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**Residual diamagnetism driven by
the superconducting fluctuations**

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<http://kft.umcs.lublin.pl/doman/lectures>

Outline

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Preliminaries

/ Cooper pairing & Higgs mechanism /

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Motivation

/ pre-pairing for BE condensation /

Outline



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Technical remarks

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Results

Outline



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Results

\Rightarrow *Bogoliubov quasiparticles above T_c*

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Results

- Bogoliubov quasiparticles above T_c

⇒ Diamagnetism above T_c

Outline



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Technical remarks



Results

- *Bogoliubov quasiparticles above T_c*
- *Diamagnetism above T_c*



Summary

Preliminaries

Superconducting state

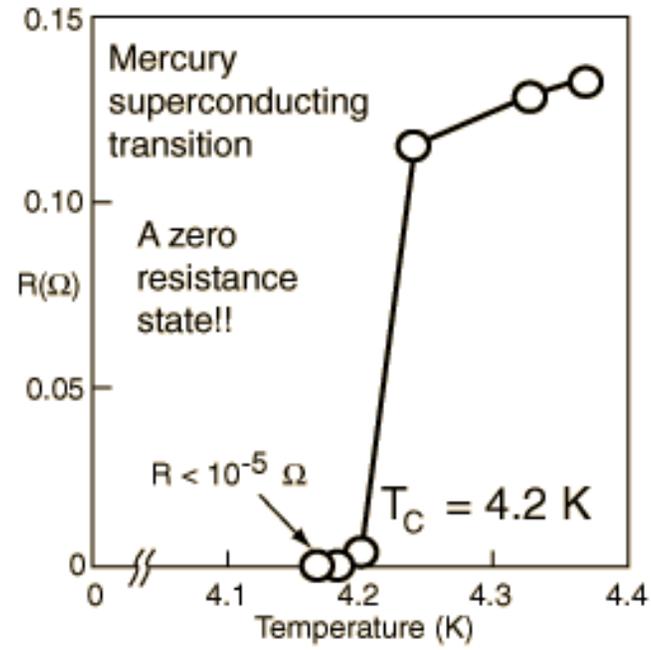
– properties

Superconducting state

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ideal d.c. conductance

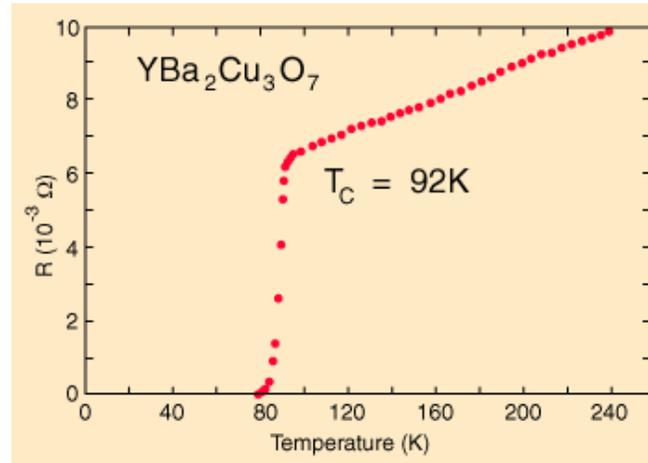


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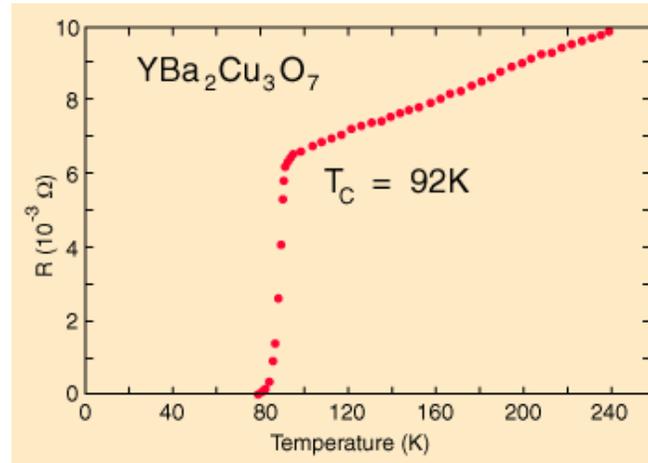


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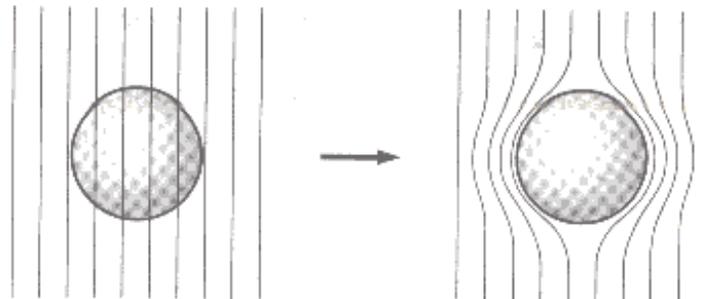


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/perfect screening of the external magnetic field/



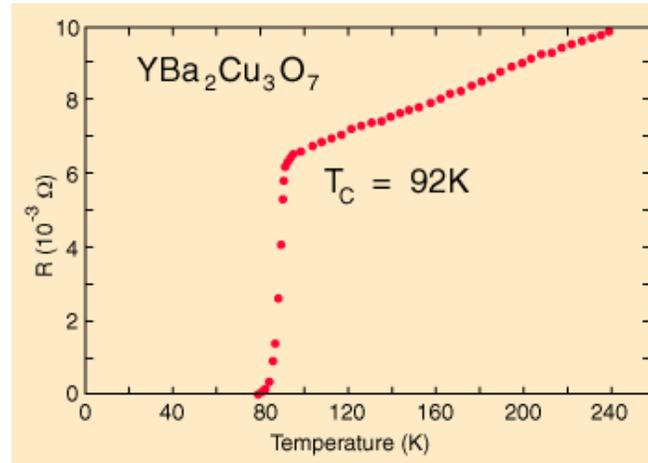
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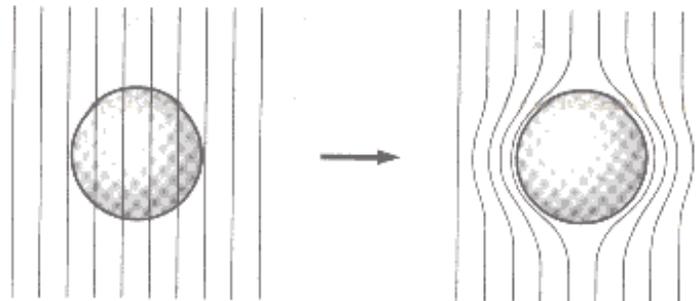


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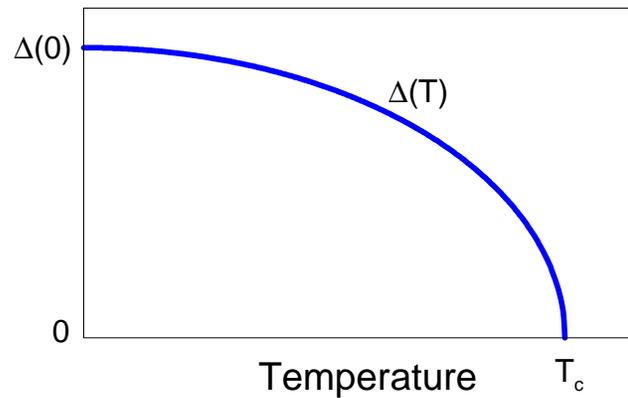
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Appearance of fermion pairs usually goes hand in hand with **superconductivity/superfluidity** but it needn't be the rule.

Conventional superconductors

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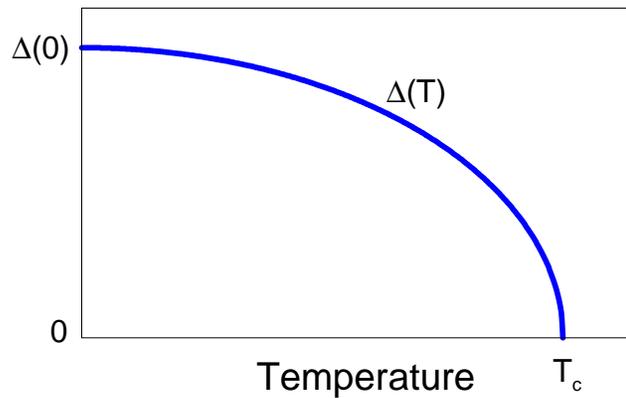
Pair formation coincides with an onset of coherence at T_c



*Pairing causes
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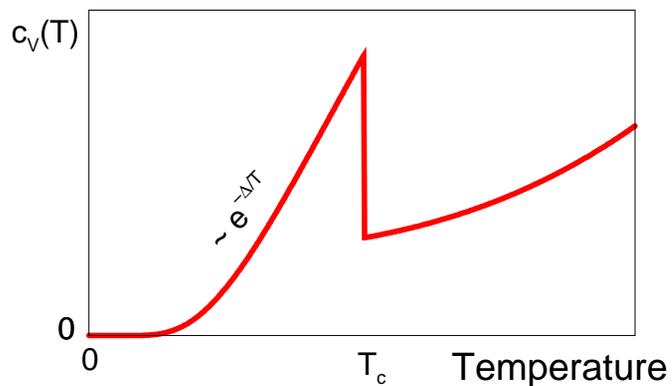
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Pairing causes the energy gap $\Delta(T)$ in a single particle spectrum

2-nd order phase transition



Below T_c there appears the order parameter (ODLRO)
 $\chi \propto \Delta(T)$

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Bose-Einstein condensate of these Cooper pairs is described by a common wave-function

$$\chi(\vec{r}, t)$$

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The order parameter

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$|\chi| \neq 0$ \longrightarrow amplitude causes the energy gap

$\nabla\theta \neq 0$ \longrightarrow phase slippage induces supercurrents

Anderson-Higgs mechanism – **outline**

The low energy excitations of the BE condensed pairs are characterized by the collective (Goldstone) phasal mode.

Anderson-Higgs mechanism

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Can any part of this mechanism survive above T_c ?

Motivation

Phase transitions – classification

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1. **closing the gap** (*BCS superconductors*)

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2. **disordering the phase** (*HTSC compounds*)

$$\lim_{T \rightarrow T_c} \langle \theta \rangle = 0$$

Historical remark

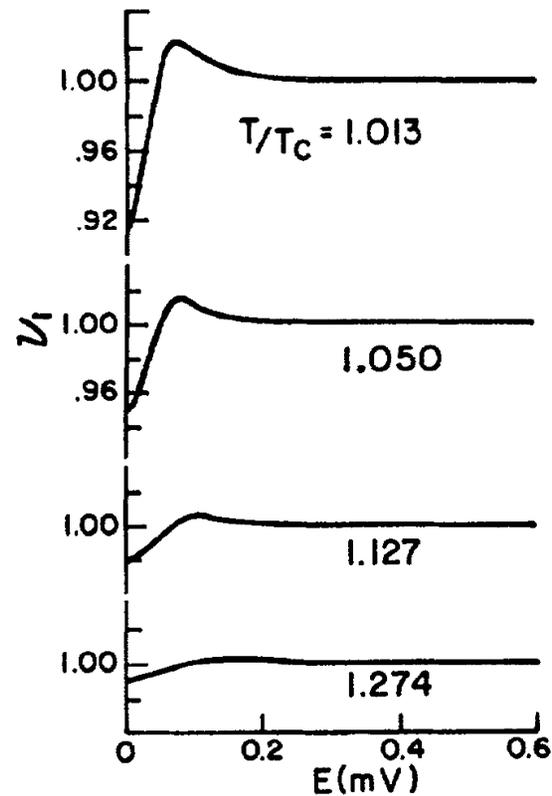
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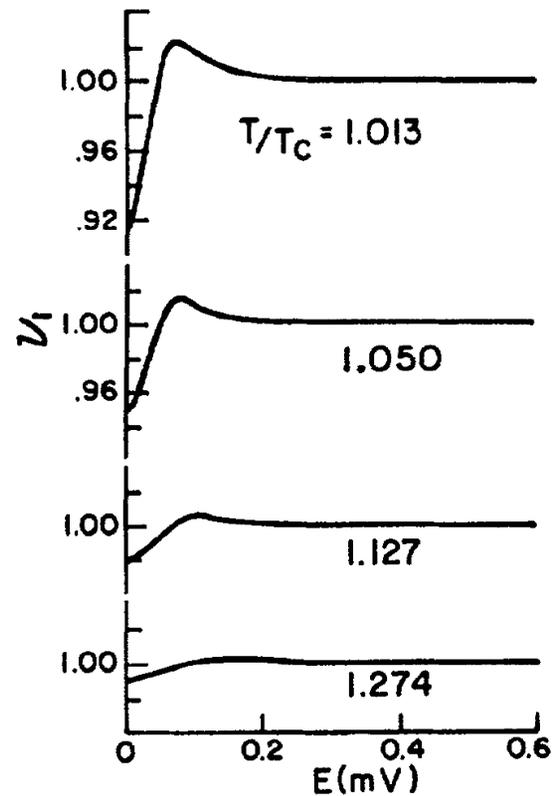
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R.W. Cohen and B. Abels, *Phys. Rev.* **168**, 444 (1968).

