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# **Majorana-type quasiparticles in nanoscopic systems**

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**/ UMCS, Lublin /**

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*Majorana returns, F. Wilczek, Nature Physics 5, 614 (2009).*

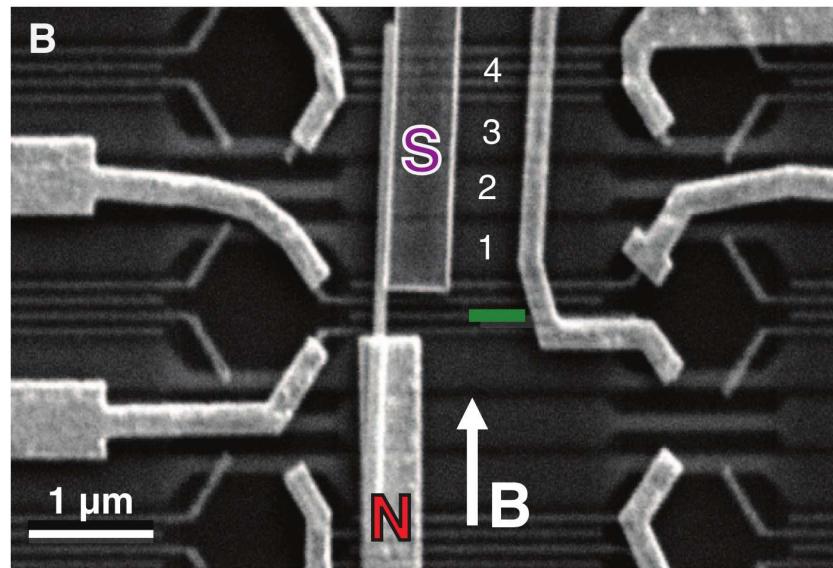
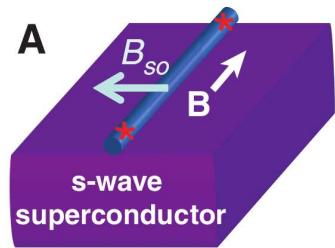
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InSb nanowire between a metal (gold) and a superconductor (Nb-Ti-N)

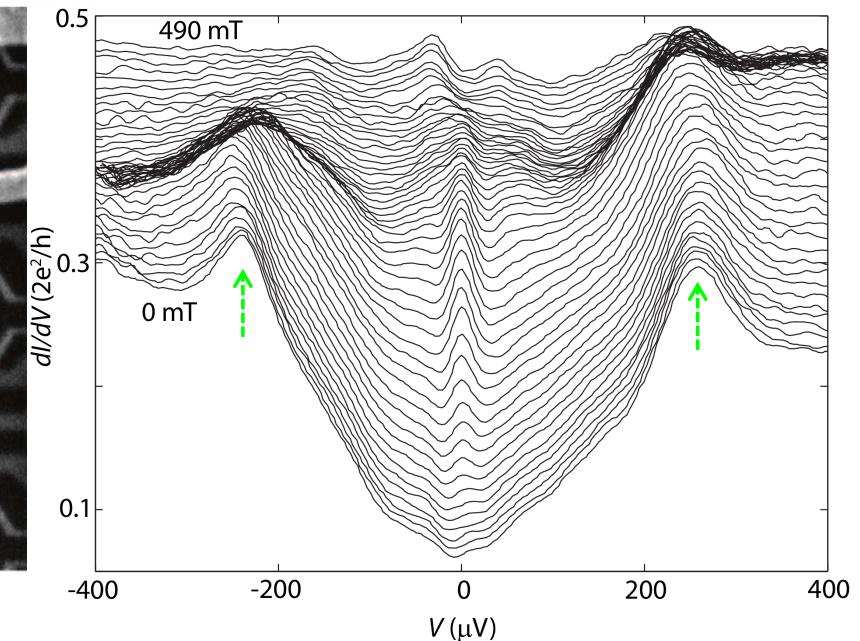
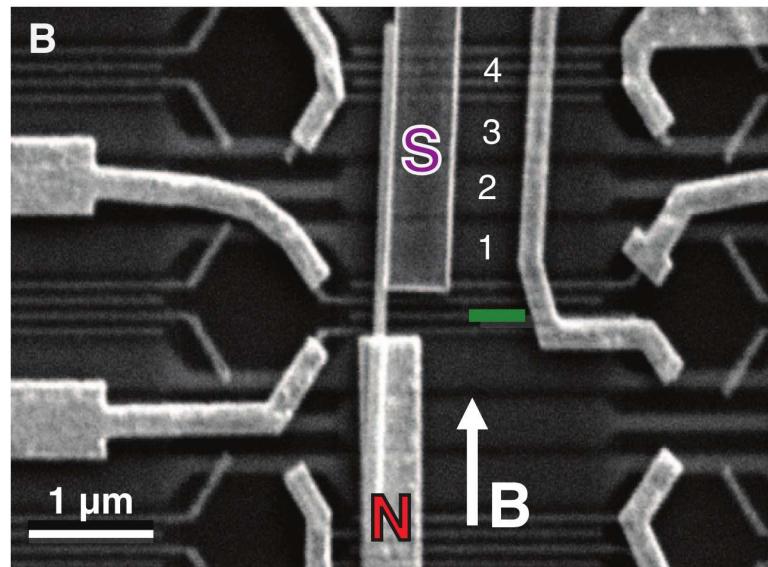
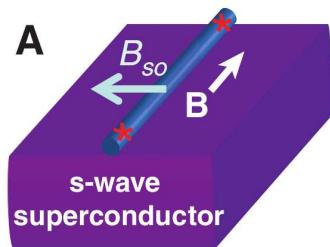


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⇒ a zero-bias enhancement due to Majorana state

V. Mourik, ..., and L.P. Kouwenhoven, Science 336, 1003 (2012).

/ Kavli Institute of Nanoscience, Delft Univ., Netherlands /

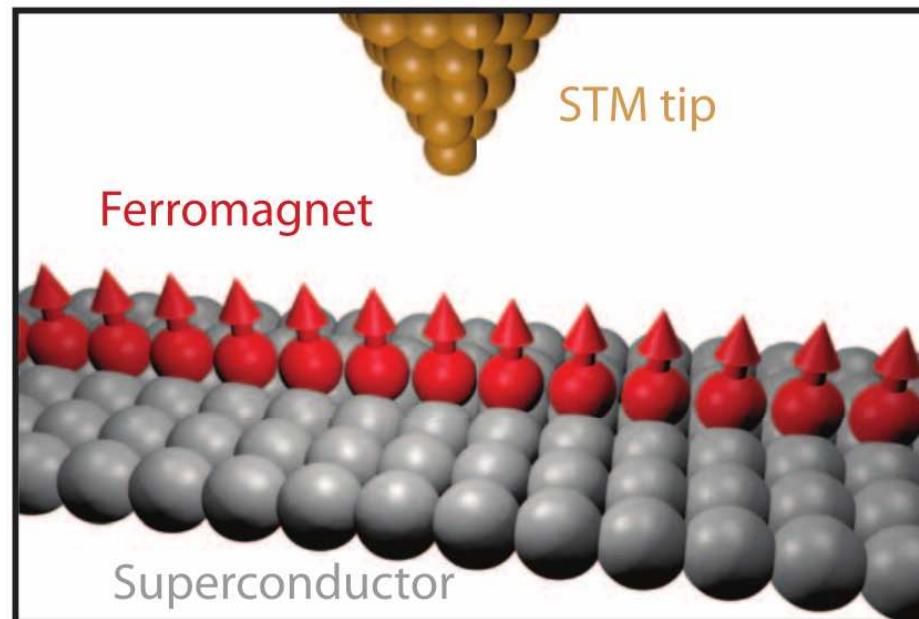
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A chain of iron atoms deposited on a surface of superconducting lead

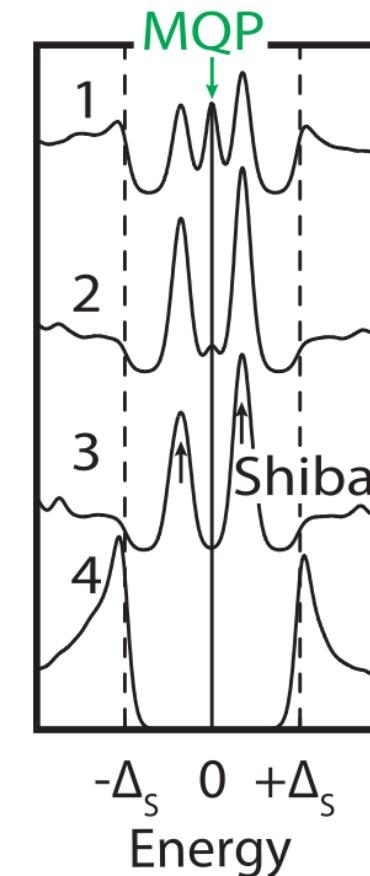
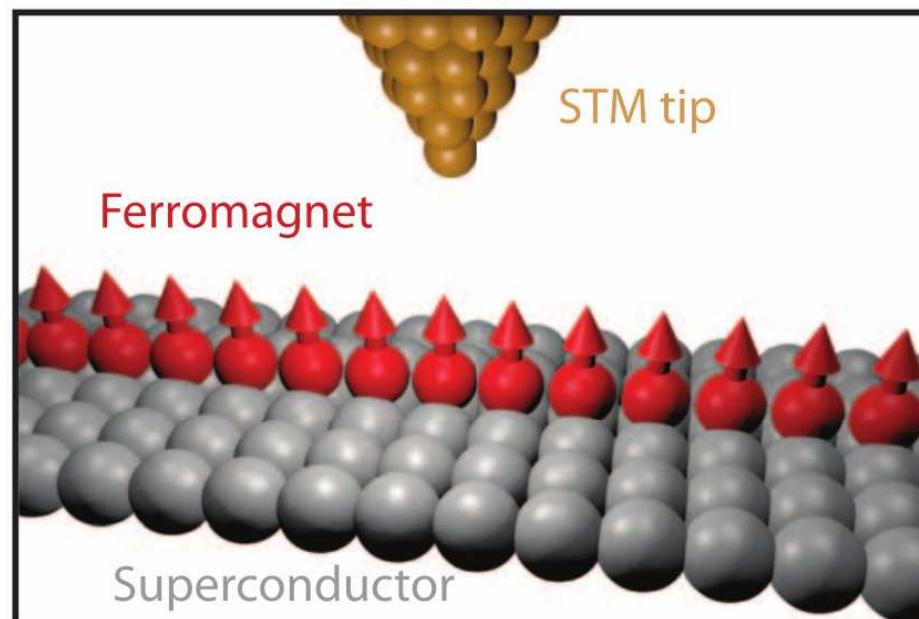


STM measurements provided evidence for:

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A chain of iron atoms deposited on a surface of superconducting lead



STM measurements provided evidence for:

⇒ Majorana bound states at the edges of a chain.

S. Nadj-Perge, ..., and A. Yazdani, Science 346, 602 (2014).

/ Princeton University, Princeton (NJ), USA /

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- Majorana vs Kondo features

## Historical remarks

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- E. Schrödinger (1926)

$$E \longrightarrow i \frac{d}{dt}$$

$$\vec{p} \longrightarrow -i\nabla$$

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- E. Majorana (1937)

particle = antiparticle

E.M. noticed that particular choice of  $\vec{\alpha}$  and  $\beta$  implies a real wave-function !

## Searching for majoranas

– in particle and nuclear physics

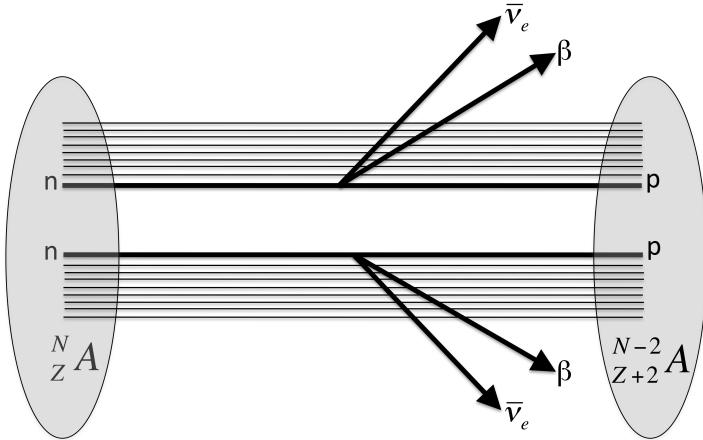
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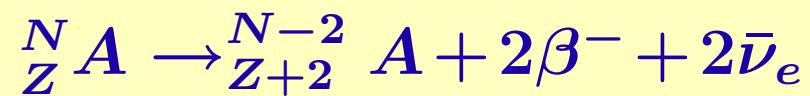
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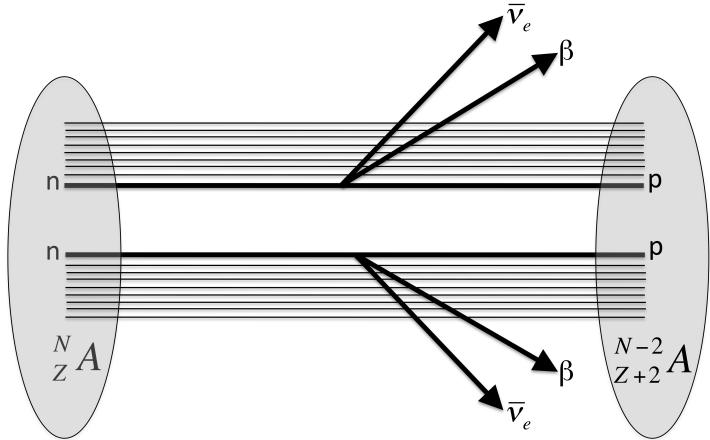


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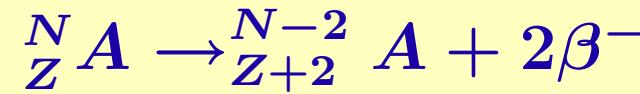
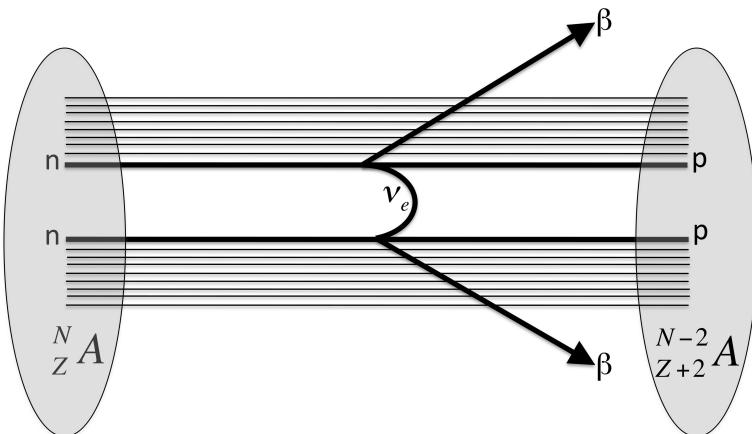


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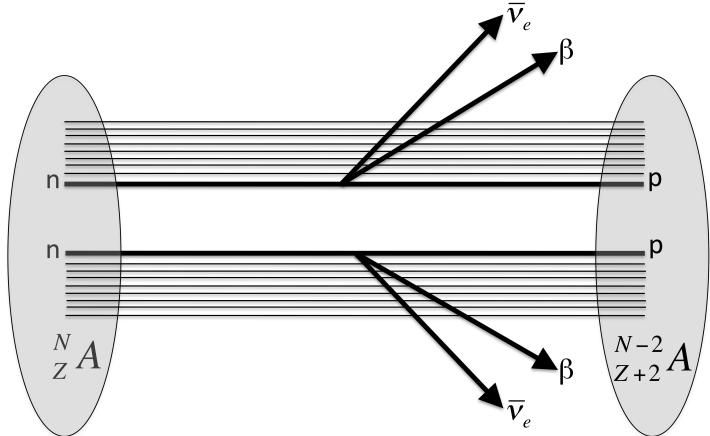
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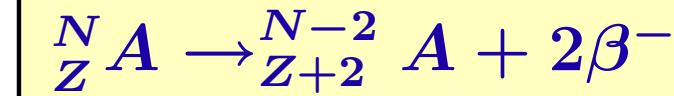
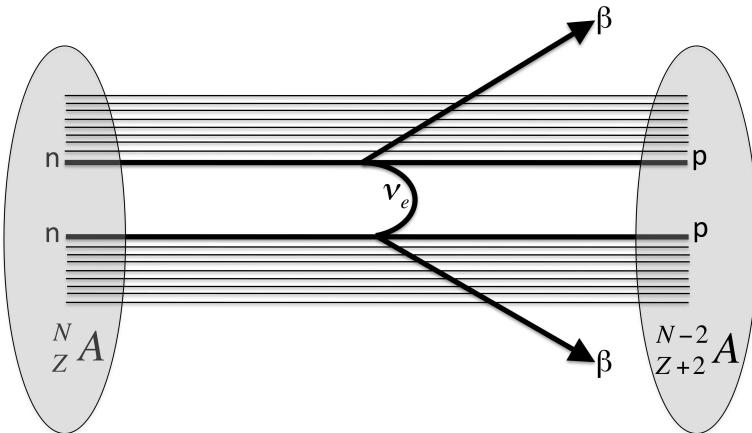
Neutrinoless decay would imply neutrinos to be majoranas.

