

*Kraków, 4 marca 2013 r.*

# **Procesy Andreeva w silnie skorelowanych układach fermionowych**

**T. Domański**

**Uniwersytet Marii Curie–Skłodowskiej  
w Lublinie**

**<http://kft.umcs.lublin.pl/doman/lectures>**

## *Outline*

## *Outline*

### 1. **Introduction**

*/ underlying idea /*

## Outline

### 1. Introduction

*/ underlying idea /*

### 2. Andreev transport via quantum dots

*/ correlations versus superconductivity /*

## Outline

### 1. Introduction

*/ underlying idea /*

### 2. Andreev transport via quantum dots

*/ correlations versus superconductivity /*

### 3. Further extensions

*/ quantum interference, dephasing, Cooper splitting, etc /*

## Outline

### 1. **Introduction**

*/ underlying idea /*

### 2. **Andreev transport via quantum dots**

*/ correlations versus superconductivity /*

### 3. **Further extensions**

*/ quantum interference, dephasing, Cooper splitting, etc /*

### 4. **Andreev spectroscopy in bulk superconductors**

*/ probing the pair coherence /*

## Outline

### 1. Introduction

*/ underlying idea /*

### 2. Andreev transport via quantum dots

*/ correlations versus superconductivity /*

### 3. Further extensions

*/ quantum interference, dephasing, Cooper splitting, etc /*

### 4. Andreev spectroscopy in bulk superconductors

*/ probing the pair coherence /*

### 5. Andreev scattering in ultracold gasses

*/ fermion vs molecular channels /*

# 1. Introduction



**Andreev reflections**

–

**the main concept**

## Andreev reflections

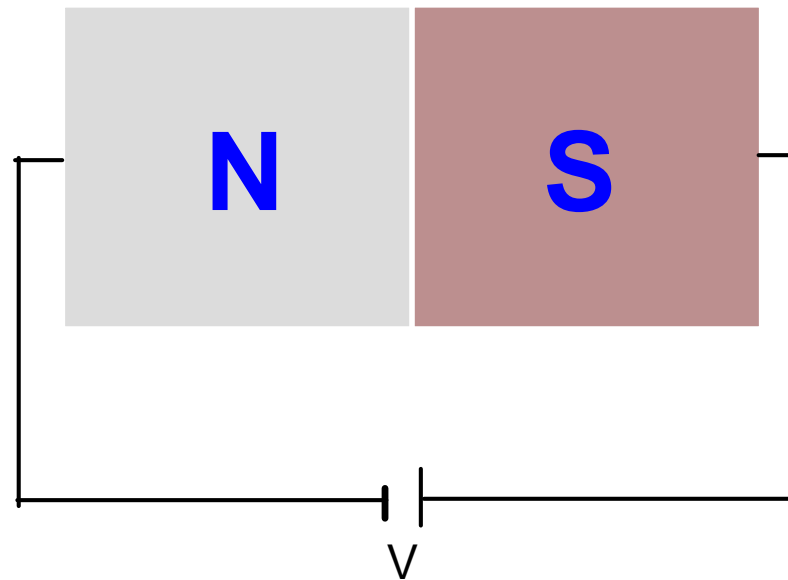
– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**

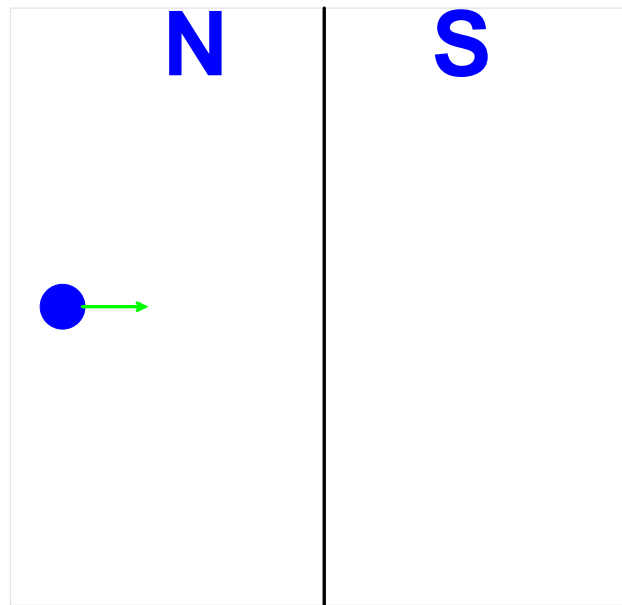


Let us restrict to the subgap regime  $|eV| \ll \Delta$  of an applied bias  $V$ .

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**

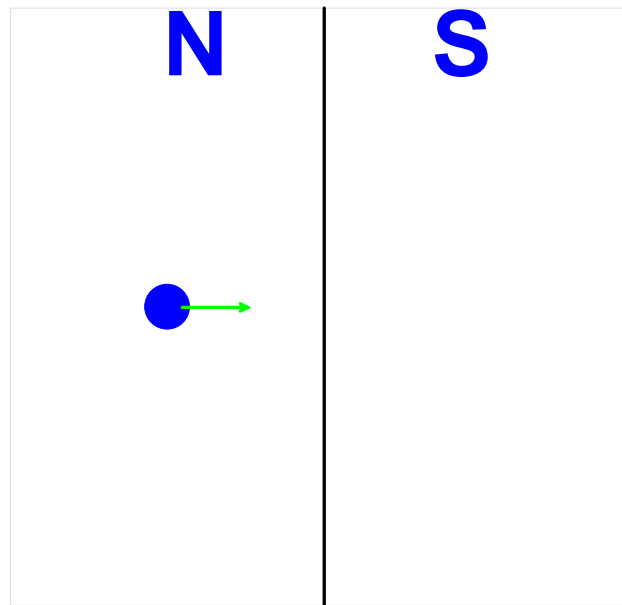


electron

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**

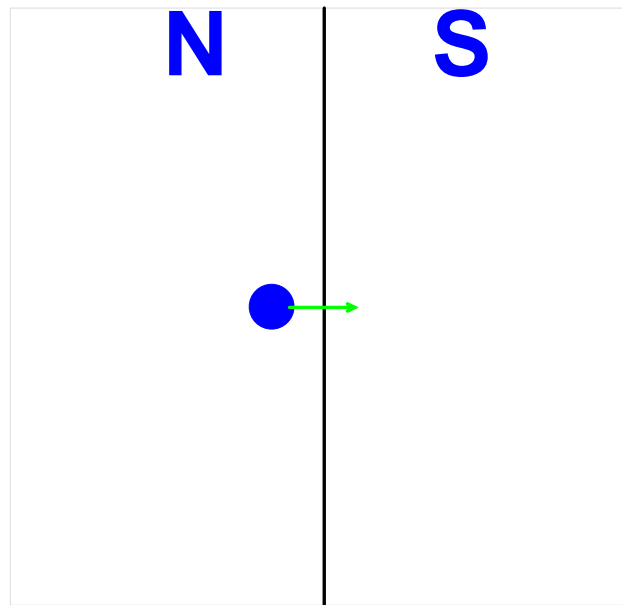


electron

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**

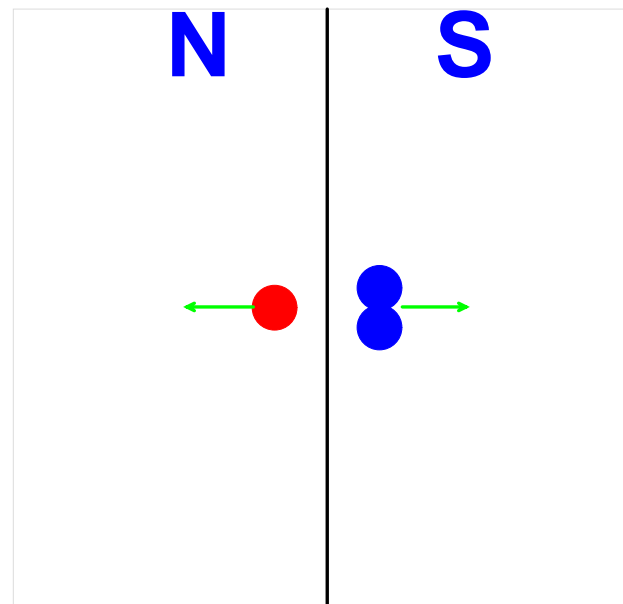


electron

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**



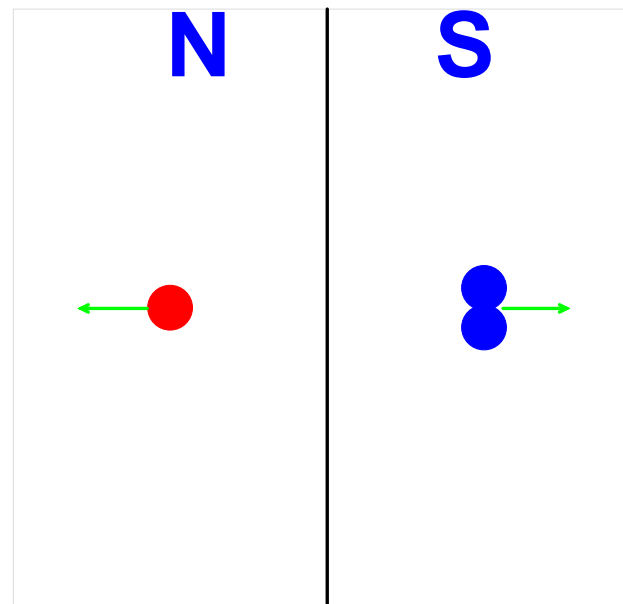
hole

Cooper pair

## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**



hole

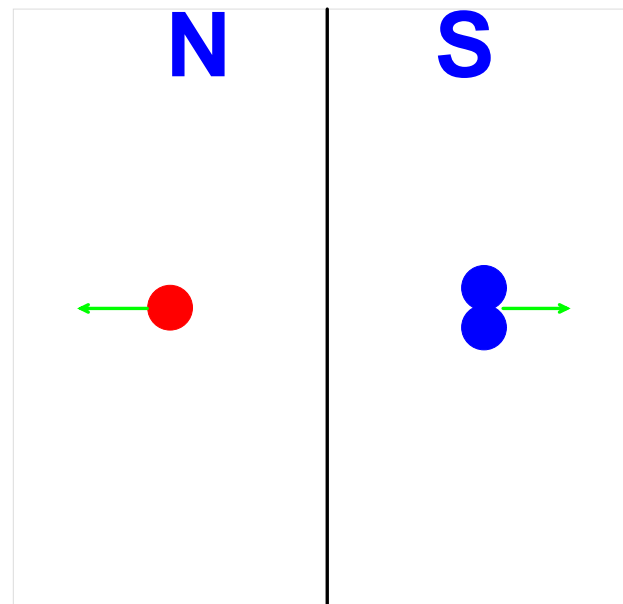
Cooper pair



## Andreev reflections

– the main concept

Let us consider the process of electron tunneling from the normal conductor **N** (e.g. metallic lead) to the superconducting electrode **S**



hole

Cooper pair

Such double-charge exchange is named the Andreev reflection (scattering).

**Andreev reflections**

–

**historical remark**

## Andreev reflections

–

## historical remark

This *anomalous* transport channel allows for a finite subgap current across N-S interface even though the single-particle transmissions are forbidden. Its original idea has been suggested by

## Andreev reflections

– historical remark

This *anomalous* transport channel allows for a finite subgap current across N-S interface even though the single-particle transmissions are forbidden. Its original idea has been suggested by



**A.F. Andreev**

/ P. Kapitza Institute, Moscow (Russia) /

*A.F. Andreev, Sov. Phys. JETP* **19**, 1228 (1964).

## **2. Andreev transport via quantum dot**

## **Physical situation**

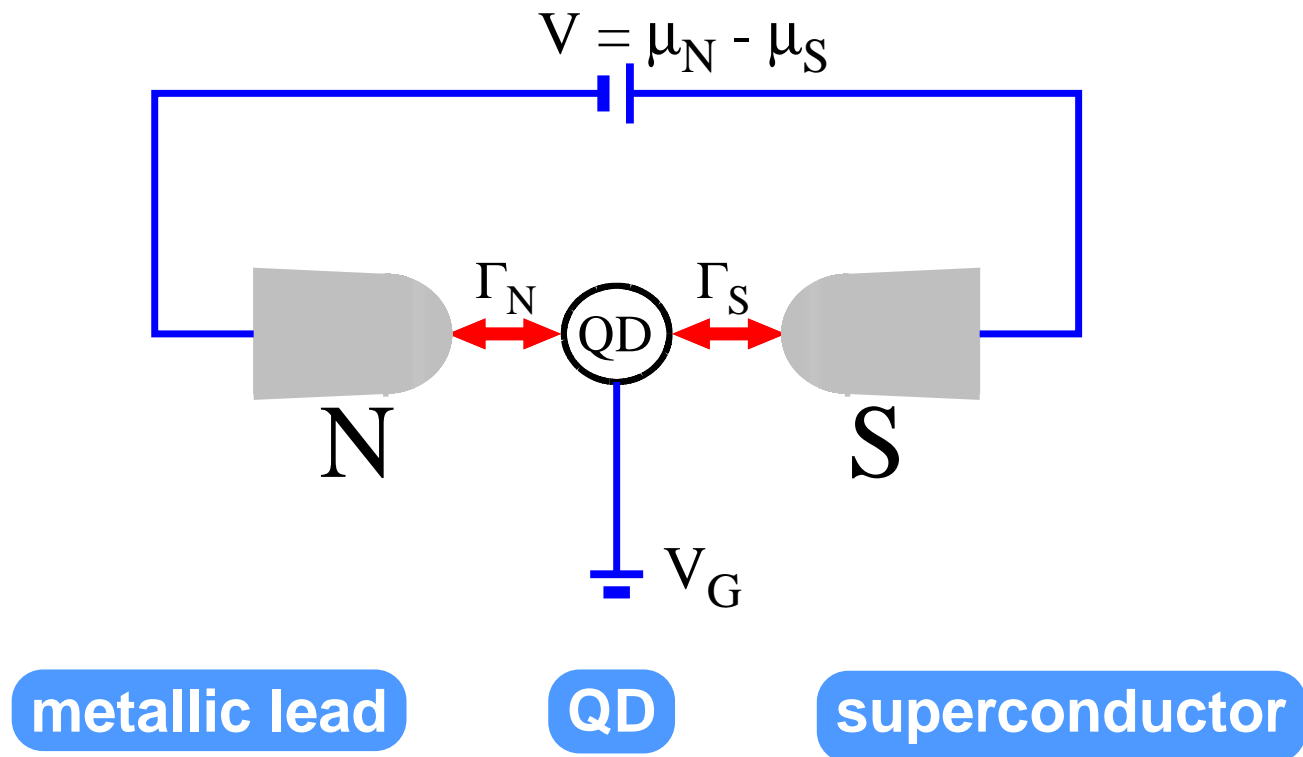
## **N-QD-S scheme**

**Let us consider the quantum dot (QD) on an interface between the external metallic (N) and superconducting (S) leads**

## Physical situation

## N-QD-S scheme

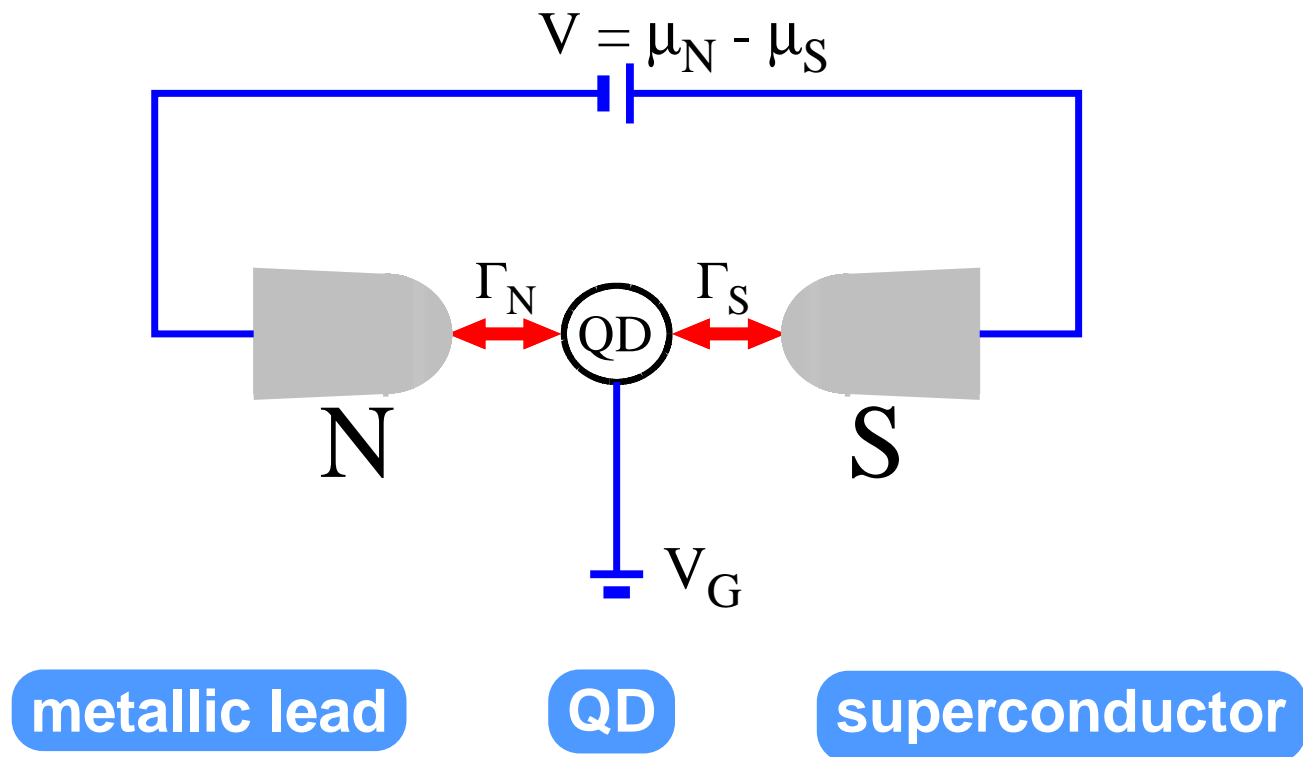
Let us consider the quantum dot (QD) on an interface between the external metallic (N) and superconducting (S) leads



## Physical situation

## N-QD-S scheme

Let us consider the quantum dot (QD) on an interface between the external metallic (N) and superconducting (S) leads



This setup can be thought of as a particular version of the SET.