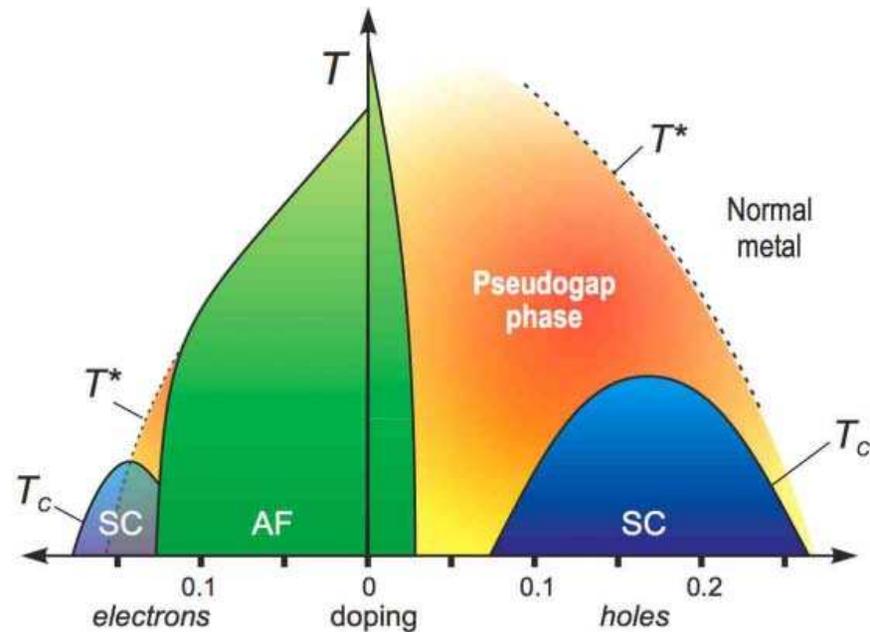
HTSC materials – phase diagram

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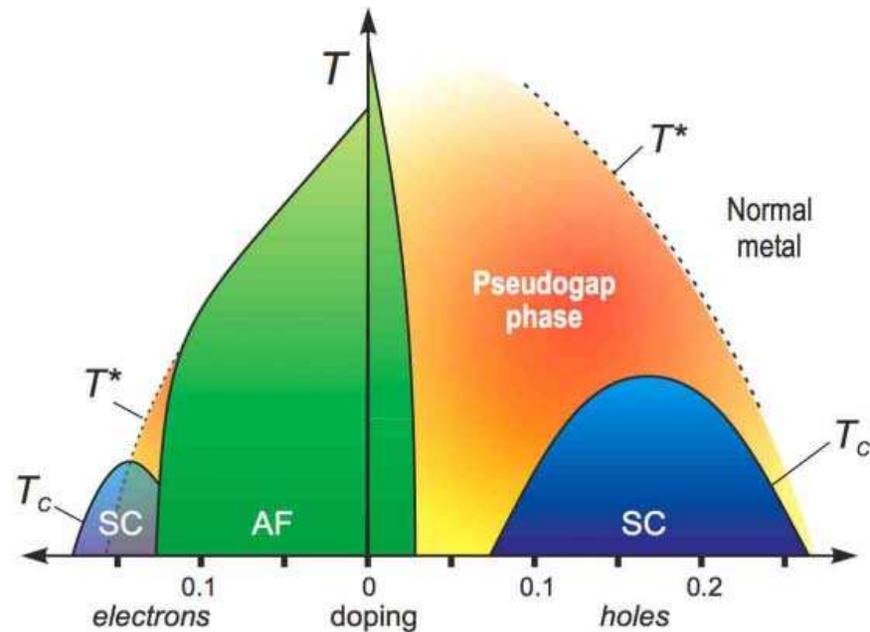


O. Fisher et al, Rev. Mod. Phys. 79, 353 (2007).

HTSC materials – phase diagram

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Unresolved problem:

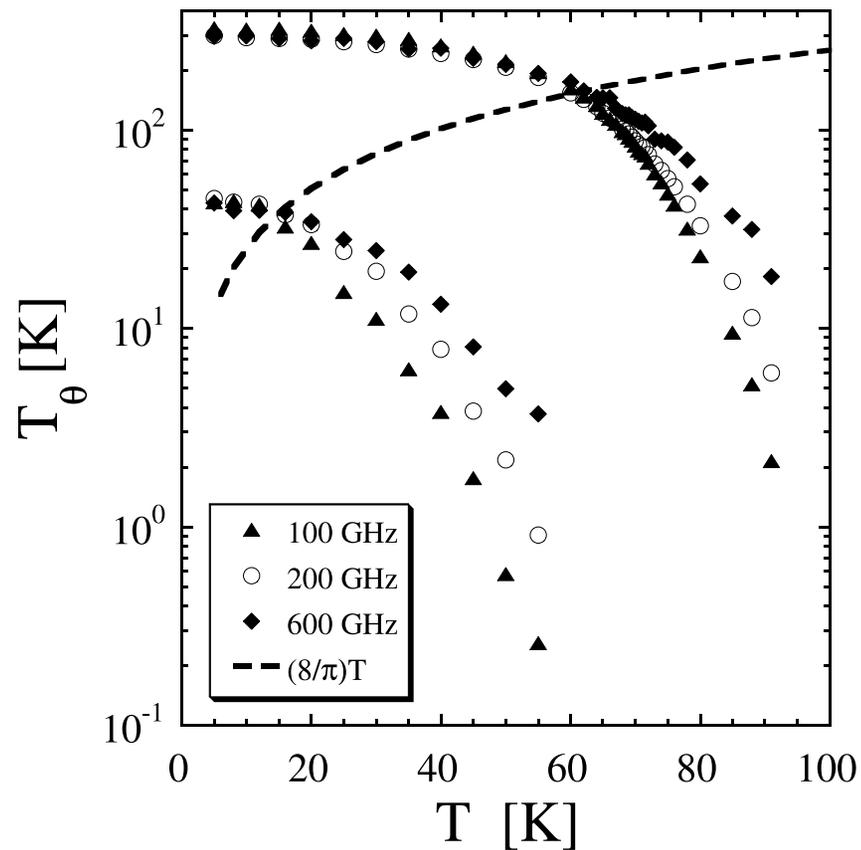
What causes the pseudogap ?

Incoherent pairs above T_c

experimental fact # 1

Incoherent pairs above T_c

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Dynamic phase-stiffness $T_\theta = \omega \operatorname{Im}\sigma(\omega, T) / \sigma_Q$
observed at the ultrafast (teraHz) external ac fields.

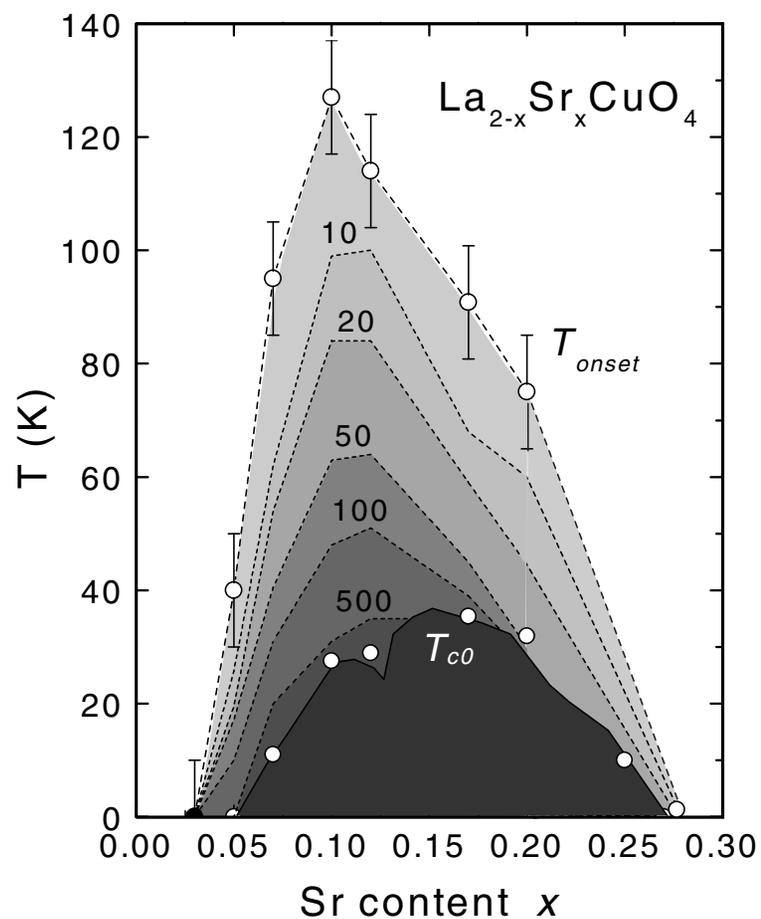
J. Corson et al, Nature 398, 221 (1999).

Incoherent pairs above T_c

experimental fact # 2

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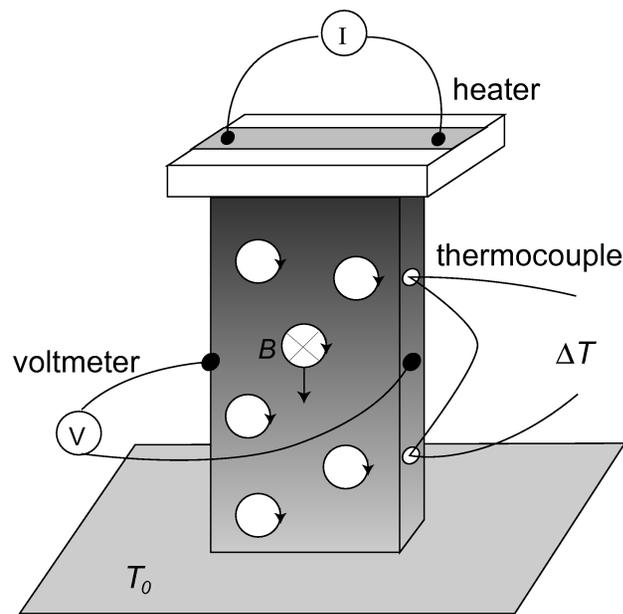
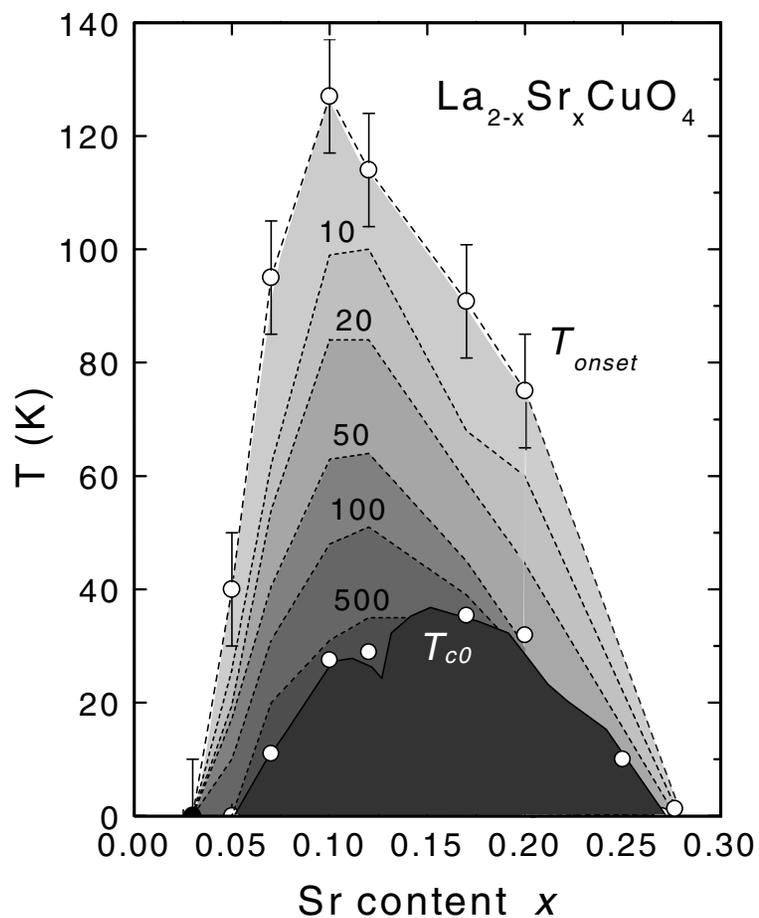


Phase slippage detected in the large Nernst effect.

