

Lublin, 16 XI 2010

Korelacje nadprzewodzące powyżej T_c

T. DOMAŃSKI

**Uniwersytet M. Curie-Skłodowskiej
w Lublinie**

Współpraca: J. Ranninger (CNRS, Grenoble)

Lublin, 16 XI 2010

**Signatures of the short-range
superconducting correlations above T_c**

T. DOMAŃSKI

**M. Curie-Skłodowska University,
Lublin, Poland**

Collaboration: J. Ranninger (CNRS, Grenoble)

Outline

Outline



Preliminaries

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology



Selected results

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology



Selected results

\Rightarrow *Bogoliubov-type modes above T_c*

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology



Selected results

- Bogoliubov-type modes above T_c

\Rightarrow Fermi arcs

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology



Selected results

- *Bogoliubov-type modes above T_c*
- *Fermi arcs*

\Rightarrow *Diamagnetism above T_c*

Outline



Preliminaries



Experimental motivation

/ pre-pairing for condensation /



Scenario & methodology



Selected results

- *Bogoliubov-type modes above T_c*
- *Fermi arcs*
- *Diamagnetism above T_c*



Conclusions

Preliminaries

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

2. **exchange of magnons**

/ superconductivity of the heavy fermion compounds /

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

2. **exchange of magnons**

/ superconductivity of the heavy fermion compounds /

3. **strong correlations**

/ high T_c superconductors /

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

2. **exchange of magnons**

/ superconductivity of the heavy fermion compounds /

3. **strong correlations**

/ high T_c superconductors /

4. **Feshbach resonance**

/ ultracold superfluid atoms /

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

2. **exchange of magnons**

/ superconductivity of the heavy fermion compounds /

3. **strong correlations**

/ high T_c superconductors /

4. **Feshbach resonance**

/ ultracold superfluid atoms /

5. **other**

/ pairing in nuclei, gluon-quark plasma /

Pairing is a common phenomenon which occurs between various kinds of fermions such as: quarks, electrons, nucleons or atoms.

The underlying **pairing mechanism** can be driven by:

1. **exchange of phonons**

/ classical superconductors, MgB_2 , diamond, ... /

2. **exchange of magnons**

/ superconductivity of the heavy fermion compounds /

3. **strong correlations**

/ high T_c superconductors /

4. **Feshbach resonance**

/ ultracold superfluid atoms /

5. **other**

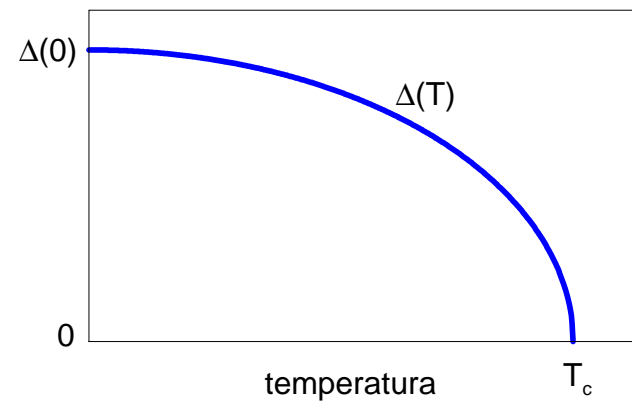
/ pairing in nuclei, gluon-quark plasma /

Very often formation of the fermion pairs goes hand in hand with **superconductivity/superfluidity** but it needs not be the rule.

Conventional superconductors – major property

Conventional superconductors – major property

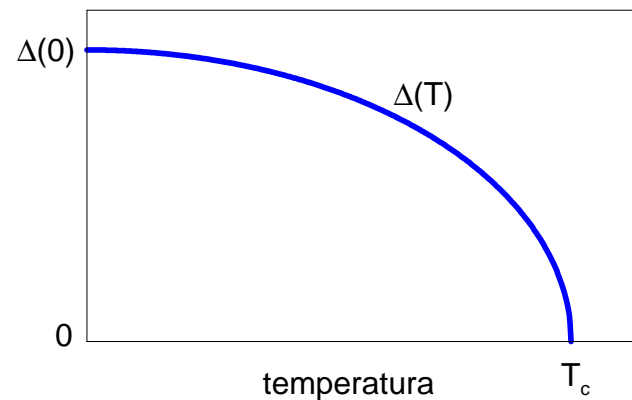
Pair formation and onset of their coherence coincide at T_c



Pairing is
responsible
for the gap
in the single
particle spectrum

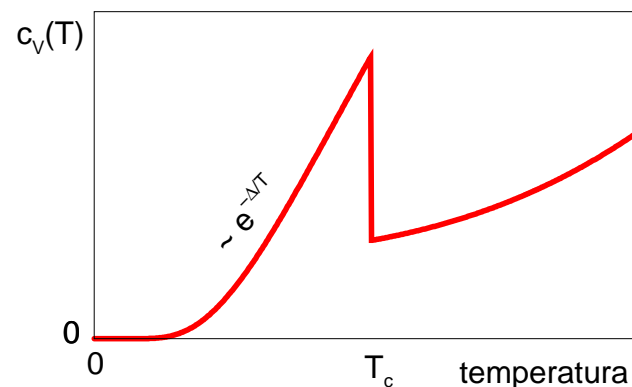
Conventional superconductors – major property

Pair formation and onset of their coherence coincide at T_c



Pairing is responsible for the gap in the single particle spectrum

The order parameter \longrightarrow 2-nd order phase transition



/ as classified by Landau /

Conventional superconductors – meaning of T_c

Conventional superconductors – meaning of T_c

At $T \rightarrow T_c$ electrons near the Fermi energy:

Conventional superconductors – meaning of T_c

At $T \rightarrow T_c$ electrons near the Fermi energy:

form the Cooper pairs

Conventional superconductors – meaning of T_c

At $T \rightarrow T_c$ electrons near the Fermi energy:

form the Cooper pairs

**and behave as a huge super-atom consisting of
 $\sim 10^{23}$ particles all gathered in an identical state.**

Conventional superconductors – meaning of T_c

At $T \rightarrow T_c$ electrons near the Fermi energy:

form the Cooper pairs

and behave as a huge super-atom consisting of
 $\sim 10^{23}$ particles all gathered in an identical state.

This is BE condensate of Cooper pairs !

Theoretical issues – generalities

Theoretical issues – generalities

The order parameter

$$\chi(\vec{r}_i, \vec{r}_j) \equiv \langle \hat{c}_{\downarrow}(\vec{r}_i) \hat{c}_{\uparrow}(\vec{r}_j) \rangle$$

Theoretical issues – generalities

The order parameter $\chi(\vec{r}_i, \vec{r}_j) \equiv \langle \hat{c}_\downarrow(\vec{r}_i) \hat{c}_\uparrow(\vec{r}_j) \rangle$

is in general a complex quantity

$$\chi = |\chi| e^{i\theta}$$

Theoretical issues – generalities

The order parameter $\chi(\vec{r}_i, \vec{r}_j) \equiv \langle \hat{c}_\downarrow(\vec{r}_i) \hat{c}_\uparrow(\vec{r}_j) \rangle$

is in general a complex quantity

$$\chi = |\chi| e^{i\theta}$$

with the following physical implications:

Theoretical issues – generalities

The order parameter $\chi(\vec{r}_i, \vec{r}_j) \equiv \langle \hat{c}_\downarrow(\vec{r}_i) \hat{c}_\uparrow(\vec{r}_j) \rangle$

is in general a complex quantity

$$\chi = |\chi| e^{i\theta}$$

with the following physical implications:

$|\chi| \neq 0 \longrightarrow$ amplitude causes the energy gap

Theoretical issues – generalities

The order parameter $\chi(\vec{r}_i, \vec{r}_j) \equiv \langle \hat{c}_\downarrow(\vec{r}_i) \hat{c}_\uparrow(\vec{r}_j) \rangle$

is in general a complex quantity

$$\chi = |\chi| e^{i\theta}$$

with the following physical implications:

$|\chi| \neq 0 \longrightarrow$ amplitude causes the energy gap

$\nabla\theta \neq 0 \longrightarrow$ phase slippage causes supercurrents

Phase transitions – classification

Phase transitions – classification

Vanishing of the complex order parameter

$$\chi = |\chi| e^{i\theta}$$

Phase transitions – classification

Vanishing of the complex order parameter

$$\chi = |\chi| e^{i\theta}$$

can be achieved at T_c either:

Phase transitions – classification

Vanishing of the complex order parameter

$$\chi = |\chi| e^{i\theta}$$

can be achieved at T_c either:

by closing the gap(BCS superconductors)

$$\lim_{T \rightarrow T_c} |\chi| = 0$$

Phase transitions – classification

Vanishing of the complex order parameter

$$\chi = |\chi| e^{i\theta}$$

can be achieved at T_c either:

by closing the gap(BCS superconductors)

$$\lim_{T \rightarrow T_c} |\chi| = 0$$

or disordering the phase(the HTSC compounds)

$$\lim_{T \rightarrow T_c} \langle \theta \rangle = 0$$

Historical remark

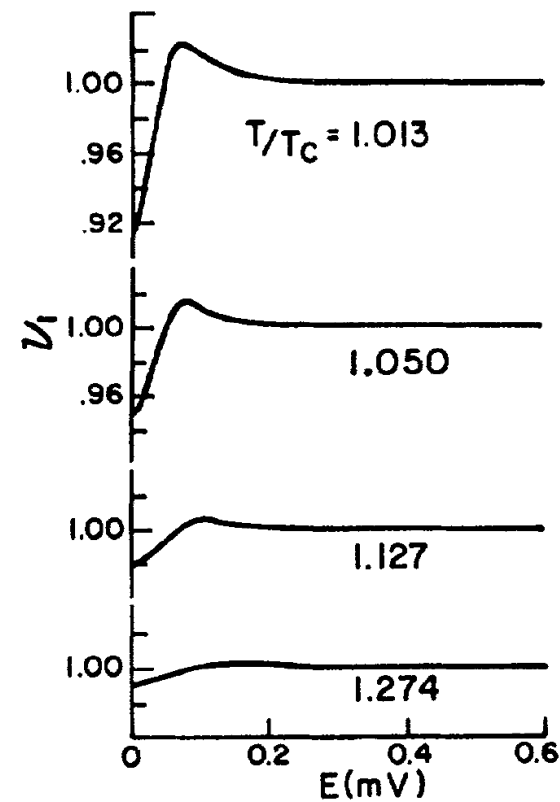
Historical remark

The first empirical observation of the sc fluctuations above T_c has been seen in **granular aluminium**.

Historical remark

The first empirical observation of the sc fluctuations above T_c has been seen in **granular aluminium**.

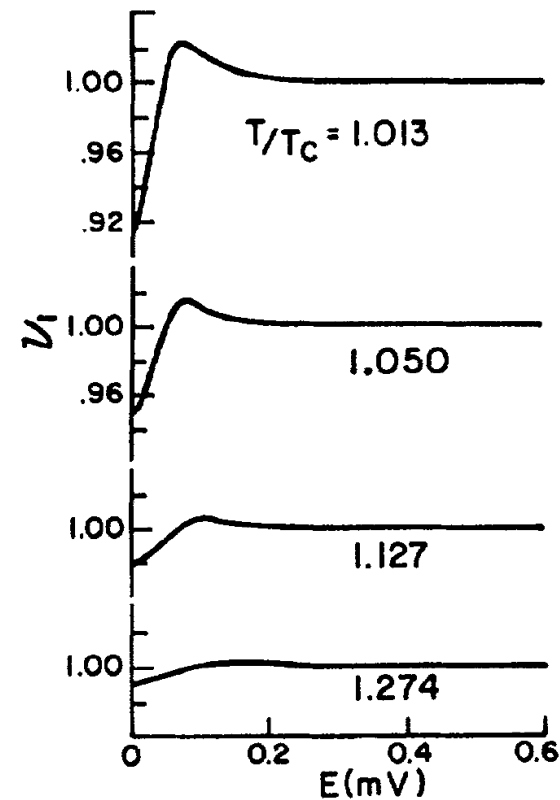
*Tunneling
conductance
revealed
a small
pseudogap
above T_c .*



Historical remark

The first empirical observation of the sc fluctuations above T_c has been seen in **granular aluminium**.

*Tunneling
conductance
revealed
a small
pseudogap
above T_c .*

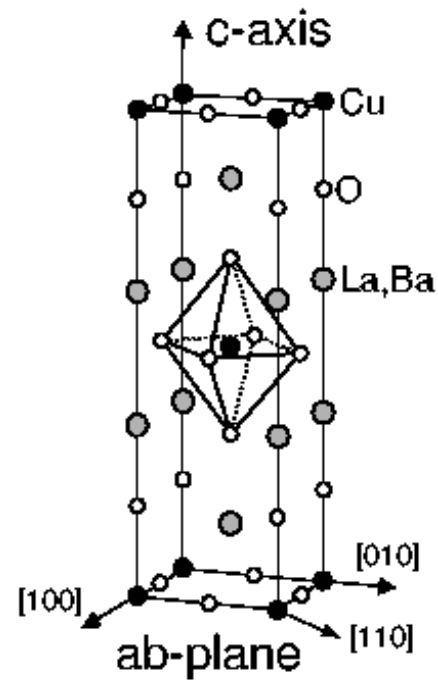


*R.W. Cohen and B. Abels, Phys. Rev. **168**, 444 (1968).*

Experimental motivation

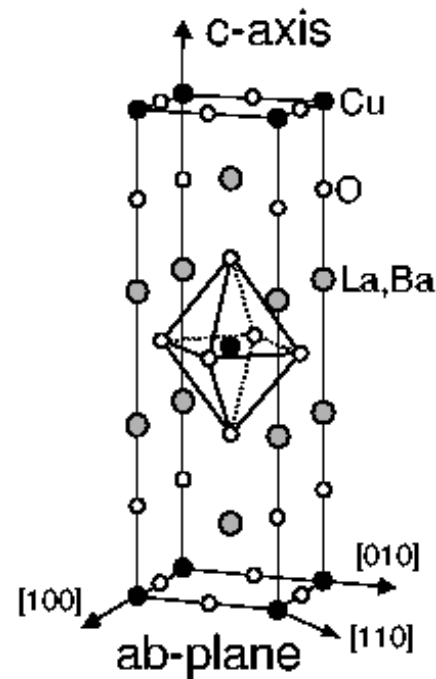
HTSC materials – structure

The parent compounds are quasi-2D Mott insulators



HTSC materials – structure

The parent compounds are quasi-2D Mott insulators



Important remark:

