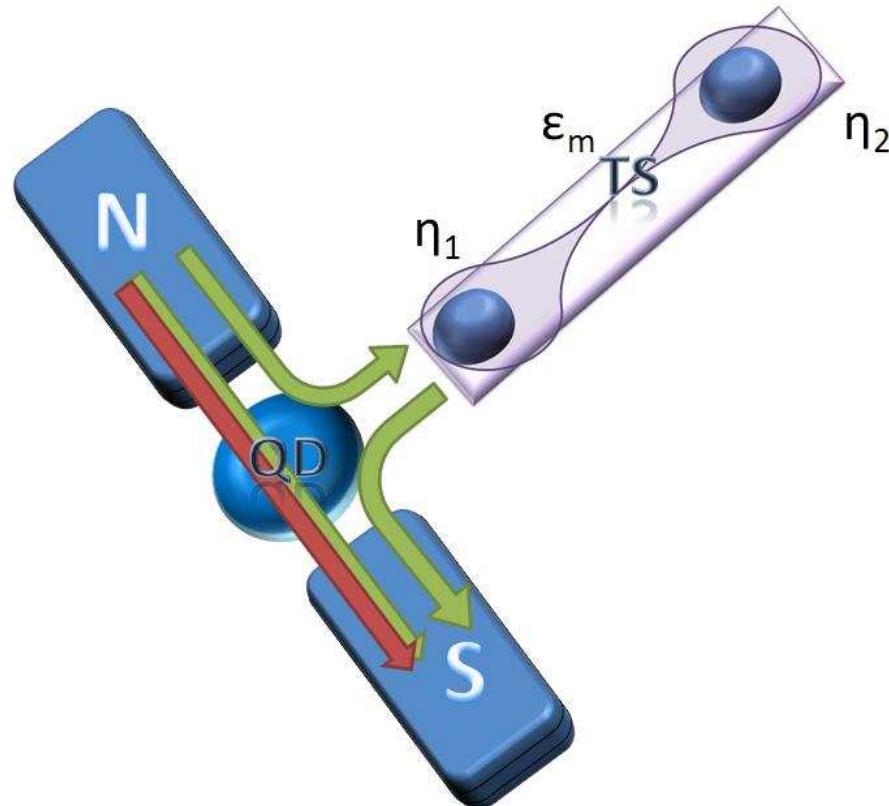
M. T. Deng,^{1,2} S. Vaitiekėnas,^{1,3} E. B. Hansen,¹ J. Danon,^{1,4} M. Leijnse,^{1,5} K. Flensberg,¹ J. Nygård,¹ P. Krogstrup,¹ C. M. Marcus^{1*}

Science 354, 1557 (2016).



Fractionality of Majoranas

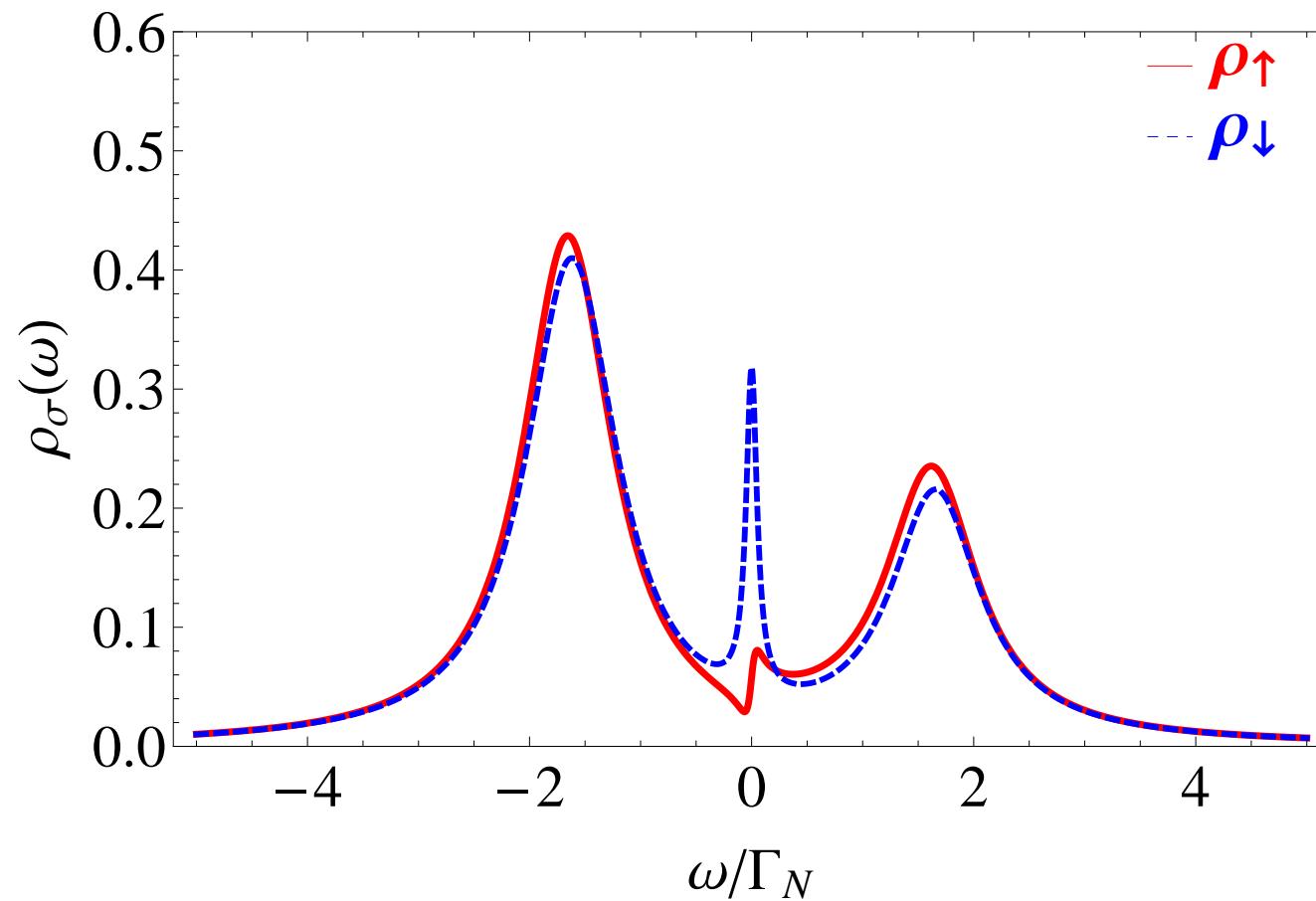
Let's consider QD with the side-attached Majorana quasiparticle.