Y. Wang et al, *Science* **299**, 86 (2003).

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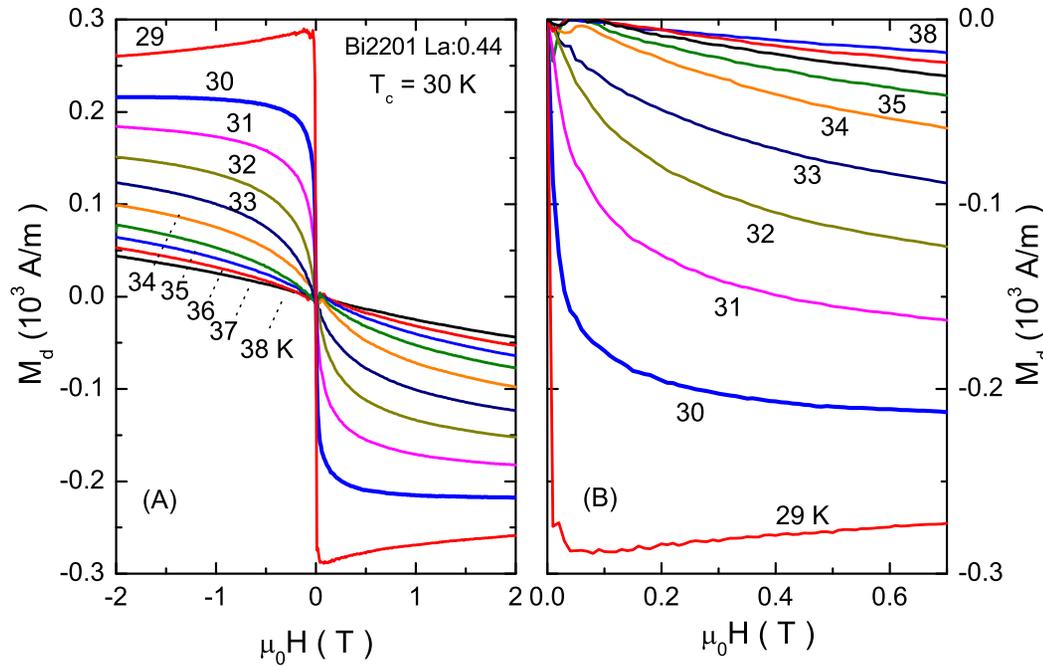
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Incoherent pairs above T_c

experimental fact # 3

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*The Van Vleck
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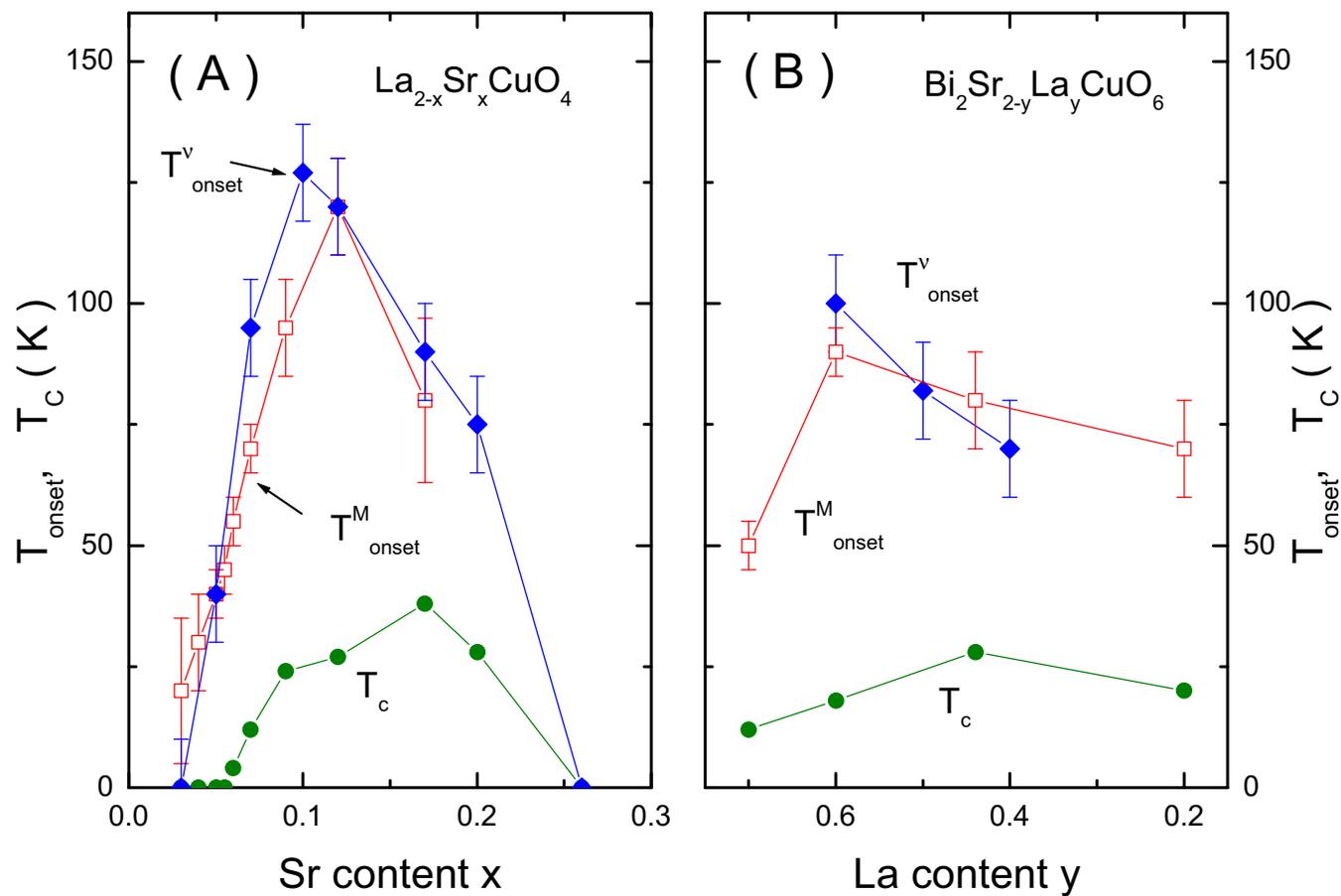
$T_c = 30K$

**Enhanced diamagnetic response revealed above T_c
by the ultrahigh precision torque magnetometry.**

*L. Li et al and N.P. Ong, Phys. Rev. B **81**, 054510 (2010).*

Incoherent pairs above T_c

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T^{ν} – onset of the Nernst effect
 T^M – onset of the diamagnetism

L. Li et al and N.P. Ong, Phys. Rev. B **81**, 054510 (2010).

Incoherent pairs above T_c

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Josephson-like features seen above T_c in the tunneling

N. Bergeal et al, Nature Phys. 4, 608 (2008).

Incoherent pairs above T_c

... continued

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⇒ **spectroscopic fingerprints of the Bogoliubov QPs seen by the unique octet patterns which survive up to $1.5T_c$**

J. Lee, ... and J.C. Davis, Science 325, 1099 (2009).

Technical remarks

Boson-Fermion model

[in a lattice representation]

$$\begin{aligned}\hat{H} &= \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ &+ \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right]\end{aligned}$$

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In the Lagrangian language we obtain this kind of physics upon applying the Hubbard-Stratonovich transformation !

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} &= \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^\dagger \hat{b}_{\mathbf{q}} \\ &+ \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^\dagger \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

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This BF scenario has been considered by various groups:

J. Ranninger *with coworkers* Grenoble

R. Micnas, S. Robaszkiewicz Poznań

T.D. Lee *with coworkers* New York

V.B. Geshkenbein, L.B. Ioffe, A.I. Larkin

E. Altman & A. Auerbach Technion

A. Griffin *with coworkers* Toronto

K. Levin *with coworkers* Chicago

and many others.

Outline of the procedure

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For studying the many-body effects we construct the sequence

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T. Domański and J. Ranninger, Phys. Rev. B 63, 134505 (2001).

Continuous unitary transformation

– algorithm

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where

$$\hat{\eta}(l) = \frac{d\hat{S}(l)}{dl} \hat{S}^\dagger(l) = -\hat{\eta}^\dagger(l).$$

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$$\hat{H} = \hat{H}_{diag} + \hat{H}_{off}$$

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For more details see for instance:

S. Kehrein, Springer Tracts in Modern Physics **217**, (2006);

F. Wegner, J. Phys. A: Math. Gen. **39**, 8221 (2006).

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Similar ideas have been also earlier developed also in the field of **control theory** under the names:

★ "double bracket flow"

R.W. Brockett, Lin. Alg. and its Appl. 146, 79 (1991).

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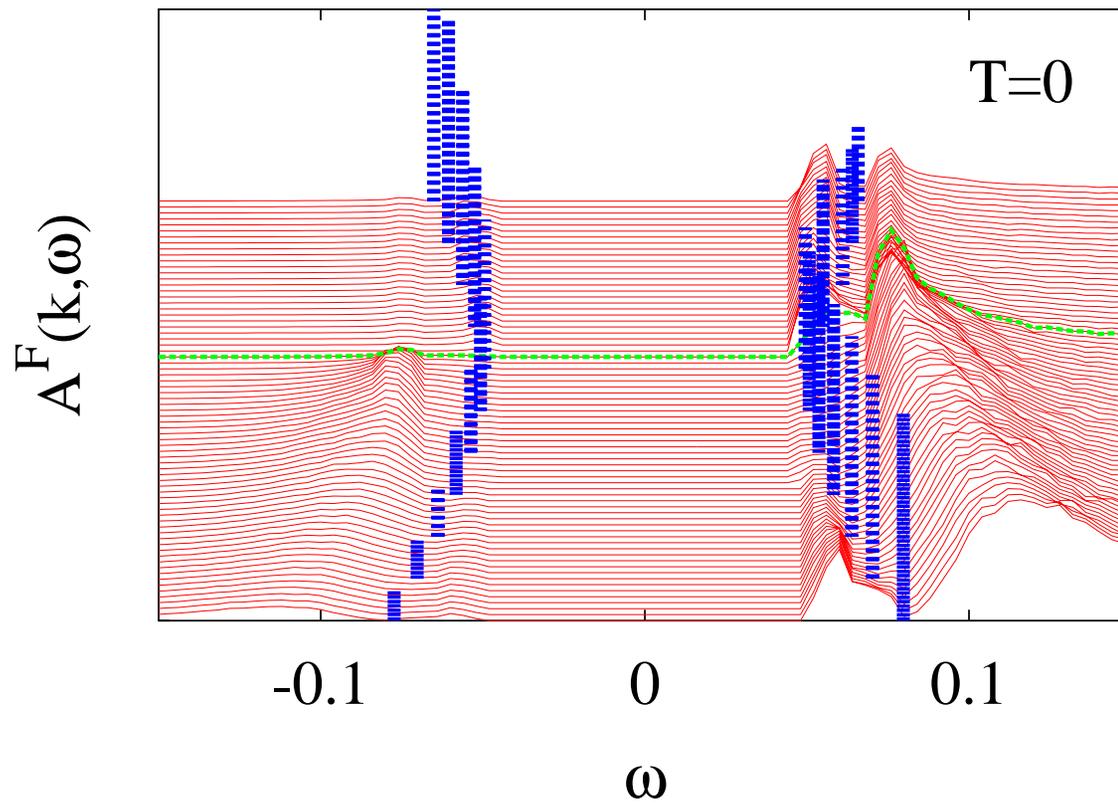
S.R. White, J. Chem. Phys. 117, 7472 (2002).