Spatial extent of the pairs is very short $\xi_{ab} \simeq 5 \text{ \AA}$

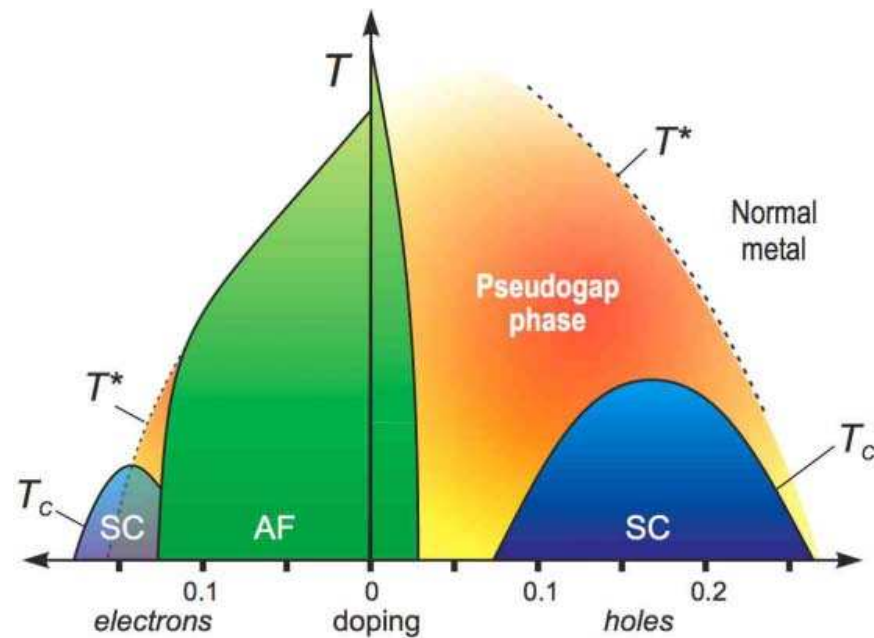
HTSC materials – effect of doping

Superconductivity appears upon doping by

HTSC materials – effect of doping

Superconductivity appears upon doping by

electrons or *holes*

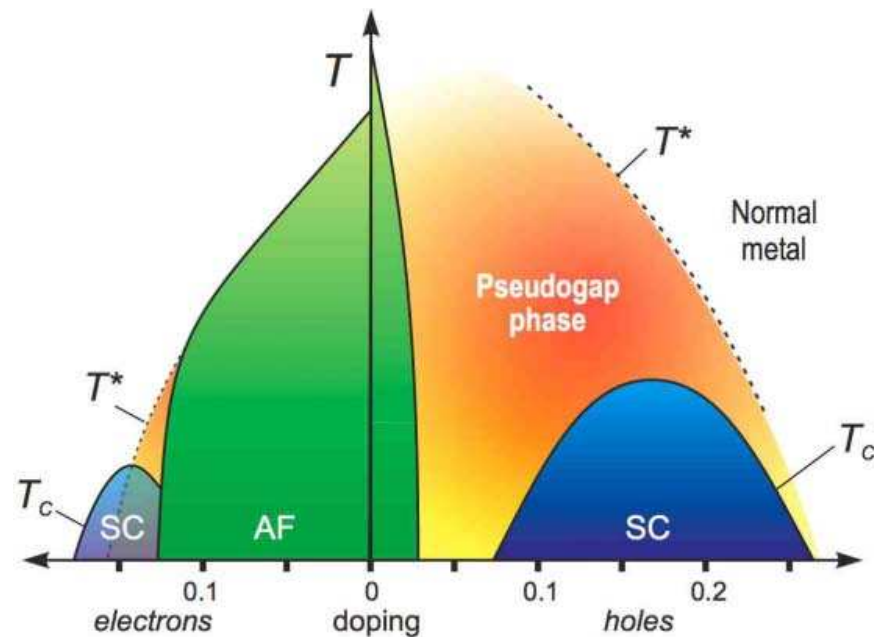


O. Fisher et al, Rev. Mod. Phys. 79, 353 (2007).

HTSC materials – effect of doping

Superconductivity appears upon doping by

electrons or *holes*



O. Fisher et al, *Rev. Mod. Phys.* **79**, 353 (2007).

Important remark:

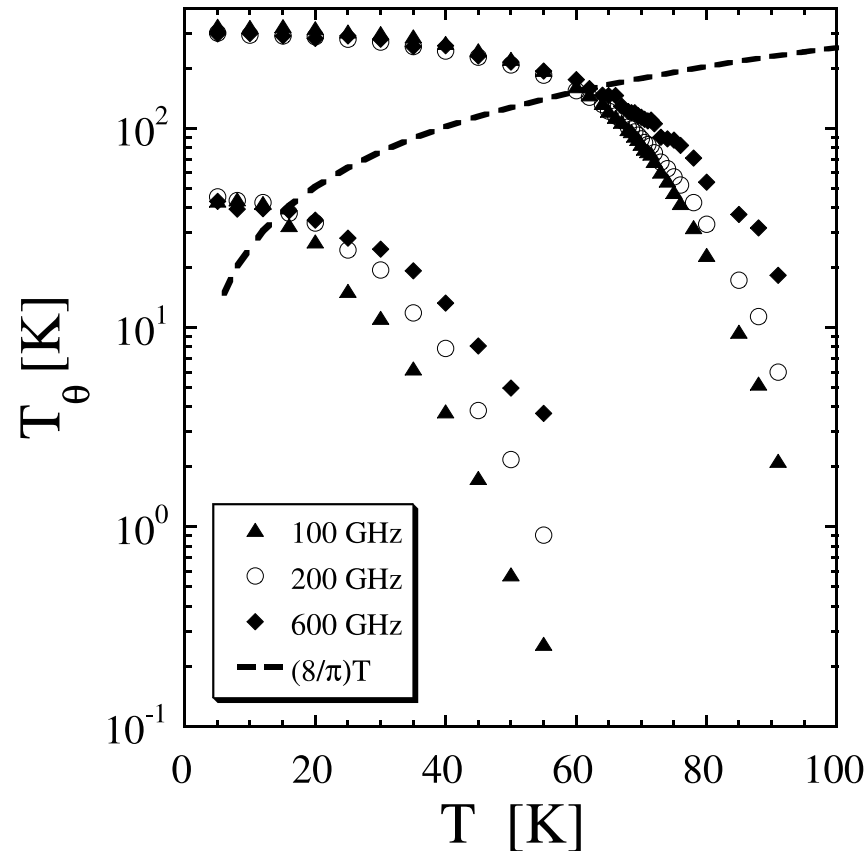
What is an origin of the pseudogap ?

Incoherent pairs above T_c

experimental fact # 1

Incoherent pairs above T_c

experimental fact # 1



Dynamic phase-stiffness $T_\theta = \omega \text{Im}\sigma(\omega, T) / \sigma_Q$
observed at the ultrafast (teraHz) external ac fields.

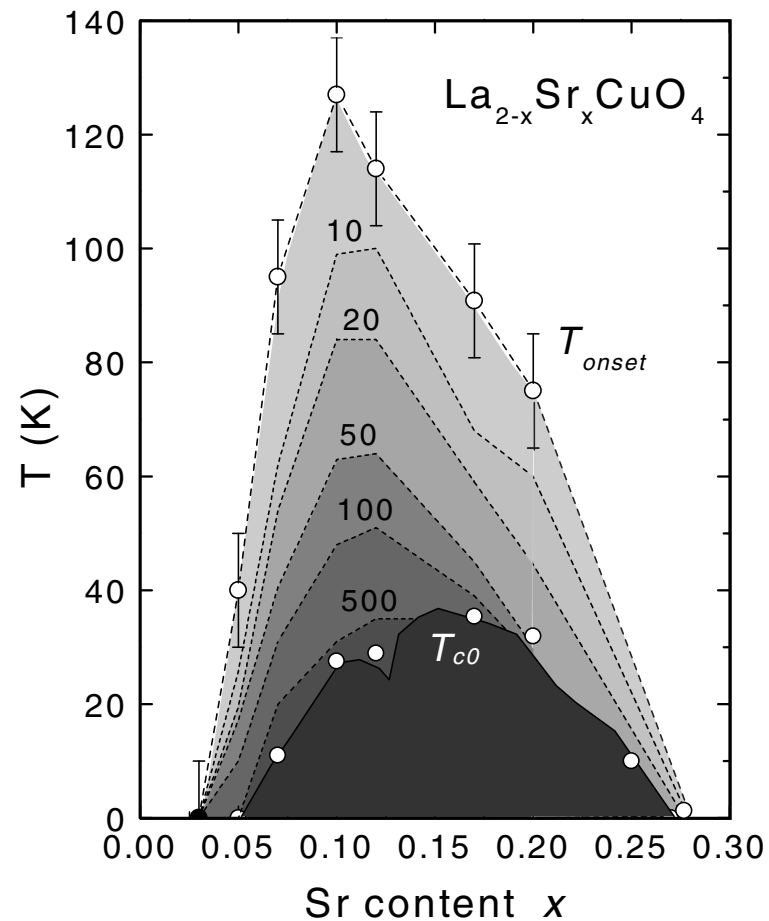
J. Corson et al, Nature 398, 221 (1999).

Incoherent pairs above T_c

experimental fact # 2

Incoherent pairs above T_c

experimental fact # 2



Phase slippage detected in the large Nernst effect.

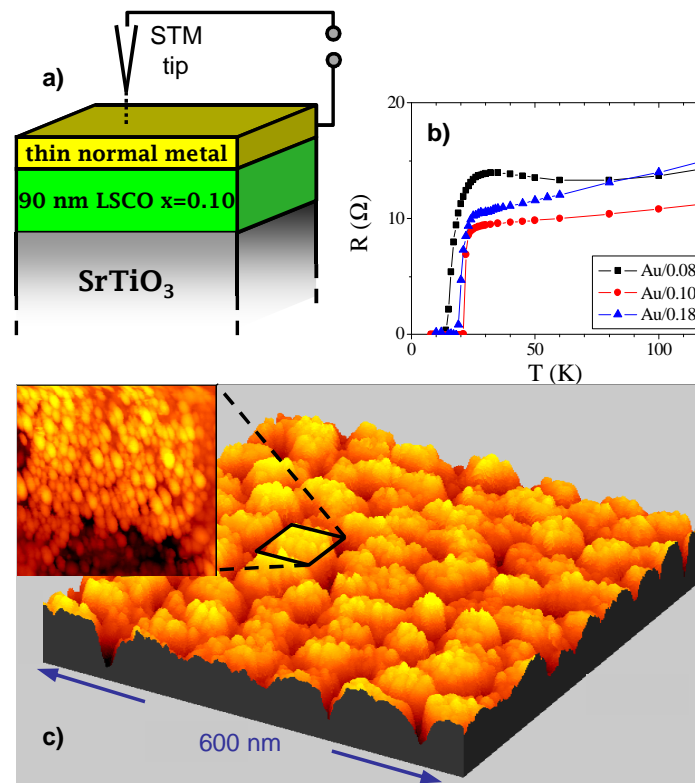
Y. Wang et al, *Science* **299**, 86 (2003).

Incoherent pairs above T_c

experimental fact # 3

Incoherent pairs above T_c

experimental fact # 3



O. Yuli et al, *Phys. Rev. Lett.* **103**, 197003 (2009).

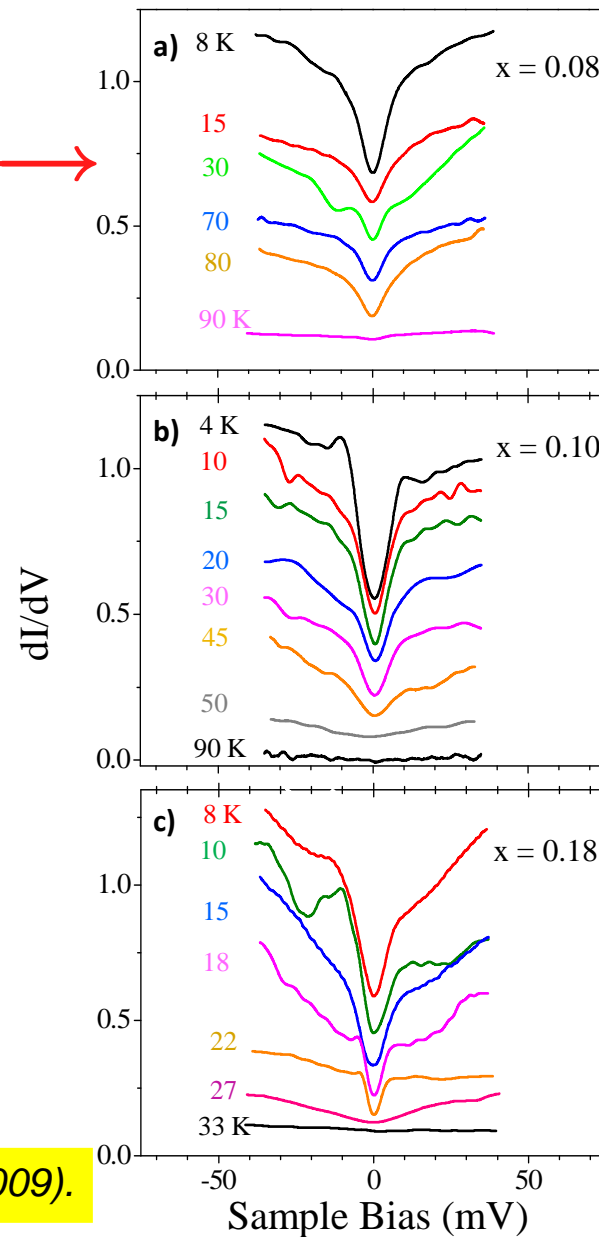
Pseudogap induced above T_c in the ultrathin metallic slab deposited on $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$.

Incoherent pairs above T_c

experimental fact # 3

Measured LDOS

Proximity
induced
pseudogap
has been
observed at
temperatures
far above T_c .