*J. Barański, A. Kobiałka and T. Domański, J. Phys.: Condens. Matt. **29**, 075603 (2017).*

Fractionality of Majoranas

Fractional (Fano-type) interference in the QD spectrum.



J. Barański, A. Kobiak and T. Domański, J. Phys.: Condens. Matt. **29**, 075603 (2017).

Summary on Majorana qps:

– message to take home

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$$\hat{\gamma}_{i,n}^\dagger = \hat{\gamma}_{i,n}$$

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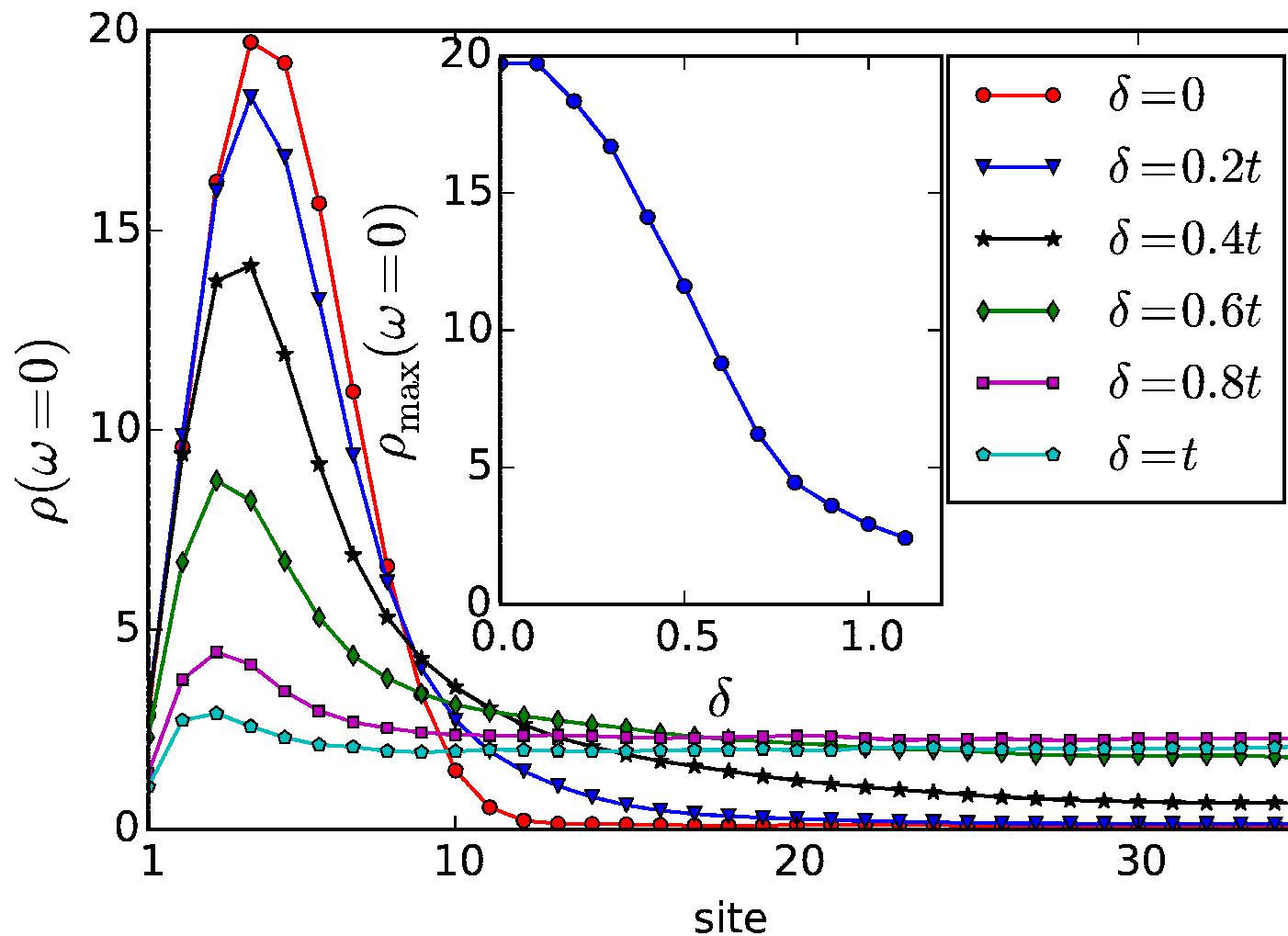
⇒ should be immune to decoherence

... yet be cautious about that !

Majorana states

- effect of disorder

Majorana quasiparticles are not truly immune to disorder !



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