Results :

1. Bogoliubov quasiparticles above T_c

Effective spectrum: BF model

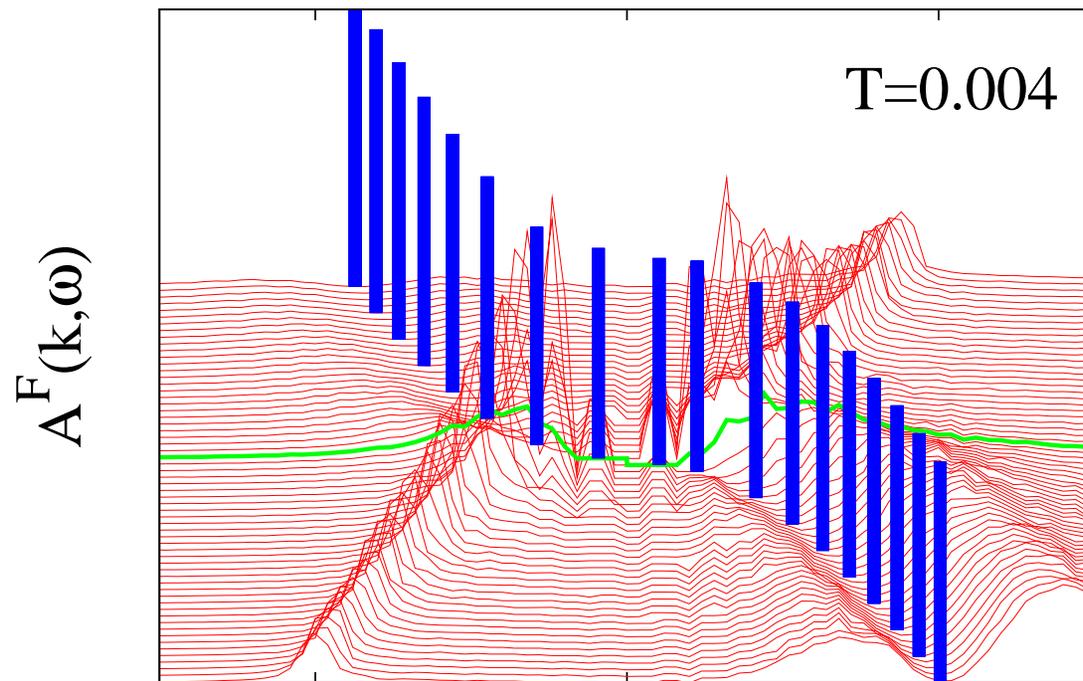
$$T < T_c$$



*T. Domański and J. Ranninger, Phys. Rev. Lett. **91**, 255301 (2003).*

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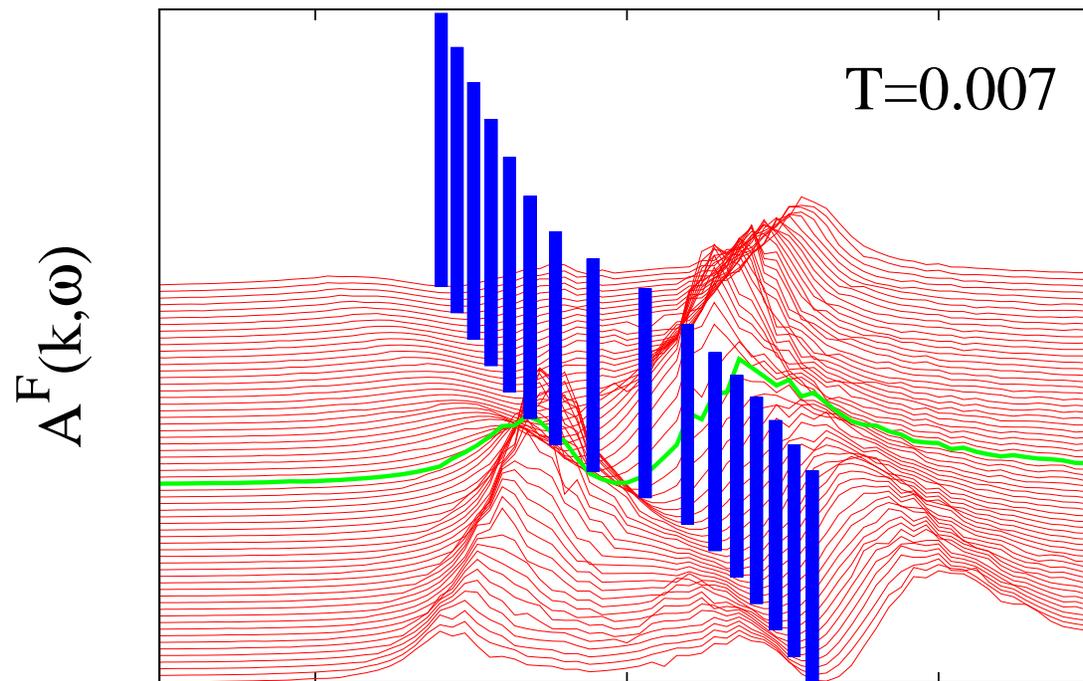
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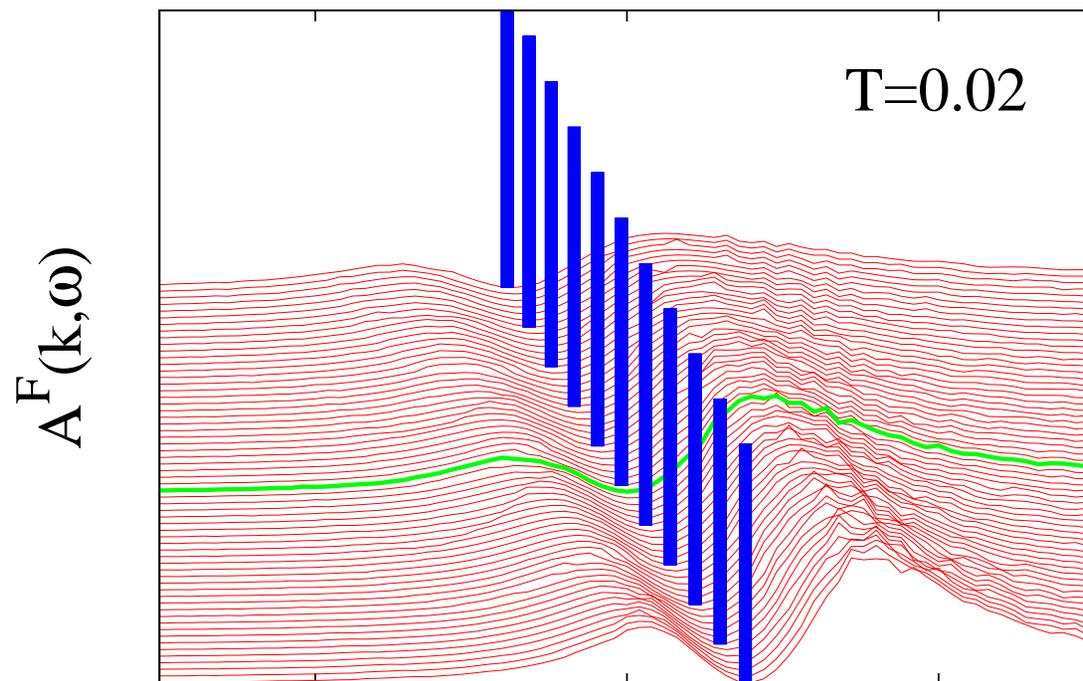
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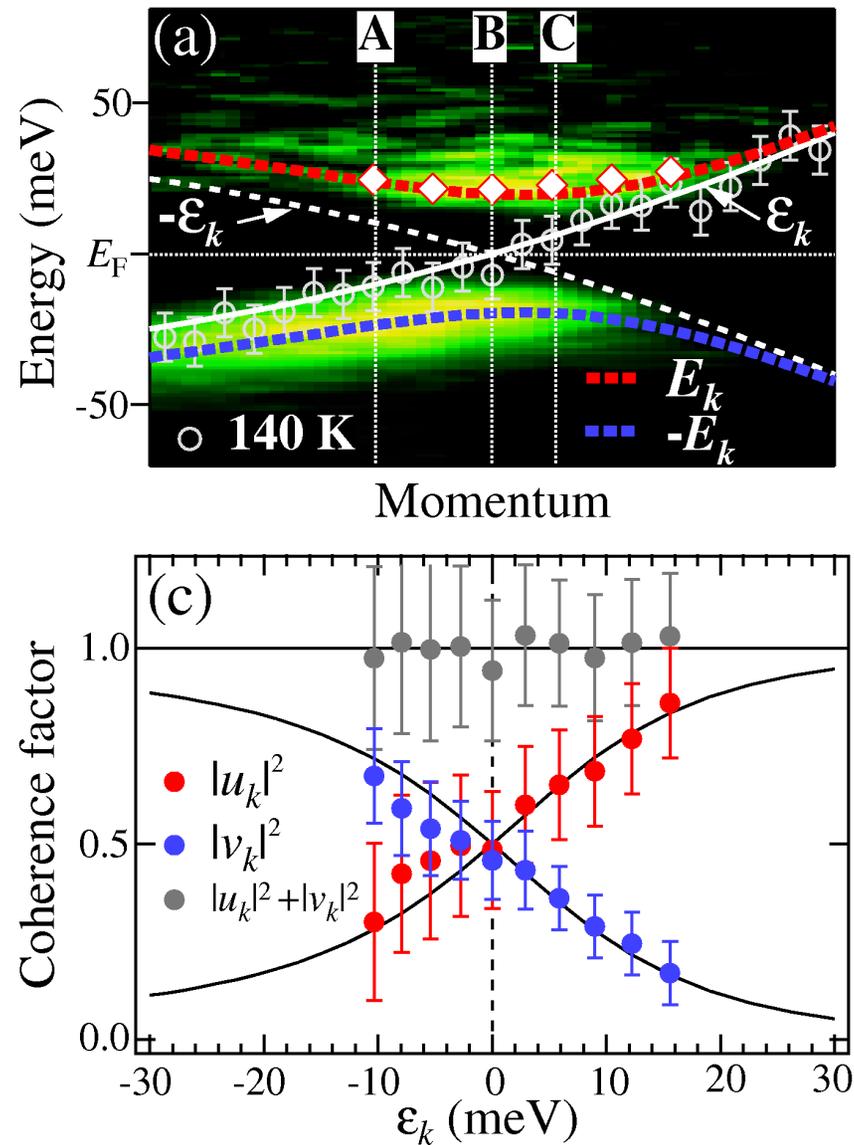
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Experimental data for $T < T_c$