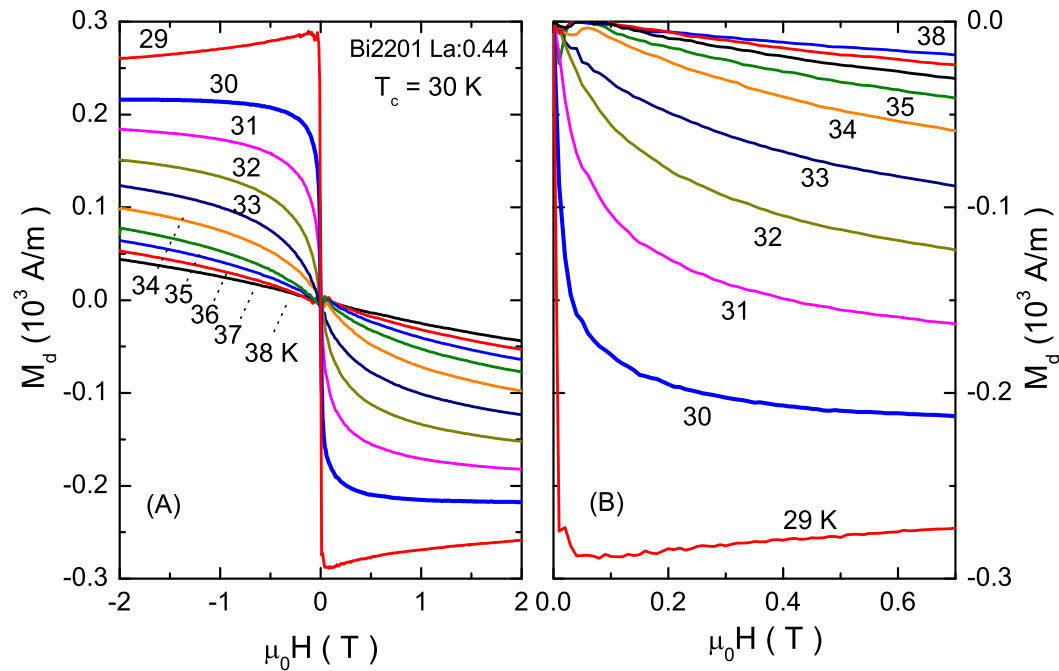
O. Yuli et al, *PRL* **103**, 197003 (2009).

Incoherent pairs above T_c

experimental fact # 4

Incoherent pairs above T_c

experimental fact # 4

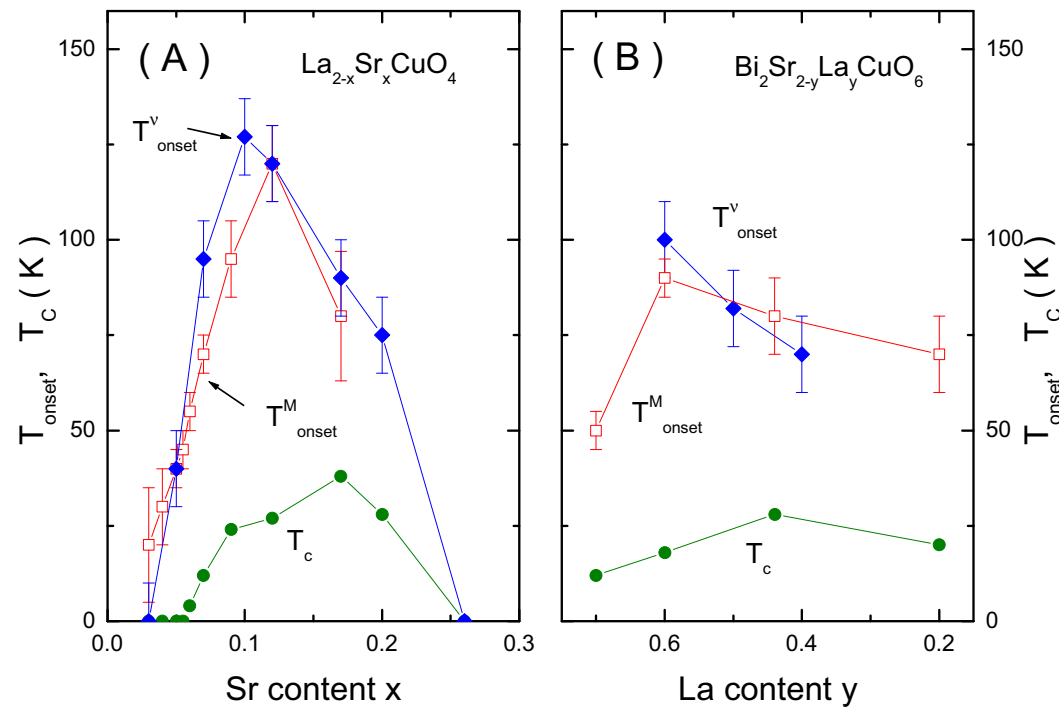


**Onset of the diamagnetic response
revealed by torque magnetometry.**

*L. Li, ... and N.P. Ong, Phys. Rev. B **81**, 054510 (2010).*

Incoherent pairs above T_c

experimental fact # 4



**Onset of the diamagnetic response
revealed by torque magnetometry.**

*L. Li, ... and N.P. Ong, Phys. Rev. B **81**, 054510 (2010).*

Incoherent pairs above T_c

... *continued*

Incoherent pairs above T_c

... continued



Josephson-like features seen above T_c in the tunneling

N. Bergeal et al, Nature Phys. 4, 608 (2008).

Incoherent pairs above T_c

... continued

⇒ **Josephson-like features seen above T_c in the tunneling**

N. Bergeal et al, Nature Phys. 4, 608 (2008).

⇒ **smooth evolution of the electronic spectrum observed by ARPES near the superconductor–insulator transition**

U. Chatterjee et al, Nature Phys. 5, 1456 (2009).

Incoherent pairs above T_c

... continued

⇒ **Josephson-like features seen above T_c in the tunneling**

N. Bergeal et al, Nature Phys. 4, 608 (2008).

⇒ **smooth evolution of the electronic spectrum observed by ARPES near the superconductor–insulator transition**

U. Chatterjee et al, Nature Phys. 5, 1456 (2009).

⇒ **spectroscopic fingerprints of the Bogoliubov QPs seen by the unique octet patterns which survive up to $1.5T_c$**

J. Lee, ... and J.C. Davis, Science 325, 1099 (2009).

II. Model & methodology

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right]\end{aligned}$$

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right]\end{aligned}$$

$$\vec{R}_l = (\vec{r}_i + \vec{r}_j)/2$$

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right]\end{aligned}$$
$$\vec{R}_l = (\vec{r}_i + \vec{r}_j)/2$$

describes a two-component system consisting of:

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right] \quad \vec{R}_l = (\vec{r}_i + \vec{r}_j)/2\end{aligned}$$

describes a two-component system consisting of:

$\hat{c}_{i\sigma}^{(\dagger)}$ itinerant fermions (e.g. holes near the Mott insulator)

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right] \quad \vec{R}_l = (\vec{r}_i + \vec{r}_j)/2\end{aligned}$$

describes a two-component system consisting of:

$\hat{c}_{i\sigma}^{(\dagger)}$ itinerant fermions (e.g. holes near the Mott insulator)

$\hat{b}_l^{(\dagger)}$ immobile local pairs (RVB defines them on the bonds)

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right]\end{aligned}$$

$\vec{R}_l = (\vec{r}_i + \vec{r}_j)/2$

describes a two-component system consisting of:

$\hat{c}_{i\sigma}^{(\dagger)}$ itinerant fermions (e.g. holes near the Mott insulator)

$\hat{b}_l^{(\dagger)}$ immobile local pairs (RVB defines them on the bonds)

interacting via:

Boson-Fermion model

$$\begin{aligned}\hat{H} = & \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l \\ & + \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right] \quad \vec{R}_l = (\vec{r}_i + \vec{r}_j)/2\end{aligned}$$

describes a two-component system consisting of:

$\hat{c}_{i\sigma}^{(\dagger)}$ itinerant fermions (e.g. holes near the Mott insulator)

$\hat{b}_l^{(\dagger)}$ immobile local pairs (RVB defines them on the bonds)

interacting via:

$\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.}$ (the Andreev-type scattering)

Boson-Fermion model

$$\hat{H} = \sum_{i,j,\sigma} (t_{ij} - \mu \delta_{i,j}) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_l \left(E_l^{(B)} - 2\mu \right) \hat{b}_l^\dagger \hat{b}_l$$

$$+ \sum_{i,j} g_{ij} \left[\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.} \right] \quad \vec{R}_l = (\vec{r}_i + \vec{r}_j)/2$$

describes a two-component system consisting of:

$\hat{c}_{i\sigma}^{(\dagger)}$ itinerant fermions (e.g. holes near the Mott insulator)

$\hat{b}_l^{(\dagger)}$ immobile local pairs (RVB defines them on the bonds)

interacting via:

$\hat{b}_l^\dagger \hat{c}_{i,\downarrow} \hat{c}_{j,\uparrow} + \text{h.c.}$ (the Andreev-type scattering)

Isotropic form of this model has been introduced 25 year ago by

J. Ranninger and S. Robaszkiewicz, Physica B **135**, 468 (1985).

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} = & \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^\dagger \hat{b}_{\mathbf{q}} \\ & + \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^\dagger \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} = & \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^\dagger \hat{b}_{\mathbf{q}} \\ & + \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^\dagger \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

BF scenario has been considered in the context of HTSC by

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} = & \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^{\dagger} \hat{b}_{\mathbf{q}} \\ & + \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^{\dagger} \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

BF scenario has been considered in the context of HTSC by

R. Micnas, J. Ranninger, S. Robaszkiewicz, Rev. Mod. Phys. 62, 113 (1990);

R. Friedberg and T.D. Lee, Phys. Rev. B 40, 423 (1989);

Ch.P. Enz, Phys. Rev. B 54, 3589 (1996);

V.B. Geshkenbein, L.B. Ioffe, A.I. Larkin, Phys. Rev. B 55, 3173 (1997);

...

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} = & \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^{\dagger} \hat{b}_{\mathbf{q}} \\ & + \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^{\dagger} \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

BF scenario has been considered in the context of HTSC by

R. Micnas, J. Ranninger, S. Robaszkiewicz, Rev. Mod. Phys. 62, 113 (1990);

R. Friedberg and T.D. Lee, Phys. Rev. B 40, 423 (1989);

Ch.P. Enz, Phys. Rev. B 54, 3589 (1996);

V.B. Geshkenbein, L.B. Ioffe, A.I. Larkin, Phys. Rev. B 55, 3173 (1997);

...

and ultracold atoms interacting with the Feshbach resonance by

Boson-Fermion model

[in the momentum space]

$$\begin{aligned}\hat{H} = & \sum_{\mathbf{k}\sigma} (\varepsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + \sum_{\mathbf{q}} \left(E^{(B)} - 2\mu \right) \hat{b}_{\mathbf{q}}^{\dagger} \hat{b}_{\mathbf{q}} \\ & + \frac{1}{\sqrt{N}} \sum_{\mathbf{k},\mathbf{q}} g_{\mathbf{k},\mathbf{q}} \left[\hat{b}_{\mathbf{q}}^{\dagger} \hat{c}_{\mathbf{k},\downarrow} \hat{c}_{\mathbf{q}-\mathbf{k},\uparrow} + \text{h.c.} \right]\end{aligned}$$

BF scenario has been considered in the context of HTSC by

R. Micnas, J. Ranninger, S. Robaszkiewicz, Rev. Mod. Phys. 62, 113 (1990);
R. Friedberg and T.D. Lee, Phys. Rev. B 40, 423 (1989);
Ch.P. Enz, Phys. Rev. B 54, 3589 (1996);
V.B. Geshkenbein, L.B. Ioffe, A.I. Larkin, Phys. Rev. B 55, 3173 (1997);
...

and ultracold atoms interacting with the Feshbach resonance by

M. Holland *et al*, Phys. Rev. Lett. 87, 120406 (2001);
Y. Ohashi, A. Griffin, Phys. Rev. Lett. 89, 130402 (2002);
R.A. Duine and H.T.C. Stoof, Phys. Rep. 396, 115 (2004);
Q. Chen, J. Stajic, S. Tan and K. Levin, Phys. Rep. 412, 1 (2005);
...

Outline of the procedure

Outline of the procedure

For studying the many-body effects we construct the sequence

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

F. Wegner (1994); K.G. Wilson (1994) - inventors of this RG-like scheme

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

F. Wegner (1994); K.G. Wilson (1994) - inventors of this RG-like scheme

Hamiltonian at $l = 0$

$$\hat{H}_F + \hat{H}_B + \hat{V}_{BF}$$

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

F. Wegner (1994); K.G. Wilson (1994) - inventors of this RG-like scheme

Hamiltonian at $0 < l < \infty$

$$\hat{H}_F(l) + \hat{H}_B(l) + \hat{V}_{BF}(l)$$

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

F. Wegner (1994); K.G. Wilson (1994) - inventors of this RG-like scheme

Hamiltonian at $l = \infty$

$$\hat{H}_F(\infty) + \hat{H}_B(\infty) + 0$$

Outline of the procedure

For studying the many-body effects we construct the sequence of unitary transformations

$$\hat{H} \longrightarrow \hat{H}(l_1) \longrightarrow \hat{H}(l_2) \longrightarrow \dots \longrightarrow \hat{H}(\infty)$$

decoupling the boson from fermion degrees of freedom.