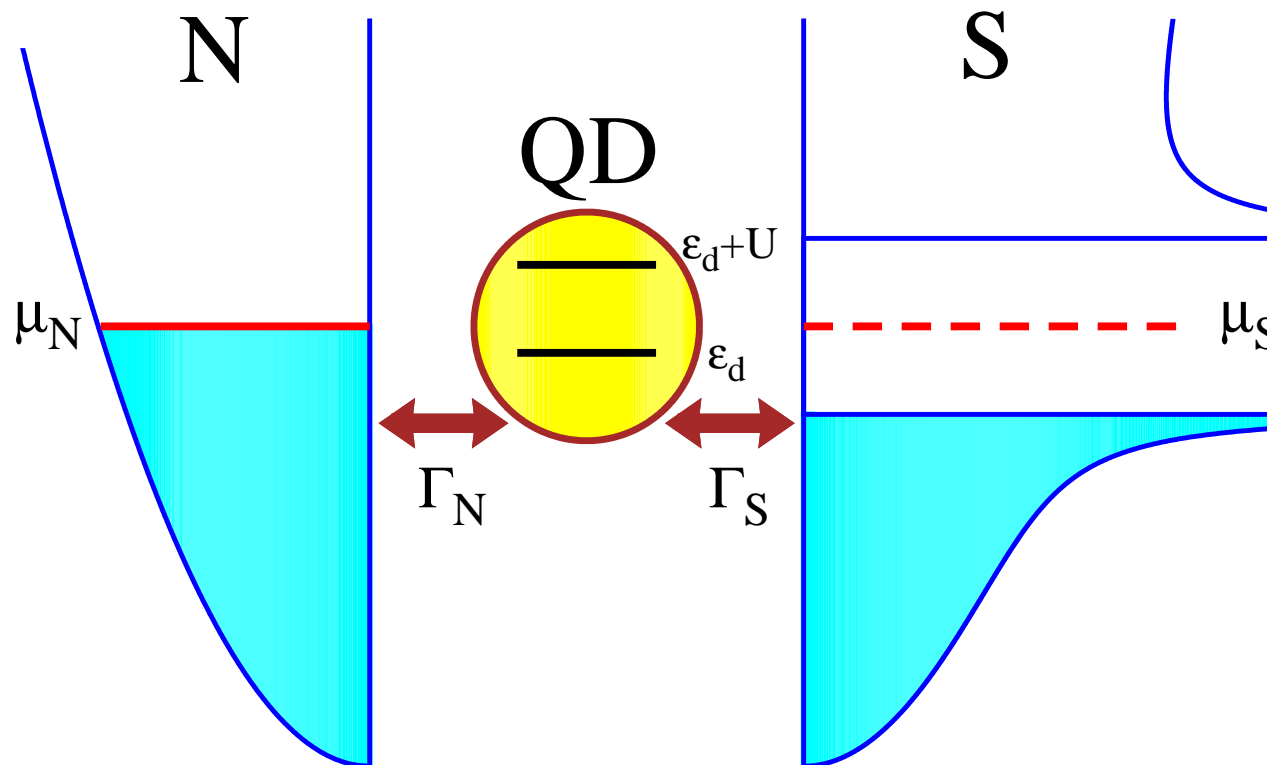
**Physical situation** – energy spectrum

## **Physical situation** – energy spectrum

**Components of the N-QD-S heterostructure have the following spectra**

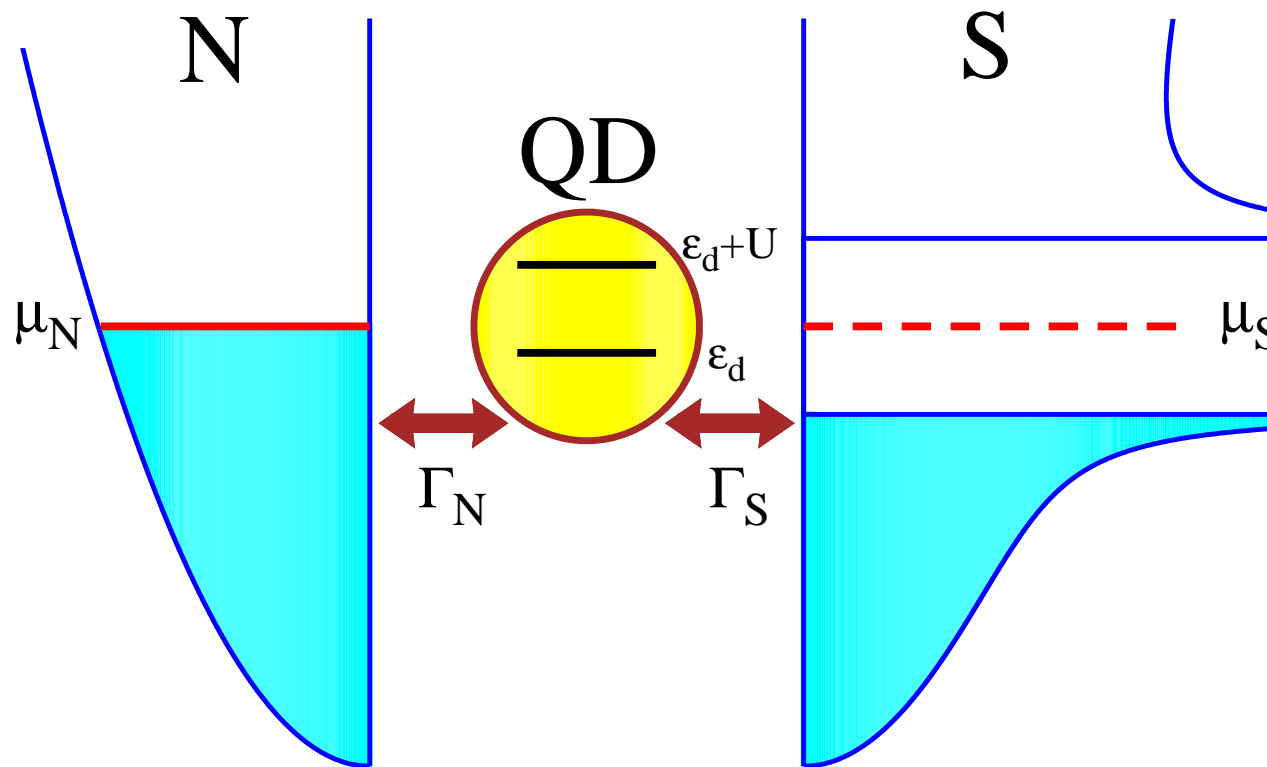
## Physical situation – energy spectrum

Components of the N-QD-S heterostructure have the following spectra



## Physical situation – energy spectrum

Components of the N-QD-S heterostructure have the following spectra



External bias  $eV = \mu_N - \mu_S$  induces the current(s) through QD.

# Microscopic model

## The correlation effects

## Microscopic model

The correlation effects

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

## Microscopic model

The correlation effects

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

are expected to affect the transport properties of the system

## Microscopic model

The correlation effects

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

are expected to affect the transport properties of the system

$$\begin{aligned} \hat{H} = & \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} + \hat{H}_N + \hat{H}_S \\ & + \sum_{\mathbf{k}, \sigma} \sum_{\beta=N, S} \left( V_{\mathbf{k}\beta} \hat{d}_{\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma\beta} + V_{\mathbf{k}\beta}^* \hat{c}_{\mathbf{k}\sigma, \beta}^{\dagger} \hat{d}_{\sigma} \right) \end{aligned}$$



## Microscopic model

The correlation effects

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

are expected to affect the transport properties of the system

$$\begin{aligned} \hat{H} = & \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} + \hat{H}_N + \hat{H}_S \\ & + \sum_{k,\sigma} \sum_{\beta=N,S} \left( V_{k\beta} \hat{d}_{\sigma}^{\dagger} \hat{c}_{k\sigma\beta} + V_{k\beta}^* \hat{c}_{k\sigma,\beta}^{\dagger} \hat{d}_{\sigma} \right) \end{aligned}$$

where

$$\hat{H}_N = \sum_{k,\sigma} (\epsilon_{k,N} - \mu_N) \hat{c}_{k\sigma N}^{\dagger} \hat{c}_{k\sigma N}$$

## Microscopic model

The correlation effects

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow}$$

are expected to affect the transport properties of the system

$$\begin{aligned} \hat{H} &= \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} + \hat{H}_N + \hat{H}_S \\ &+ \sum_{k,\sigma} \sum_{\beta=N,S} \left( V_{k\beta} \hat{d}_{\sigma}^{\dagger} \hat{c}_{k\sigma\beta} + V_{k\beta}^* \hat{c}_{k\sigma,\beta}^{\dagger} \hat{d}_{\sigma} \right) \end{aligned}$$

where

$$\hat{H}_S = \sum_{k,\sigma} (\epsilon_{k,S} - \mu_S) \hat{c}_{k\sigma S}^{\dagger} \hat{c}_{k\sigma S} - \sum_k \left( \Delta \hat{c}_{k\uparrow S}^{\dagger} \hat{c}_{k\downarrow S}^{\dagger} + \text{h.c.} \right)$$

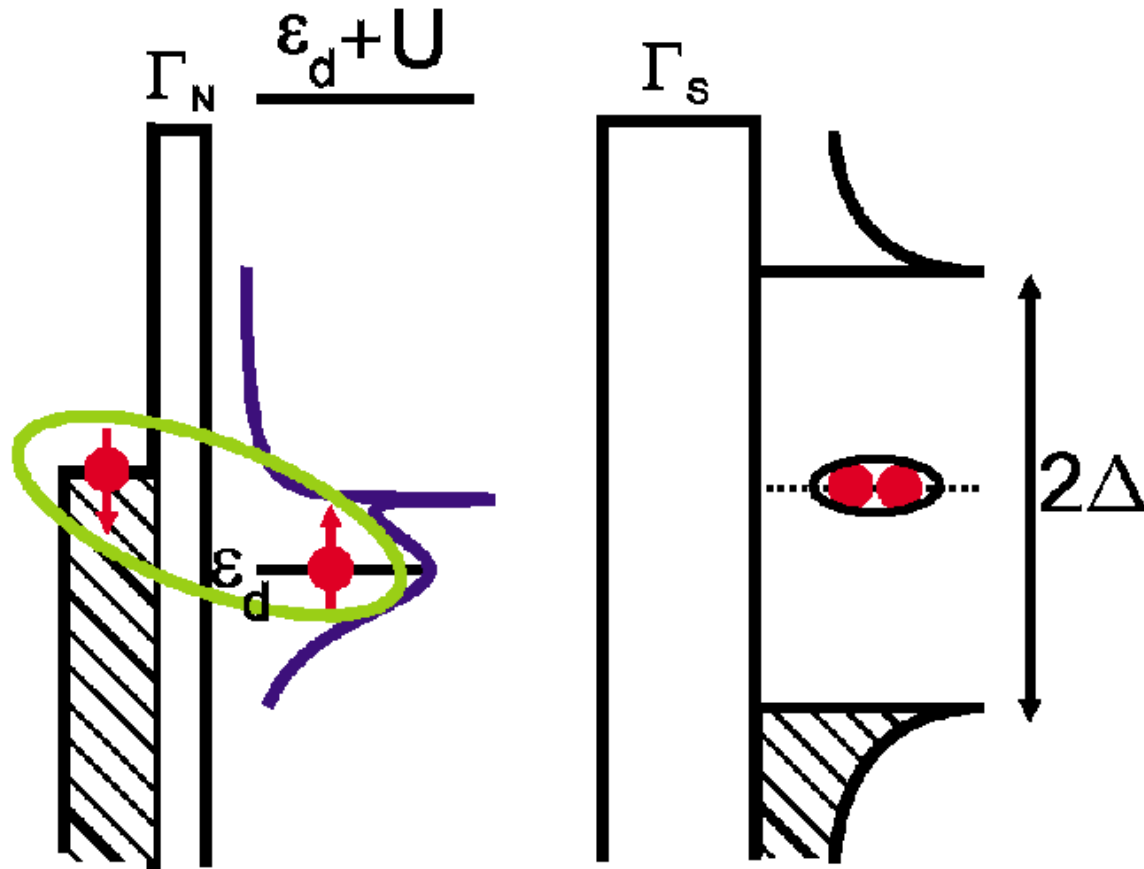
**Relevant problems :** **issue # 1**

## **Relevant problems :** issue # 1

Hybridization of QD to the metallic lead is responsible for:

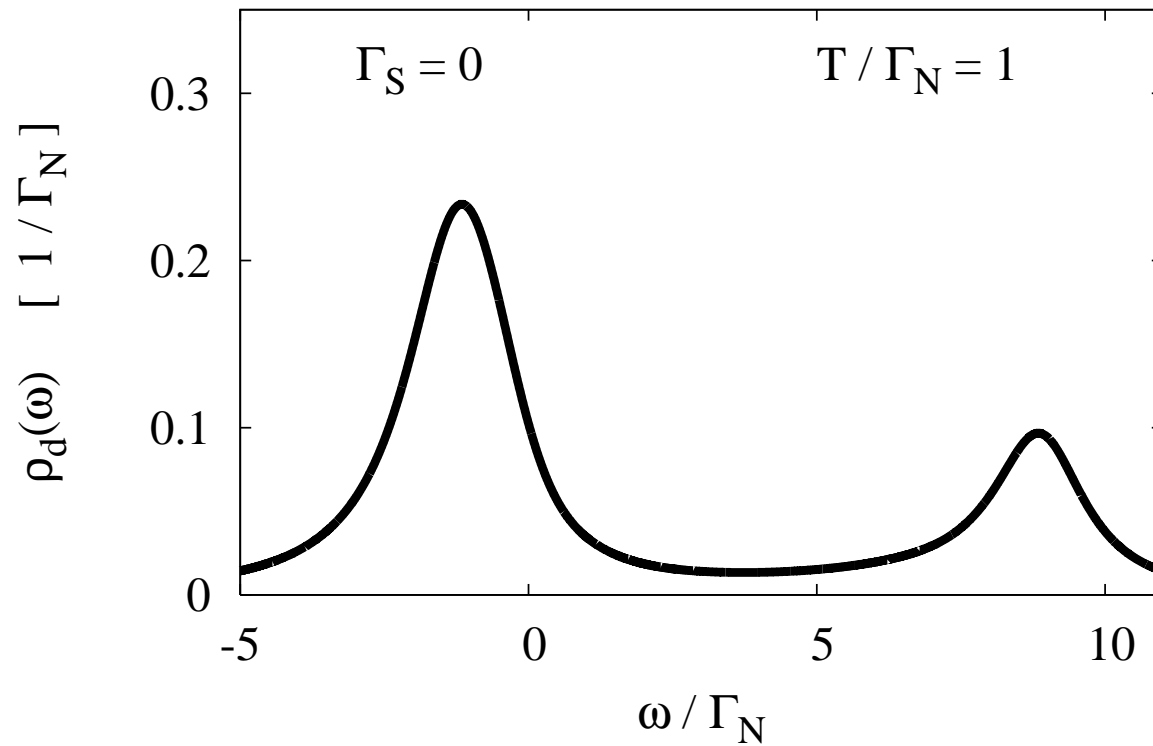
## Relevant problems : issue # 1

Hybridization of QD to the metallic lead is responsible for:



## Relevant problems : issue # 1

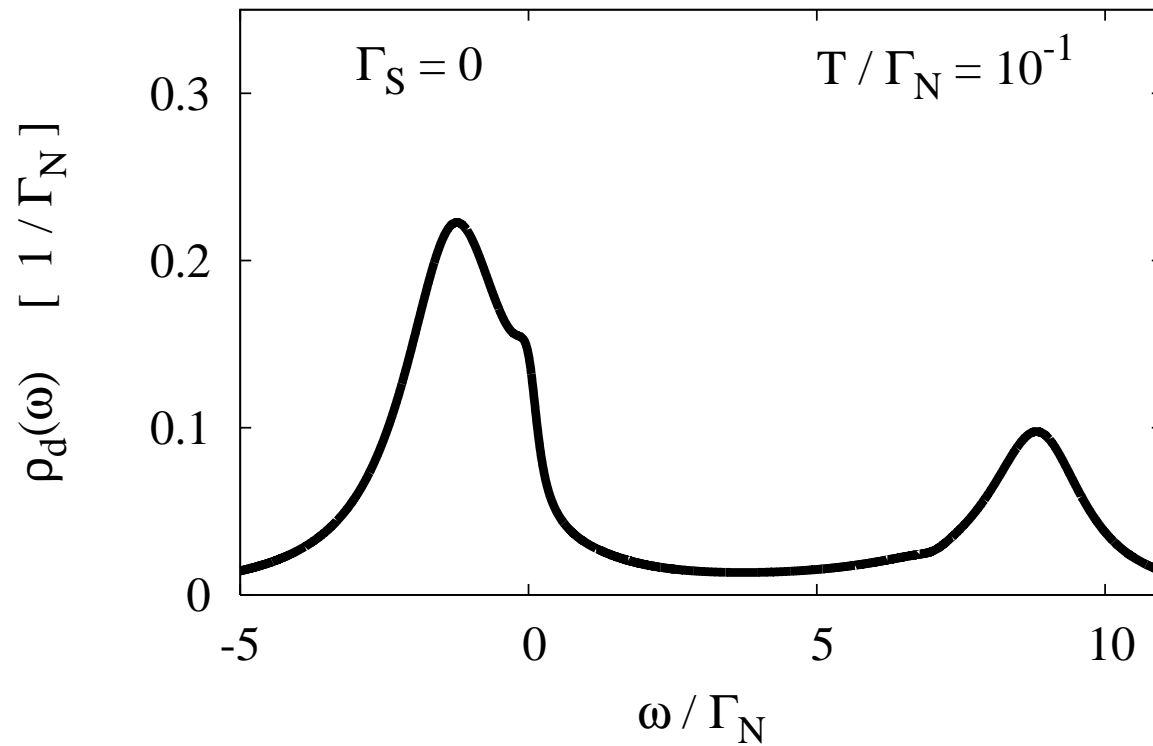
Hybridization of QD to the metallic lead is responsible for:



a broadening of QD levels

## Relevant problems : issue # 1

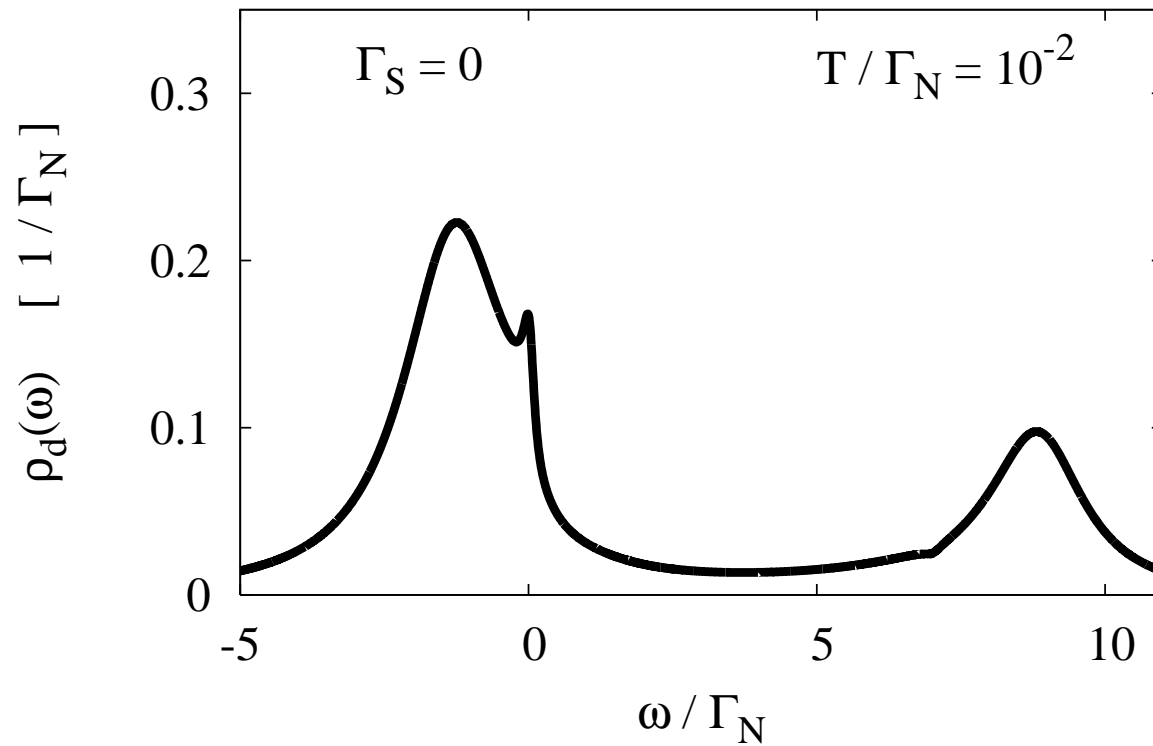
Hybridization of QD to the metallic lead is responsible for:



a broadening of QD levels and ...

## Relevant problems : issue # 1

Hybridization of QD to the metallic lead is responsible for:

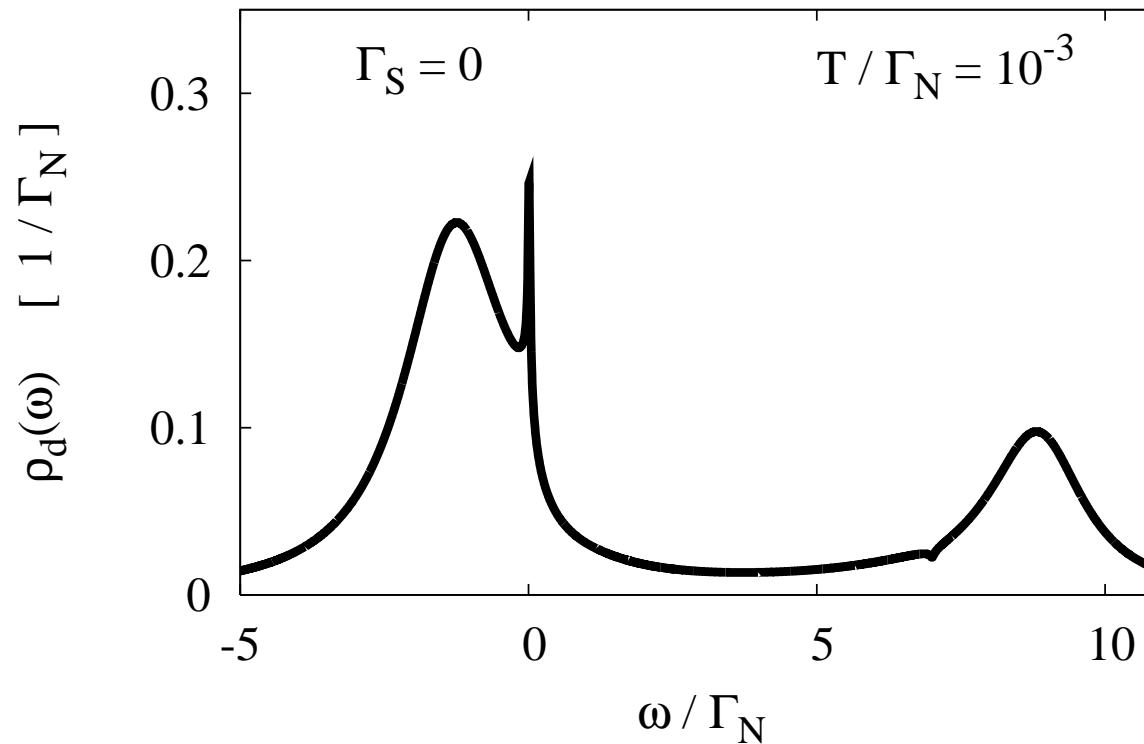


a broadening of QD levels and ...



## Relevant problems : issue # 1

Hybridization of QD to the metallic lead is responsible for:

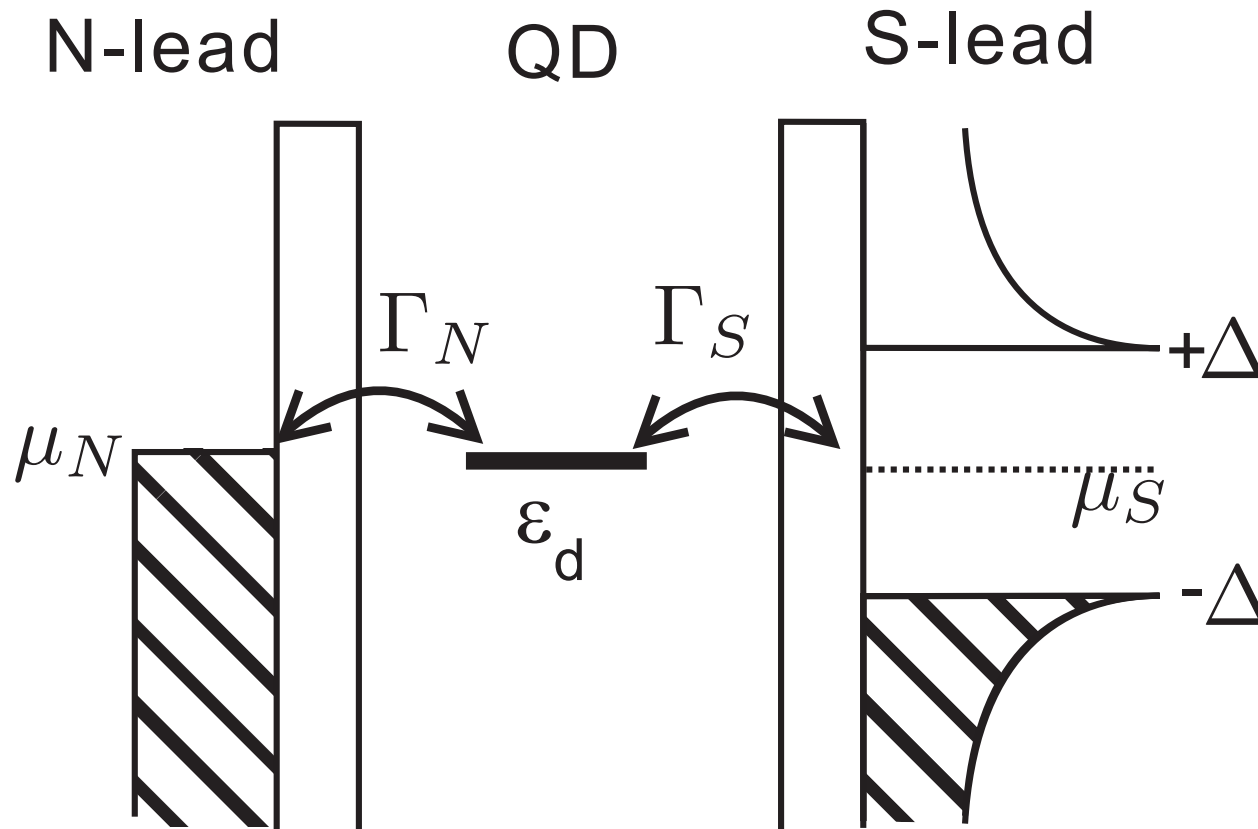


- ★ a broadening of QD levels and
- ★ appearance of the Kondo resonance below  $T_K$ .