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Neutrinoless decay would imply neutrinos to be majoranas.

Does it really occur ?

# Search for majorana quasiparticles – in solids

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- ★ Many-body effects, however, can induce emergent quasiparticles / concept 'More is different' emphasized by P.W. Anderson (1972) / Examples: phonons, polarons, magnons, spinons, holons etc.
- ★ Formally, any Dirac fermion can be majoranized via a canonical transformation to the Majorana basis .

# Majoranization

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- Usual (Dirac) fermions obey the anticommutation relations

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- $c_j^{(\dagger)}$  can be recast in terms of majoranas

$$\begin{aligned}\hat{c}_j &\equiv (\hat{\gamma}_{j1} + i\hat{\gamma}_{j2}) / 2 \\ \hat{c}_j^\dagger &\equiv (\hat{\gamma}_{j1} - i\hat{\gamma}_{j2}) / 2\end{aligned}$$

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- These  $\gamma_{i,n}$  operators obey unconventional algebra

$$\begin{aligned}\{\hat{\gamma}_{i,n}, \hat{\gamma}_{j,m}^\dagger\} &= 2\delta_{i,j}\delta_{n,m} \\ \hat{\gamma}_{i,n}^\dagger &= \hat{\gamma}_{i,n}\end{aligned}$$

creation = annihilation !

## Majoranization

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$$\begin{aligned}\hat{\beta}_{k\uparrow} &\equiv u_k \hat{c}_{k\uparrow} + v_k \hat{c}_{-k\downarrow}^\dagger \\ \hat{\beta}_{-k\downarrow}^\dagger &\equiv -v_k \hat{c}_{k\uparrow} + u_k \hat{c}_{-k\downarrow}^\dagger\end{aligned}$$

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- At the Fermi level the BCS coefficients  $u_{k_F} = v_{k_F} = 1/\sqrt{2}$ , thus

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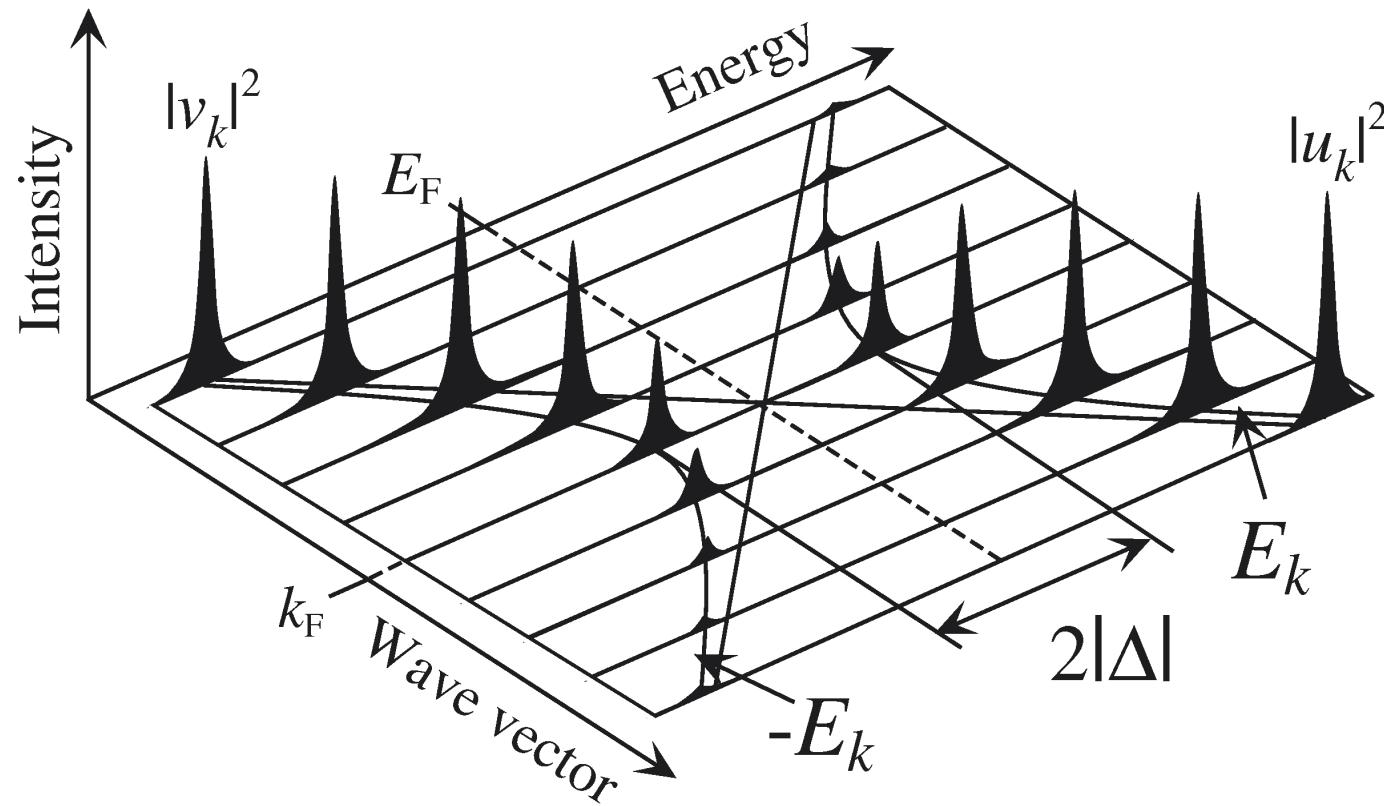
OK, but they must be zero-energy modes

# **Quasiparticles**

## **– of usual superconductors**

## Quasiparticles

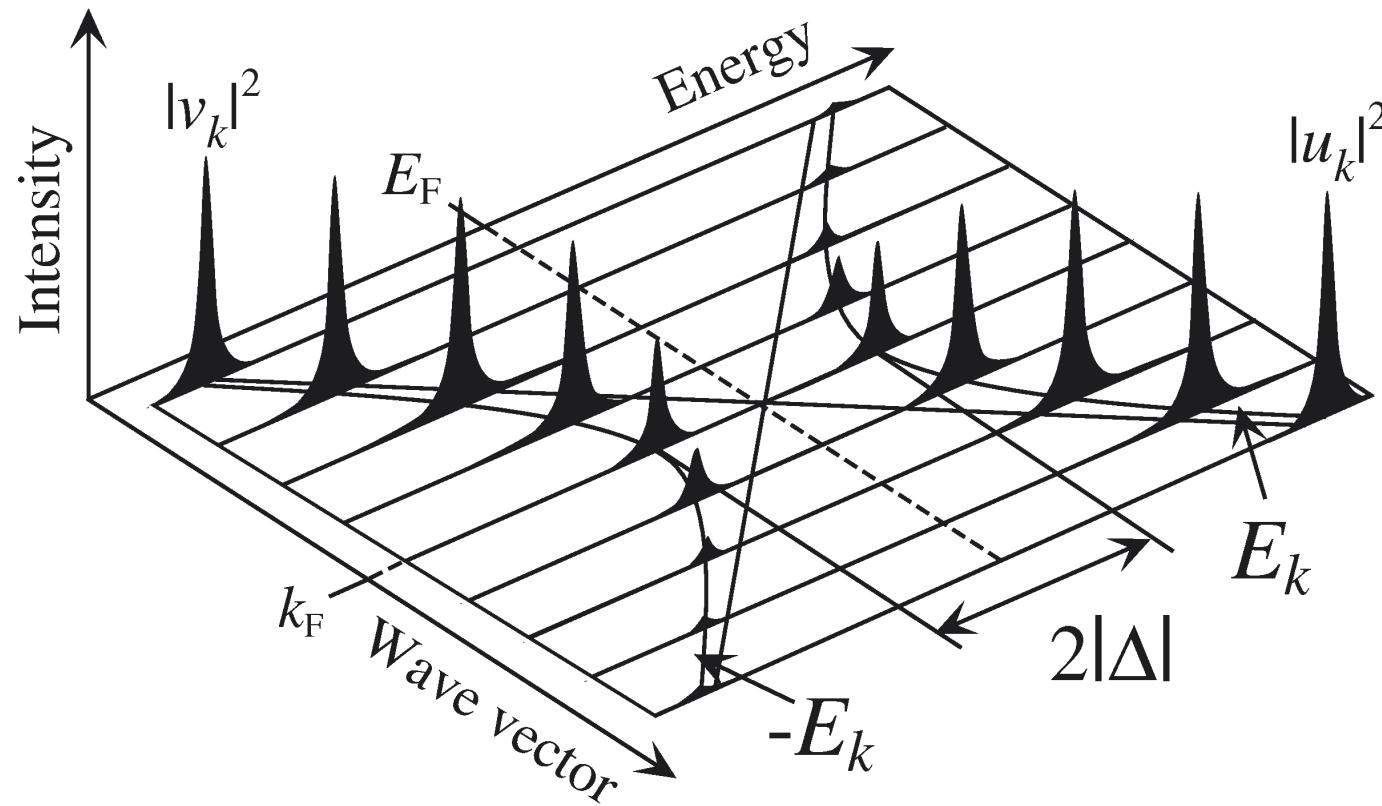
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To obtain true majoranas we need the zero energy ( $E_k = 0$ ) quasiparticles !

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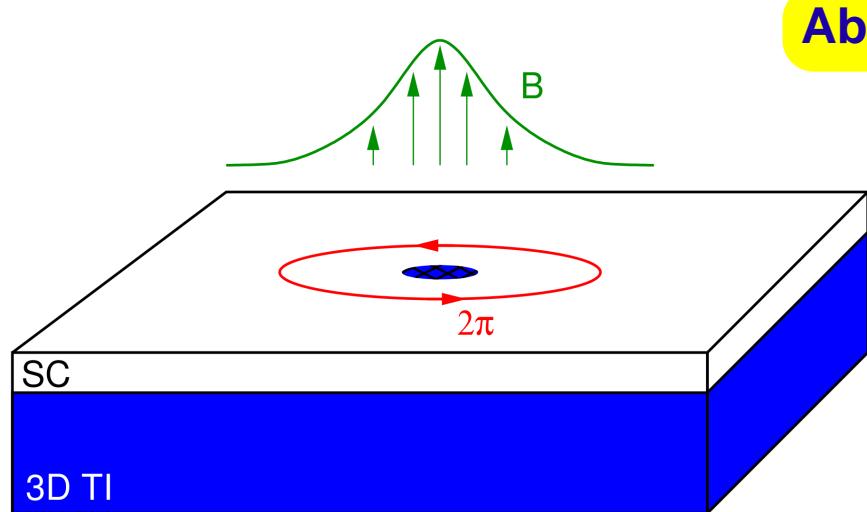
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- ★ Un-paired majoranas need time-reversal symmetry breaking , for instance, this can be obtained in p-wave superconductors.
- ★ Majorana quasiparticles obey non-Abelian (anyon) statistics, thus might be usefull for quantum computation.

## Theoretical proposals

- Fu-Kane model (2008)

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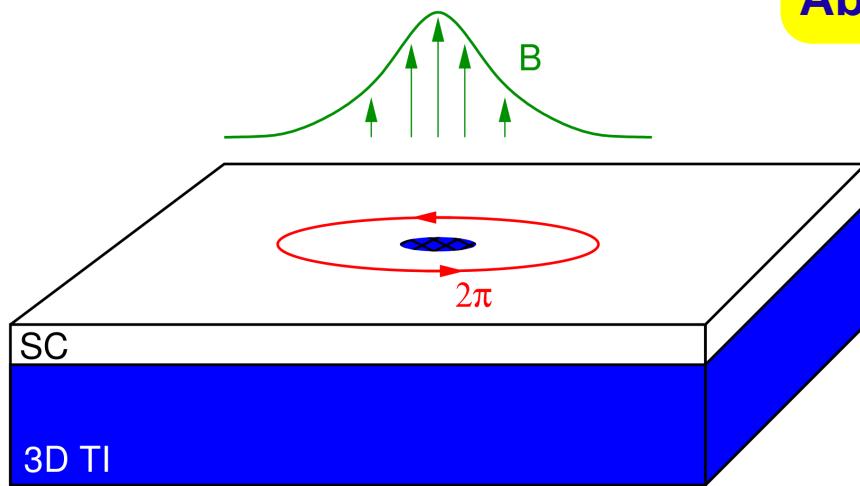
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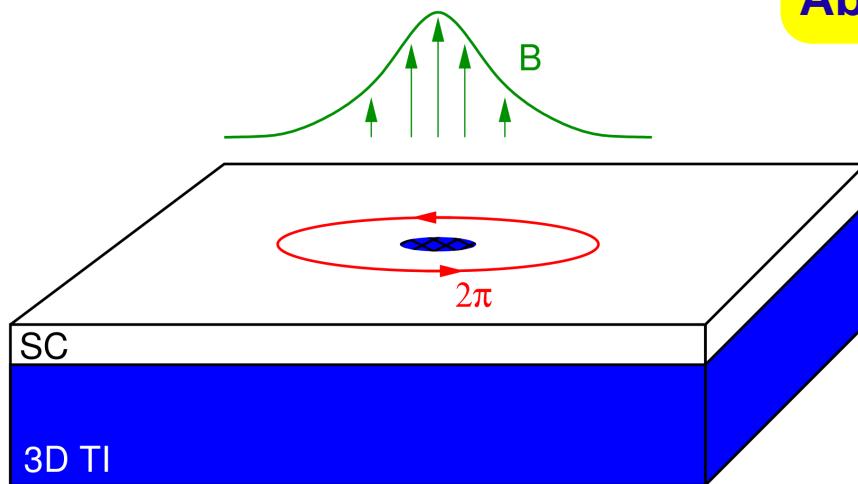
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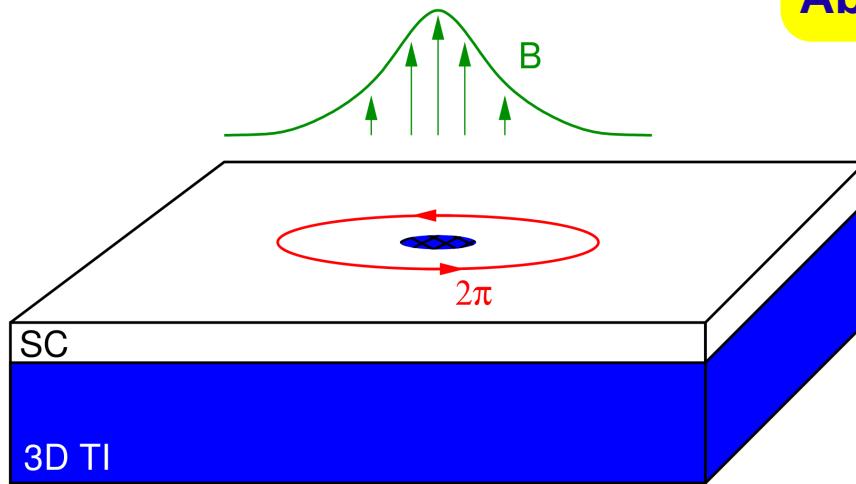
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Superconductivity is induced in the surface state of 3D topological insulator