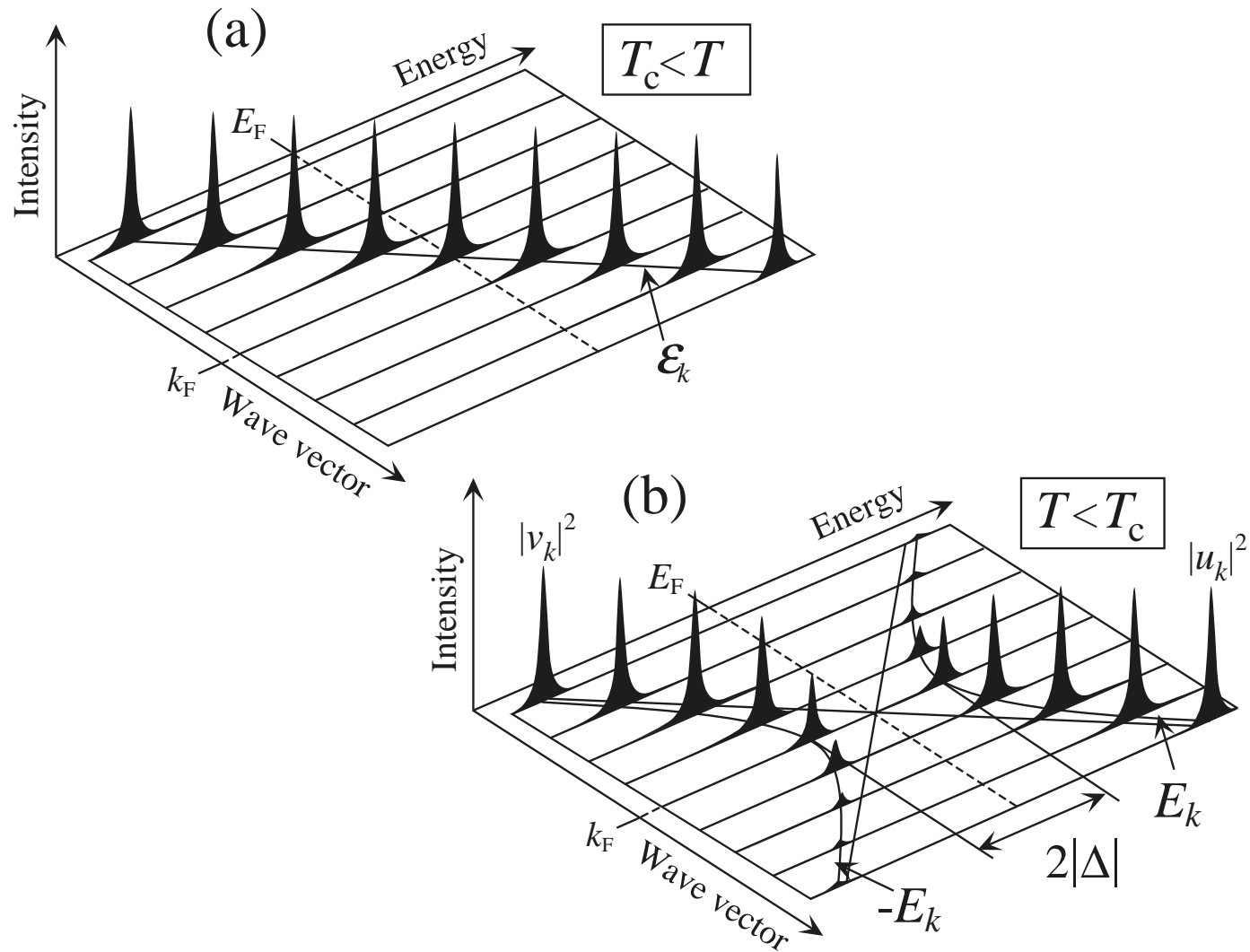
F. Wegner (1994); K.G. Wilson (1994) - inventors of this RG-like scheme

Hamiltonian at $l = \infty$

$$\hat{H}_F(\infty) + \hat{H}_B(\infty) + 0$$

*T. Domański and J. Ranninger, Phys. Rev. B **63**, 134505 (2001).*

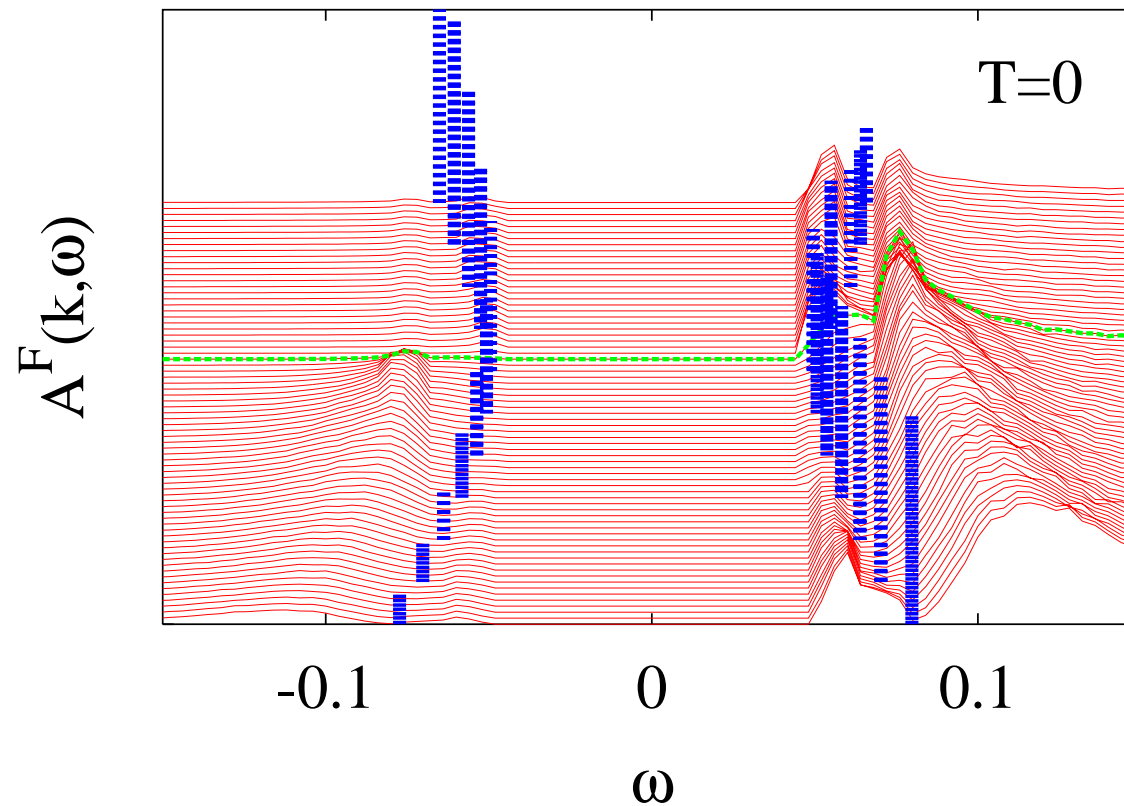
III. Selected results



Single particle spectra of conventional superconductors consist of the Bogoliubov branches separated around E_F by $2\Delta_{sc}$ (the fluctuation effects are neglected).