**Relevant problems :** **issue # 2**

## Relevant problems : issue # 2

### Hybridization of QD to the superconducting lead



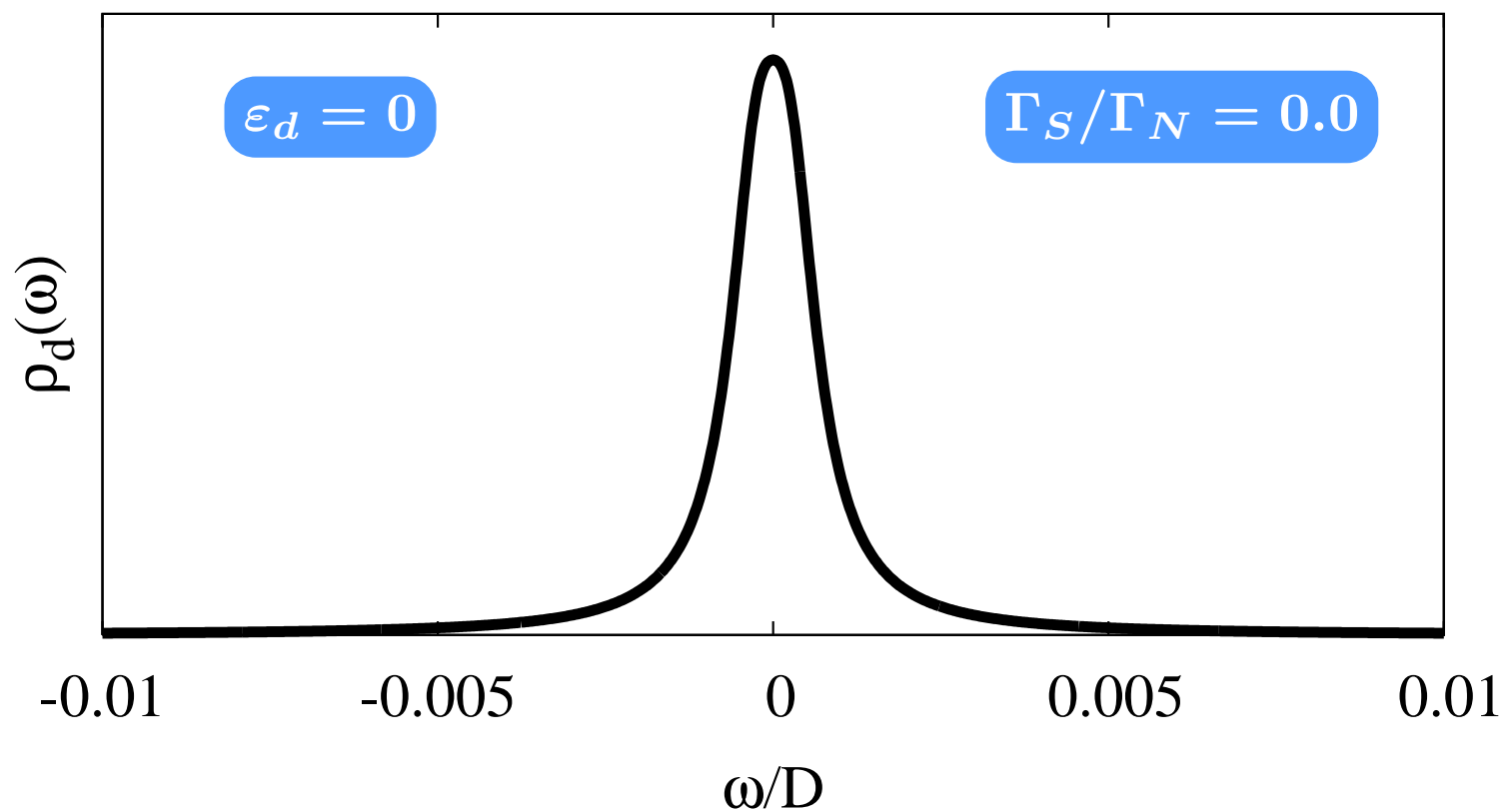
## Relevant problems : issue # 2

Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

## Relevant problems : issue # 2

Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

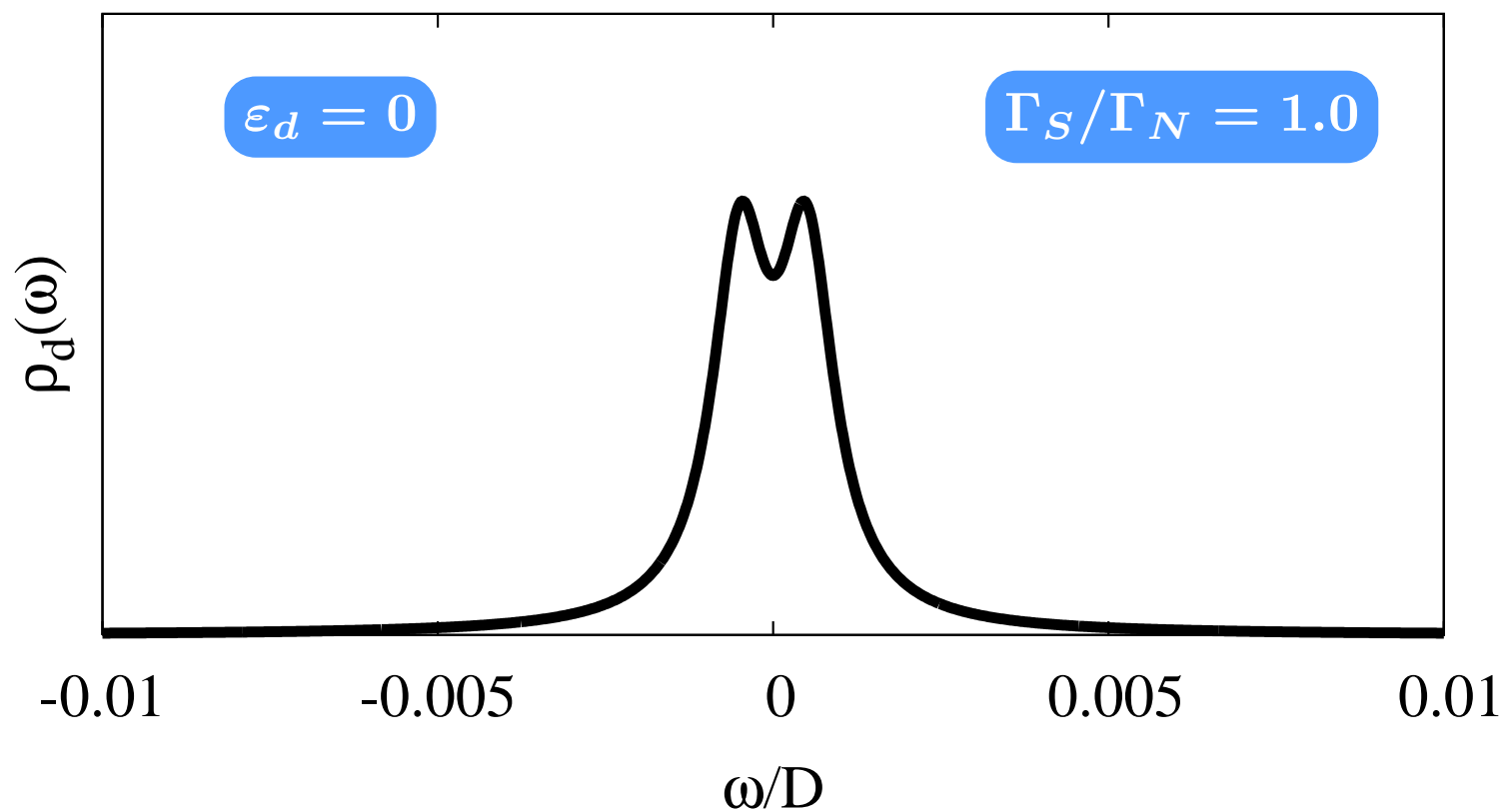
$$U = 0$$



## Relevant problems : issue # 2

Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

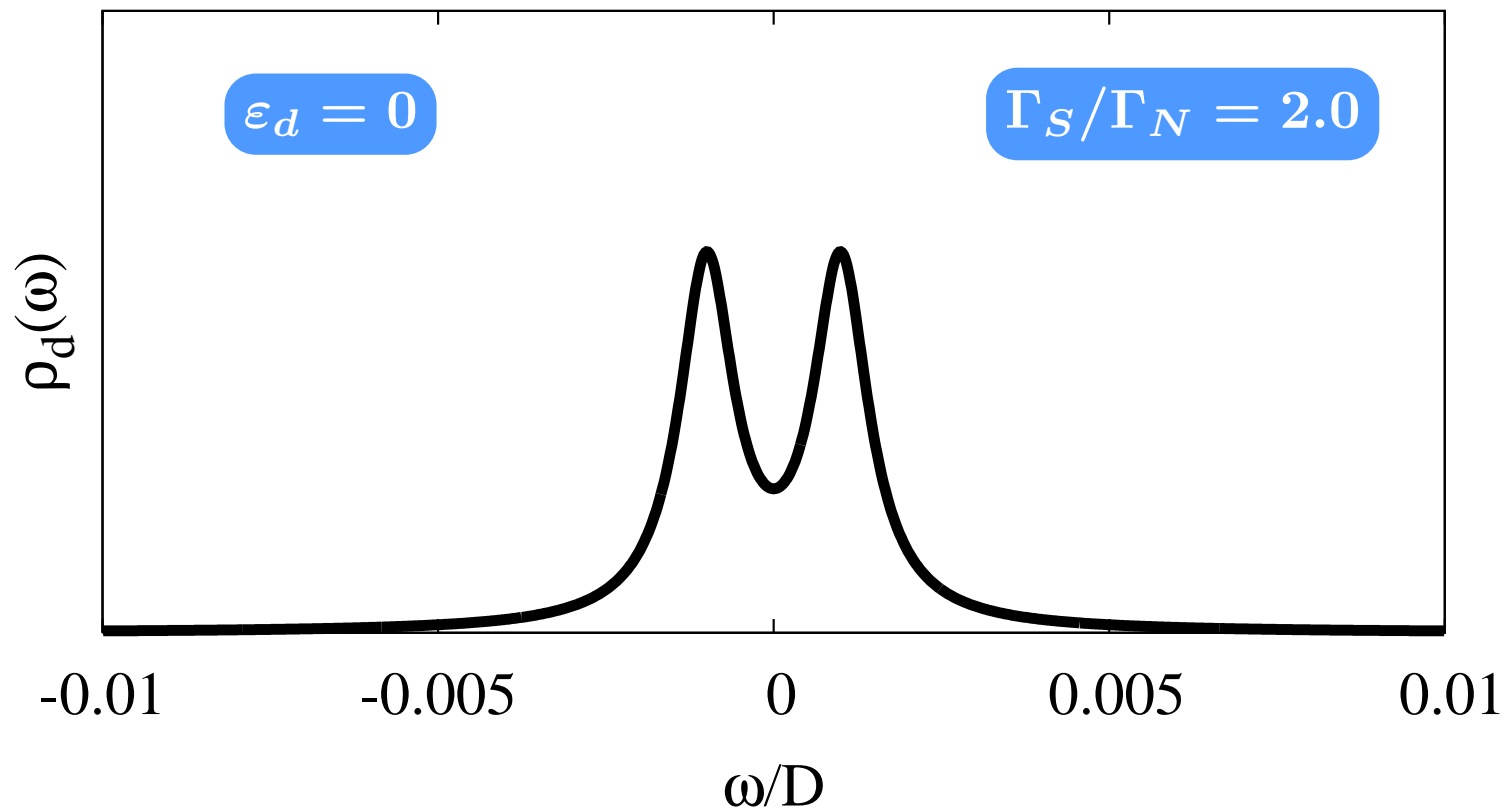
$$U = 0$$



## Relevant problems : issue # 2

Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

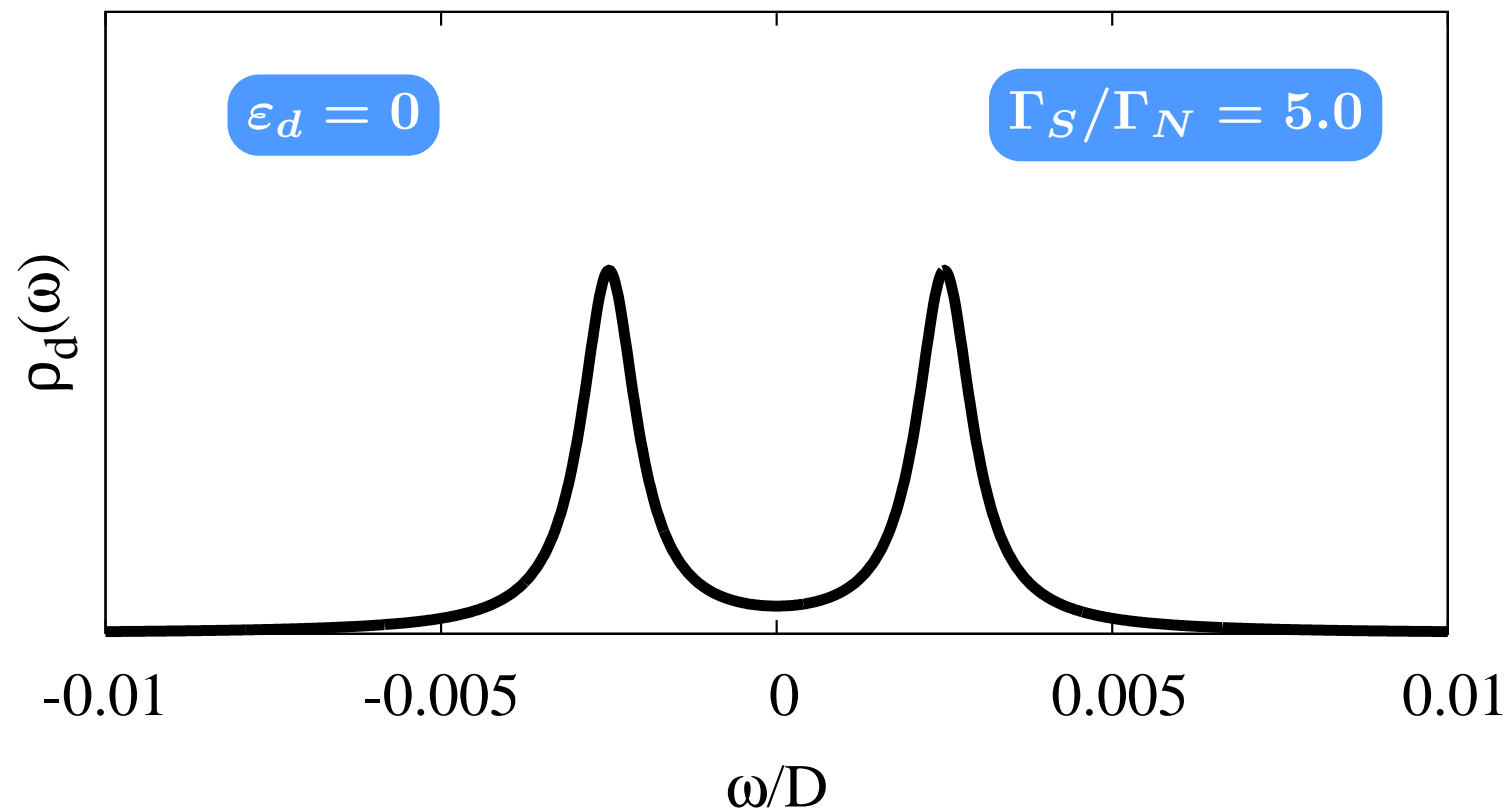
$$U = 0$$



## Relevant problems : issue # 2

Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

$$U = 0$$

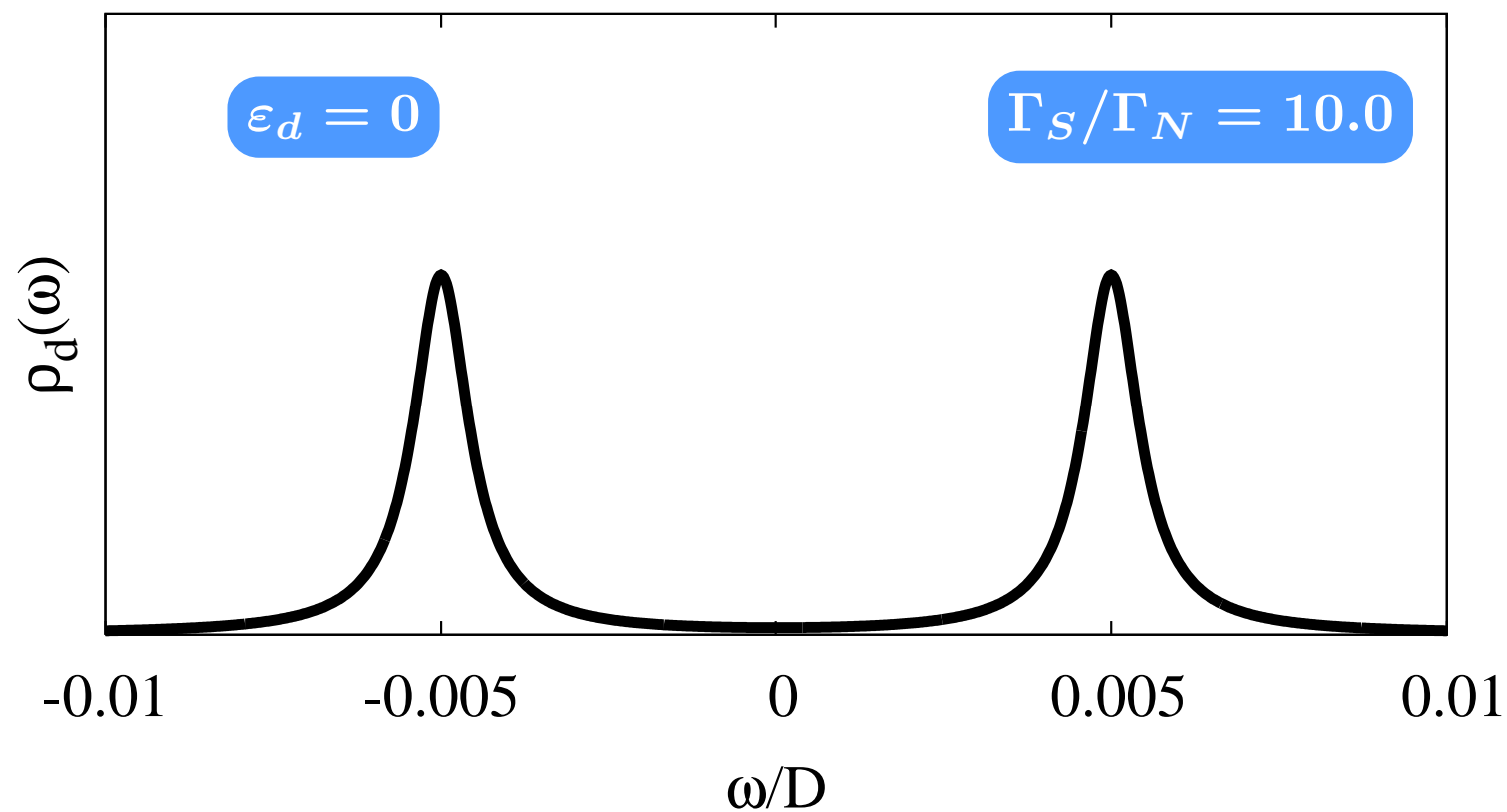




## Relevant problems : issue # 2

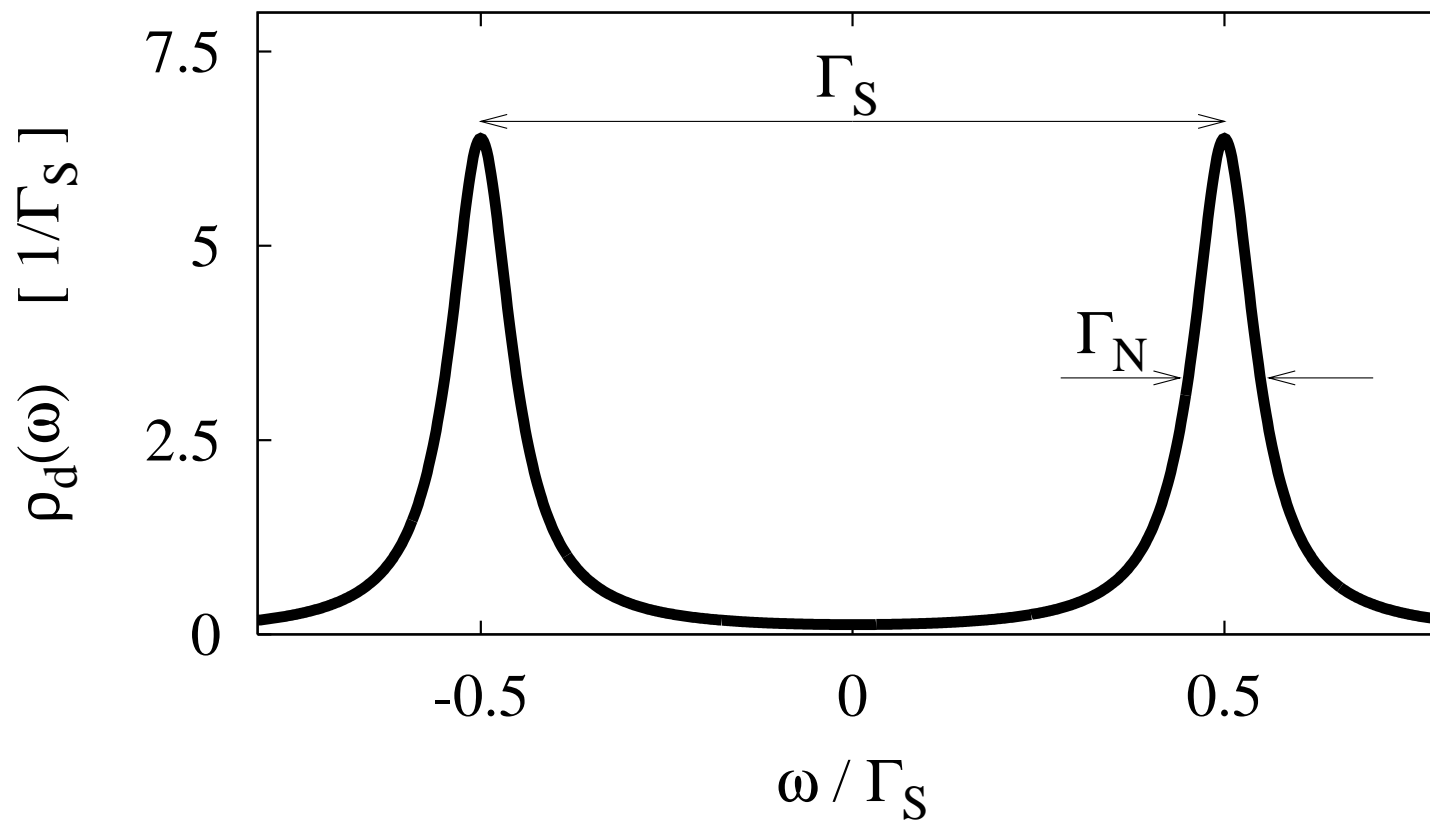
Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).