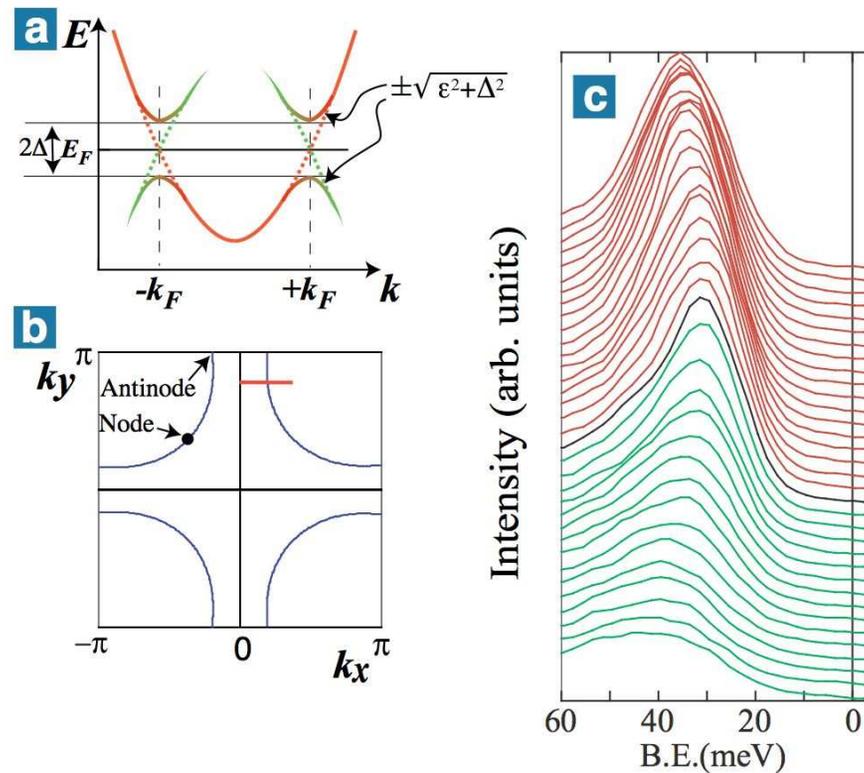


H. Matsui, T. Sato, and T. Takahashi et al, *Phys. Rev. Lett.* **90**, 217002 (2003).

Evidence for Bogoliubov QPs above T_c

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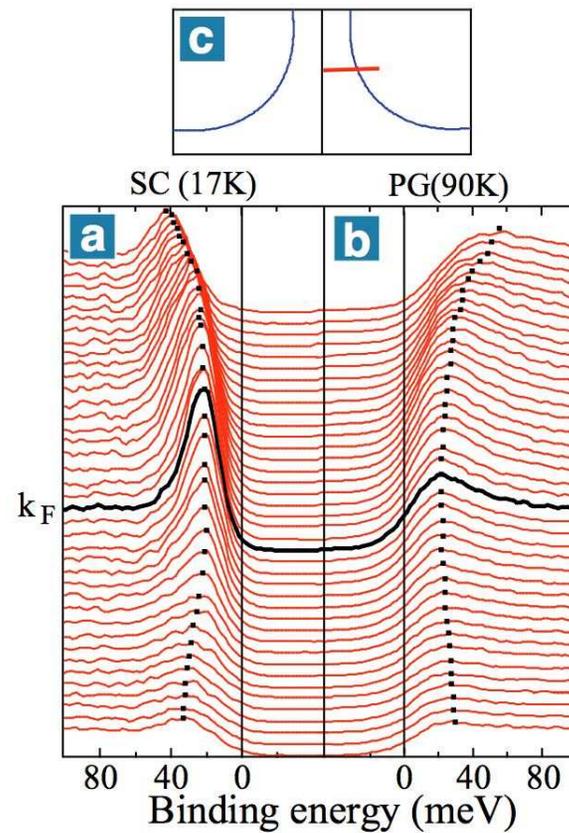


Results for: $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$

A. Kanigel et al, *Phys. Rev. Lett.* **101**, 137002 (2008).

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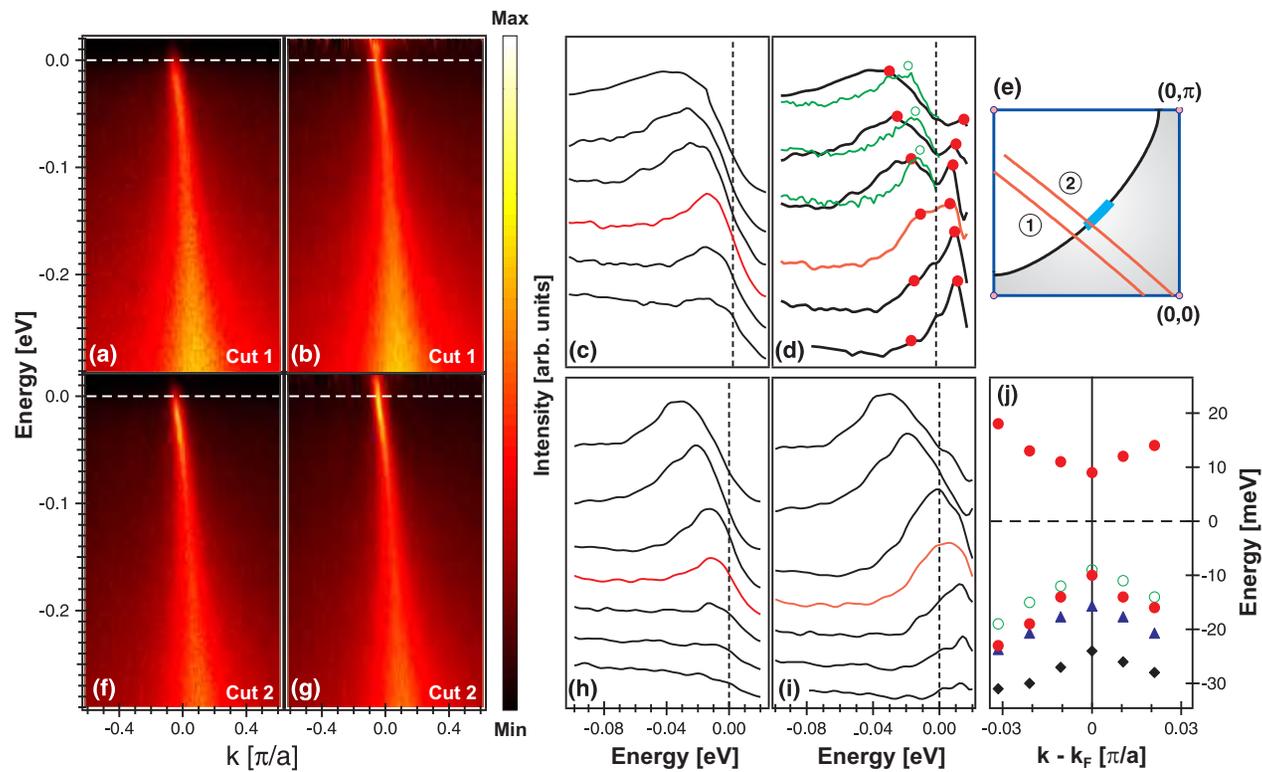


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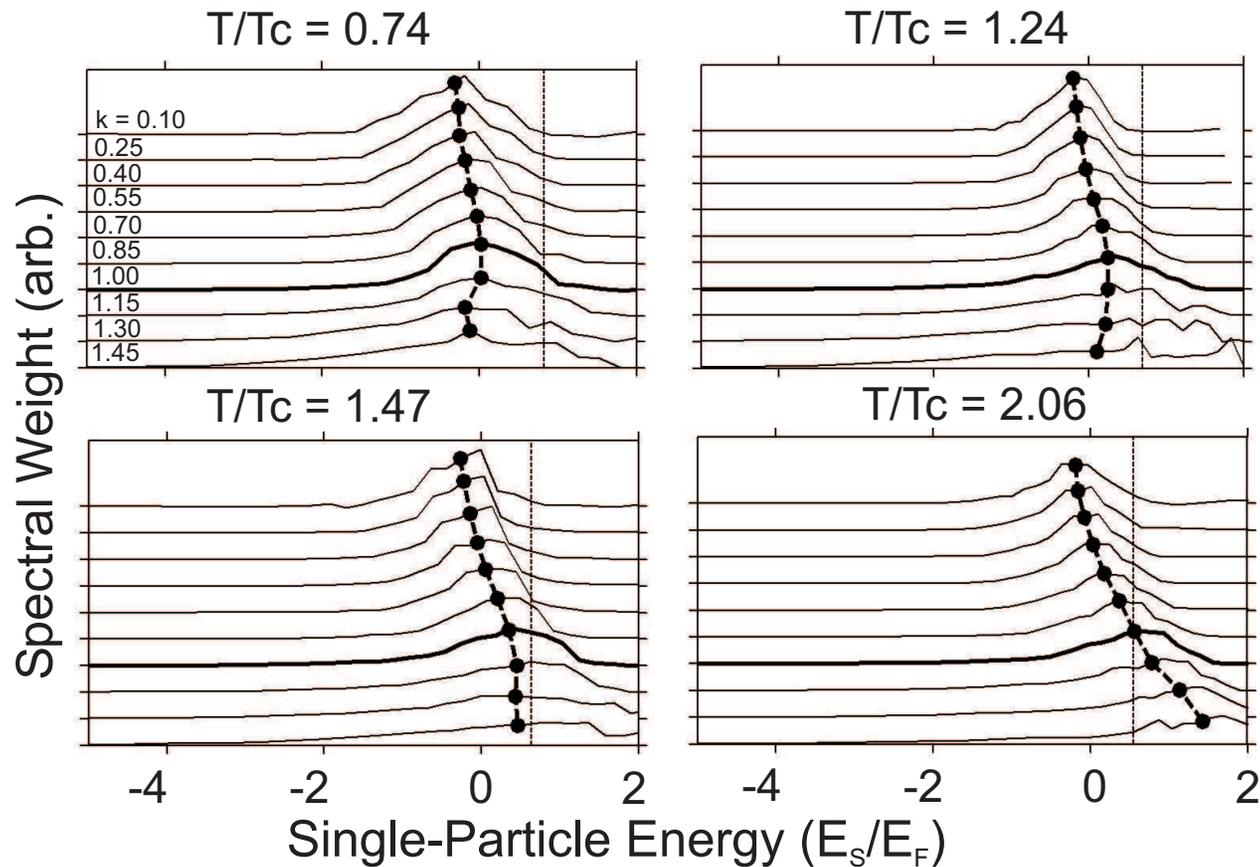


Results for: $\text{La}_{1.895}\text{Sr}_{0.105}\text{CuO}_4$

M. Shi et al, Eur. Phys. Lett. 88, 27008 (2009).

Evidence for Bogoliubov QPs above T_c

D. Jin group (Boulder, USA)



Results for: ultracold ^{40}K atoms

J.P. Gaebler et al, Nature Phys. 6, 569 (2010).

Results :

2. Diamagnetism above T_c

Correlation functions

For studying the diamagnetic response (in the Kubo formalism) we have to determine the current-current correlation function

$$- \hat{T}_\tau \langle \hat{j}_q(\tau) \hat{j}_{-q}(0) \rangle$$

with statistical averaging defined as

$$\langle \dots \rangle = \text{Tr} \left\{ e^{-\beta \hat{H}} \dots \right\} / \text{Tr} \left\{ e^{-\beta \hat{H}} \right\}$$

and $\beta^{-1} = k_B T$.