Effective spectrum: BF model

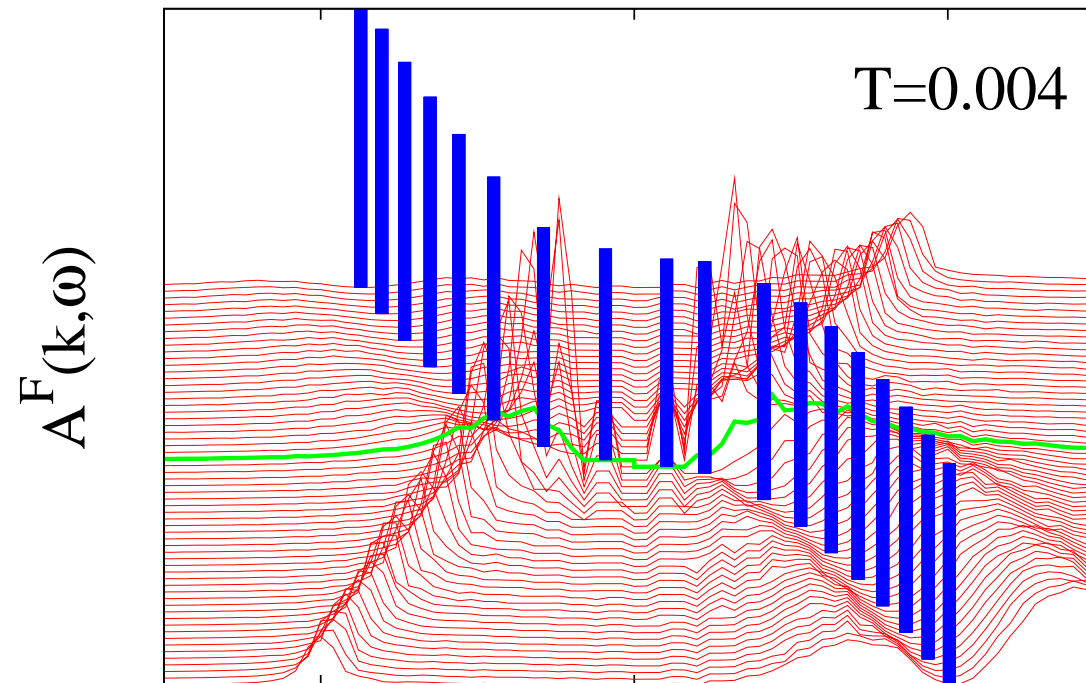
$$T < T_c$$



*T. Domański and J. Ranninger, Phys. Rev. Lett. **91**, 255301 (2003).*

Effective spectrum: BF model

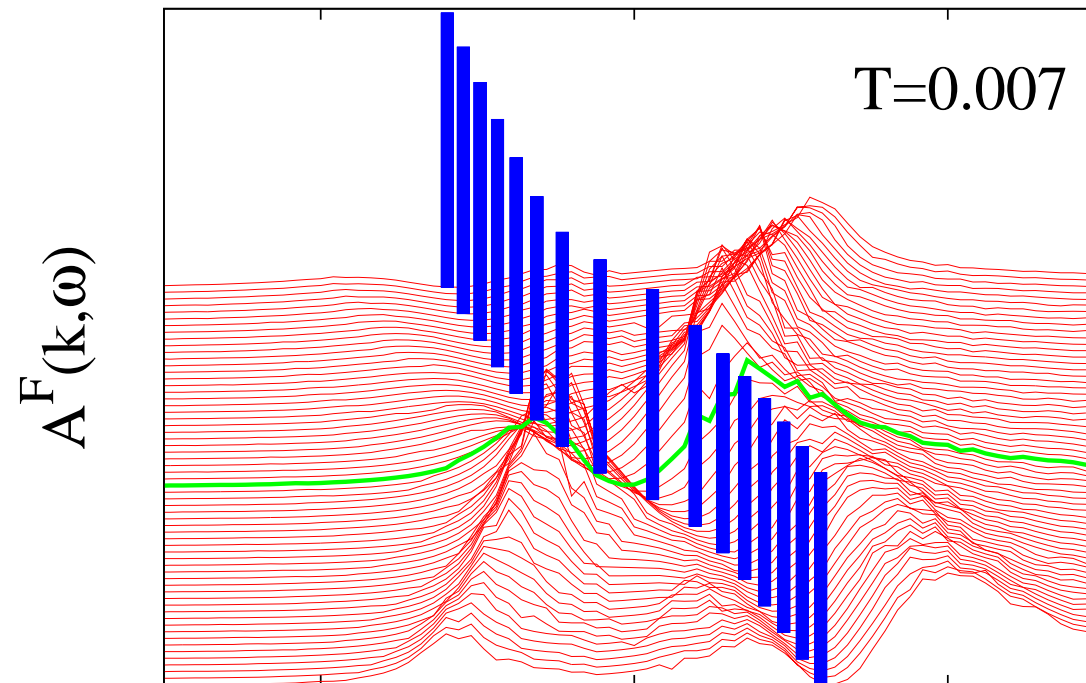
$$T_c < T < T^*$$



*T. Domański and J. Ranninger, Phys. Rev. Lett. **91**, 255301 (2003).*

Effective spectrum: BF model

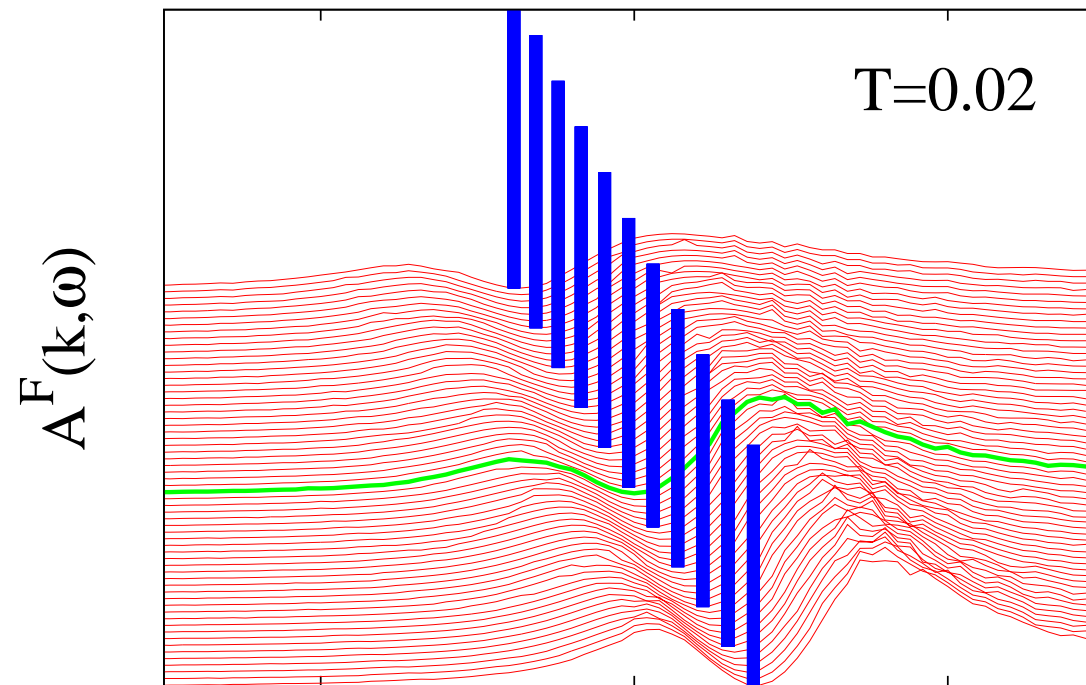
$$T_c < T < T^*$$



*T. Domański and J. Ranninger, Phys. Rev. Lett. **91**, 255301 (2003).*

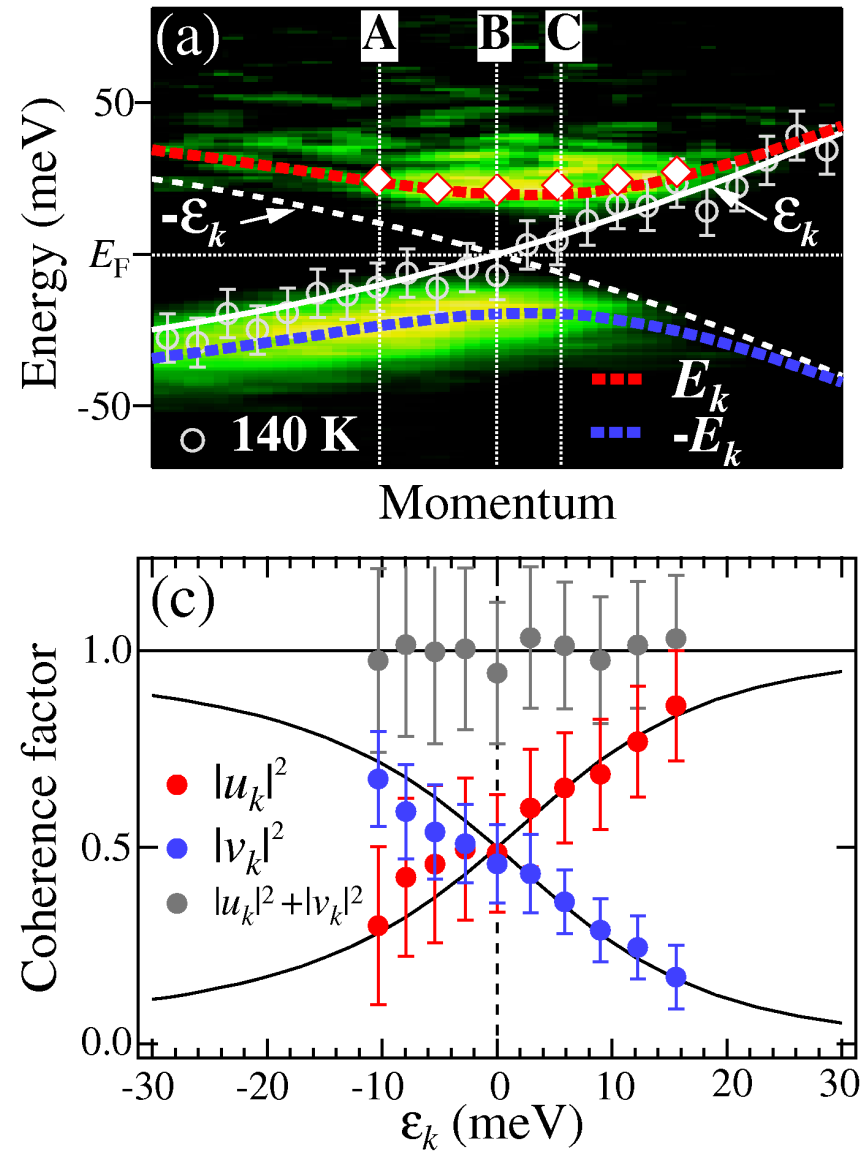
Effective spectrum: BF model

$$T > T^*$$



*T. Domański and J. Ranninger, Phys. Rev. Lett. **91**, 255301 (2003).*

Experimental data for $T < T_c$



H. Matsui, T. Sato, and T. Takahashi et al, *Phys. Rev. Lett.* **90**, 217002 (2003).

Date: Tue, 27 Feb 2007 19:05:55 +0900

From: Hiroaki Matsui <h.matsui@arpes.phys.tohoku.ac.jp>

To: Tadeusz Domanski <doman@kft.umcs.lublin.pl>

Dear Dr. Domanski,

...

We completely agree with you on that detecting the normal state BQP in the UD cuprates has a huge potential impact on the pseudogap problem. As you know, this kind of measurement is not very easy because the ARPES peak is broad in UD at anti-node and high-temperature. We do not have the data at present, but we are trying to realize such an experiment by selecting the conditions.

Thank you very much for contacting us.

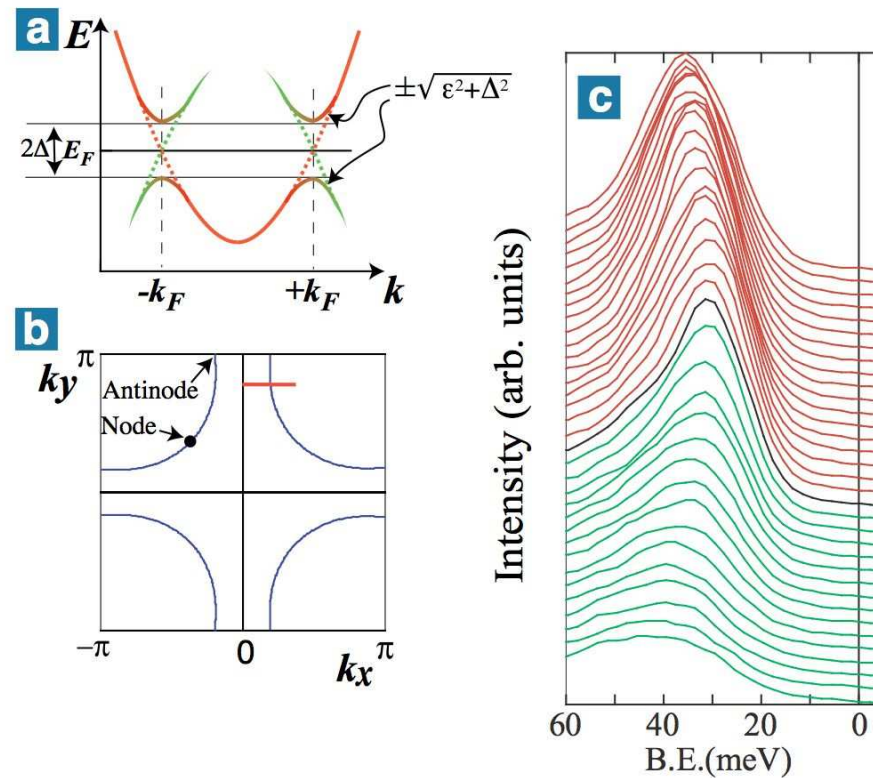
Sincerely yours,

H. Matsui

Evidence for Bogoliubov QPs above T_c

Evidence for Bogoliubov QPs above T_c

J. Campuzano group (Chicago, USA)

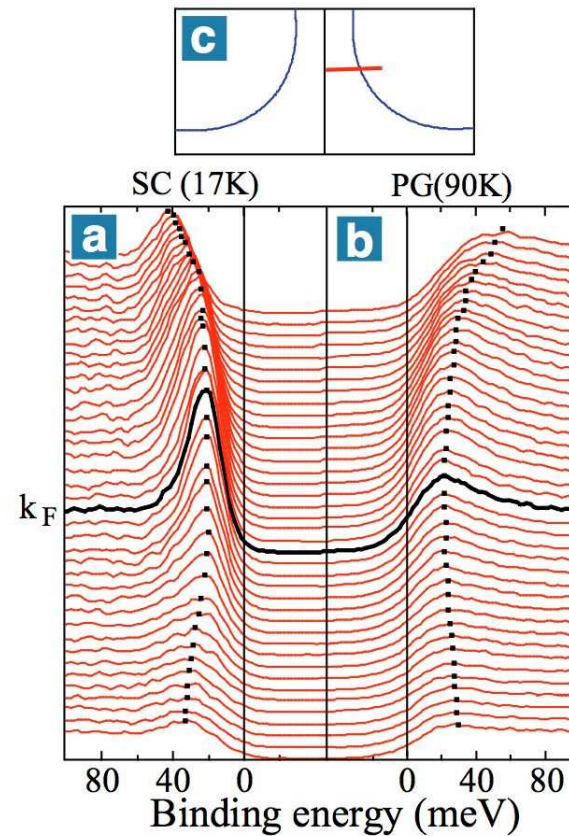


Results for: $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$

A. Kanigel et al, *Phys. Rev. Lett.* **101**, 137002 (2008).

Evidence for Bogoliubov QPs above T_c

J. Campuzano group (Chicago, USA)

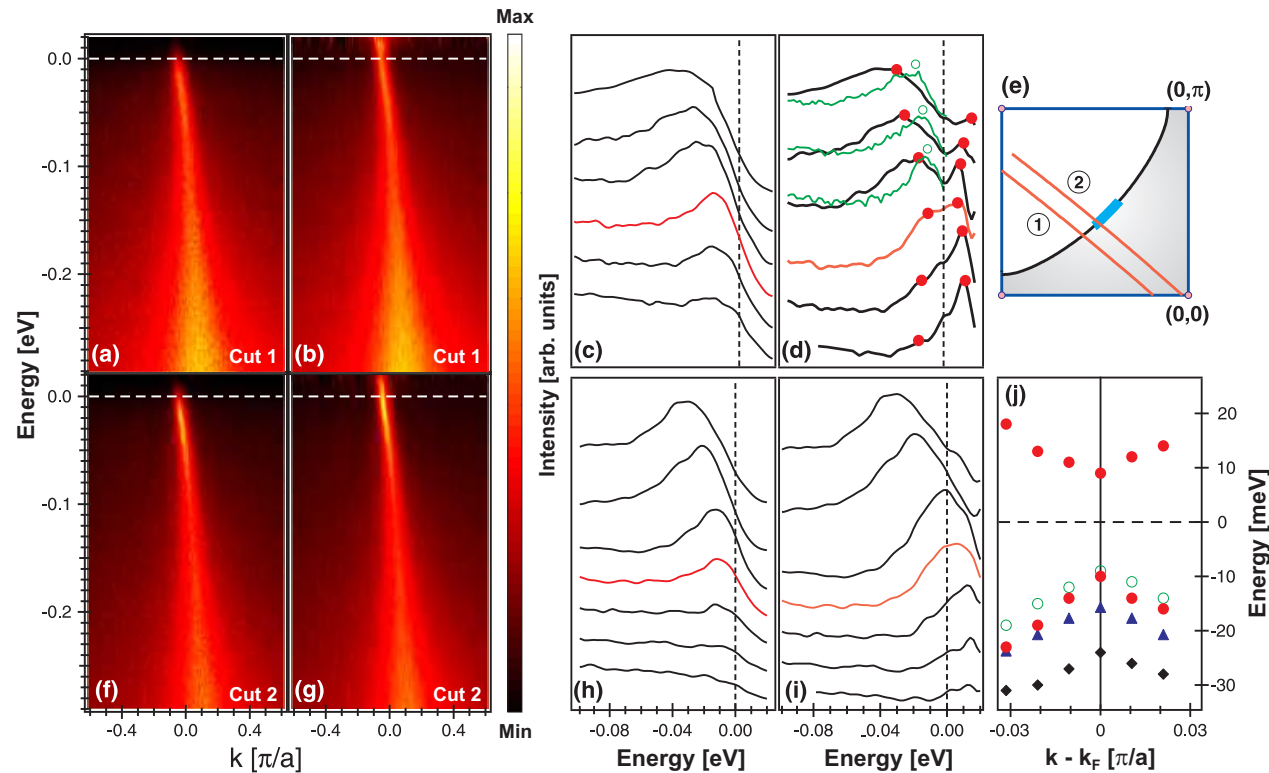


Results for: $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$

A. Kanigel et al, *Phys. Rev. Lett.* **101**, 137002 (2008).

Evidence for Bogoliubov QPs above T_c

PSI group (Villigen, Switzerland)

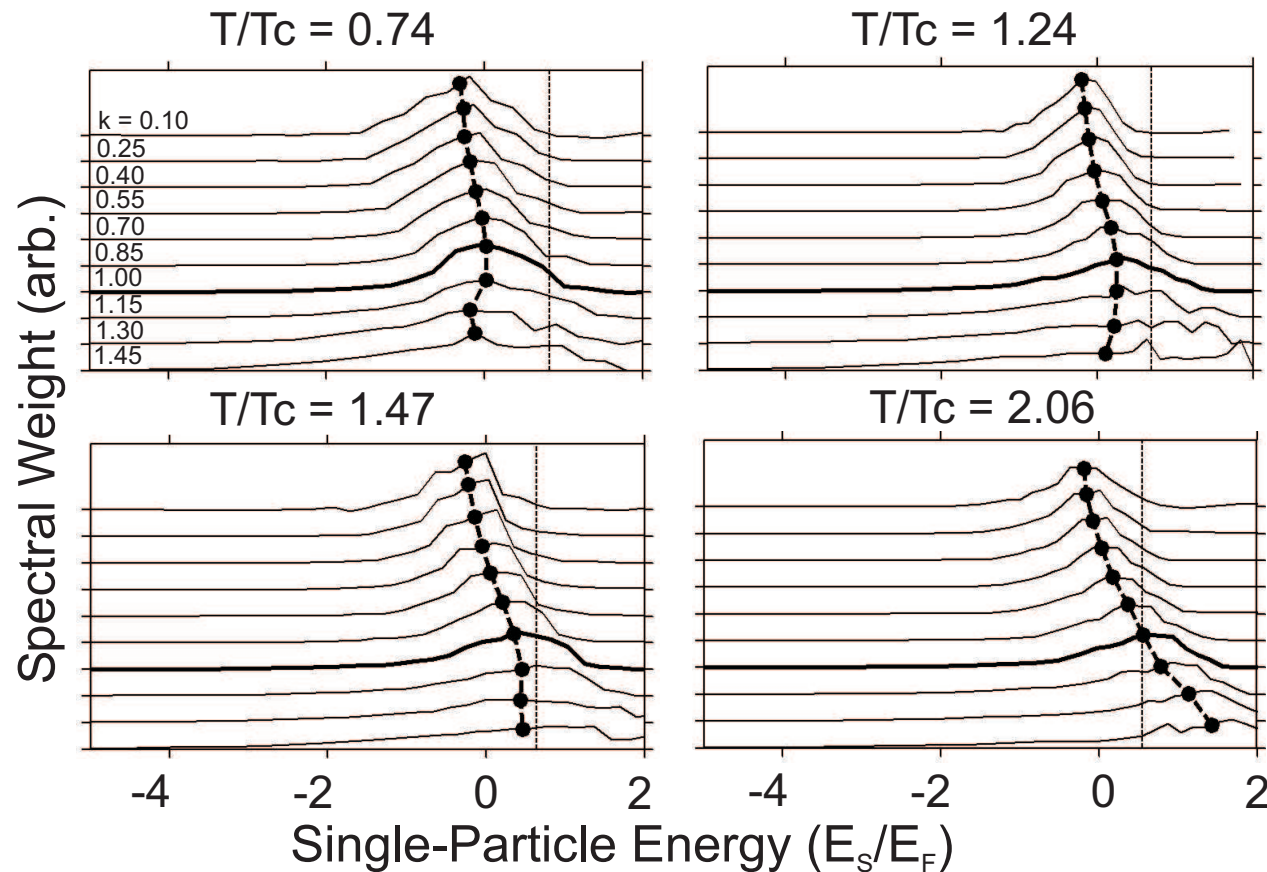


Results for: $\text{La}_{1.895}\text{Sr}_{0.105}\text{CuO}_4$

M. Shi et al, Eur. Phys. Lett. 88, 27008 (2009).

Evidence for Bogoliubov QPs above T_c

D. Jin group (Boulder, USA)



Results for: ultracold ^{40}K atoms

J.P. Gaebler et al, Nature Phys. 6, 569 (2010).