Possible examples:      SC  $\Rightarrow$  Pb, Nb      3D TI  $\Rightarrow$  Bi<sub>2</sub>Se<sub>3</sub>, Bi<sub>2</sub>Te<sub>3</sub>

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The zero-mode Majorana quasiparticle

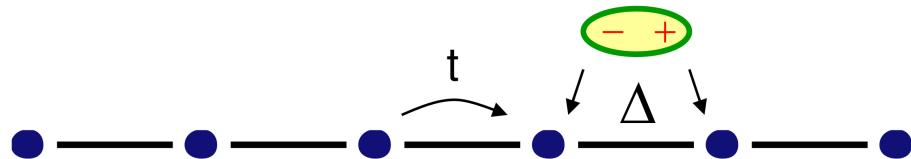
$$\hat{\gamma} = \frac{i}{\sqrt{2}} \int d^2\vec{r} \left[ e^{i(\alpha/2 - \pi/4)} \hat{c}_{\vec{r}\downarrow} - e^{-i(\alpha/2 - \pi/4)} \hat{c}_{\vec{r}\downarrow}^\dagger \right] f(r)$$

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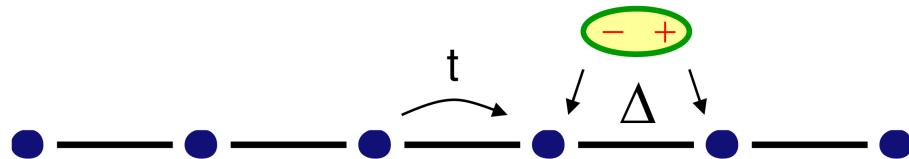
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**p-wave pairing of spinless 1D fermions**

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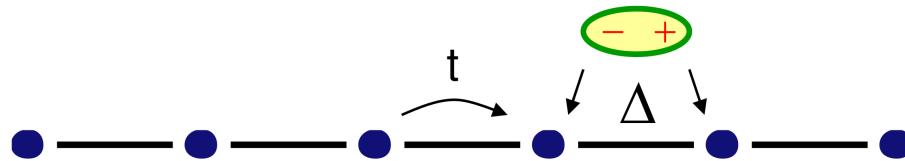


p-wave pairing of spinless 1D fermions

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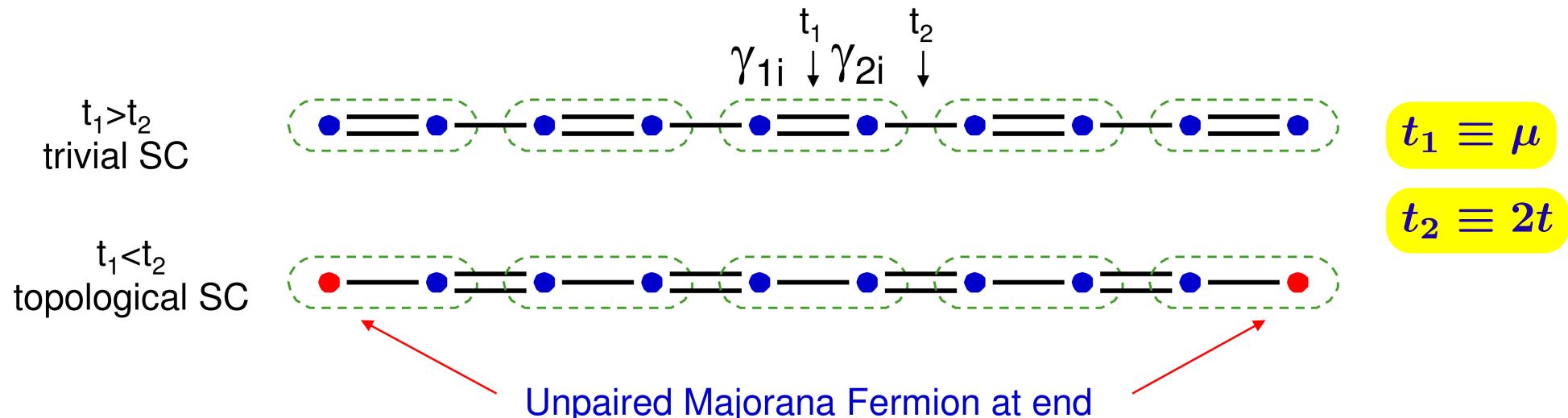
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This toy-model is exactly soluble in Majorana basis. Two special cases:



## Physical realizations

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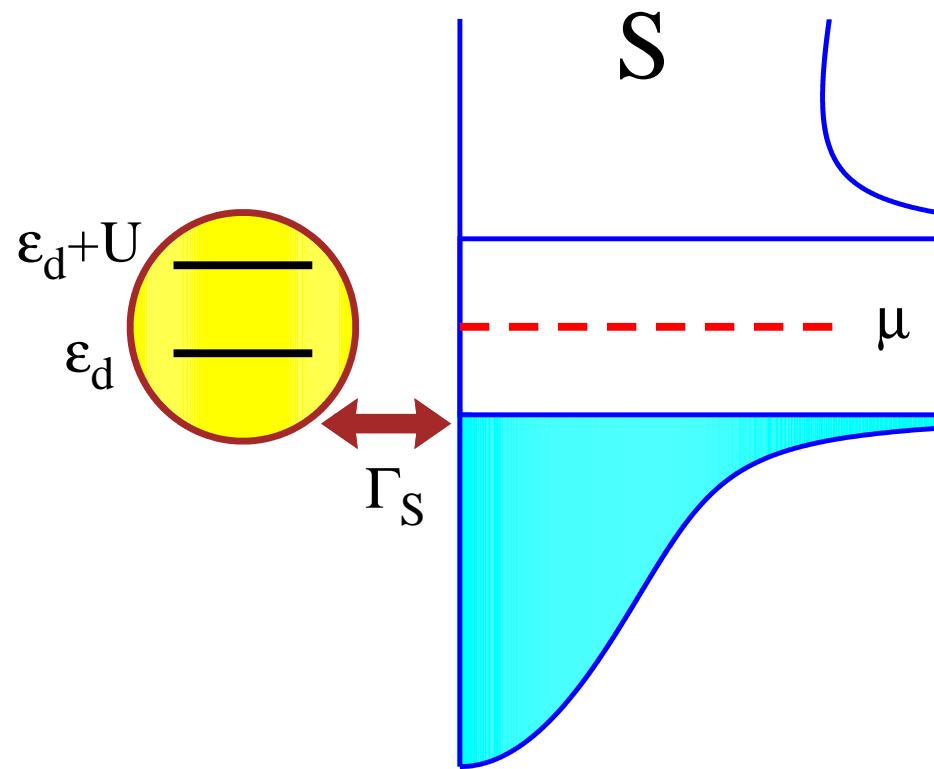
Choy *et al* (2011); Martin & Morpugo (2012); Nadj-Perge *et al* (2013)

- superconductor-double quantum dot

A.R. Wright, M. Veldhorst, Phys. Rev. Lett. (2013)

**Quantum impurity (dot or wire)  
coupled to a superconductor**

## Electronic spectrum



## Microscopic model

## Anderson-type Hamiltonian

Quantum impurity (dot)

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

coupled with a superconductor

$$\begin{aligned} \hat{H} &= \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} + \hat{H}_S \\ &+ \sum_{\mathbf{k}, \sigma} \left( V_{\mathbf{k}} \hat{d}_{\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + V_{\mathbf{k}}^* \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{d}_{\sigma} \right) \end{aligned}$$

where

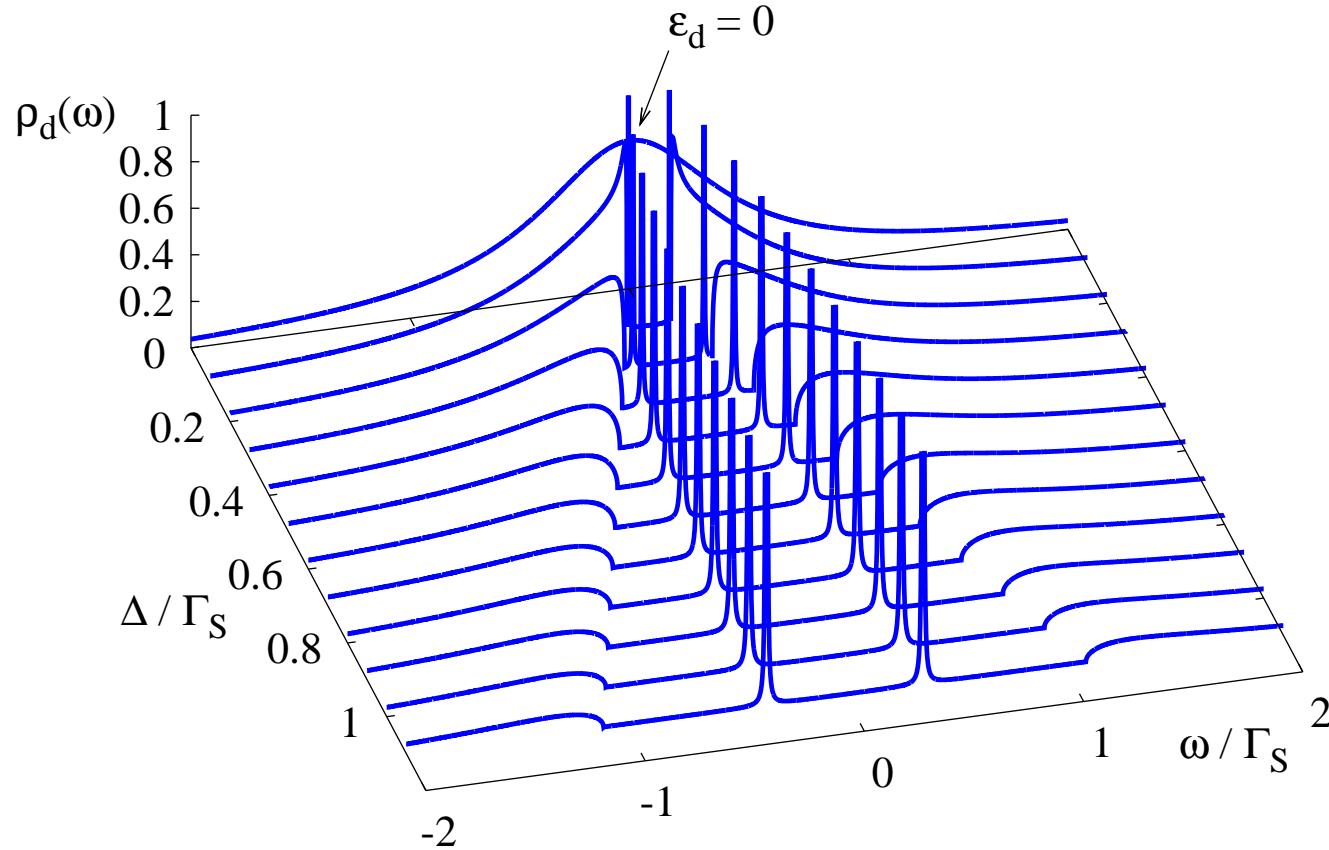
$$\hat{H}_S = \sum_{\mathbf{k}, \sigma} (\epsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} - \sum_{\mathbf{k}} \left( \Delta \hat{c}_{\mathbf{k}\uparrow}^{\dagger} \hat{c}_{\mathbf{k}\downarrow}^{\dagger} + \text{h.c.} \right)$$

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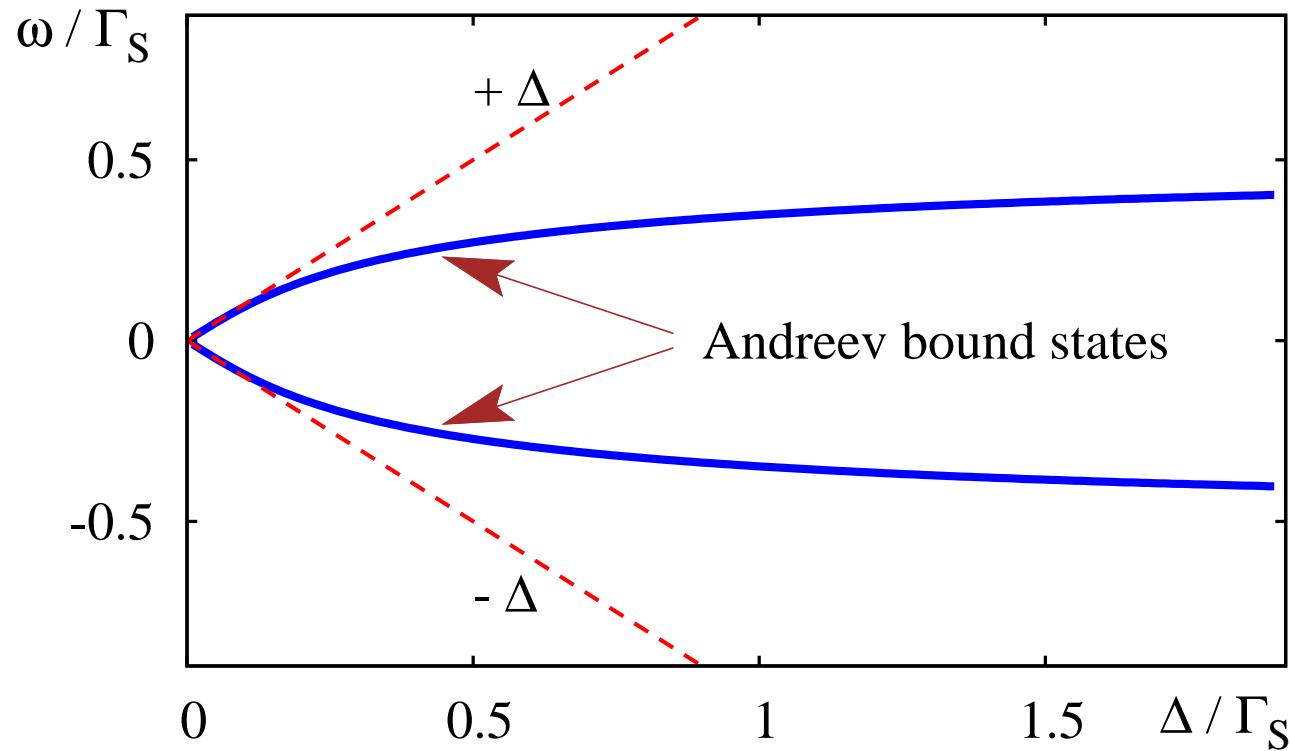


In-gap resonances: Andreev bound states.