This can be achieved using the following invariance

$$\begin{aligned} \text{Tr} \left\{ e^{-\beta \hat{H}} \hat{O} \right\} &= \text{Tr} \left\{ e^{\hat{S}(l)} e^{-\beta \hat{H}} \hat{O} e^{-\hat{S}(l)} \right\} \\ &= \text{Tr} \left\{ e^{\hat{S}(l)} e^{-\beta \hat{H}} e^{-\hat{S}(l)} e^{\hat{S}(l)} \hat{O} e^{-\hat{S}(l)} \right\} \\ &= \text{Tr} \left\{ e^{-\beta \hat{H}(l)} \hat{O}(l) \right\} \end{aligned}$$

where

$$\hat{H}(l) = e^{\hat{S}(l)} \hat{H} e^{-\hat{S}(l)}$$

$$\hat{O}(l) = e^{\hat{S}(l)} \hat{O} e^{-\hat{S}(l)}$$

Correlation functions

– some remarks

- ★ The easiest way for calculating $\langle \hat{j}_q(\tau) \hat{j}_{-q} \rangle$ is in a limit $l \longrightarrow \infty$ when $\hat{H}(\infty)$ becomes (block-)diagonal.
- ★ The operators \hat{j}_q must however undergo the continuous transformation

$$\hat{j} \longrightarrow \dots \longrightarrow \hat{j}(l) \longrightarrow \dots \longrightarrow \hat{j}(\infty)$$

- ★ obeying **the flow equation**:

$$\frac{\partial \hat{j}_q(l)}{\partial l} = \left[\hat{\eta}(l), \hat{j}_q(l) \right]$$

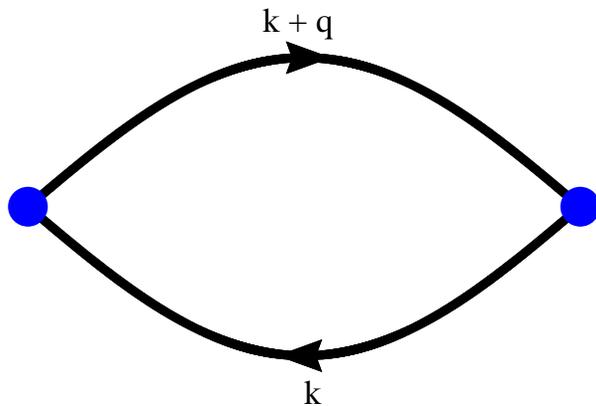
Diamagnetic response above T_c

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Main contributions to the current-current response function:

Diamagnetic response above T_c

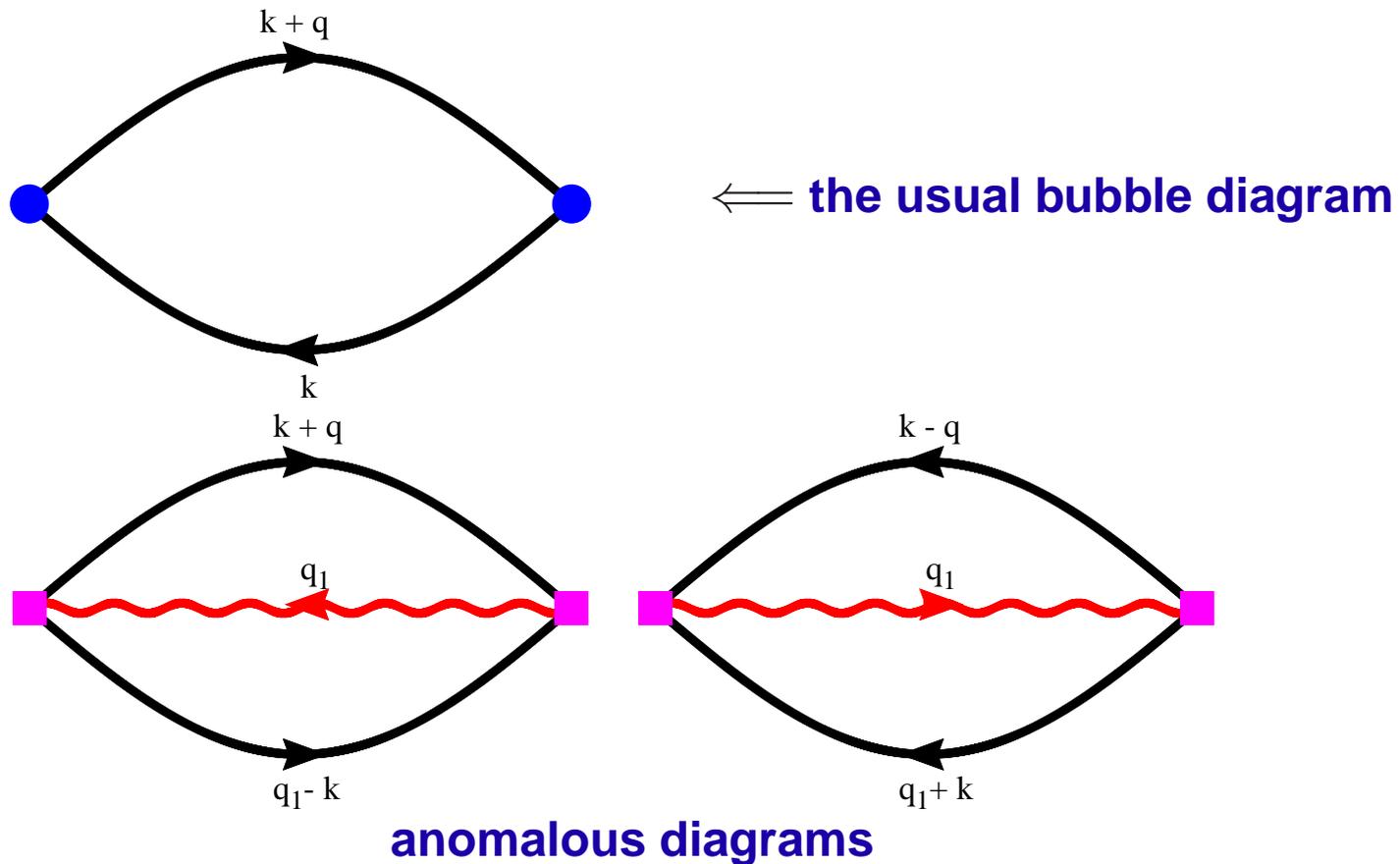
Main contributions to the current-current response function:



← the usual bubble diagram

Diamagnetic response above T_c

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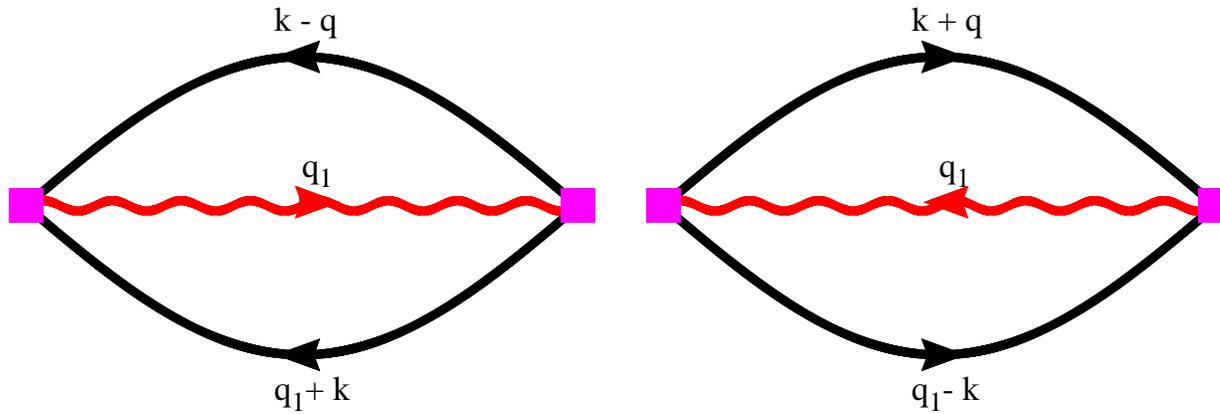
Each **vertex** has to be determined from the flow equations.

T. Domanski and J. Ranninger, (to be published).

Diamagnetic response above T_c

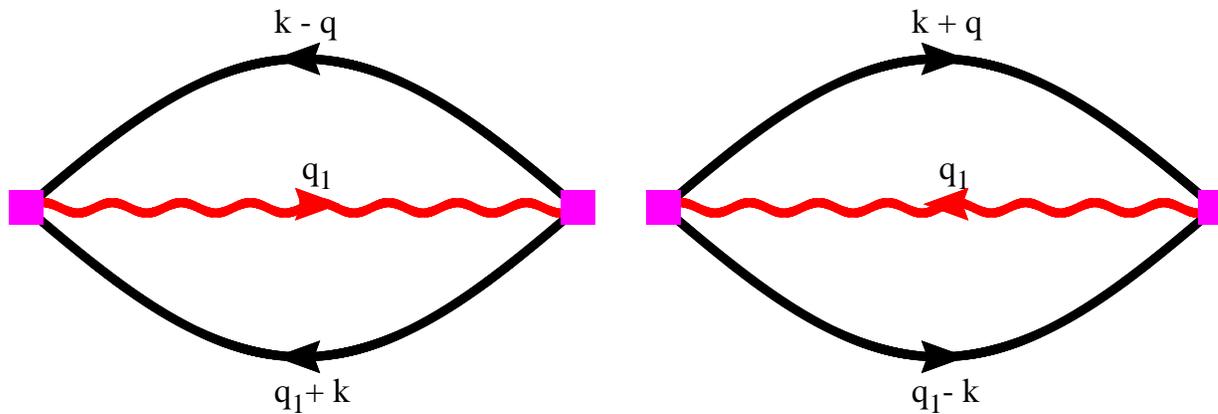
Diamagnetic response above T_c

The anomalous contributions

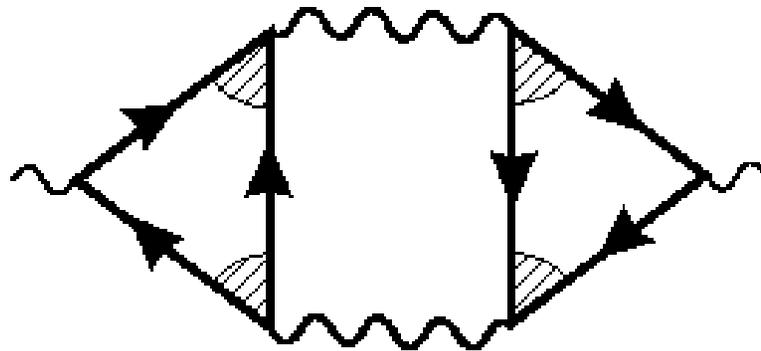


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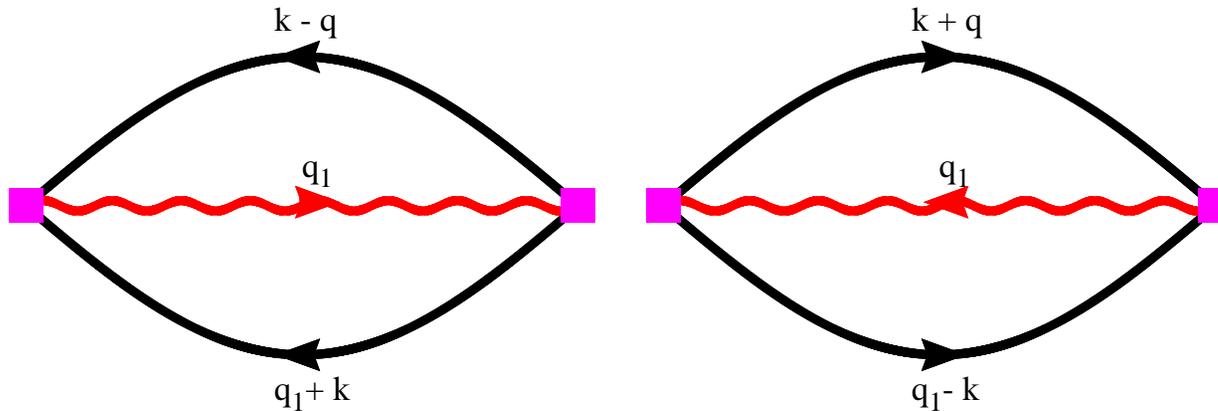


resemble the Aslamasov-Larkin diagram

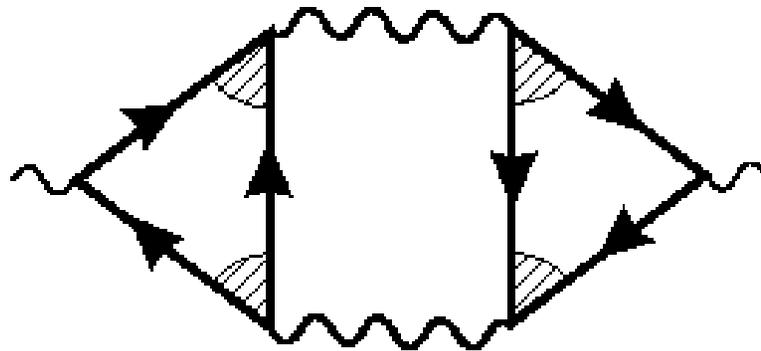


Diamagnetic response above T_c

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resemble the Aslamasov-Larkin diagram



enhancing the conductance/diamagnetism above T_c .