Date: Mon, 31 Mar 2008 13:57:05 +0300

From: Amit Kanigel <amitk@physics.technion.ac.il>

To: Tadeusz Domanski <doman@kft.umcs.lublin.pl>

Dear Prof. Domanski,

I'm really happy for your remarks. I read your paper (the PRL) and indeed found it very interesting. I must apologize and admit that I was not aware of the paper. While writing my paper I looked quite intensively for theoretical models predicting BG-like dispersion and for some reason I missed your work. Although the paper was already submitted I hope I'll have the chance to put in a reference to your work before publication.

If you have no objection, after I'll read the longer paper I might have few questions for you regarding the Boson-Fermion model.

Best regards,

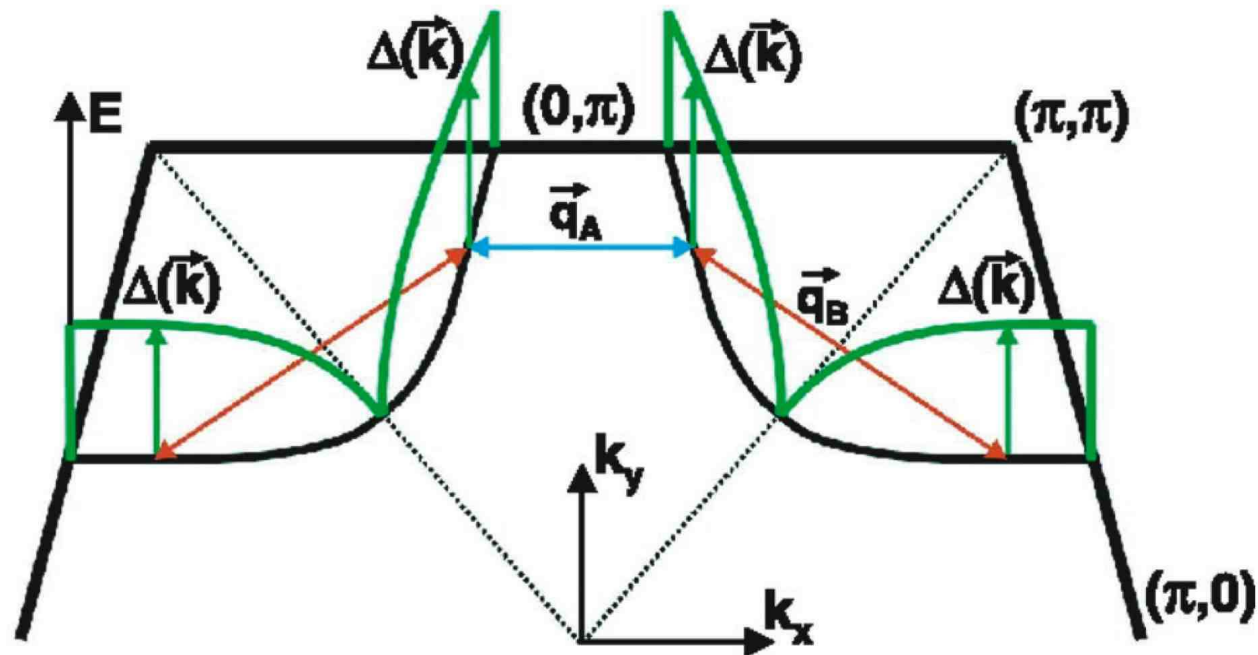
Amit

Angular dependence of the gap

Various experiments indicate that below T_c the gap in the cuprate superconductors has d -wave symmetry.

Angular dependence of the gap

Various experiments indicate that below T_c the gap in the cuprate superconductors has d -wave symmetry.



J.E. Hoffman et al, Science 297, 1148 (2002).

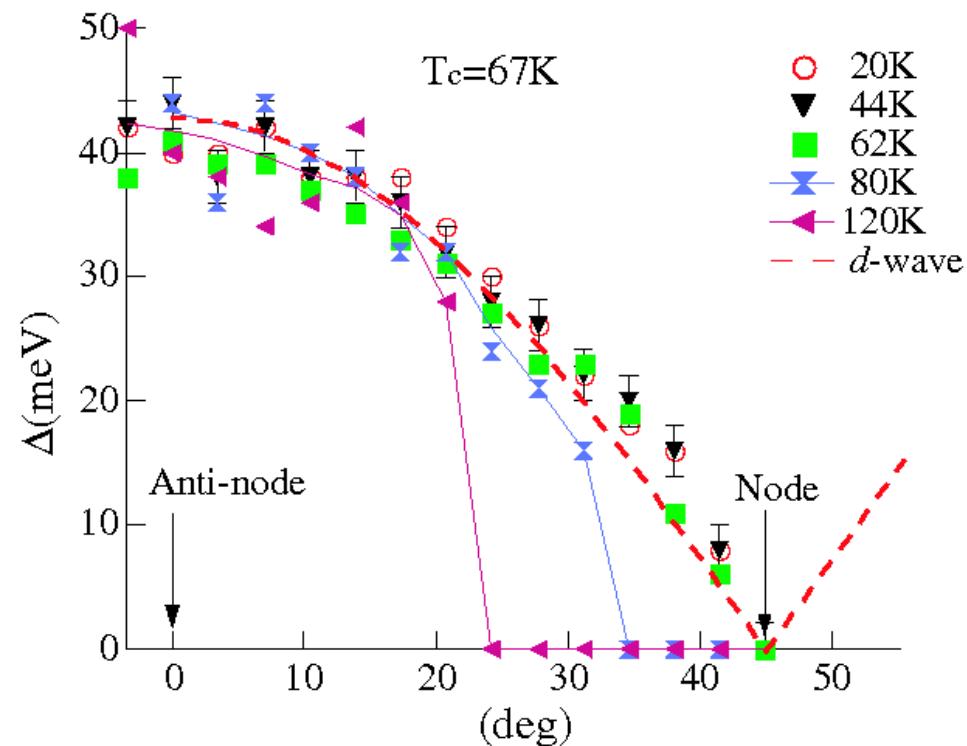
Energy gap above T_c

Energy gap above T_c

In a normal state the energy gap does survive above T_c .
Upon increasing temperature it gradually closes, starting from the nodal area where *the Fermi arcs* emerge.

Energy gap above T_c

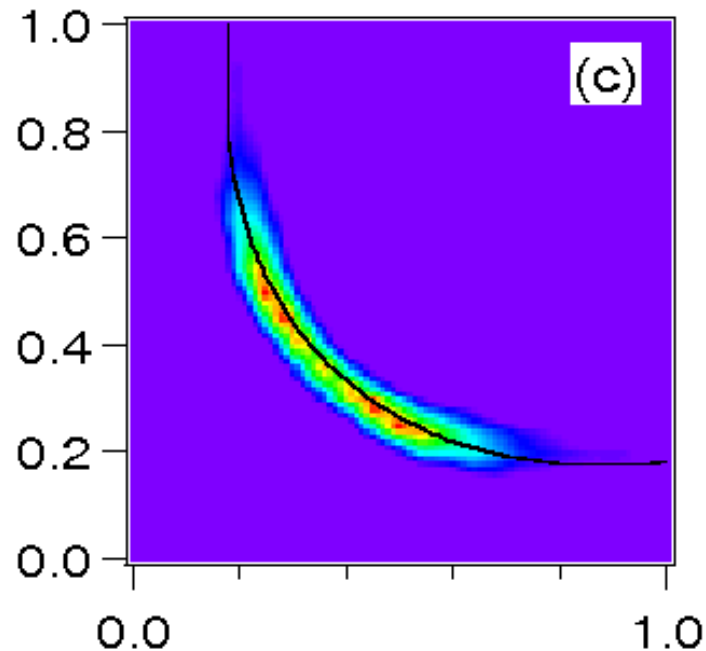
In a normal state the energy gap does survive above T_c . Upon increasing temperature it gradually closes, starting from the nodal area where *the Fermi arcs* emerge.



A. Kanigel et al, *Phys. Rev. Lett.* **99**, 157001 (2007).

Energy gap above T_c

In a normal state the energy gap does survive above T_c . Upon increasing temperature it gradually closes, starting from the nodal area where *the Fermi arcs* emerge.



Pieces of the Fermi surface near the antinodal area are missing.

"Death of a Fermi surface" K. McElroy, *Nature Physics* 2, 441 (2006) .

B-F study of the angular variation

B-F study of the angular variation

We have examined the effect of anisotropic B-F coupling

$g_{\vec{k}} = g [\cos(k_x) - \cos(k_y)]$ using the realistic dispersion

$$\varepsilon_{\vec{k}} = -2t [\cos(k_x) + \cos(k_y)] - 4t' \cos(k_x) \cos(k_y).$$

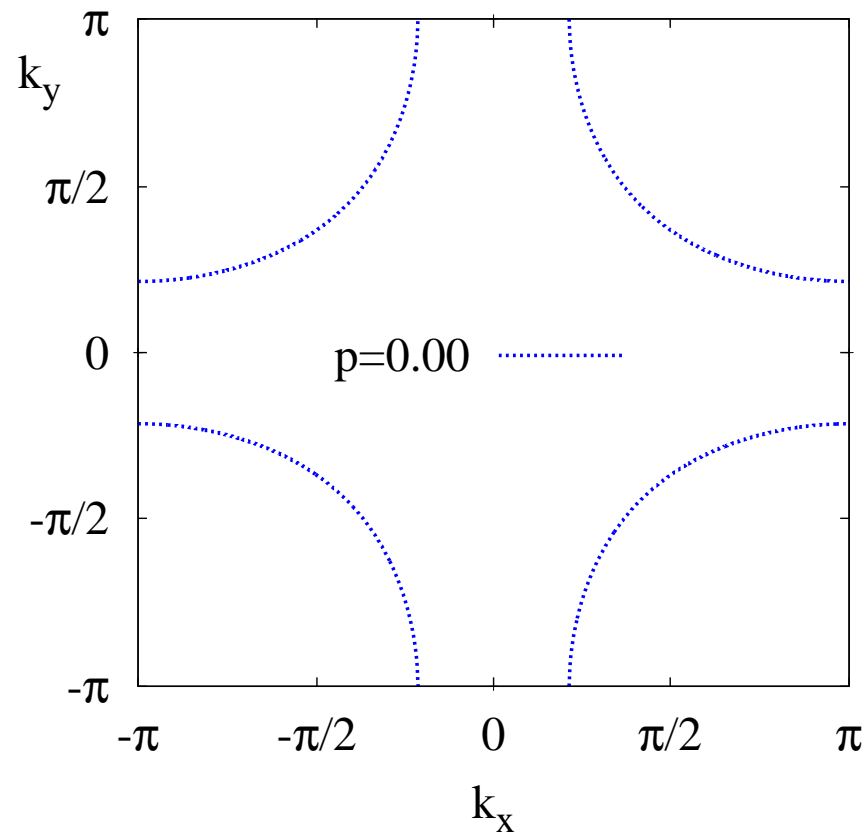
B-F study of the angular variation

We have examined the effect of anisotropic B-F coupling

$g_{\vec{k}} = g [\cos(k_x) - \cos(k_y)]$ using the realistic dispersion

$$\varepsilon_{\vec{k}} = -2t [\cos(k_x) + \cos(k_y)] - 4t' \cos(k_x) \cos(k_y).$$

$$t'/t = -0.4$$



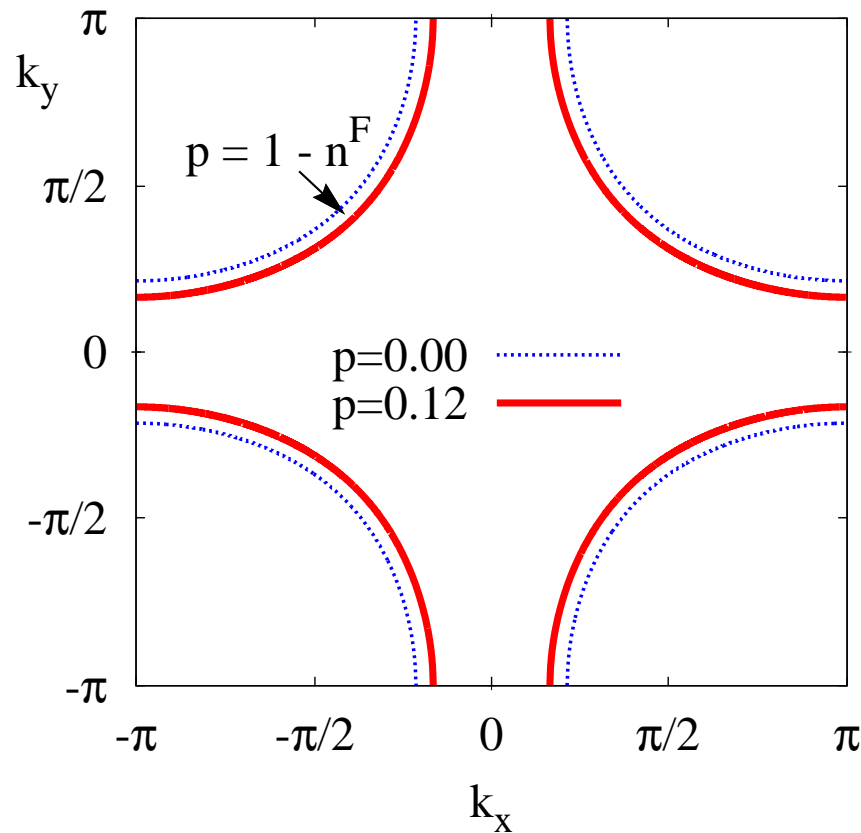
B-F study of the angular variation

We have examined the effect of anisotropic B-F coupling

$g_{\vec{k}} = g [\cos(k_x) - \cos(k_y)]$ using the realistic dispersion

$$\varepsilon_{\vec{k}} = -2t [\cos(k_x) + \cos(k_y)] - 4t' \cos(k_x) \cos(k_y).$$