J. Barański and T. Domański, J. Phys.: Condens. Matter **25**, 435305 (2013).

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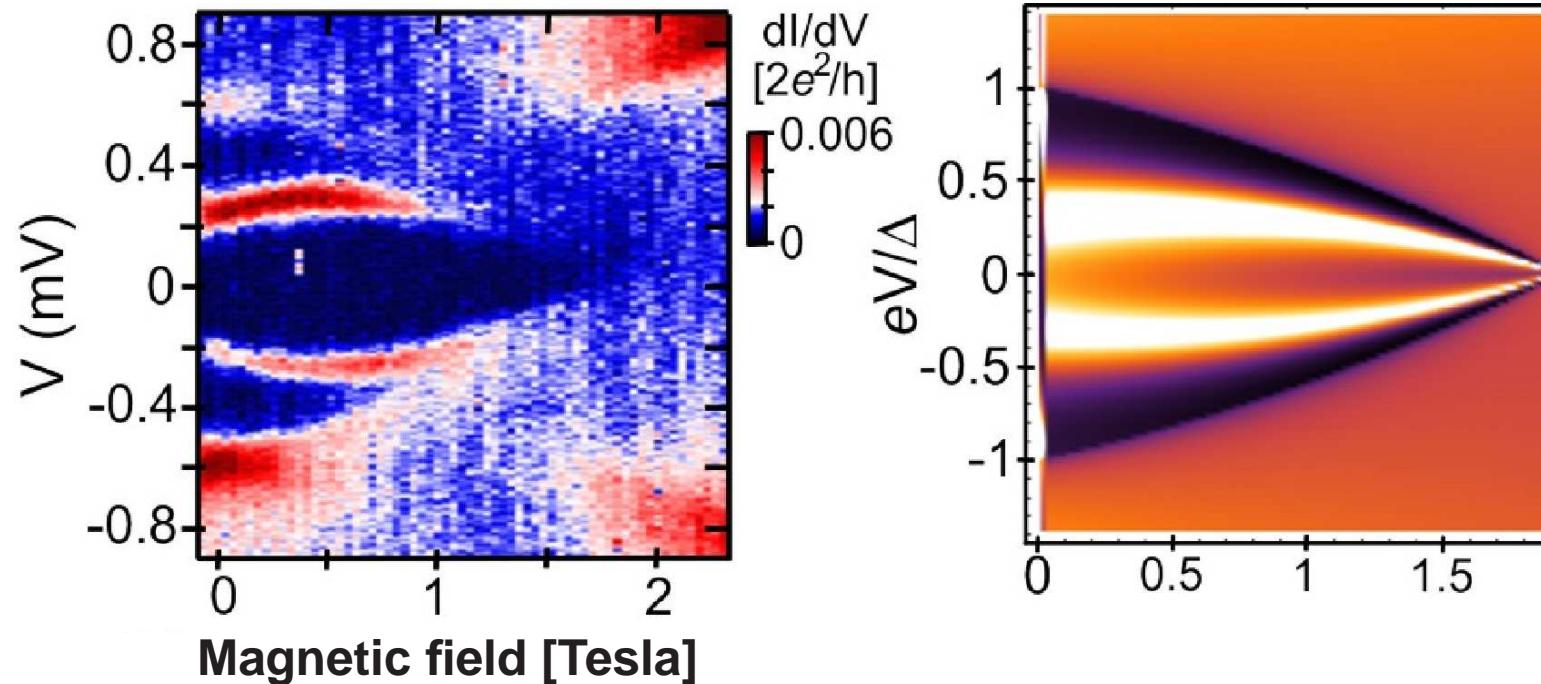


Energies of the in-gap resonances (Andreev bound states)

J. Barański and T. Domański, J. Phys.: Condens. Matter **25**, 435305 (2013).

## Subgap states

- experimental data



Differential conductance of nanotubes coupled to vanadium (S) and gold (N)

/ external magnetic field changes the magnitude of pairing gap  $\Delta(B)$  /

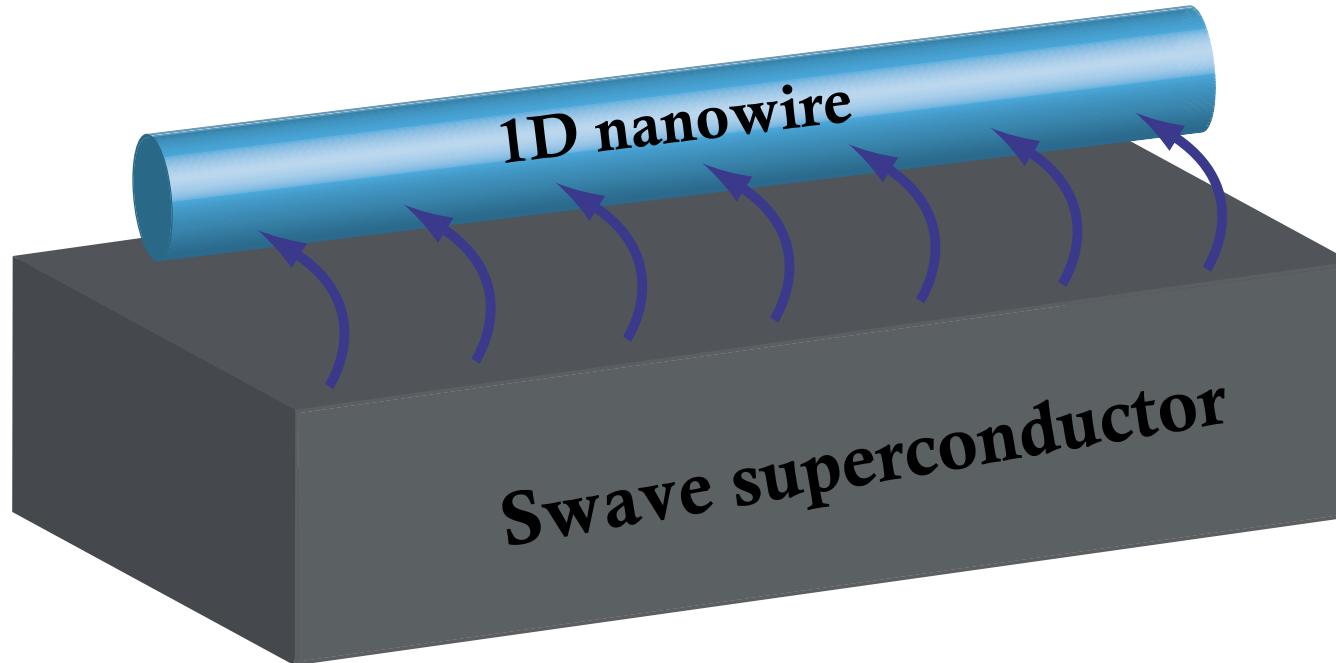
Eduardo J.H. Lee, ..., S. De Franceschi, Nature Nanotechnology 9, 79 (2014).

# Andreev vs Majorana states

## – a story of mutation

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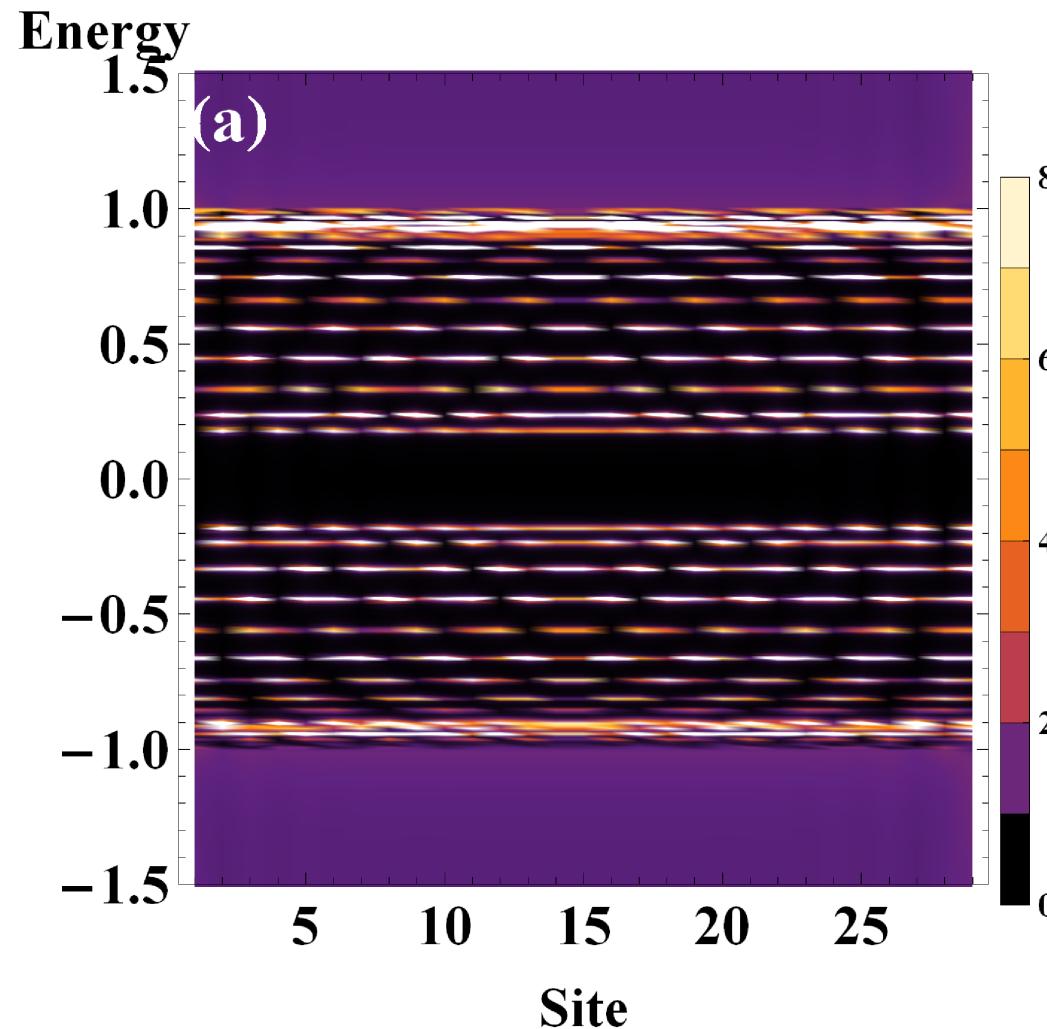


Imagine a quantum wire deposited on s-wave superconductor

*D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B 88, 165401 (2013).*

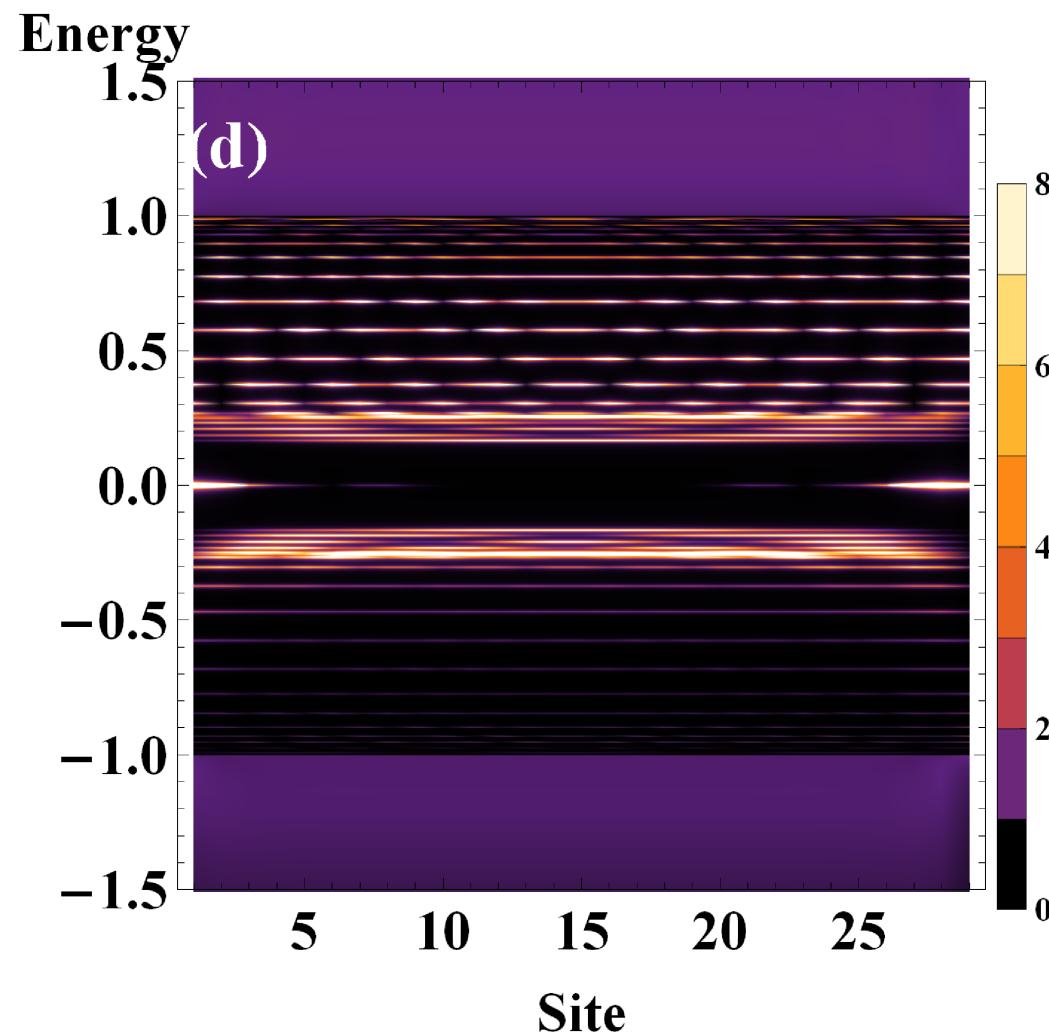
# Andreev vs Majorana states

- a story of mutation



Spectrum of a quantum wire reveals a number of Andreev states.