$$U = 0$$



## Relevant problems : issue # 2

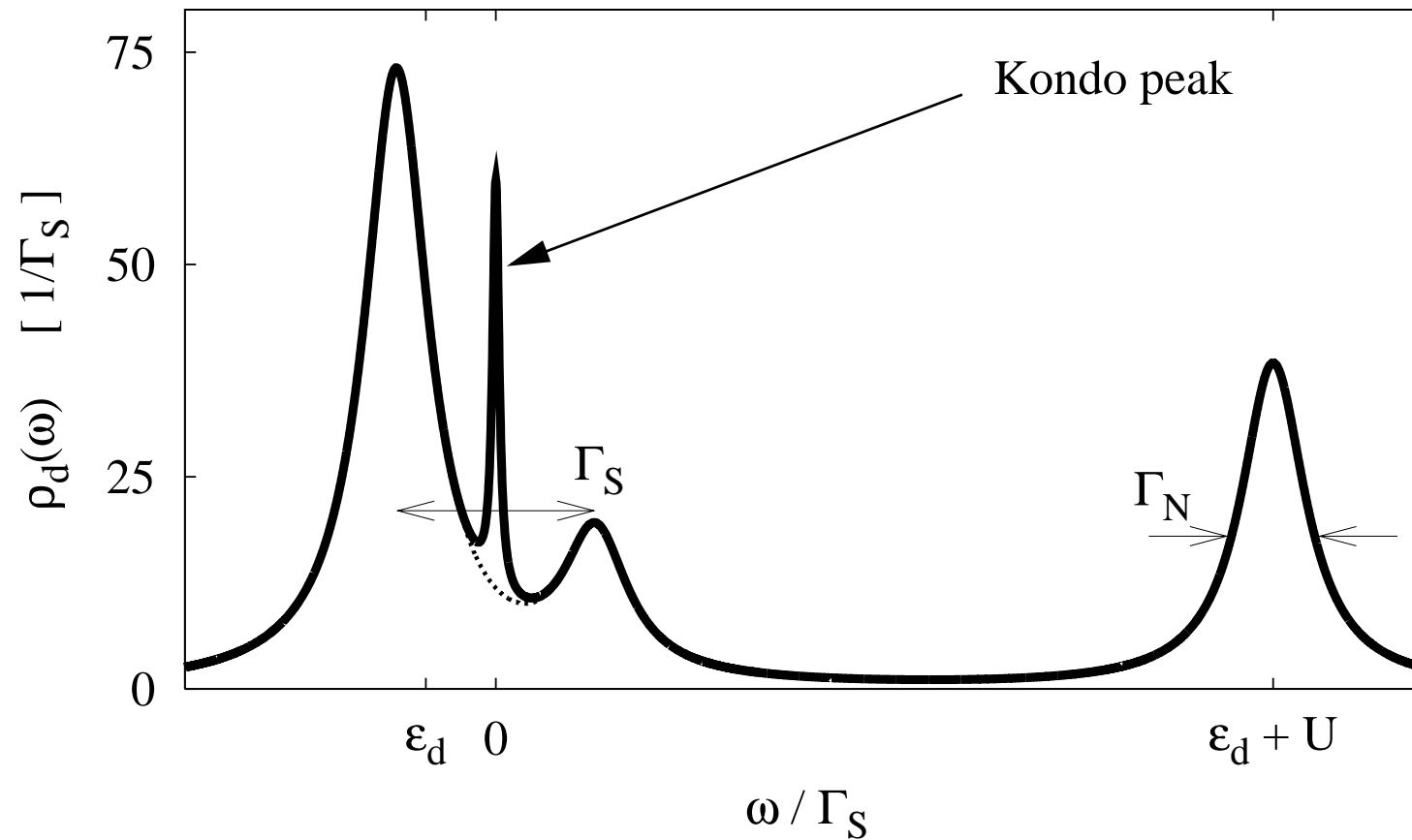
Hybridization of QD to S transmits the **on-dot pairing** (*proximity effect*).



**Relevant problems :** # 1 + 2

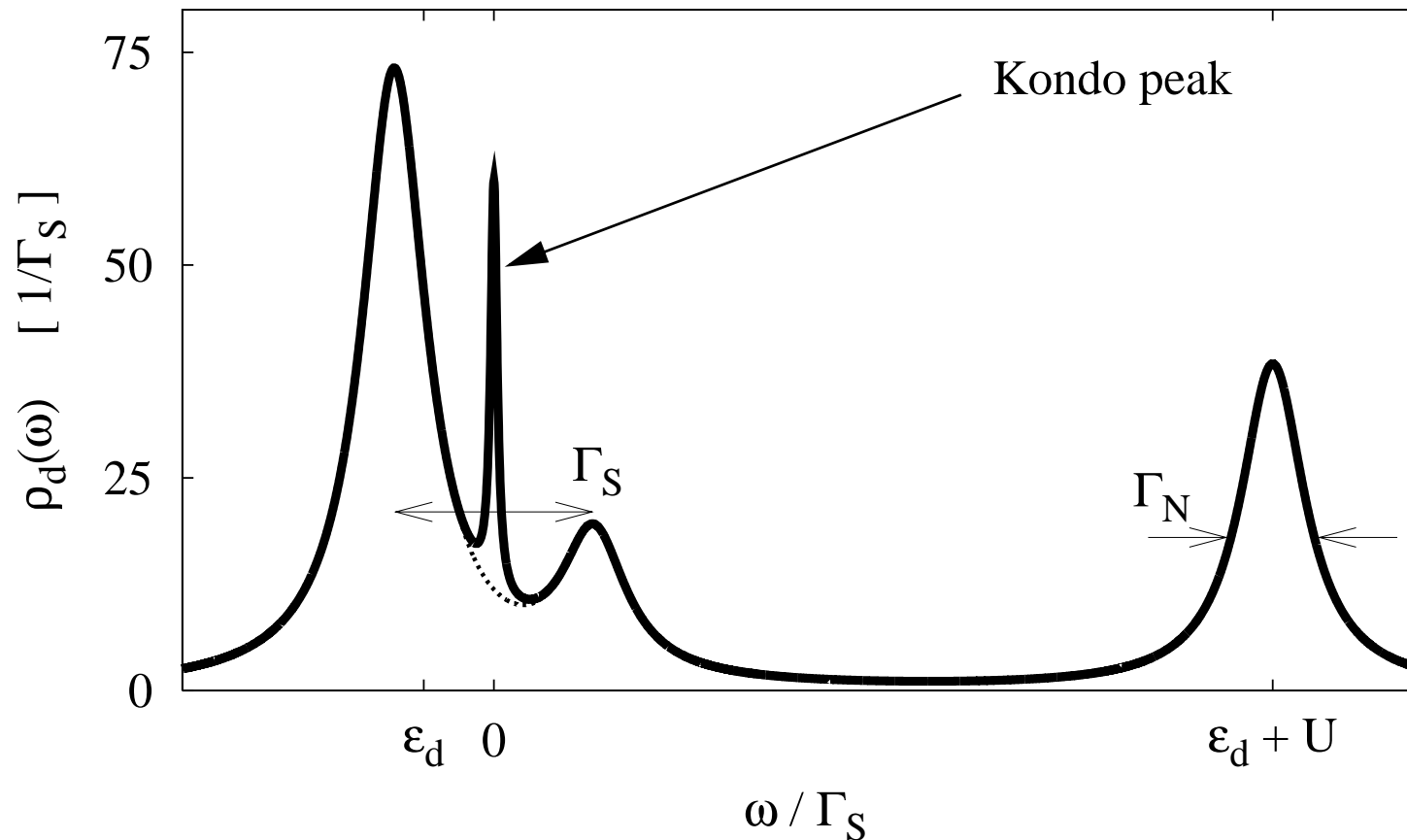
Relevant problems : # 1 + 2

Hybridizations  $\Gamma_N$  and  $\Gamma_S$  are thus effectively leading to



Relevant problems : # 1 + 2

Hybridizations  $\Gamma_N$  and  $\Gamma_S$  are thus effectively leading to



/ interplay between the Kondo effect and superconductivity /

**Questions:**

## Questions:

- ★ What relation does occur between superconductivity (transmitted onto the QD) and the Kondo effect ?

## Questions:

★ What relation does occur between superconductivity (transmitted onto the QD) and the Kondo effect ?

Do they coexist or compete ?



## Questions:

- ★ What relation does occur between superconductivity (transmitted onto the QD) and the Kondo effect ?

Do they coexist or compete ?

- ★ How do these effects show up in the charge current through N-QD-S junction ?

## Questions:

- ★ What relation does occur between superconductivity (transmitted onto the QD) and the Kondo effect ?

Do they coexist or compete ?

- ★ How do these effects show up in the charge current through N-QD-S junction ?

Are there any particular features ?

## Formal aspects

## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

$$G_d(\tau, \tau') = - \begin{pmatrix} \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\downarrow(\tau') \rangle \\ \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\downarrow(\tau') \rangle \end{pmatrix}$$

## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

$$G_d(\tau, \tau') = - \begin{pmatrix} \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\downarrow(\tau') \rangle \\ \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\downarrow(\tau') \rangle \end{pmatrix}$$

In equilibrium its Fourier transform obeys the Dyson equation

## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

$$G_d(\tau, \tau') = - \begin{pmatrix} \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\downarrow(\tau') \rangle \\ \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\downarrow(\tau') \rangle \end{pmatrix}$$

In equilibrium its Fourier transform obeys the Dyson equation

$$G_d(\omega)^{-1} = \begin{pmatrix} \omega - \varepsilon_d & 0 \\ 0 & \omega + \varepsilon_d \end{pmatrix} - \Sigma_d^0(\omega) - \Sigma_d^U(\omega)$$

## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

$$G_d(\tau, \tau') = - \begin{pmatrix} \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\downarrow(\tau') \rangle \\ \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\downarrow(\tau') \rangle \end{pmatrix}$$

In equilibrium its Fourier transform obeys the Dyson equation

$$G_d(\omega)^{-1} = \begin{pmatrix} \omega - \varepsilon_d & 0 \\ 0 & \omega + \varepsilon_d \end{pmatrix} - \Sigma_d^0(\omega) - \Sigma_d^U(\omega)$$

with

$$\Sigma_d^0(\omega) \quad \text{the selfenergy for } U = 0$$



## Formal aspects

To account for both, the proximity effect and the correlations, we have to deal with the Nambu ( $2 \times 2$  matrix) Green's function

$$G_d(\tau, \tau') = - \begin{pmatrix} \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\uparrow(\tau) \hat{d}_\downarrow(\tau') \rangle \\ \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\uparrow^\dagger(\tau') \rangle & \hat{T}_\tau \langle \hat{d}_\downarrow^\dagger(\tau) \hat{d}_\downarrow(\tau') \rangle \end{pmatrix}$$

In equilibrium its Fourier transform obeys the Dyson equation

$$G_d(\omega)^{-1} = \begin{pmatrix} \omega - \varepsilon_d & 0 \\ 0 & \omega + \varepsilon_d \end{pmatrix} - \Sigma_d^0(\omega) - \Sigma_d^U(\omega)$$

with

$$\Sigma_d^U(\omega) \quad \text{correction due to } U \neq 0.$$

## Non-equilibrium phenomena

## Non-equilibrium phenomena

The steady current  $J_L = -J_R$  is found to consist of two contributions

$$J(V) = J_1(V) + J_A(V)$$

## Non-equilibrium phenomena

The steady current  $J_L = -J_R$  is found to consist of two contributions

$$J(V) = J_1(V) + J_A(V)$$

which can be expressed by the Landauer-type formula

$$J_1(V) = \frac{2e}{h} \int d\omega \, T_1(\omega) [f(\omega + eV, T) - f(\omega, T)]$$

$$J_A(V) = \frac{2e}{h} \int d\omega \, T_A(\omega) [f(\omega + eV, T) - f(\omega - eV, T)]$$

with the transmittance

$$T_1(\omega) = \Gamma_N \Gamma_S \left( |G_{11}^r(\omega)|^2 + |G_{12}^r(\omega)|^2 - \frac{2\Delta}{|\omega|} \text{Re} G_{11}^r(\omega) G_{12}^r(\omega) \right)$$

## Non-equilibrium phenomena

The steady current  $J_L = -J_R$  is found to consist of two contributions

$$J(V) = J_1(V) + J_A(V)$$

which can be expressed by the Landauer-type formula

$$J_1(V) = \frac{2e}{h} \int d\omega \, T_1(\omega) [f(\omega + eV, T) - f(\omega, T)]$$

$$J_A(V) = \frac{2e}{h} \int d\omega \, T_A(\omega) [f(\omega + eV, T) - f(\omega - eV, T)]$$

with the transmittance

$$T_A(\omega) = \Gamma_N^2 |G_{12}(\omega)|^2$$

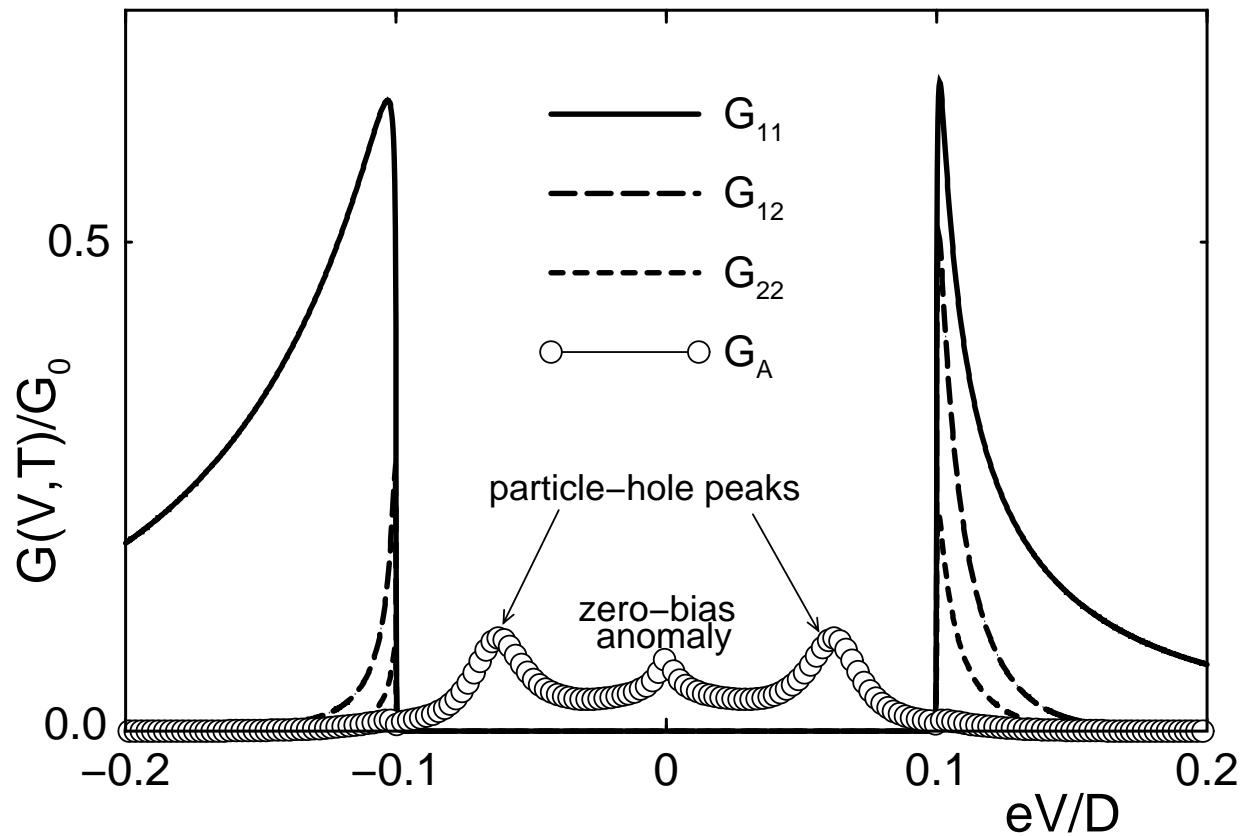
## Transport channels

## Transport channels

Qualitative features in the differential conductance  $G(V) = \frac{\partial J(V)}{\partial V}$