$$t'/t = -0.4$$



$\varepsilon_{\vec{k}_F}$ topology corresponding to the hole doping $p = 0.12$

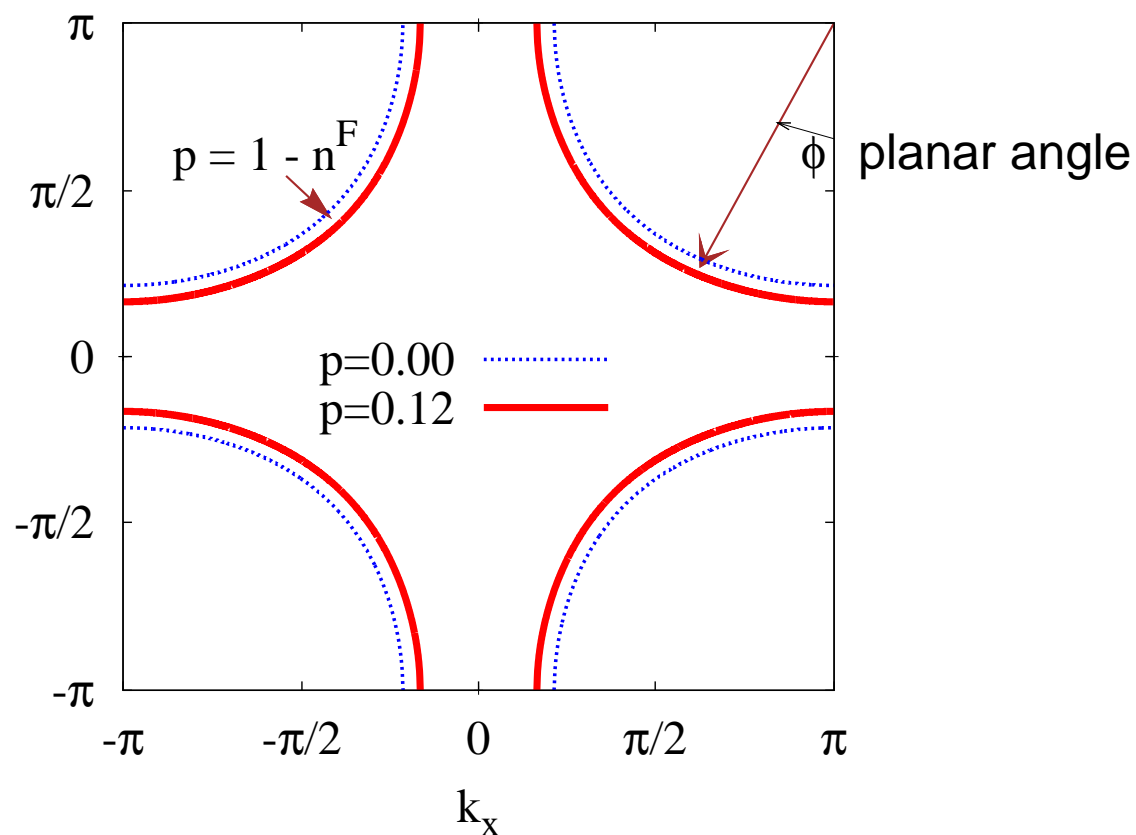
B-F study of the angular variation

We have examined the effect of anisotropic B-F coupling

$g_{\vec{k}} = g [\cos(k_x) - \cos(k_y)]$ using the realistic dispersion

$$\varepsilon_{\vec{k}} = -2t [\cos(k_x) + \cos(k_y)] - 4t' \cos(k_x) \cos(k_y).$$

$$t'/t = -0.4$$



$\varepsilon_{\vec{k}_F}$ topology corresponding to the hole doping $p = 0.12$

Emergence of the Fermi arcs

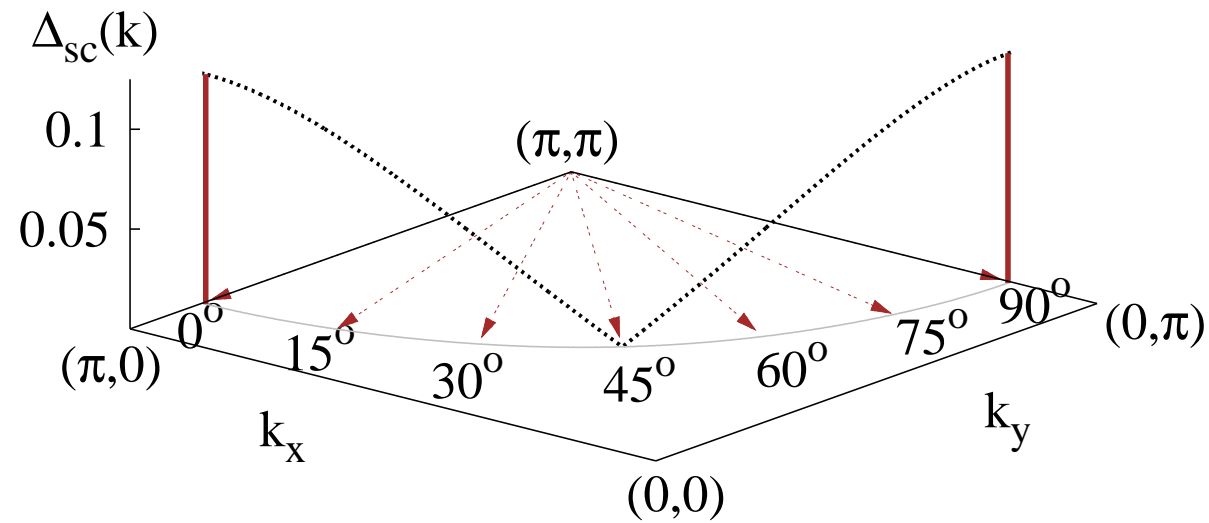
Emergence of the Fermi arcs

The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$

Emergence of the Fermi arcs

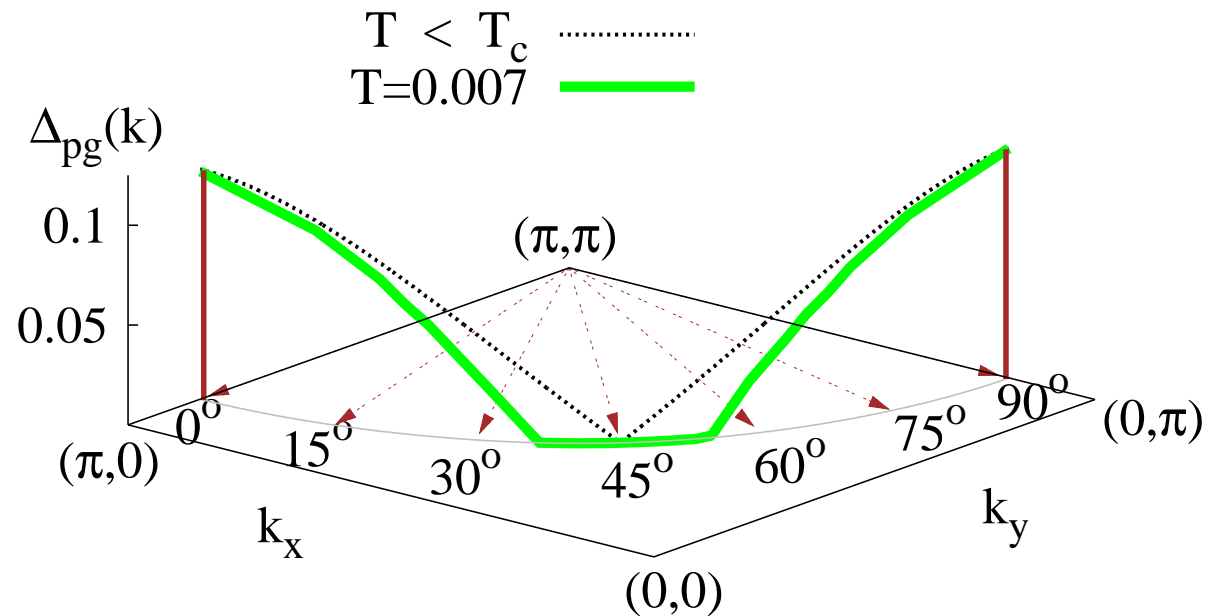
The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$

$T < T_c$



Emergence of the Fermi arcs

The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$

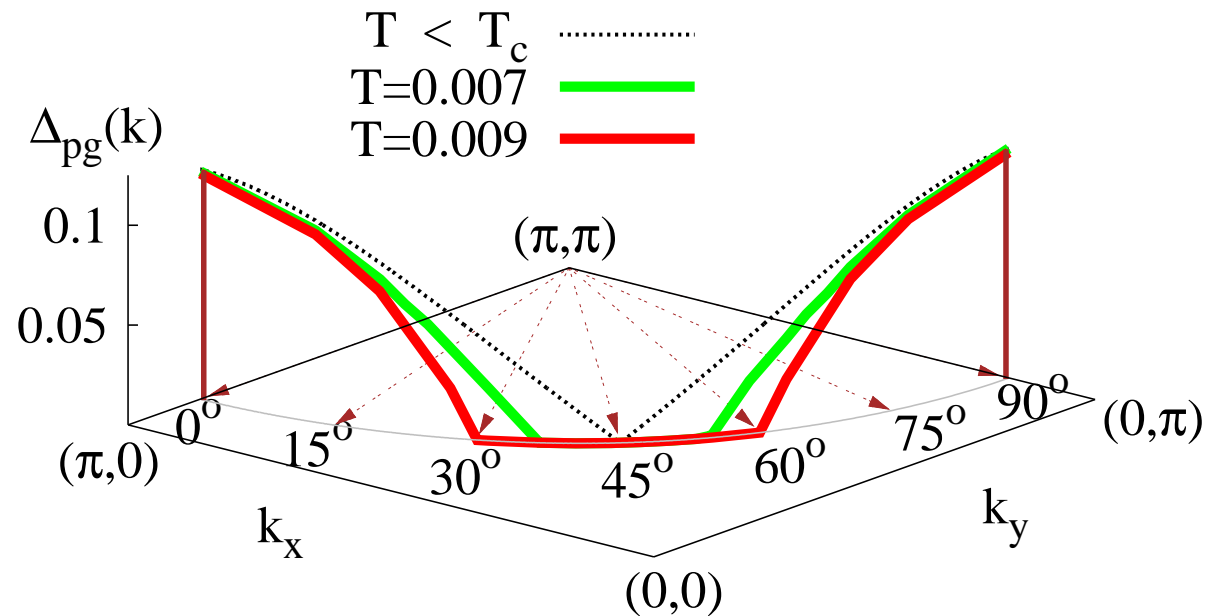


By increasing temperature:

– $\Delta_{pg}(\vec{k})$ is almost unaffected in the antinodal areas,

Emergence of the Fermi arcs

The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$

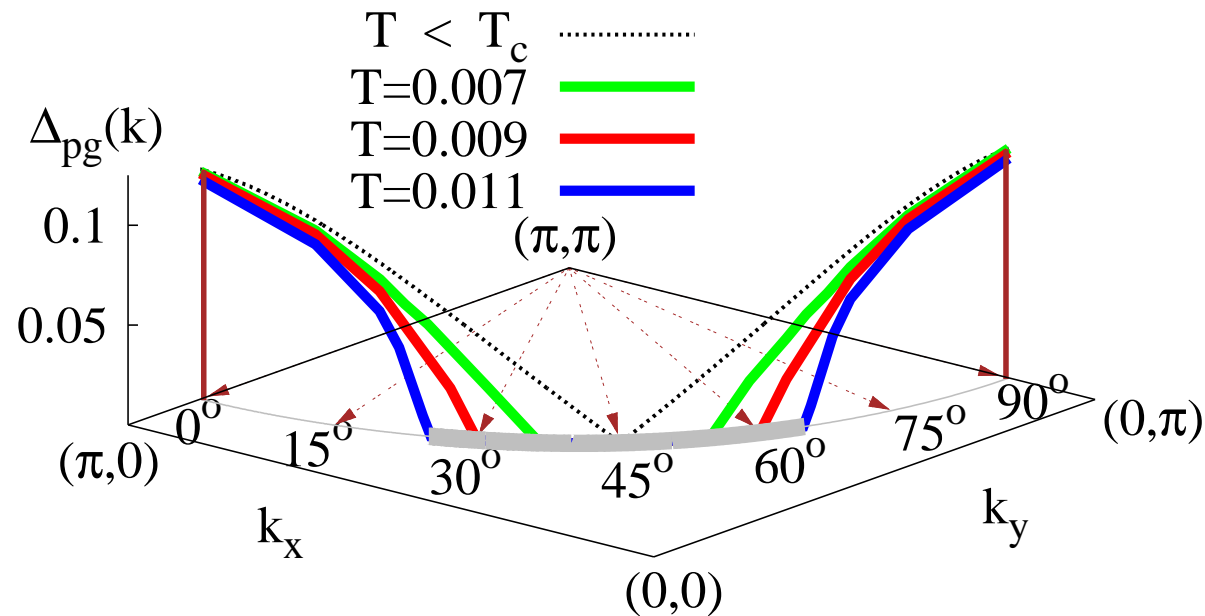


By increasing temperature:

- $\Delta_{pg}(\vec{k})$ is almost unaffected in the antinodal areas,
- Fermi surface gradually rebuilds near the nodal parts,

Emergence of the Fermi arcs

The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$

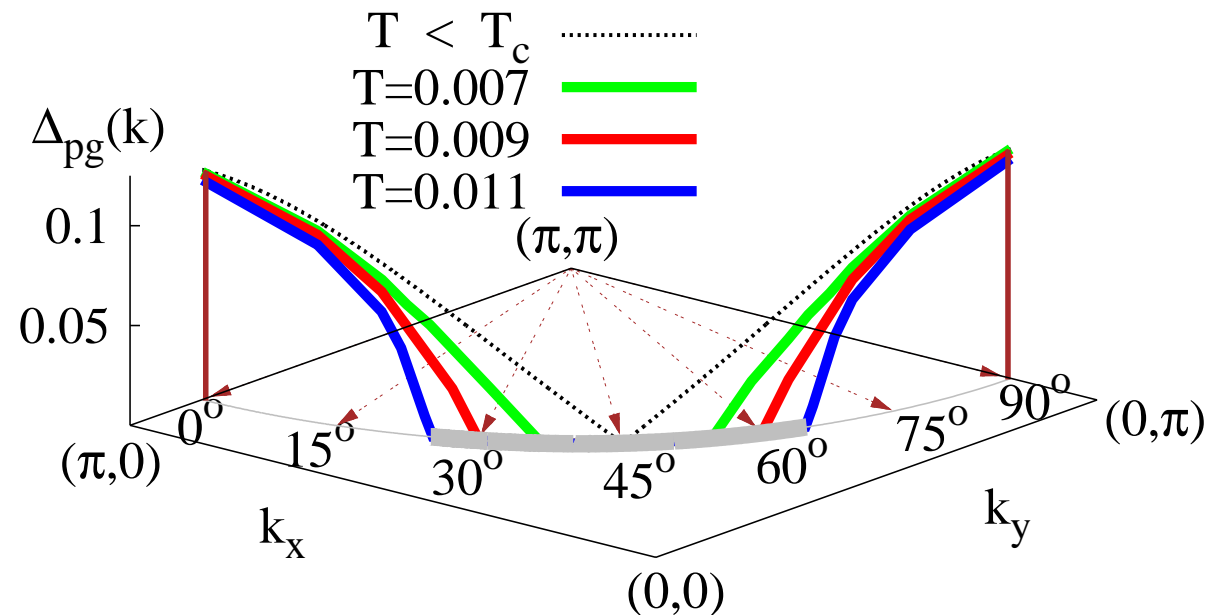


By increasing temperature:

- $\Delta_{pg}(\vec{k})$ is almost unaffected in the antinodal areas,
- Fermi surface gradually rebuilds near the nodal parts,
- length of the Fermi arc scales linearly with $T - T_c$.