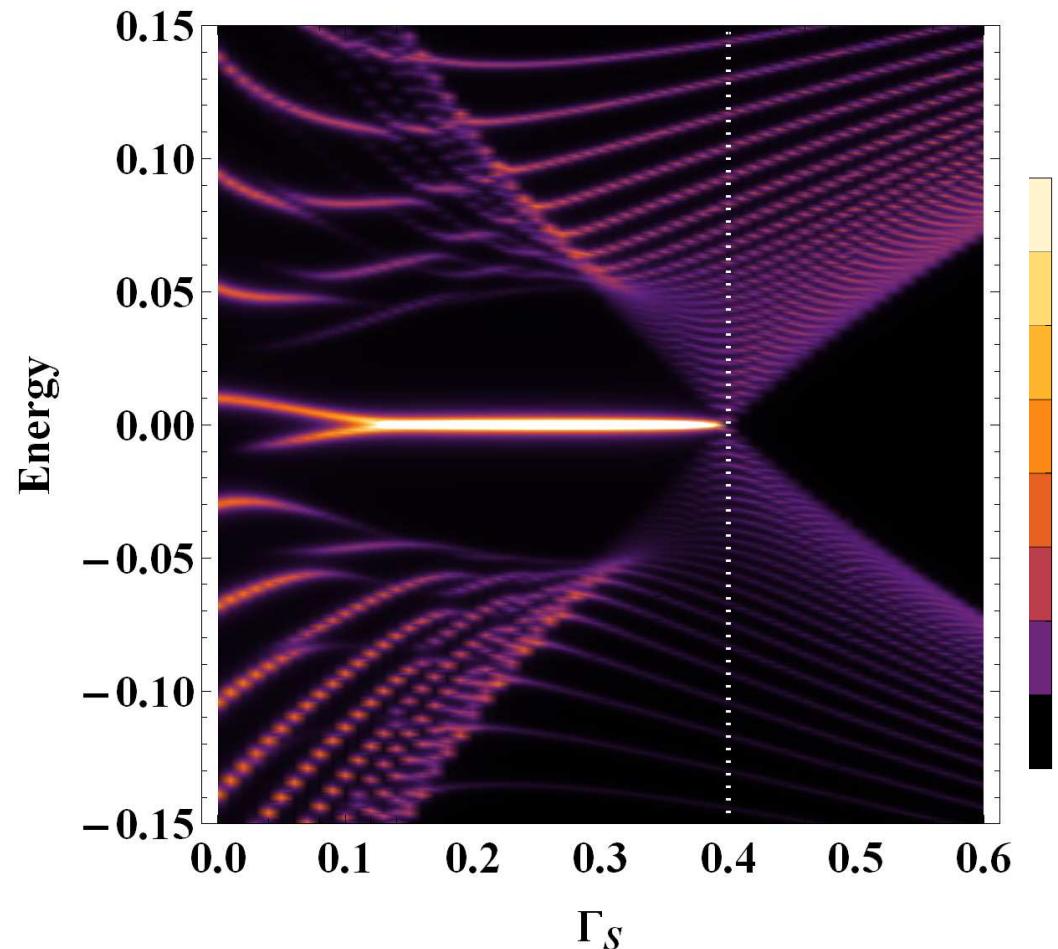
# Andreev vs Majorana states – a story of mutation



Spin-orbit coupling can induce the Majorana-type quasiparticles.

## Andreev vs Majorana states

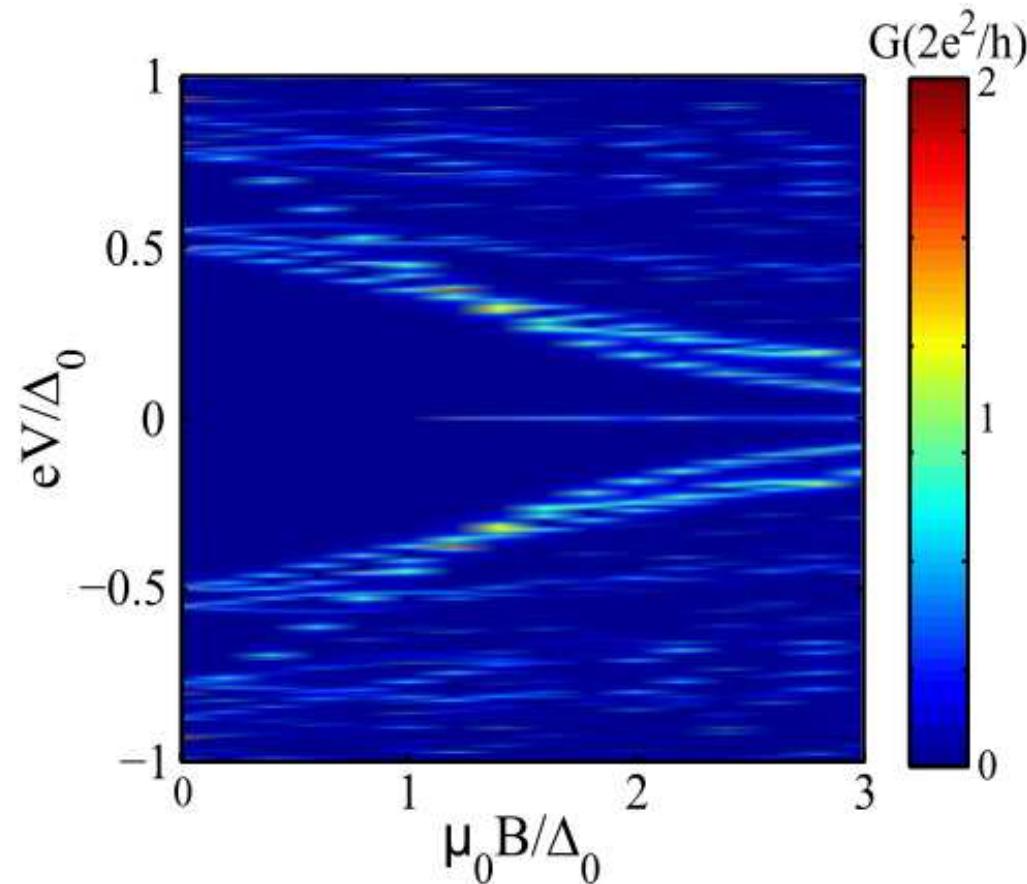
– a story of mutation



**Majorana quasiparticles appear at the edges of a quantum wire.**

*D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B **88**, 165401 (2013).*

# Andreev vs Majorana states – a story of mutation

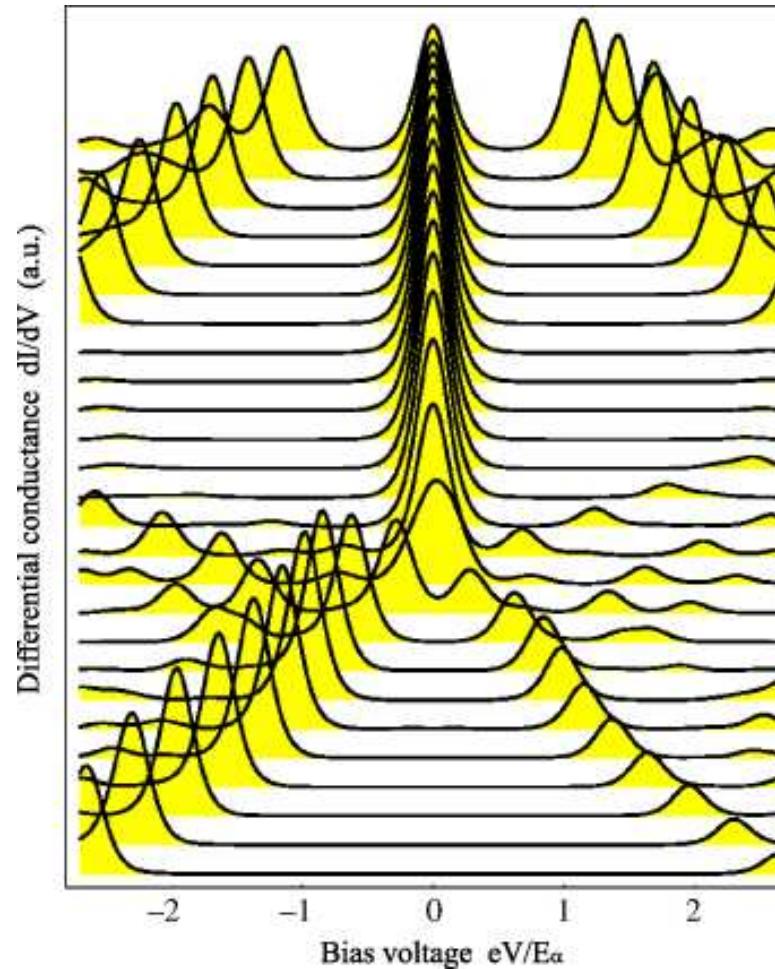


Quasiparticles at the edge of a quantum wire for varying magnetic field.

J. Liu, A.C. Potter, K.T. Law, and P.A. Lee, Phys. Rev. Lett. **109**, 267002 (2012).

# Andreev vs Majorana states

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**Quasiparticles at the edge of a quantum wire for varying magnetic field.**

*T.D. Stanescu, R.M. Lutchyn, and S. Das Sarma, Phys. Rev. B 84, 144522 (2011).*

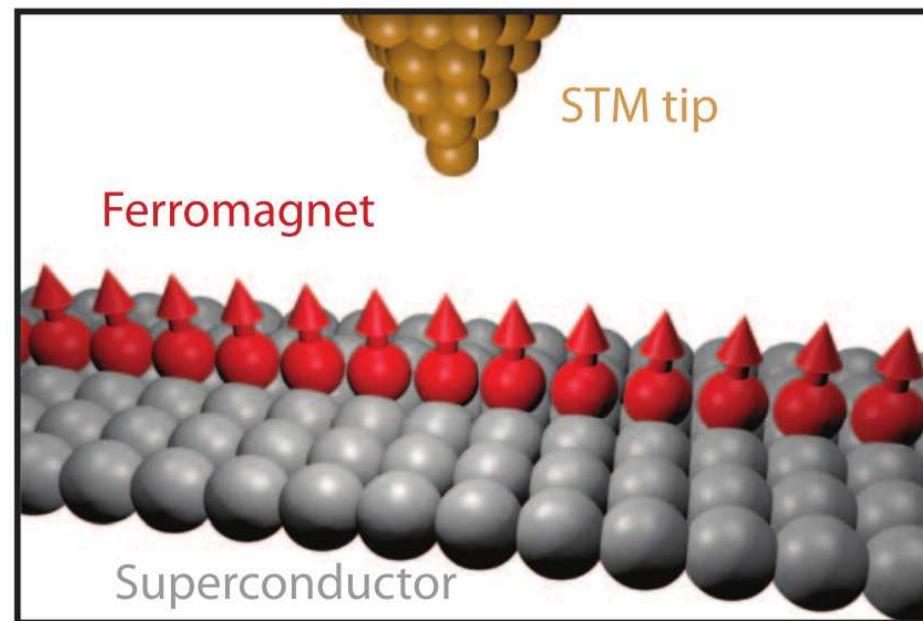
## Experimental results

### – for Majorana quasiparticles

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## – for Majorana quasiparticles

A chain of iron atoms deposited on a surface of superconducting lead



STM measurements provided evidence for:

⇒ Majorana bound states at the edges of a chain.

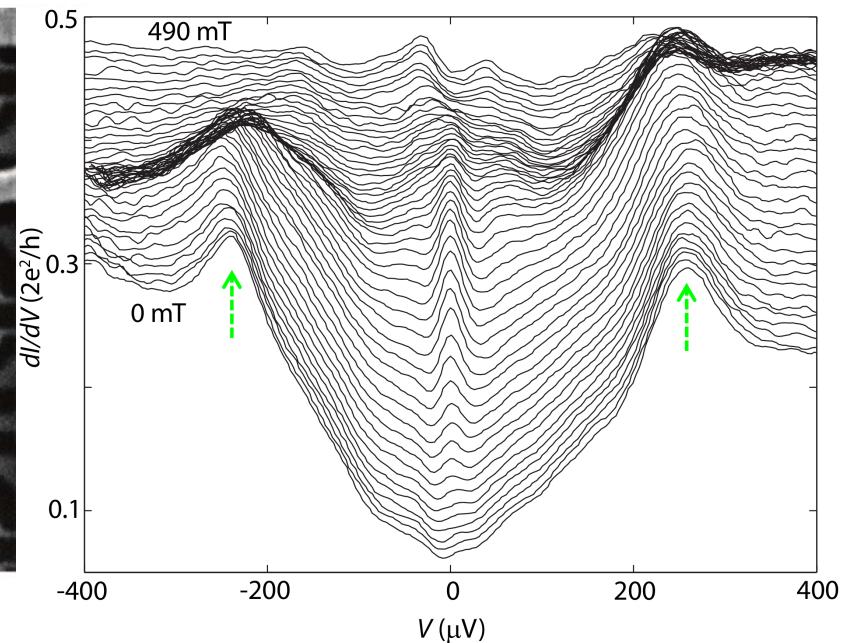
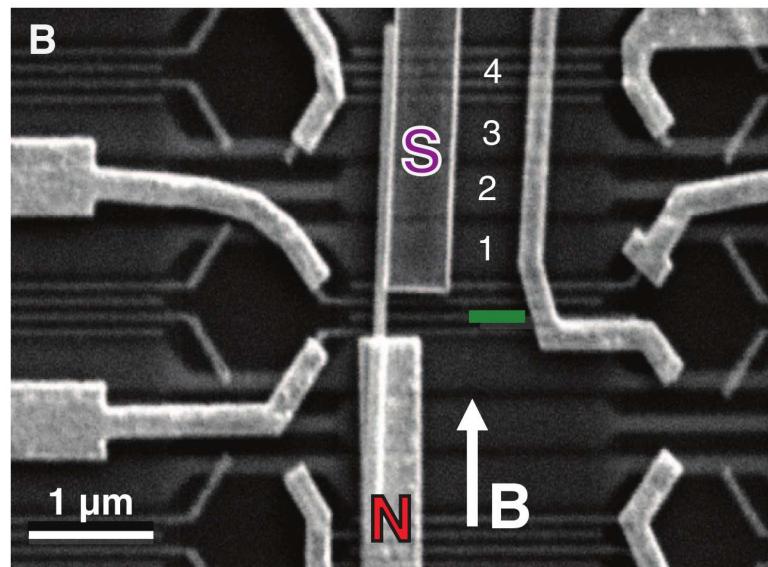
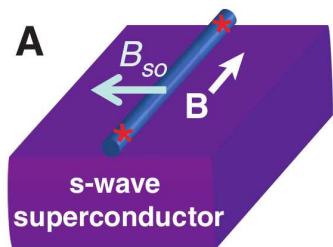
S. Nadj-Perge, ..., and A. Yazdani, Science 346, 602 (2014).