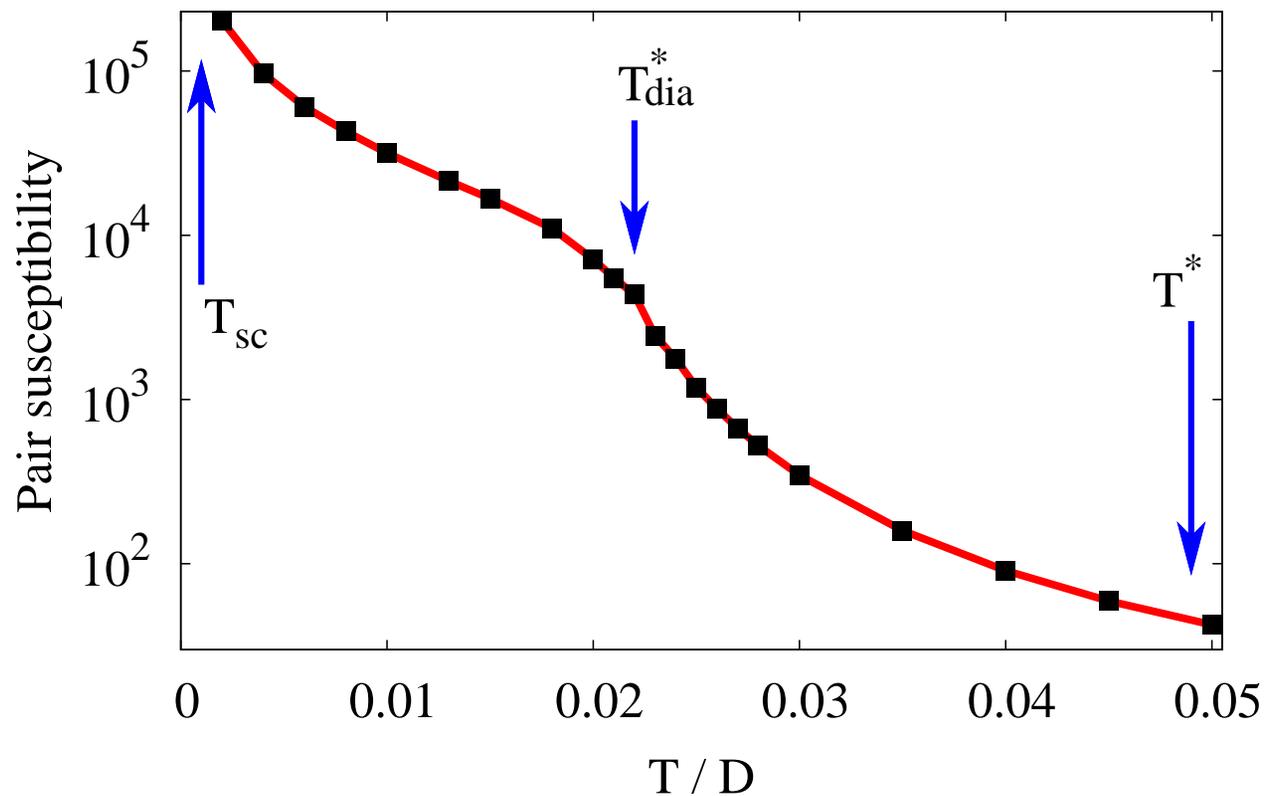
Diamagnetic response above T_c

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Residual diamagnetism originates from the collective behavior of pairs. It is closely related with increase of **pair susceptibility**, which **is enhanced at T_{dia}^*** and ultimately **diverges at T_{sc}** .

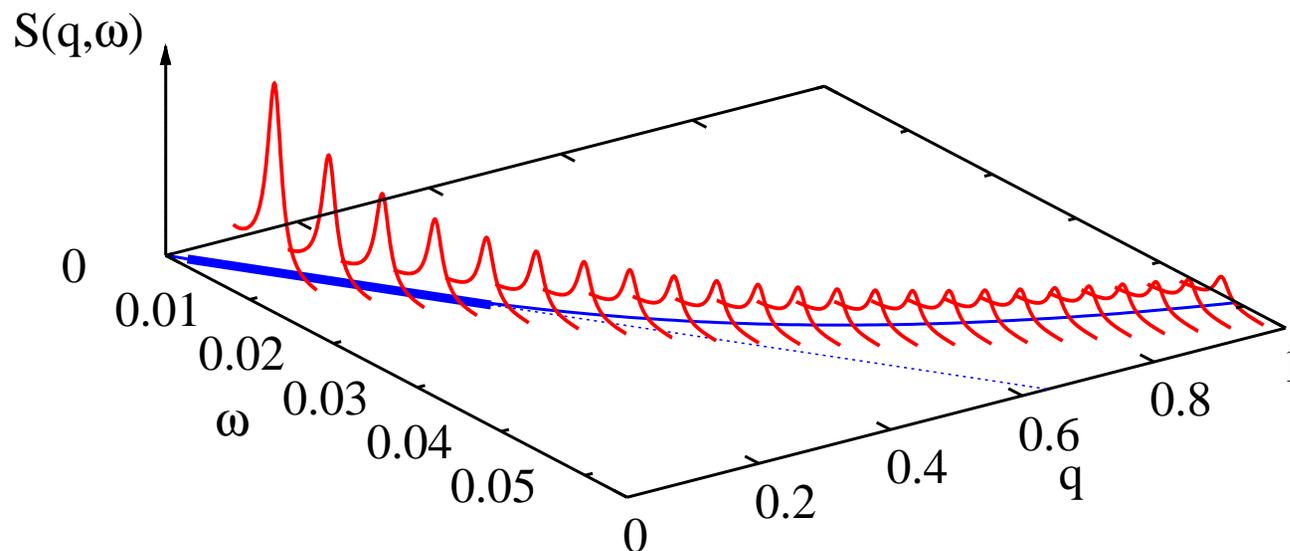
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The structure function $\frac{1}{\omega} \text{Im} \chi(q, \omega)$ showing a piece of the collective (Goldstone) branch for $q_{c1} < q < q_{c2}$.

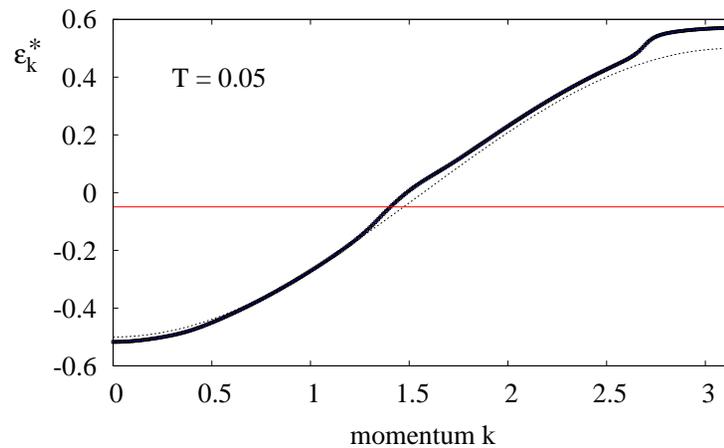
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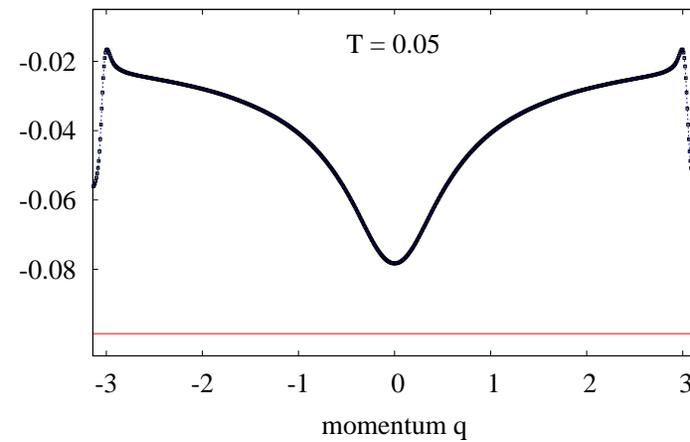
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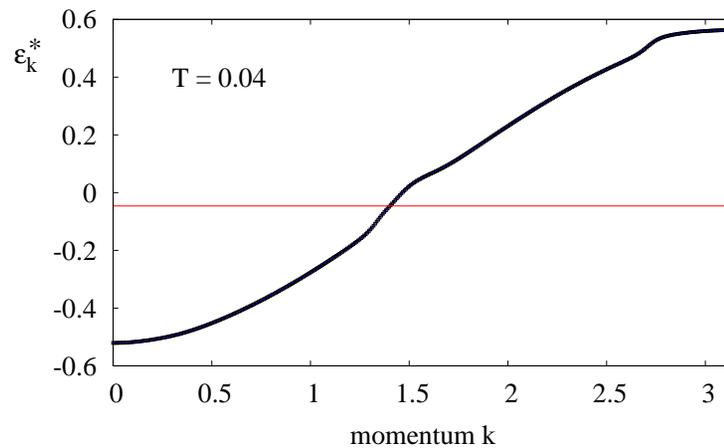
fermion dispersion



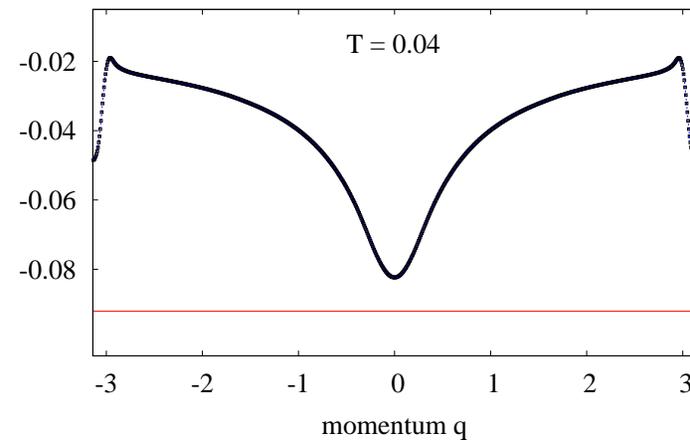
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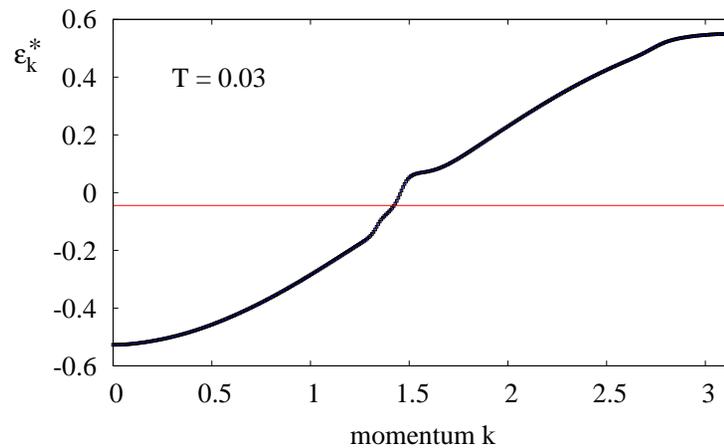
fermion dispersion