## Transport channels

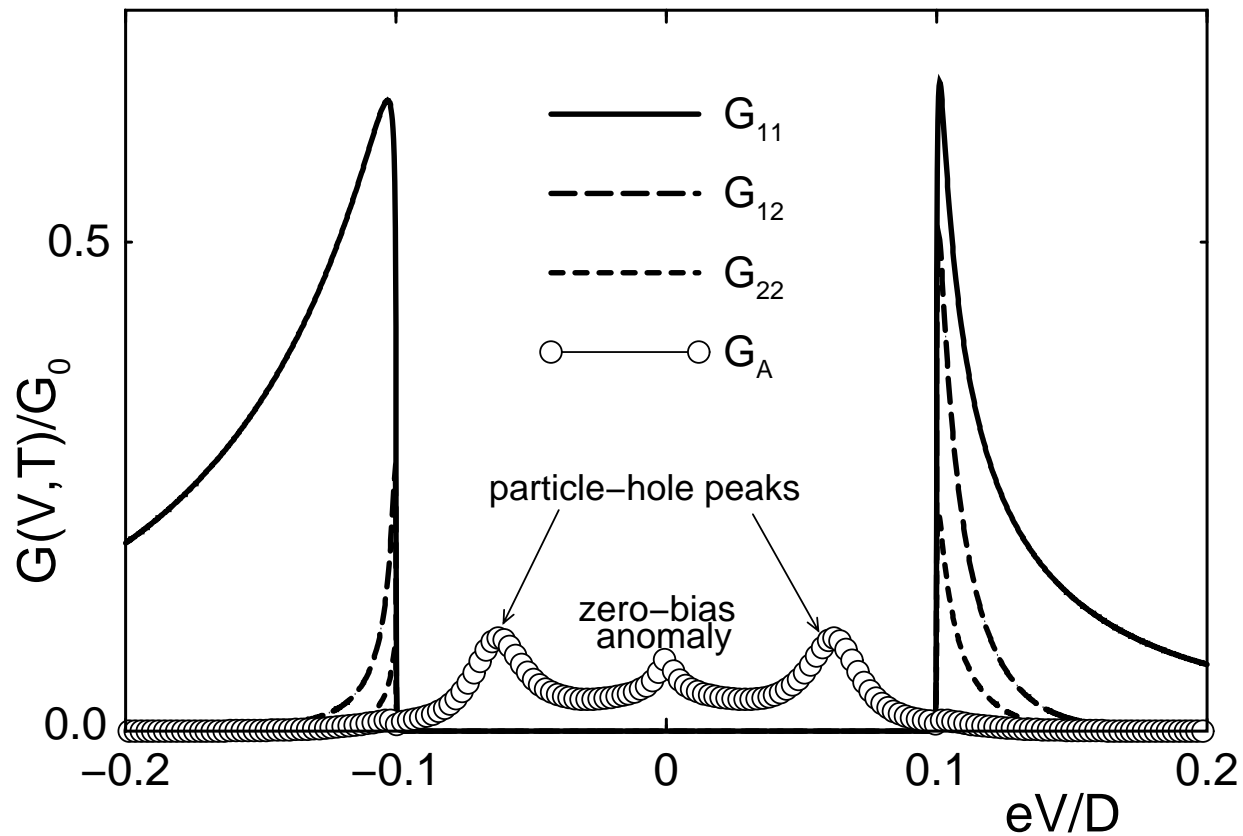
Qualitative features in the differential conductance  $G(V) = \frac{\partial J(V)}{\partial V}$





## Transport channels

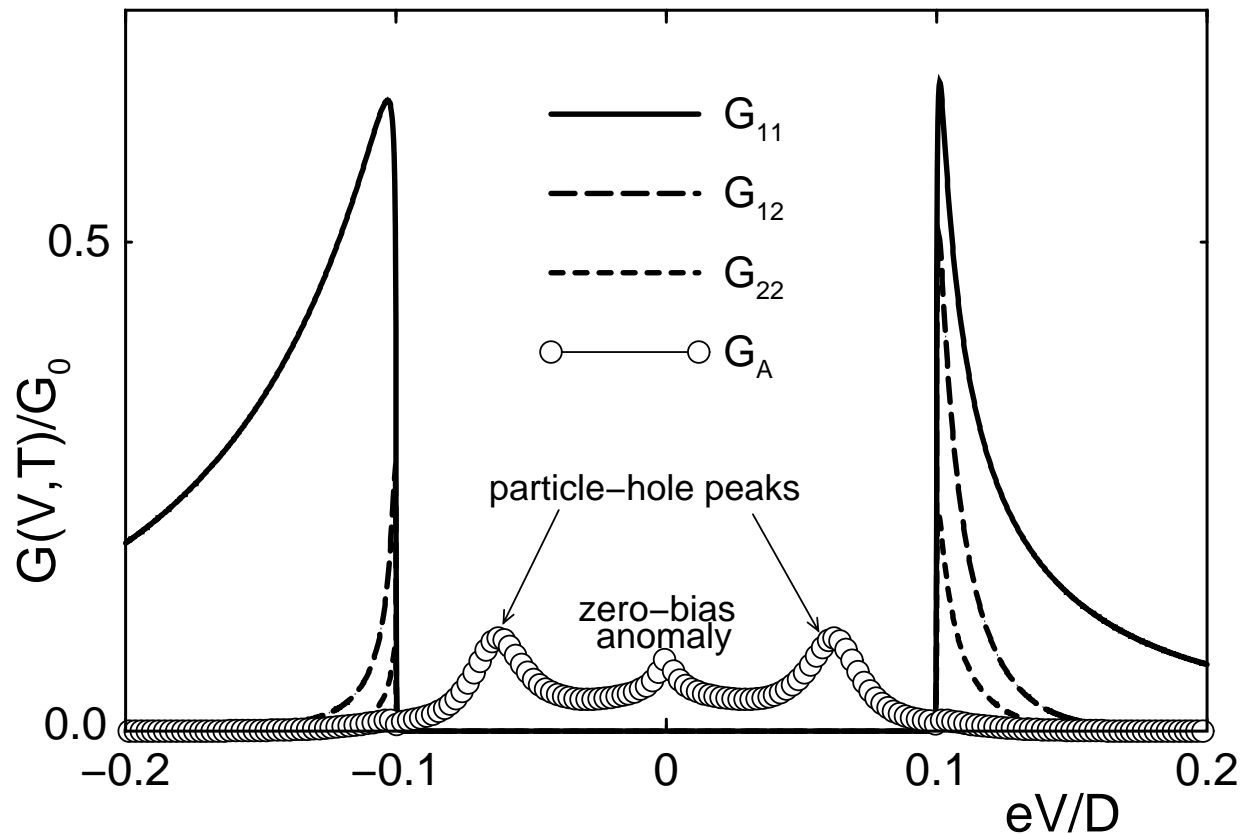
Qualitative features in the differential conductance  $G(V) = \frac{\partial J(V)}{\partial V}$



T. Domański, A. Donabidowicz, K.I. Wysokiński, PRB **76**, 104514 (2007).

## Transport channels

Qualitative features in the differential conductance  $G(V) = \frac{\partial J(V)}{\partial V}$



T. Domański, A. Donabidowicz, K.I. Wysokiński, PRB **76**, 104514 (2007).

We shall now focus on the subgap Andreev conductance.

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

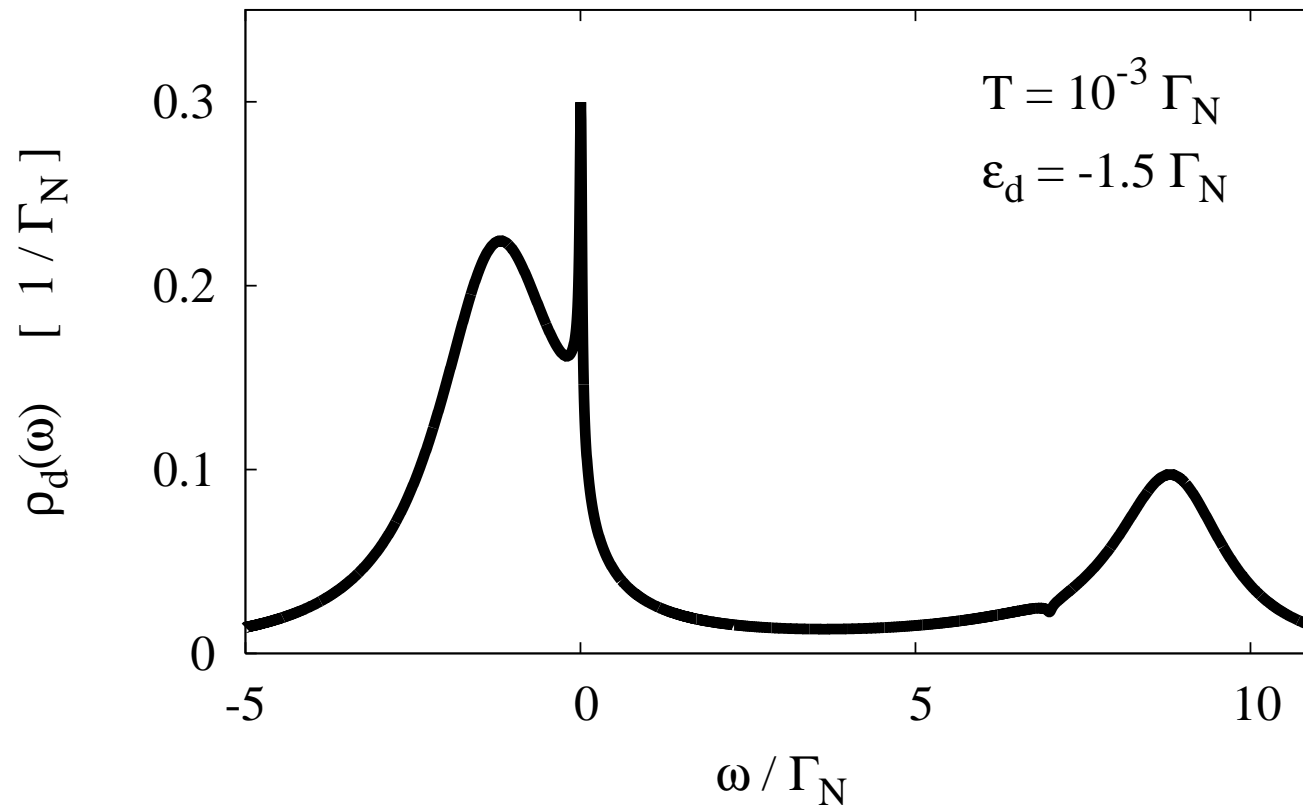
**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