Emergence of the Fermi arcs

The energy gap below T_c is $\Delta_{sc}(\vec{k}) = \Delta [\cos k_x - \cos k_y]$



By increasing temperature:

- $\Delta_{pg}(\vec{k})$ is almost unaffected in the antinodal areas,
- Fermi surface gradually rebuilds near the nodal parts,
- length of the Fermi arc scales linearly with $T - T_c$.

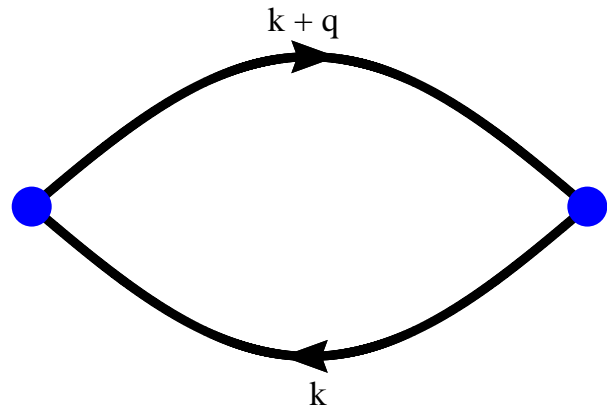
Diamagnetic response above T_c

Diamagnetic response above T_c

Main contributions to the current-current response function:

Diamagnetic response above T_c

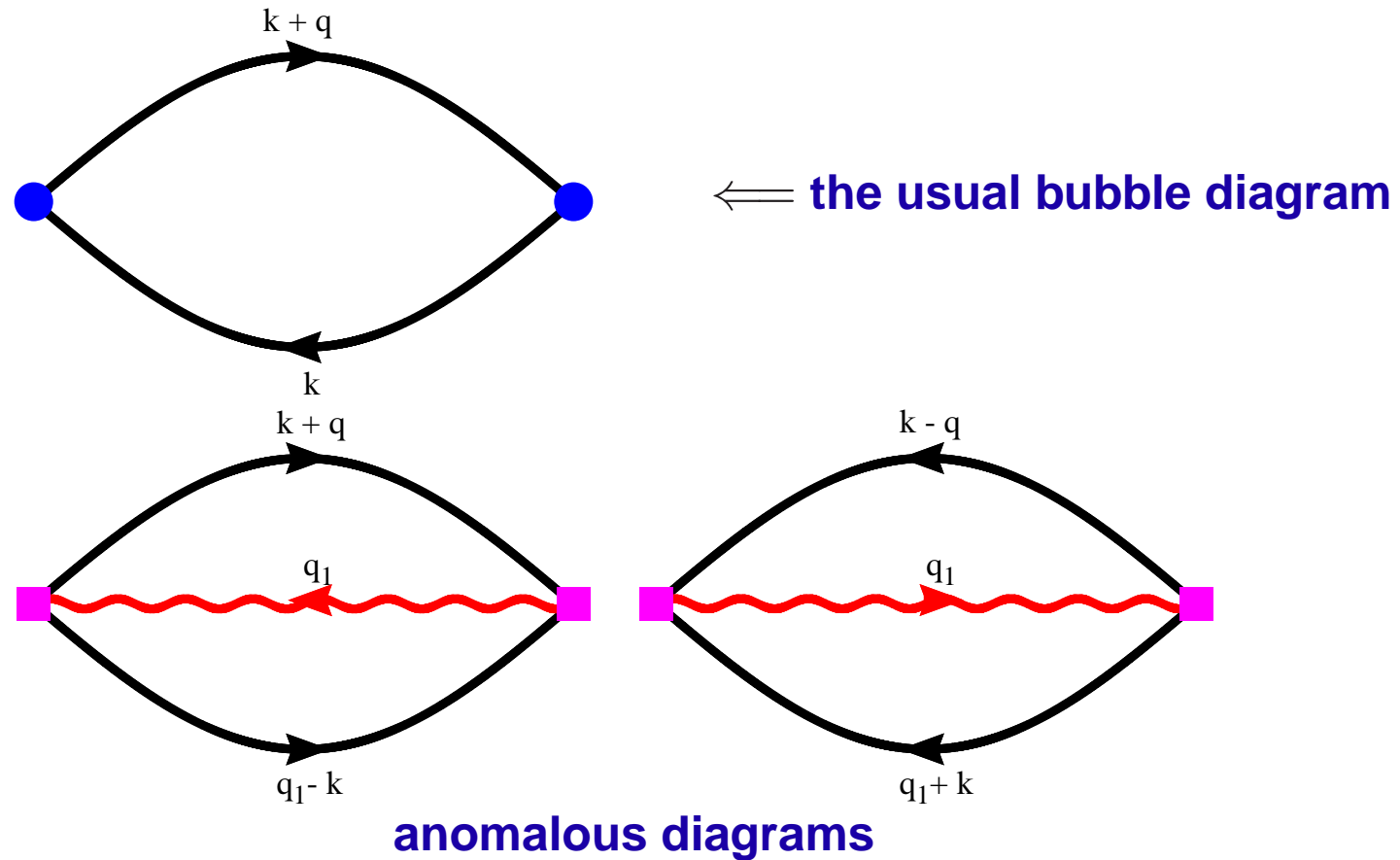
Main contributions to the current-current response function:



\Leftarrow the usual bubble diagram

Diamagnetic response above T_c

Main contributions to the current-current response function:



Each **vertex** has to be determined from the flow equations.

T. Domanski and J. Ranninger, (2010).

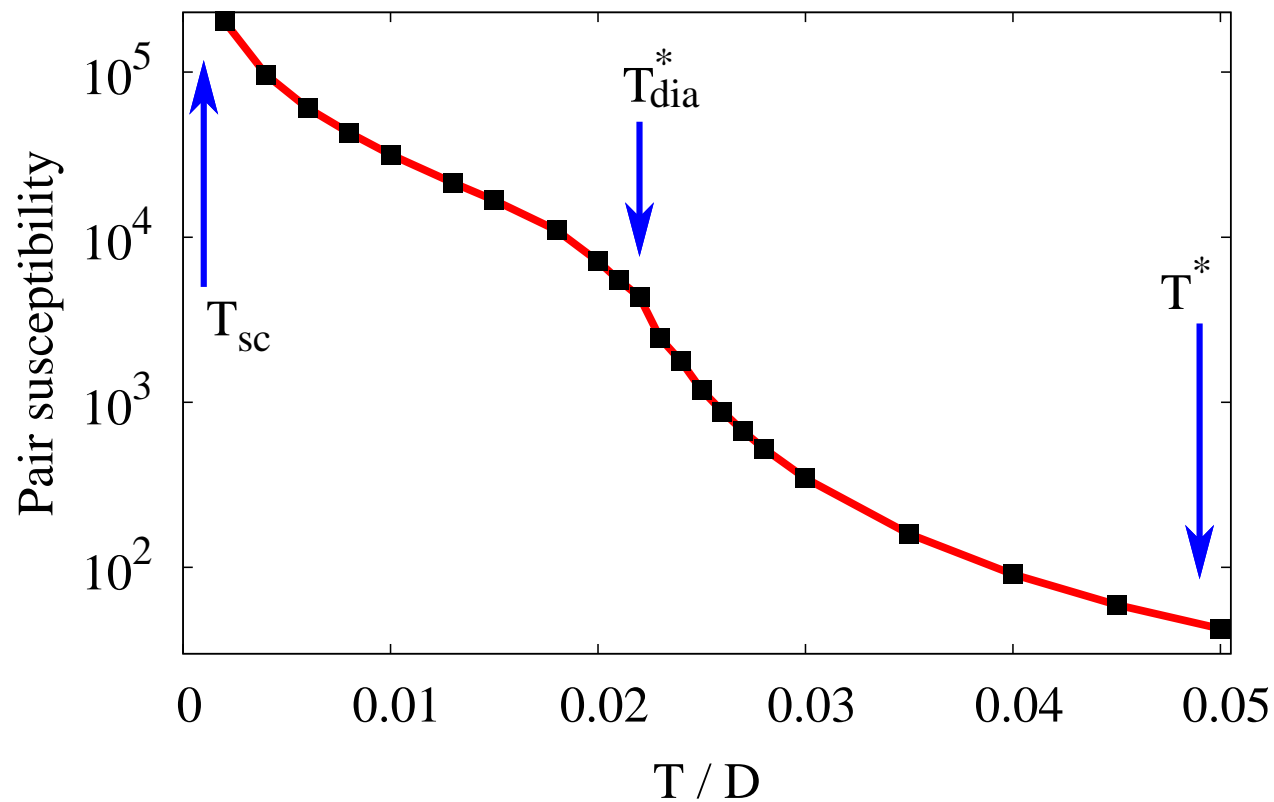
Diamagnetic response above T_c

Diamagnetic response above T_c

Residual diamagnetism appears together with enhancement of the pairing susceptibility well above T_c but safely below T^*

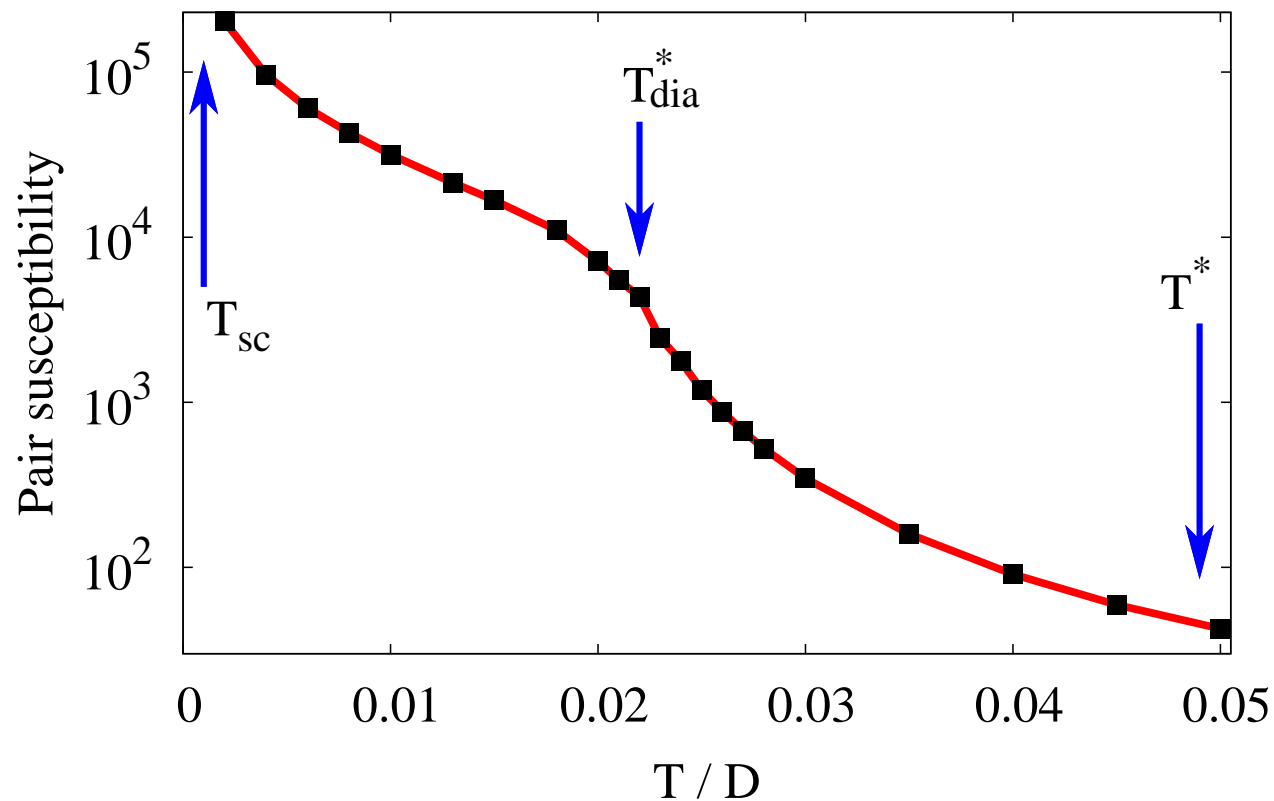
Diamagnetic response above T_c

Residual diamagnetism appears together with enhancement of the pairing susceptibility well above T_c but safely below T^*



Diamagnetic response above T_c

Residual diamagnetism appears together with enhancement of the pairing susceptibility well above T_c but safely below T^*



T. Domanski and J. Ranninger, (2010).

Conclusions

Conclusions

- 1) There is evidence for the pre-formed pairs existing above T_c

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

2) Precursor of superconductivity is seen there

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

2) Precursor of superconductivity is seen there

⇒ by the Bogoliubov-type quasiparticles

/ revealed by ARPES and FT-STM spectroscopies /

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

2) Precursor of superconductivity is seen there

⇒ by the Bogoliubov-type quasiparticles

/ revealed by ARPES and FT-STM spectroscopies /

⇒ the residual diamagnetic response

/ indicated by the torque magnetometry /

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

2) Precursor of superconductivity is seen there

⇒ by the Bogoliubov-type quasiparticles

/ revealed by ARPES and FT-STM spectroscopies /

⇒ the residual diamagnetic response

/ indicated by the torque magnetometry /
etc.

Conclusions

1) There is evidence for the pre-formed pairs existing above T_c

⇒ in the HTSC materials

/ underdoped samples /

⇒ and ultracold fermion atoms

/ nearby the Feshbach resonance /

2) Precursor of superconductivity is seen there

⇒ by the Bogoliubov-type quasiparticles

/ revealed by ARPES and FT-STM spectroscopies /

⇒ the residual diamagnetic response

/ indicated by the torque magnetometry /
etc.

<http://kft.umcs.lublin.pl/doman/lectures>