/ Princeton University, Princeton (NJ), USA /

## Experimental results

## – for Majorana quasiparticles

InSb nanowire between a metal (gold) and a superconductor (Nb-Ti-N)



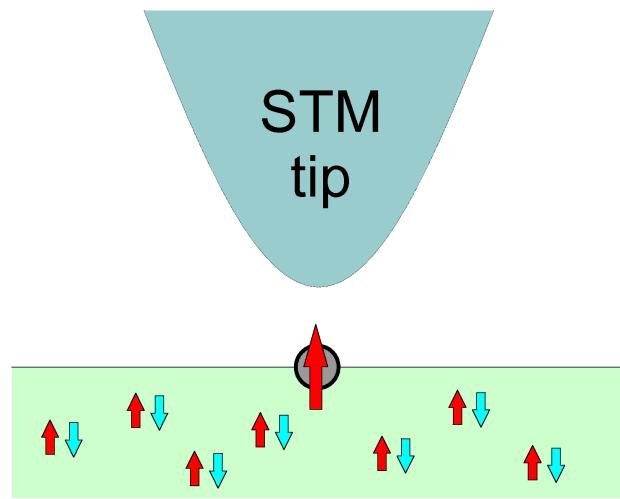
$dI/dV$  measured at 70 mK for varying magnetic field  $B$  indicated:

⇒ a zero-bias enhancement due to Majorana state

V. Mourik, ..., and L.P. Kouwenhoven, Science **336**, 1003 (2012).

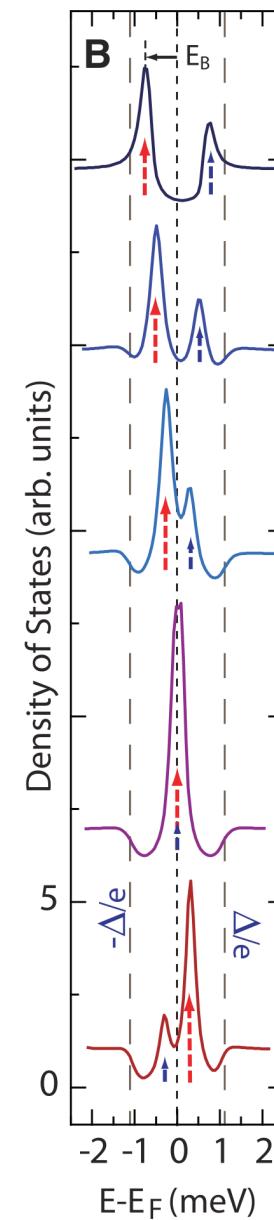
/ Kavli Institute of Nanoscience, Delft Univ., Netherlands /

## Subgap states of the bulk materials



**STM scheme (left) and the experimental data (right) obtained for Mn impurities on the superconducting Pb(111) surface.**

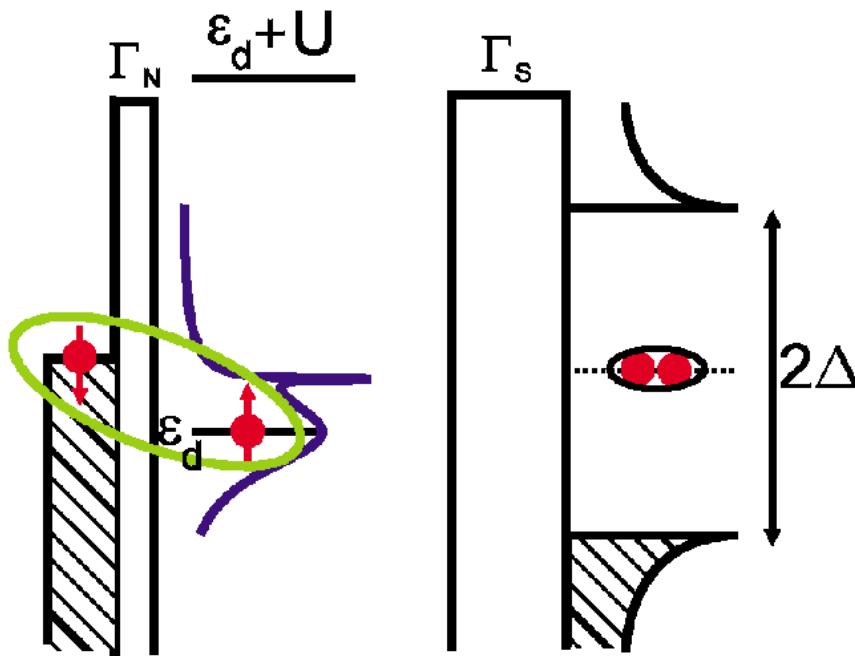
K.J. Franke, G. Schulze, and J.I. Pascual, Science **332**, 940 (2011).



## **Other related issues**

## Relevant problems : issue # 1

Coupling of the QD with a metallic electrode:

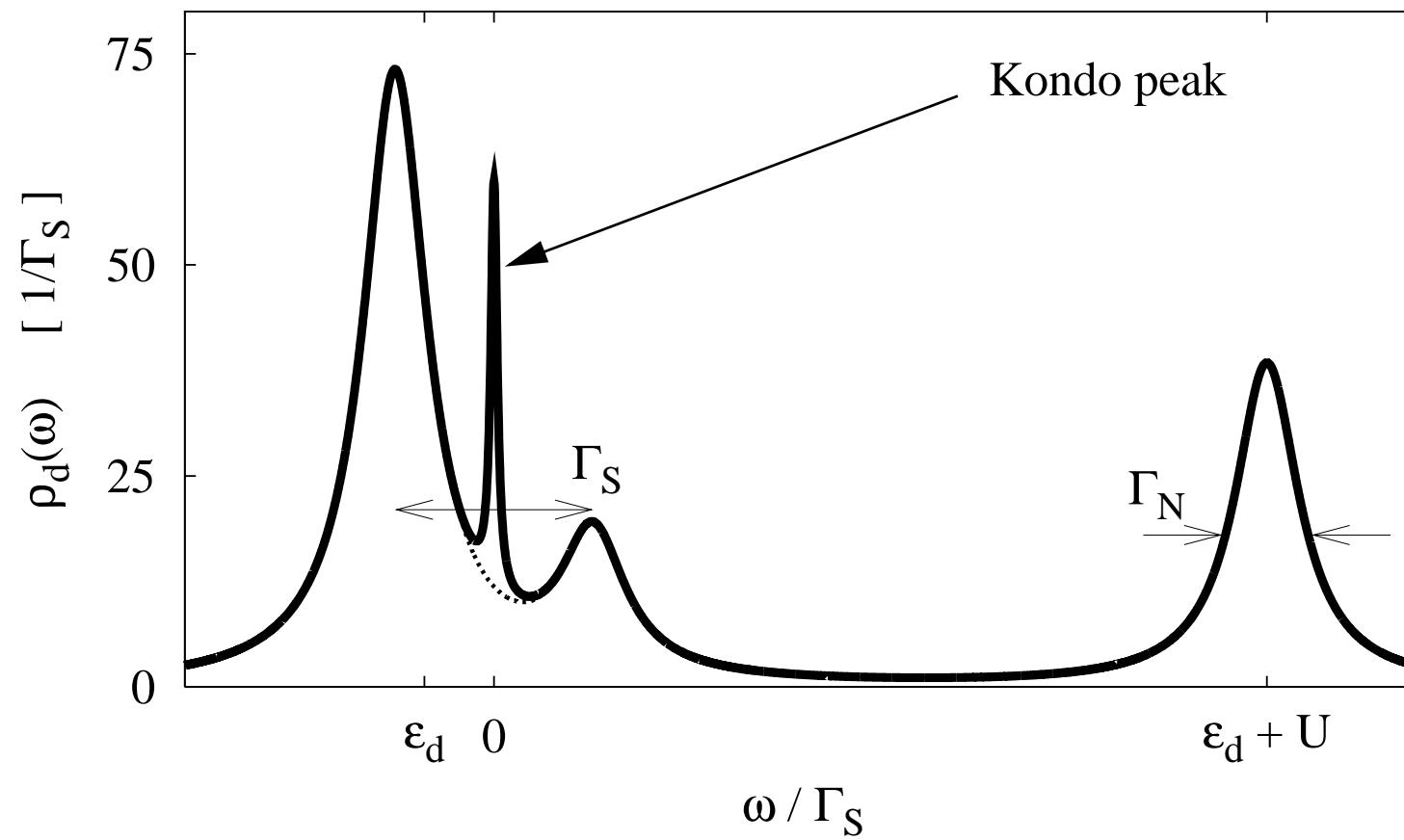


- ★ broadens the QD levels ( $\sim \Gamma_N$ )
- ★ induces the Kondo resonance (below  $T_K$ ).

## Relevant problems :

## issue # 2

Additional hybridization  $\Gamma_S$  with superconductor induces on-dot pairing



## Theoretical background:

– various many-body techniques

EOM

R. Fazio and R. Raimondi (1998)

slave bosons

P. Schwab and R. Raimondi (1999)

NCA

A.A. Clerk, V. Ambegaokar, and S. Hershfield (2000)

IPT

J.C. Cuevas, A. Levy Yeyati, and A. Martin-Rodero (2001)

constrained sb

M. Krawiec and K.I. Wysokiński (2004)

NRG

Y. Tanaka, N. Kawakami, and A. Oguri, (2007)

EOM revised

T. Domański et al, (2007)

NRG

J. Bauer, A. Oguri, and A.C. Hewson, (2007)

f-RG

C. Karrasch, A. Oguri, and V. Meden, (2008)

QMC

A. Koga, (2013)

CUT

M. Zapalska and T. Domański, (2014)

NRG

R. Žitko et al, (2015)

## Andreev conductance – theoretical results

Subgap conductance  $G_A(V)$  obtained for:

$$U = 10\Gamma_N$$

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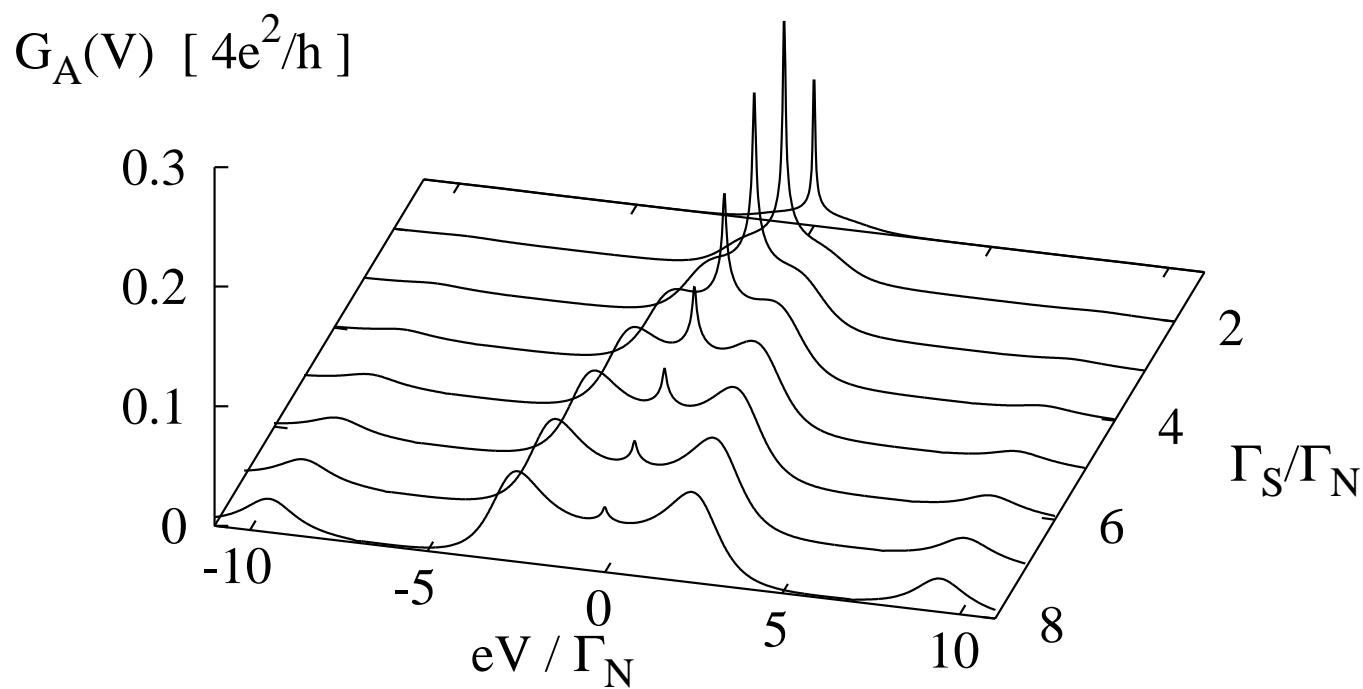
T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