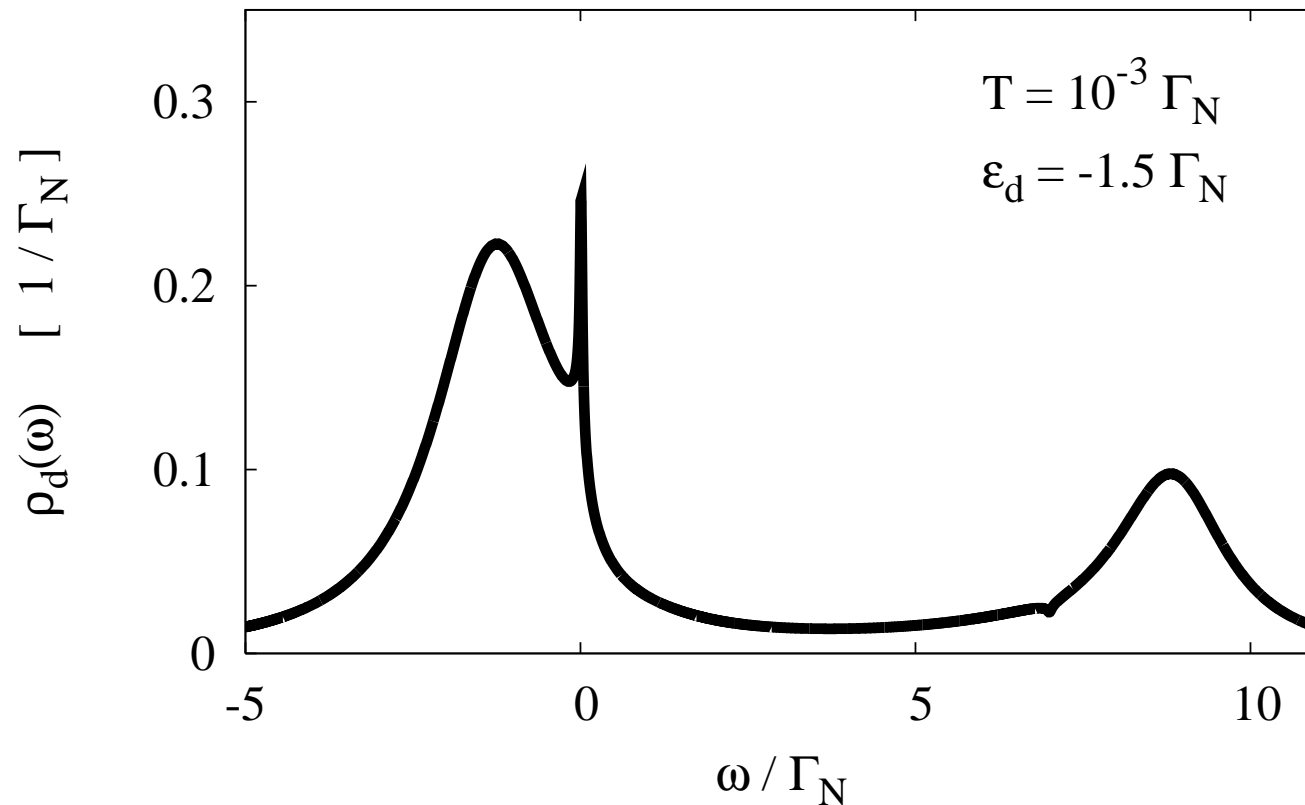


$$\Gamma_S/\Gamma_N = 0$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

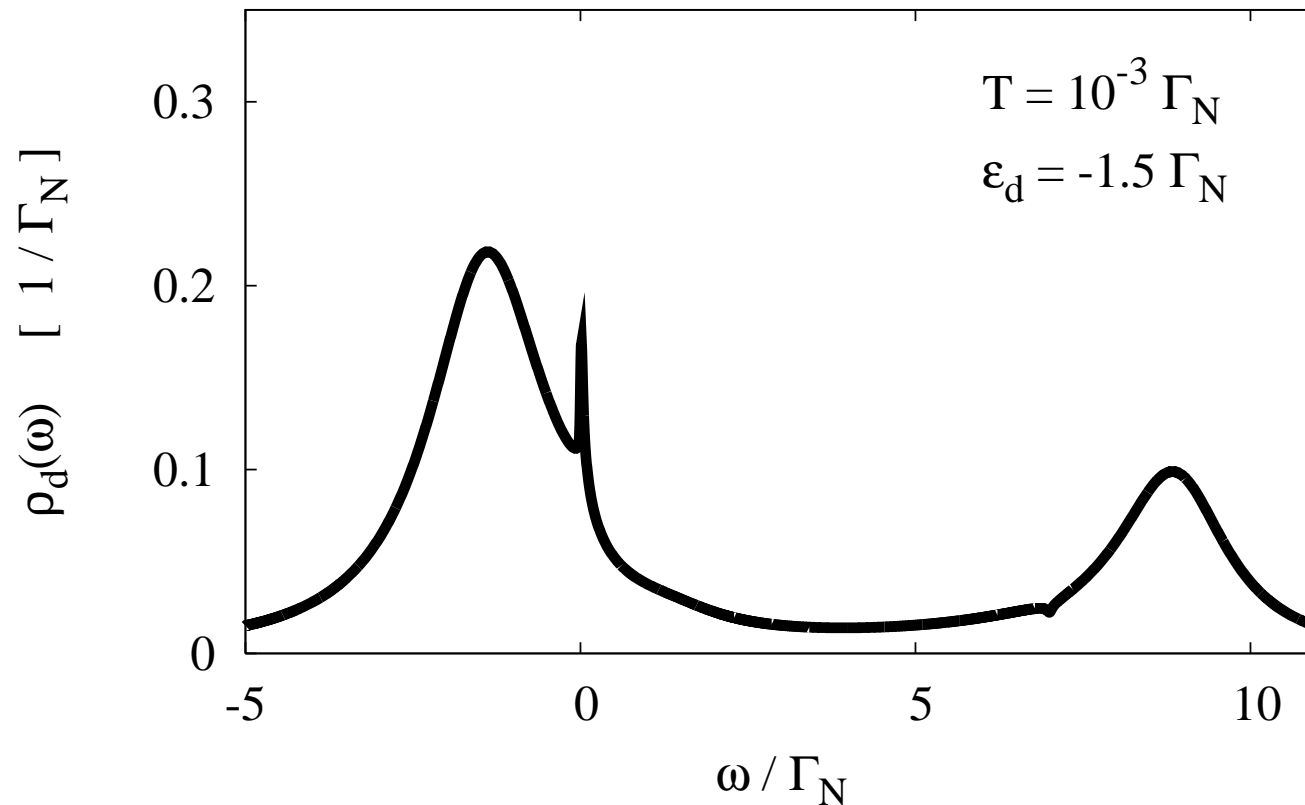


$$\Gamma_S/\Gamma_N = 1$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

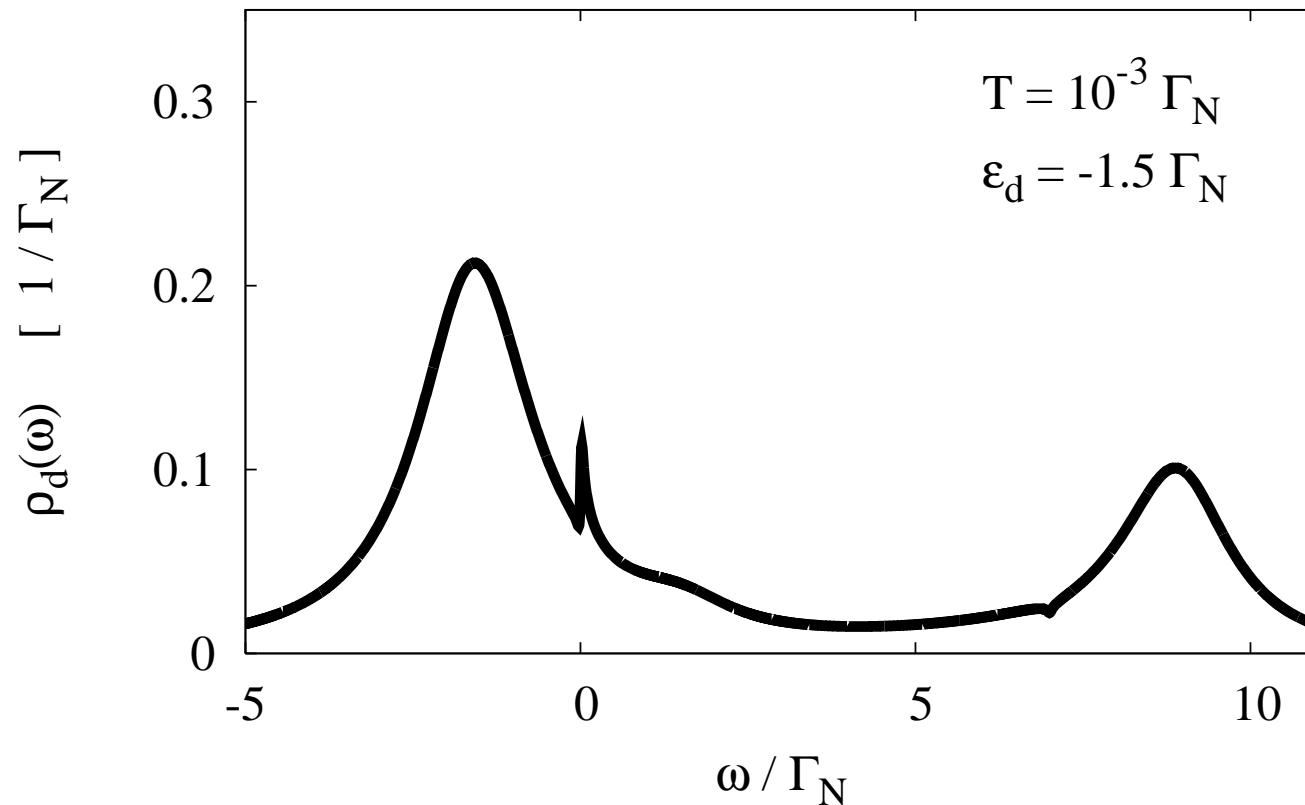


$$\Gamma_S/\Gamma_N = 2$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$



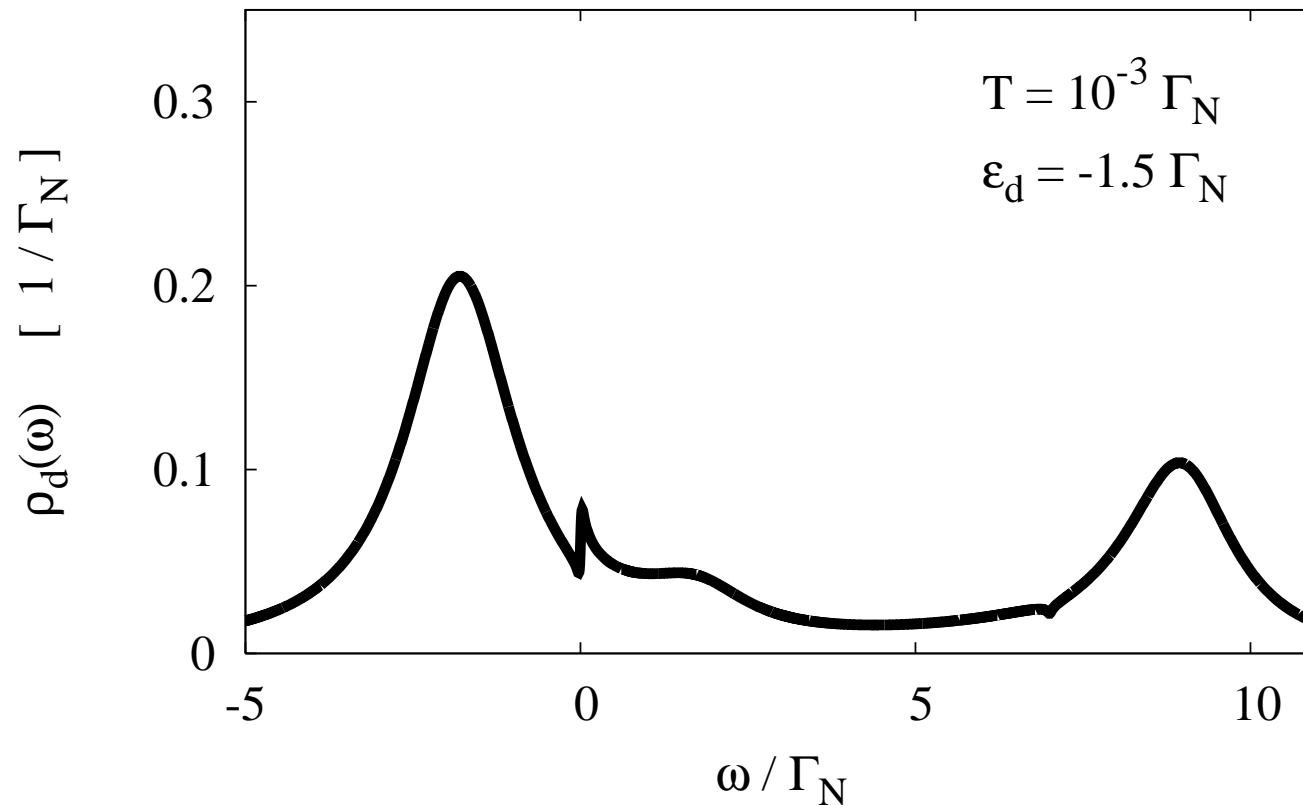
$$\Gamma_S/\Gamma_N = 3$$



## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

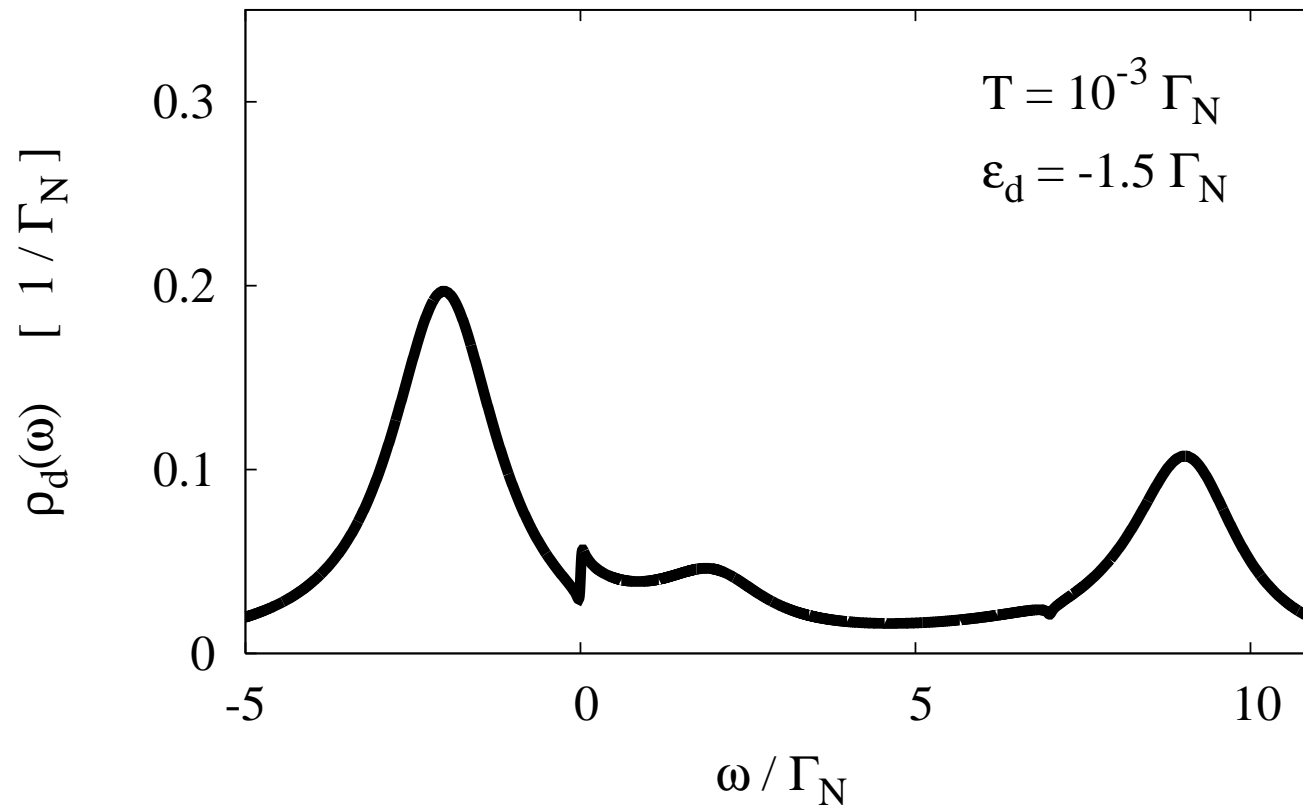


$$\Gamma_S/\Gamma_N = 4$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

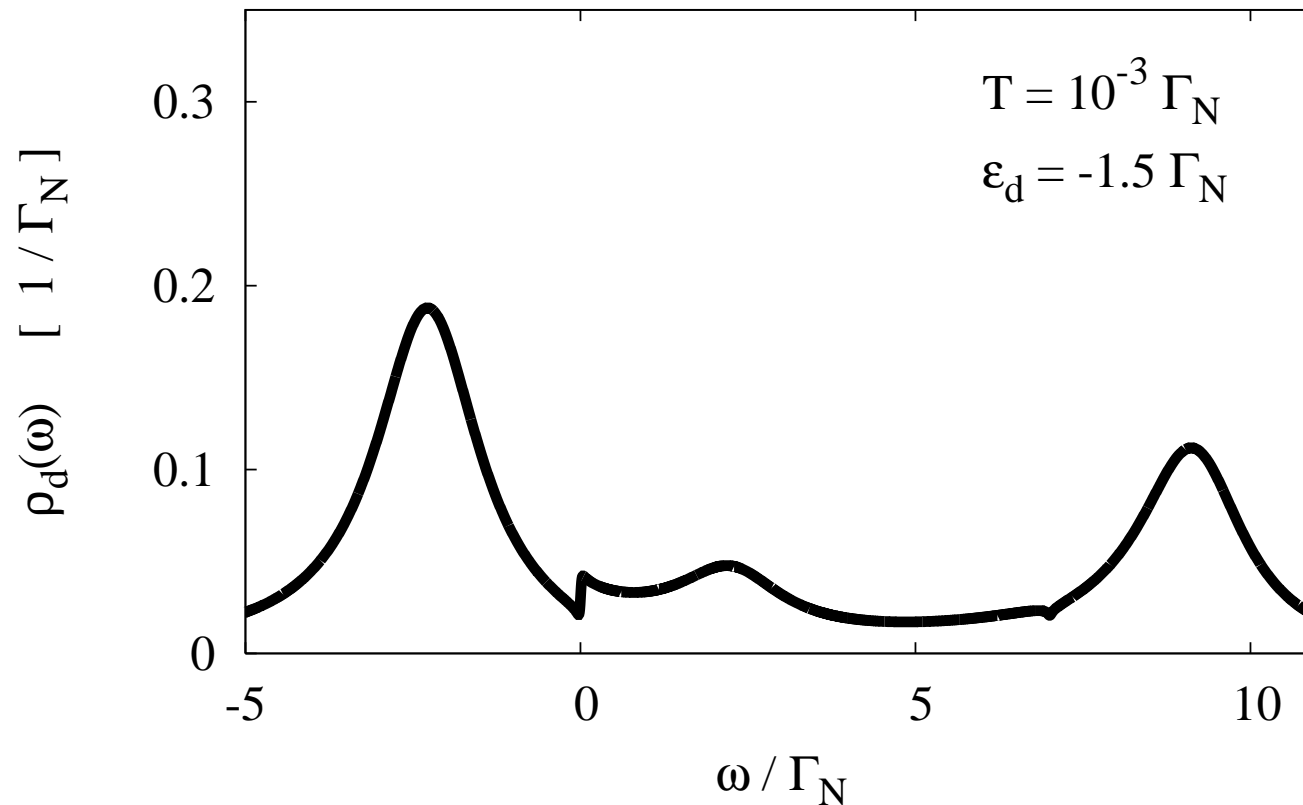


$$\Gamma_S/\Gamma_N = 5$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

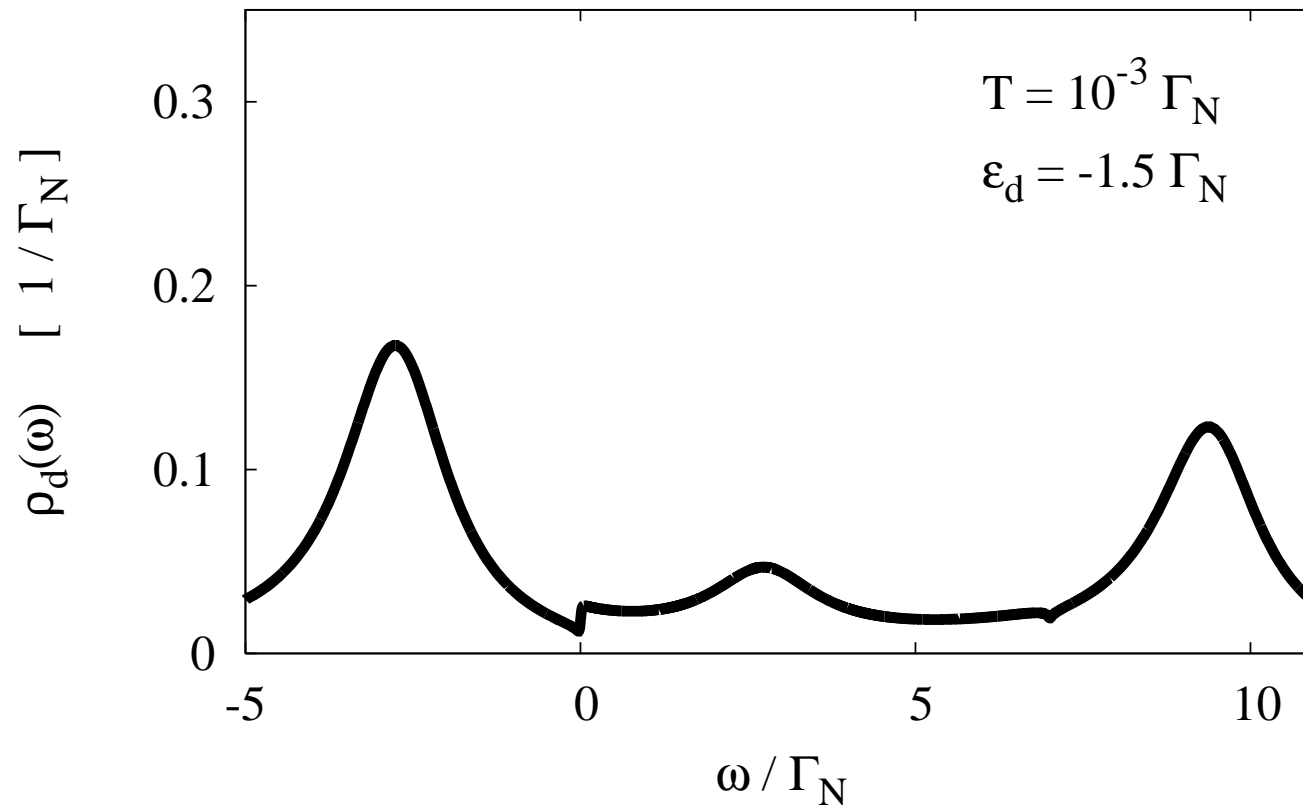


$$\Gamma_S/\Gamma_N = 6$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$

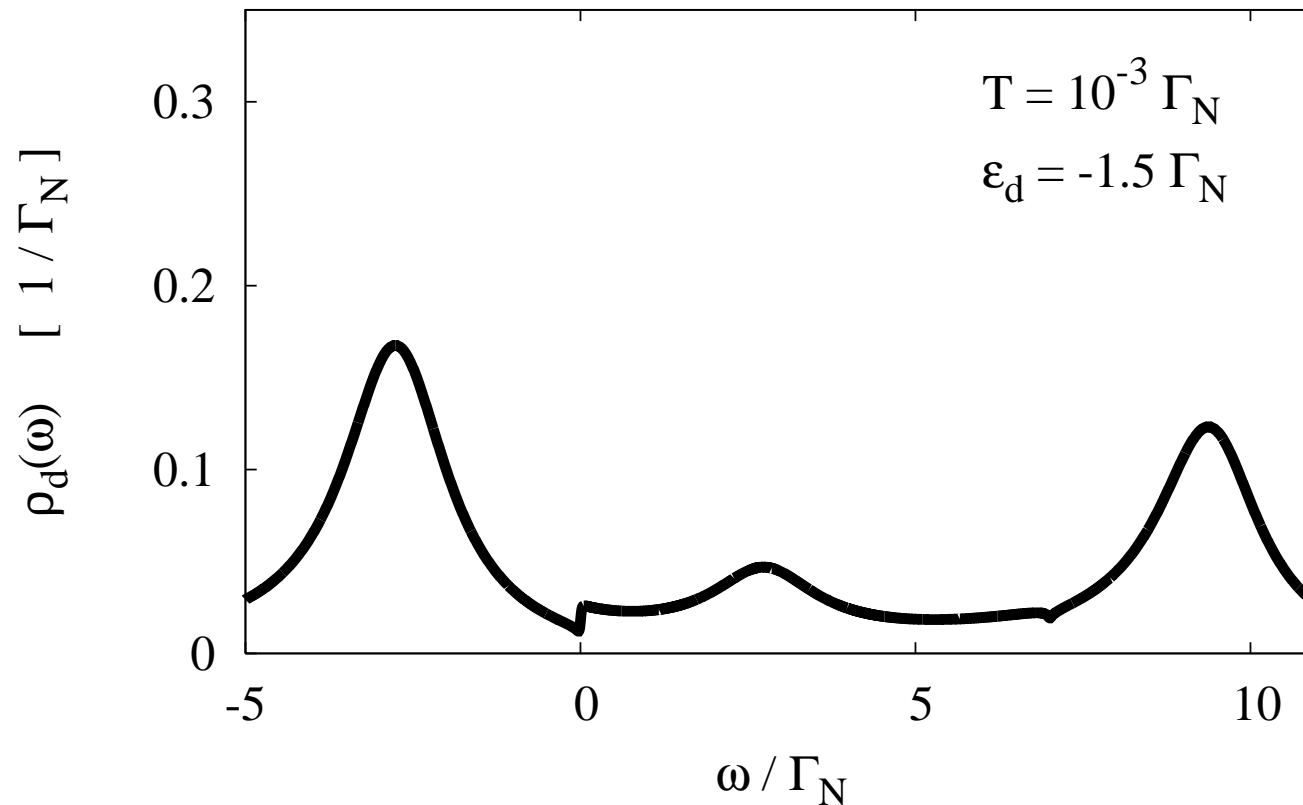


$$\Gamma_S/\Gamma_N = 8$$

## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Spectral function obtained below  $T_K$  for  $U = 10\Gamma_N$