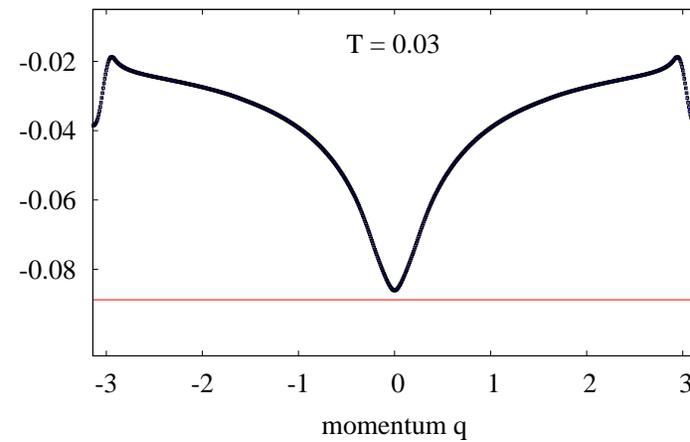
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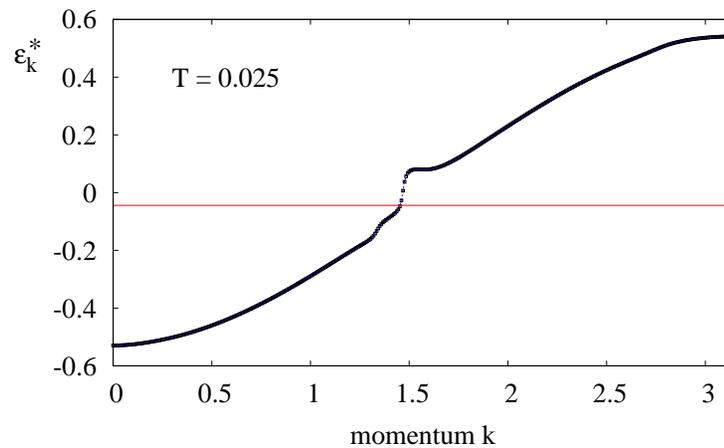
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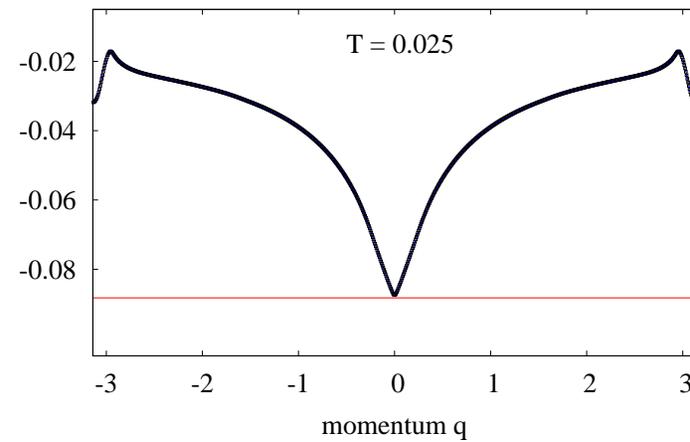
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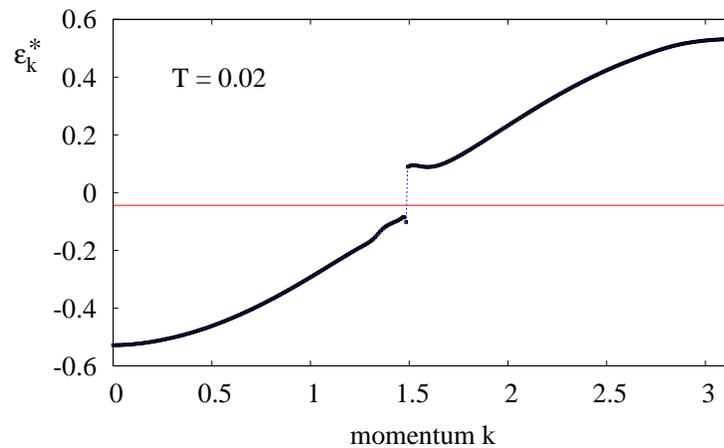
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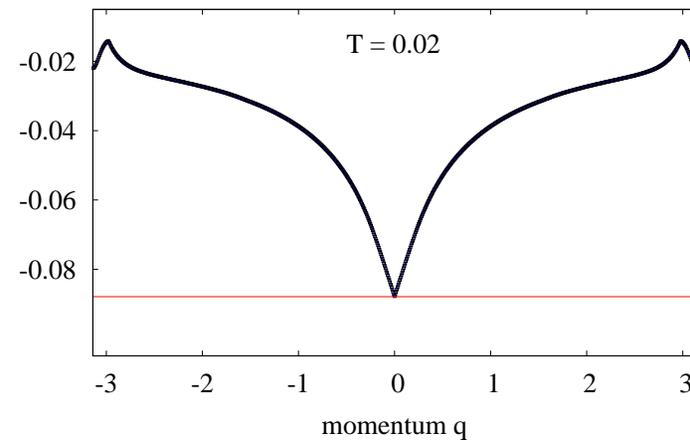
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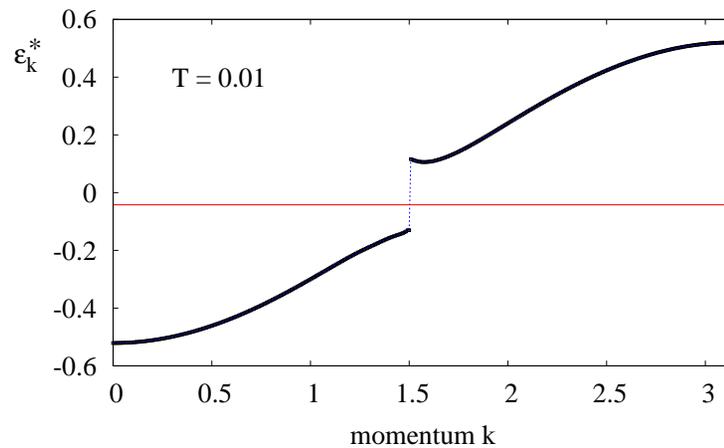
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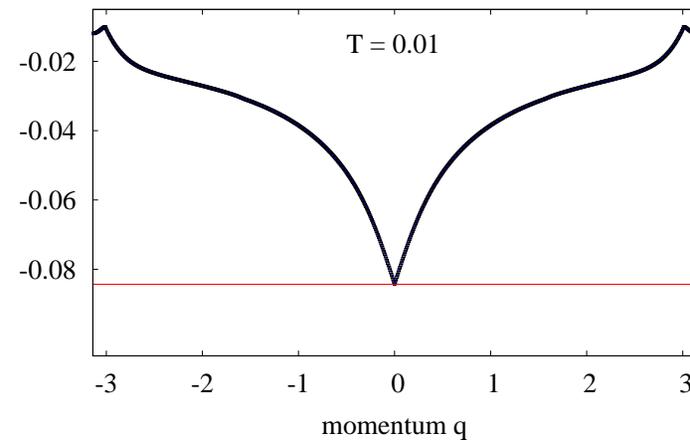
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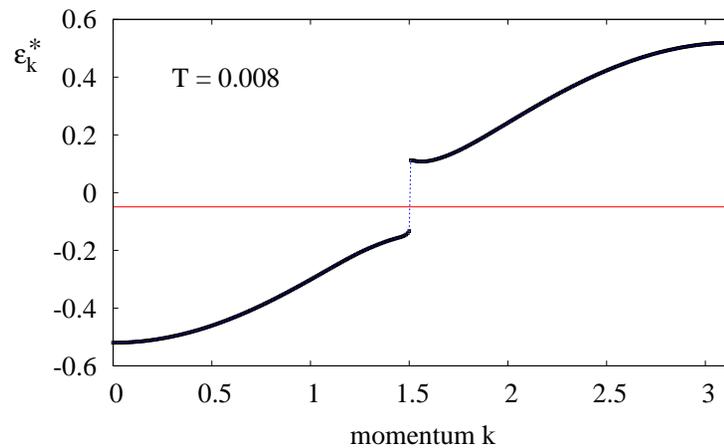
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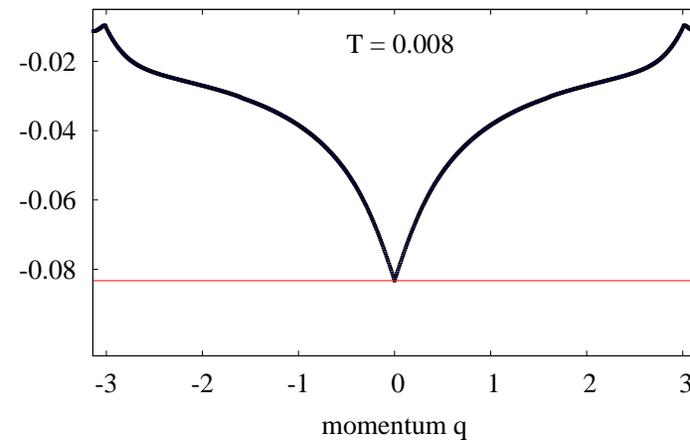
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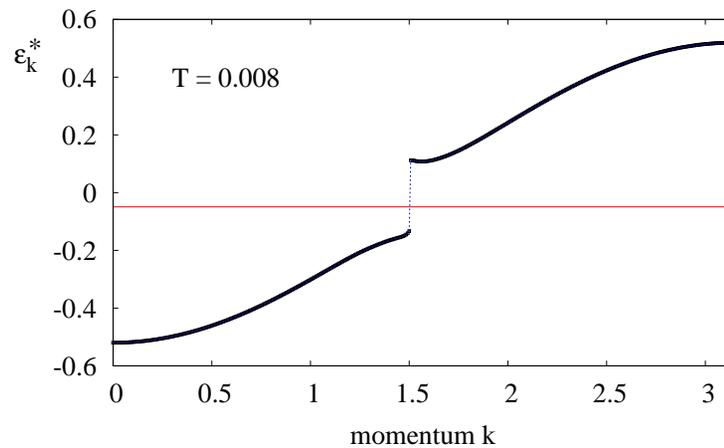
fermion dispersion



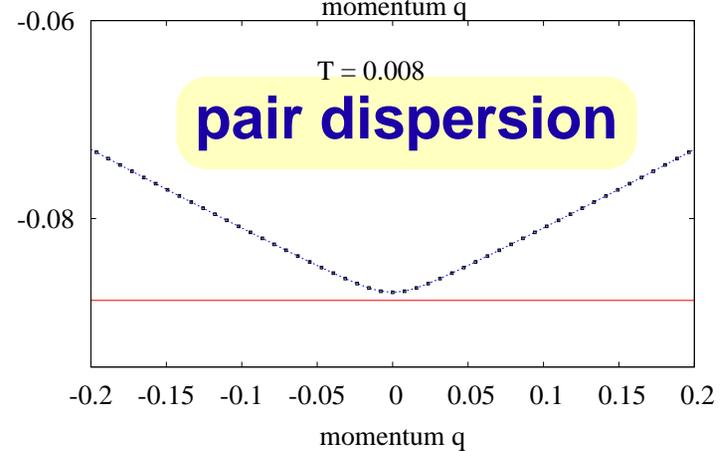
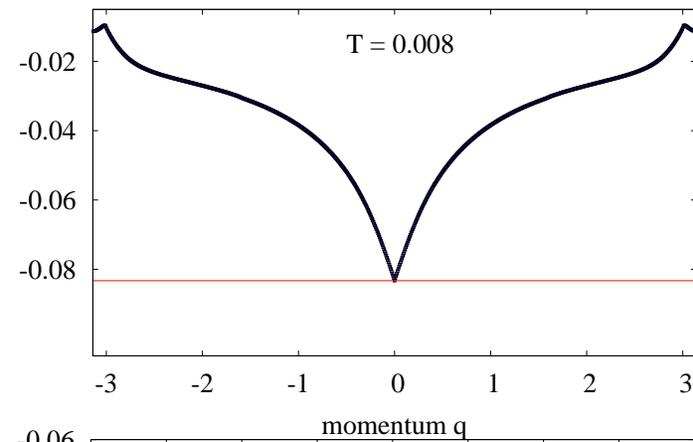
pair dispersion

Onset of diamagnetism above T_c

Onset of the diamagnetism coincides with appearance of the collective features in the fermion/boson spectrum.



fermion dispersion



pair dispersion

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Thank you.