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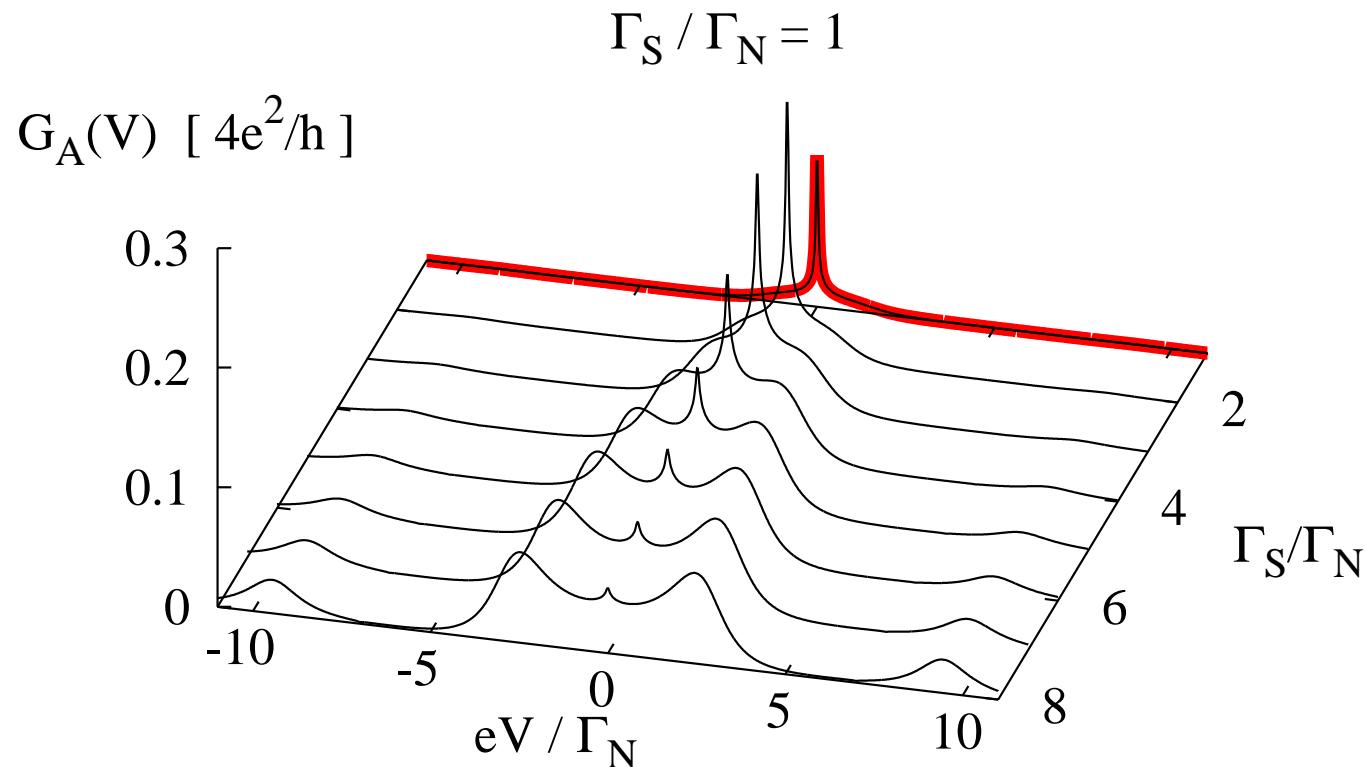
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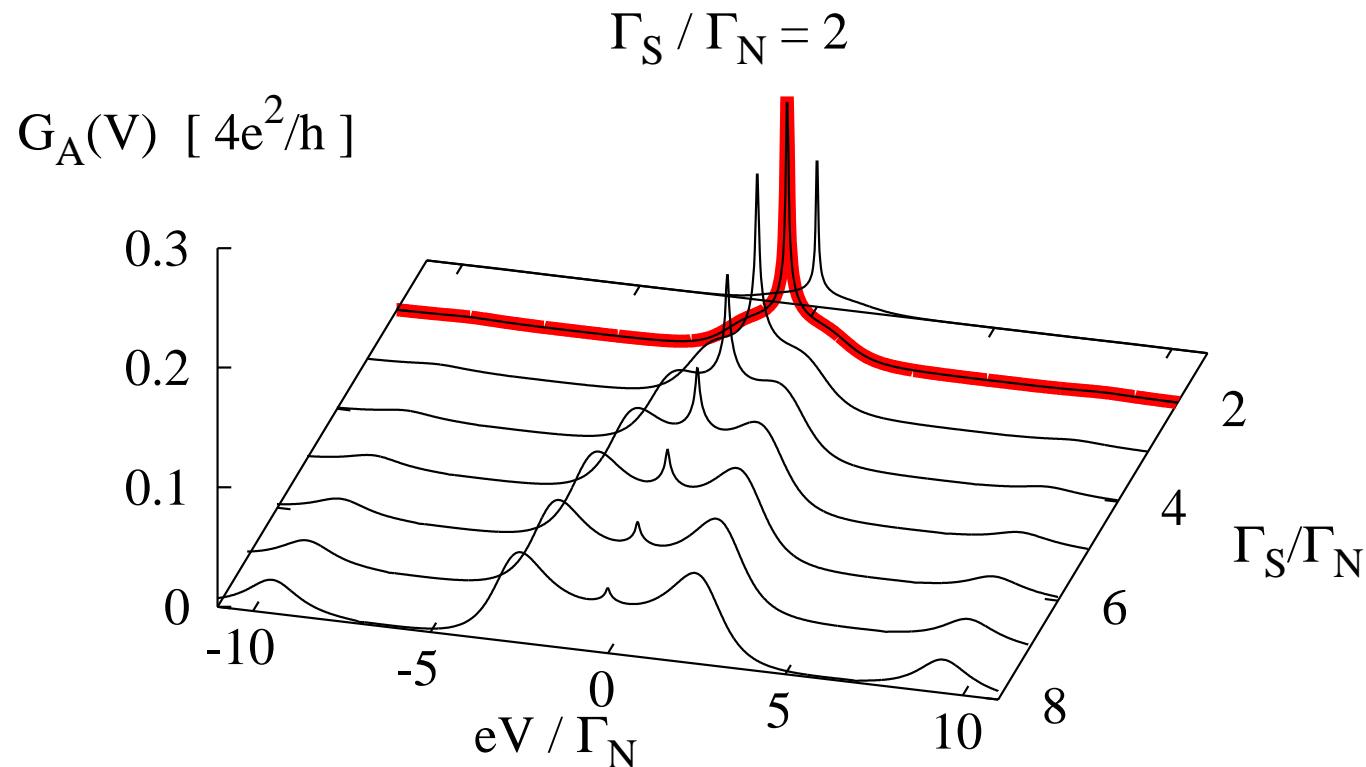
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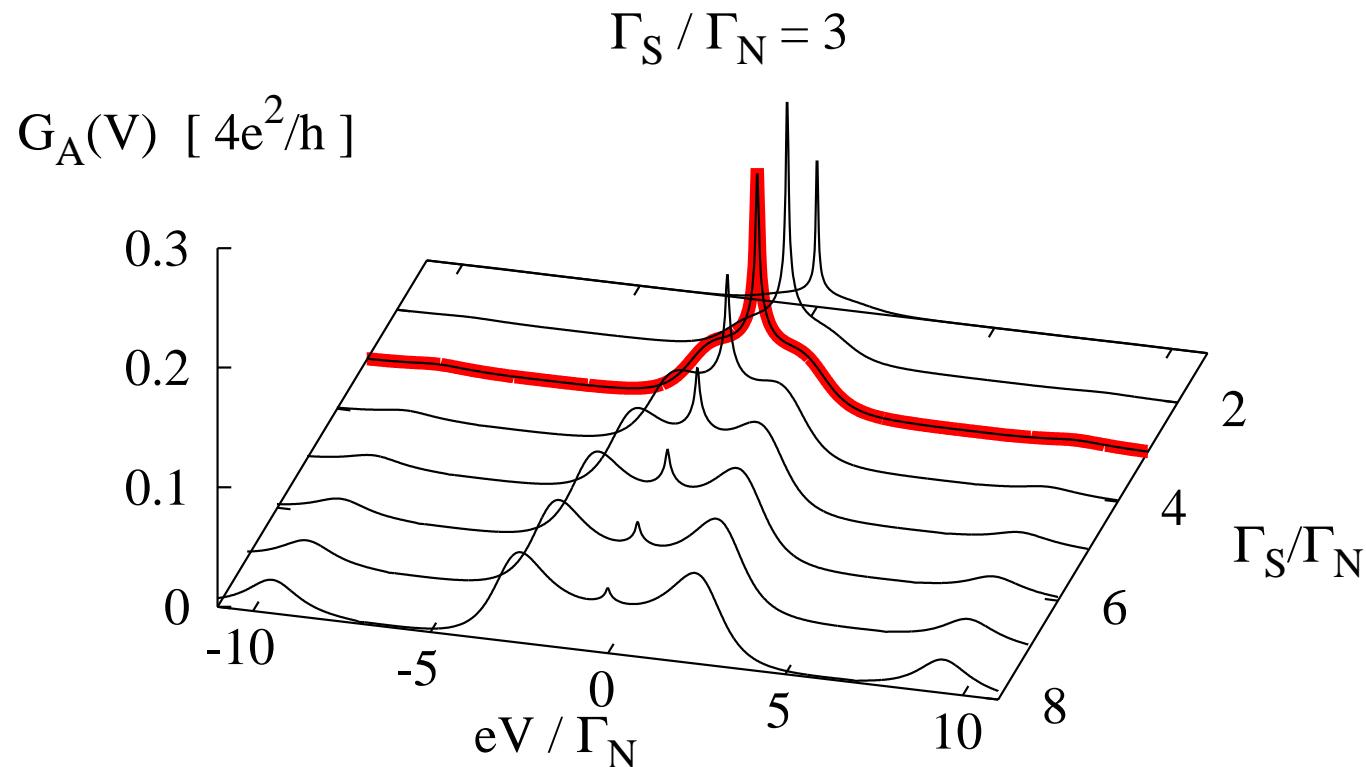
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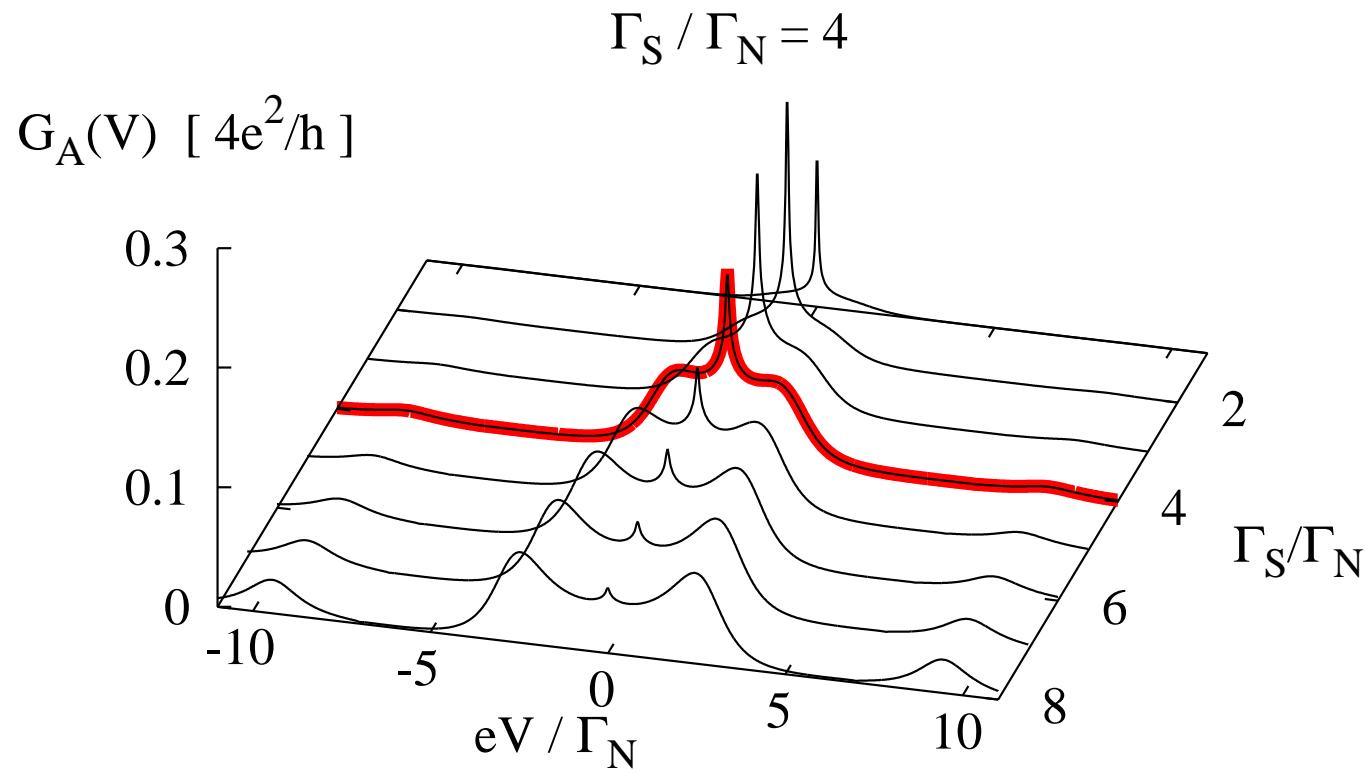
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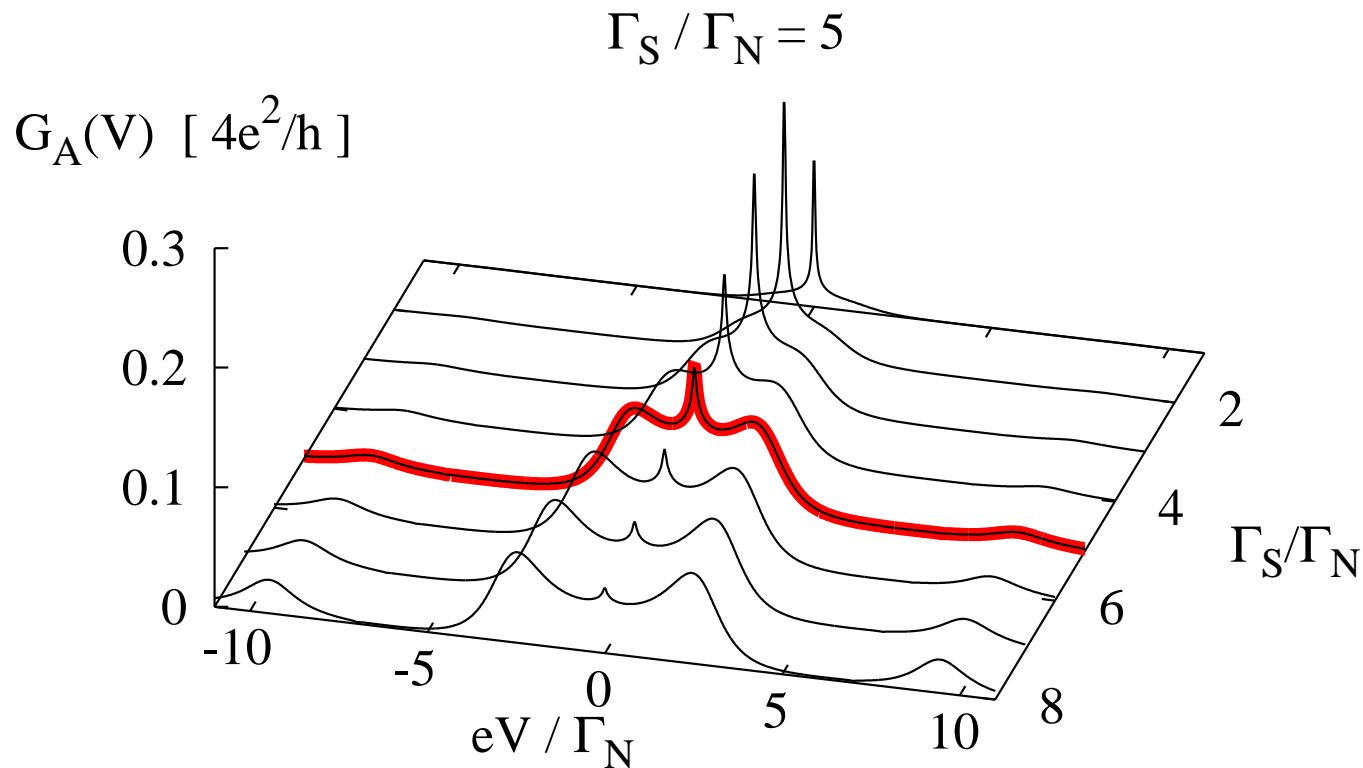
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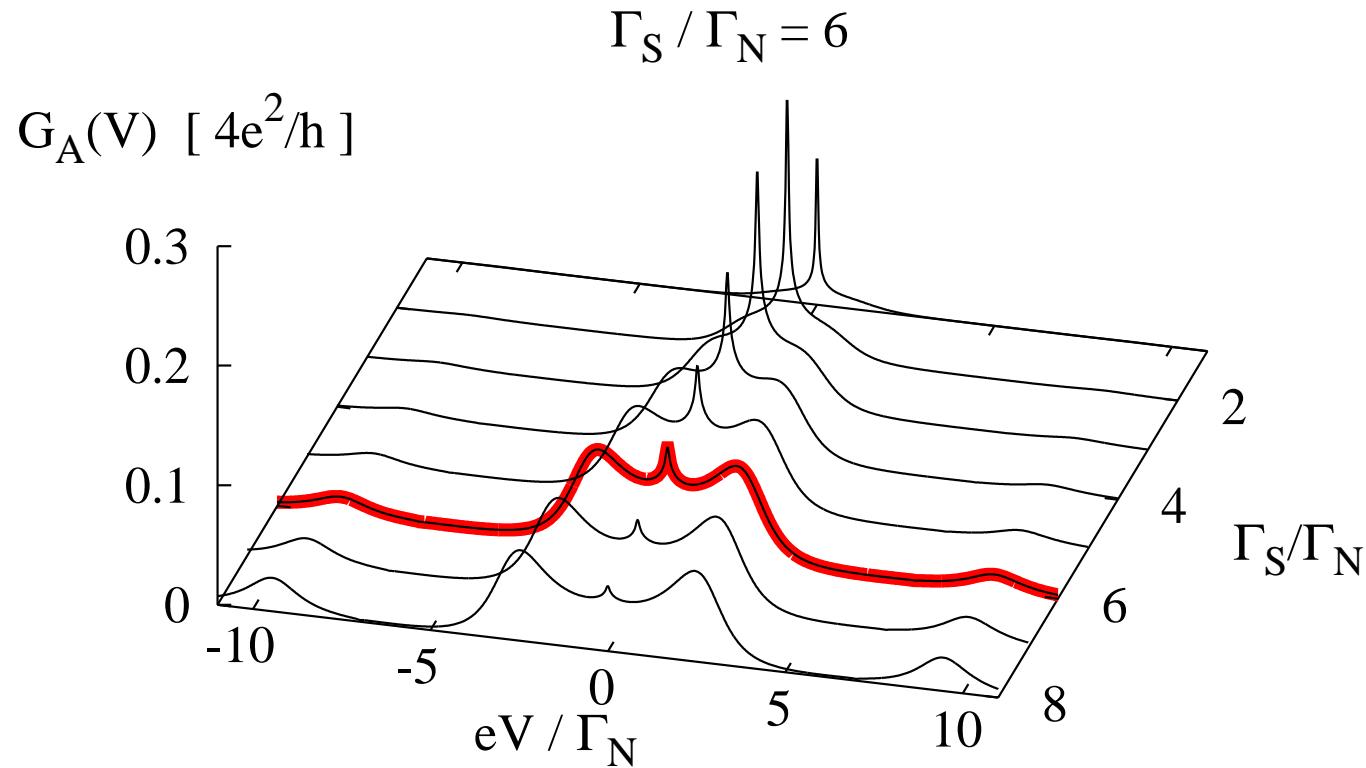
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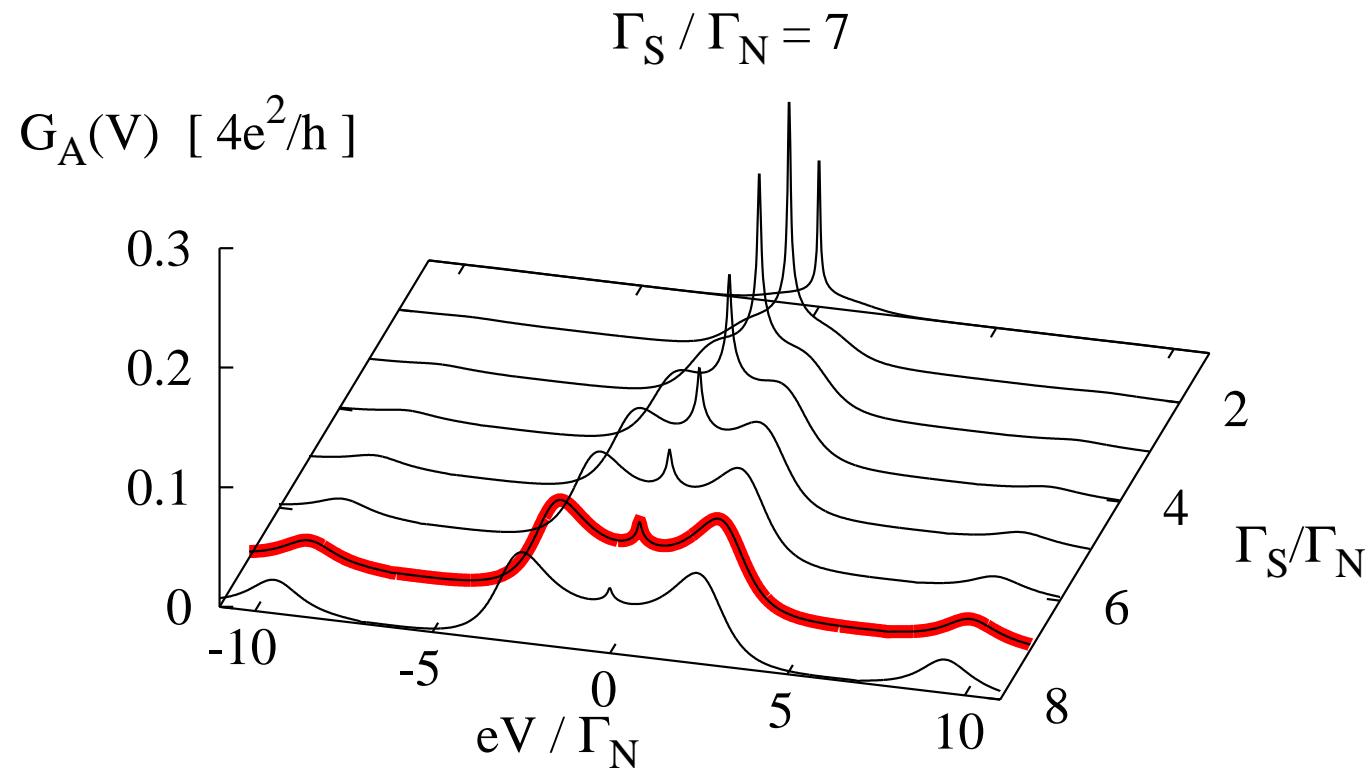
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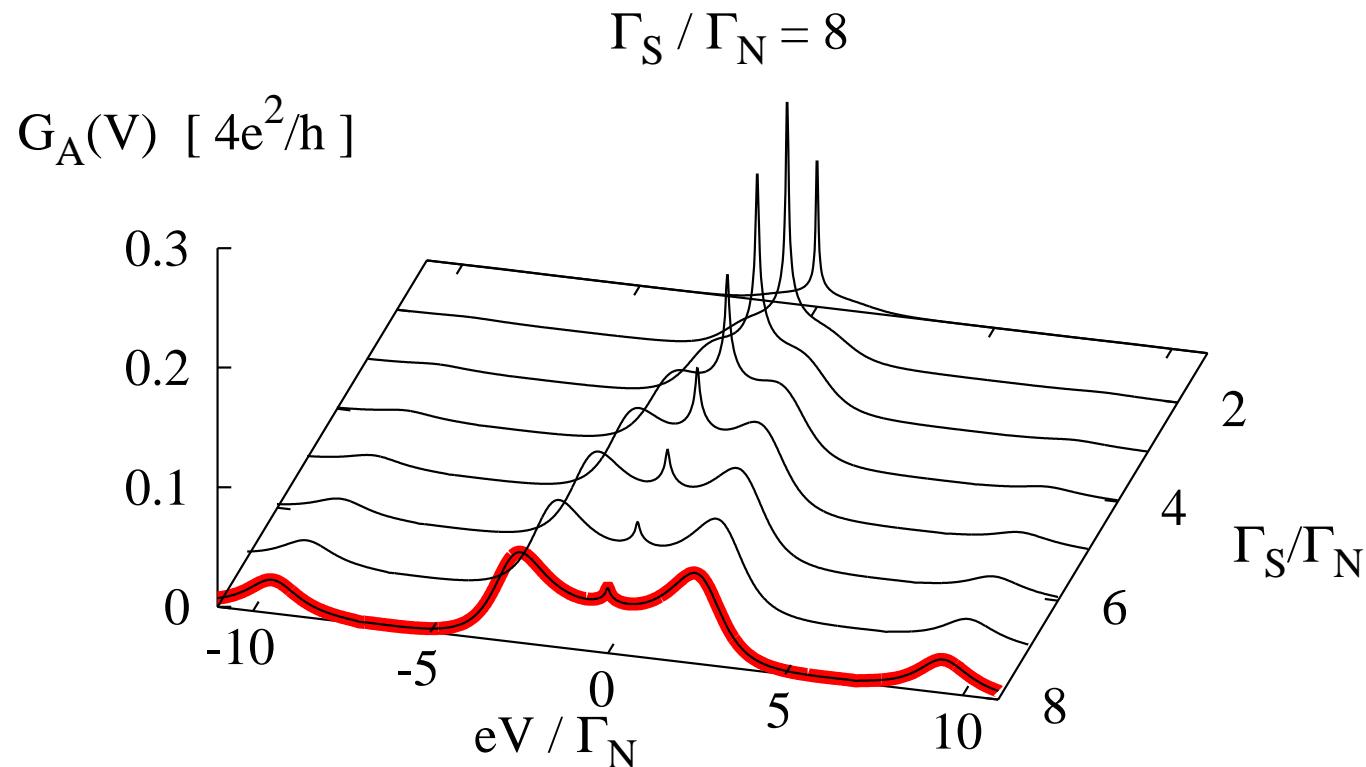
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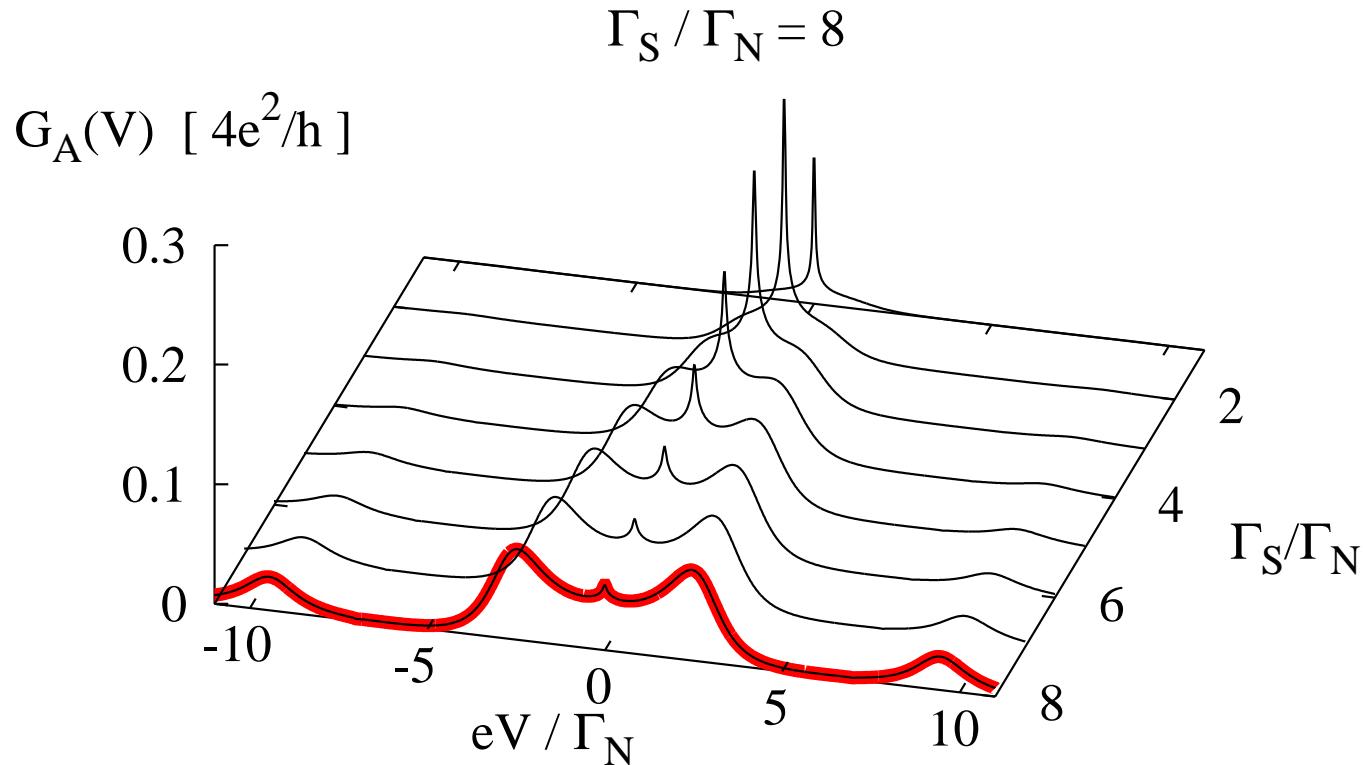
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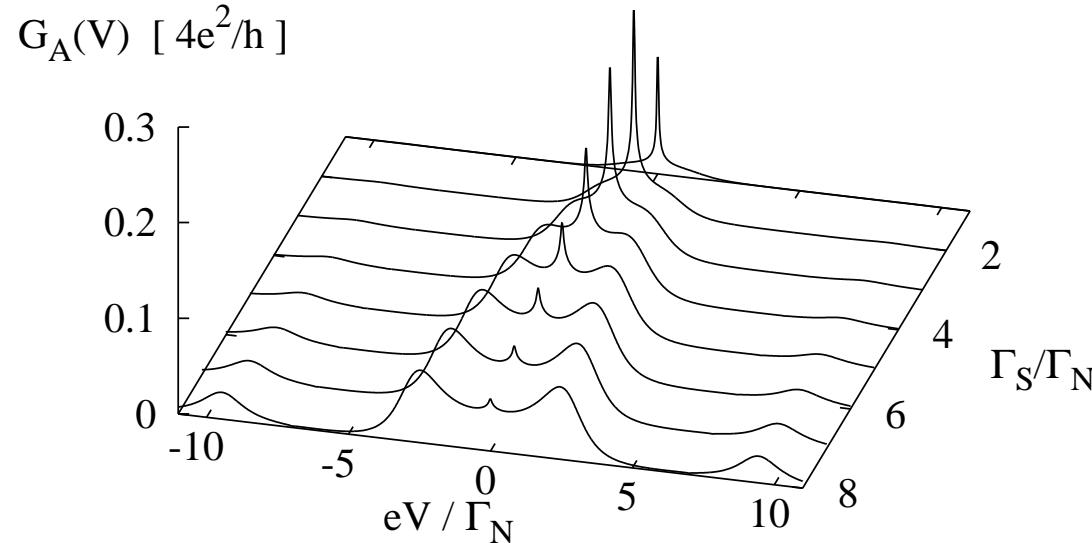
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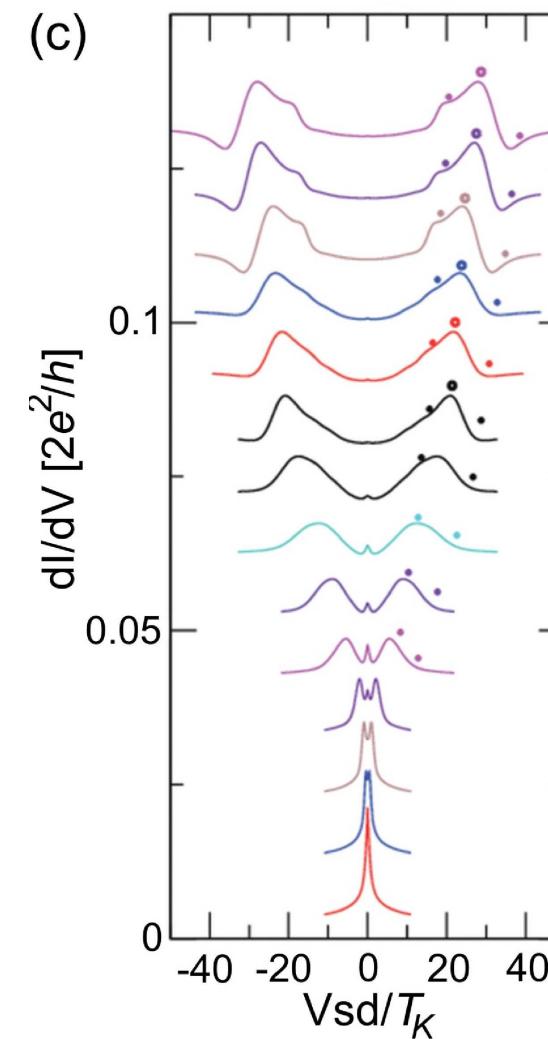
T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

Kondo state enhances the zero-bias conductance !

# Kondo effect vs induced pairing



**Theory:** T. Domański, A. Donabidowicz, PRB **78**, 073105 (2008).



**Experiment:** E.J.H. Lee *et al*, Phys. Rev. Lett. **109**, 186802 (2012).

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