Superconductivity suppresses the Kondo resonance

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

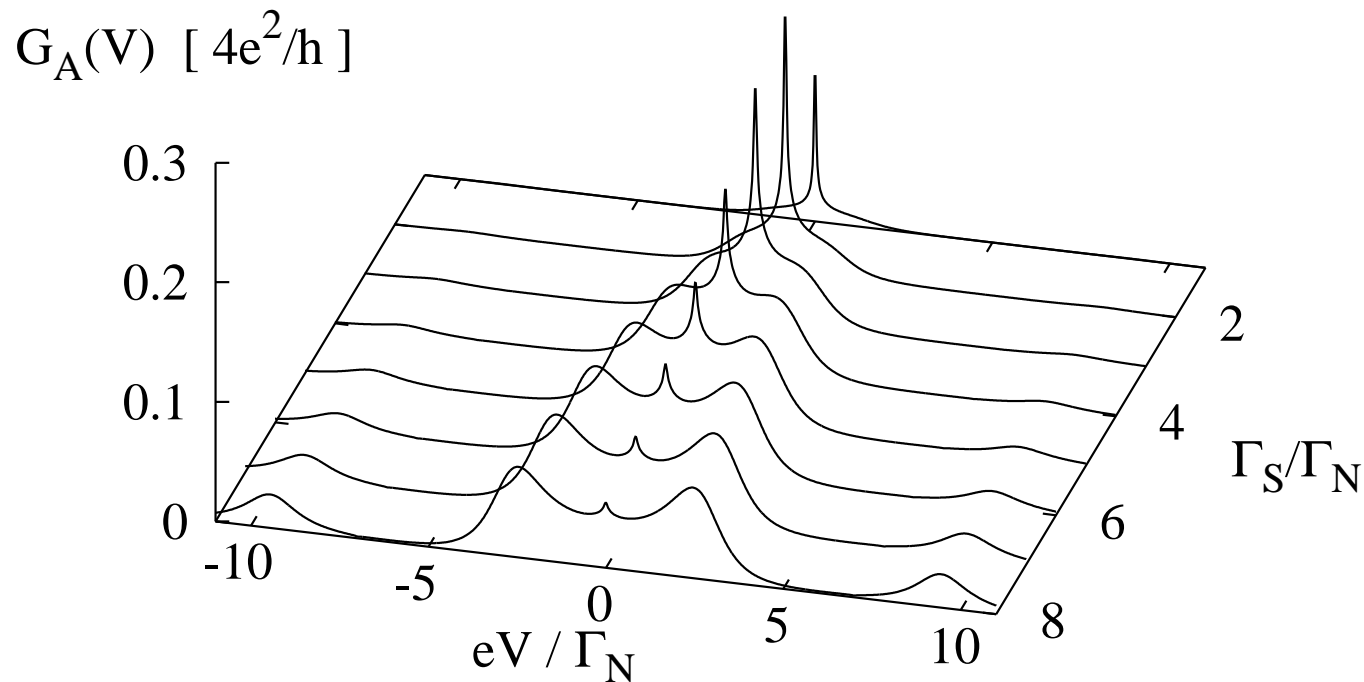
$$U = 10\Gamma_N$$

T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).



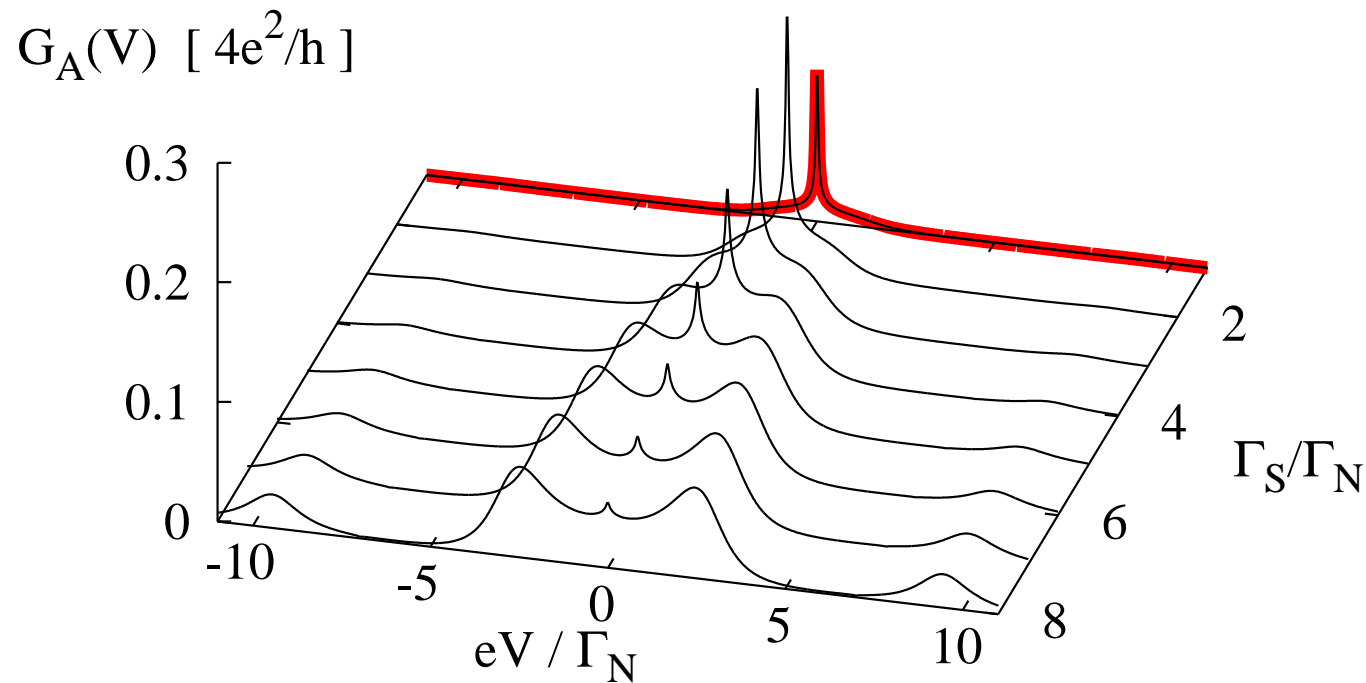
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 1$$



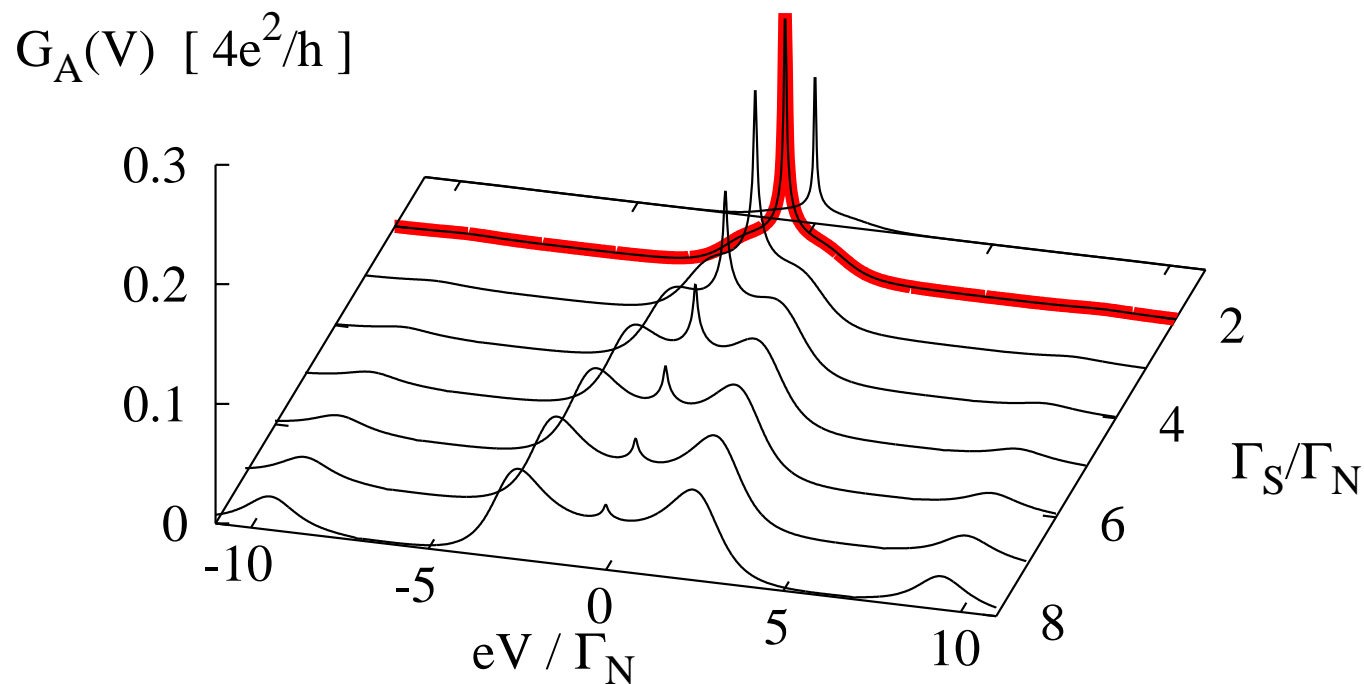
T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 2$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

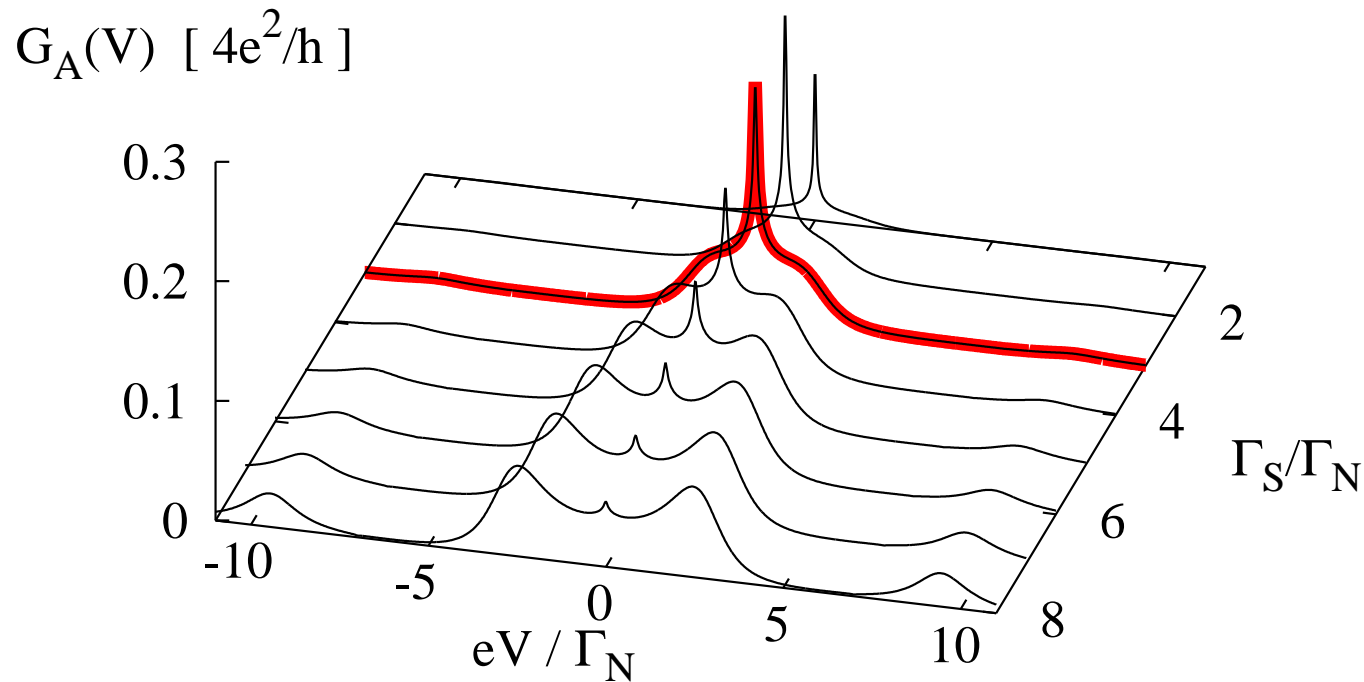
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 3$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

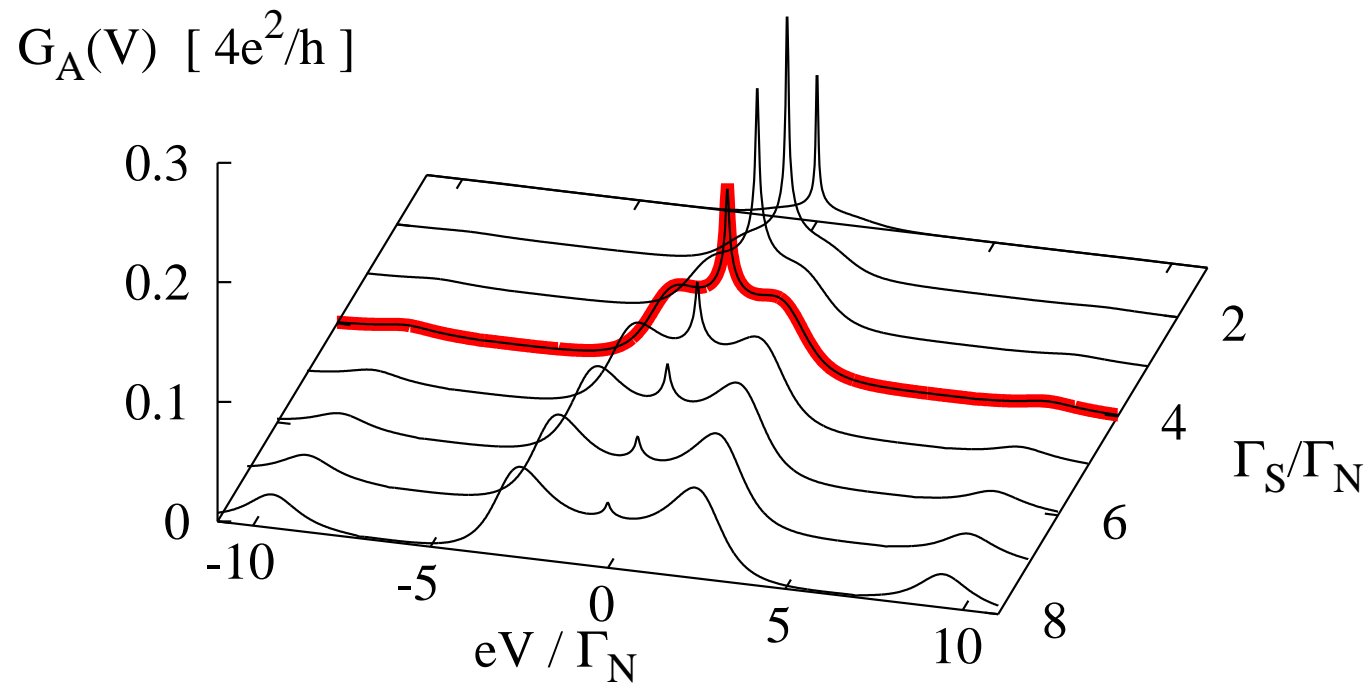
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 4$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

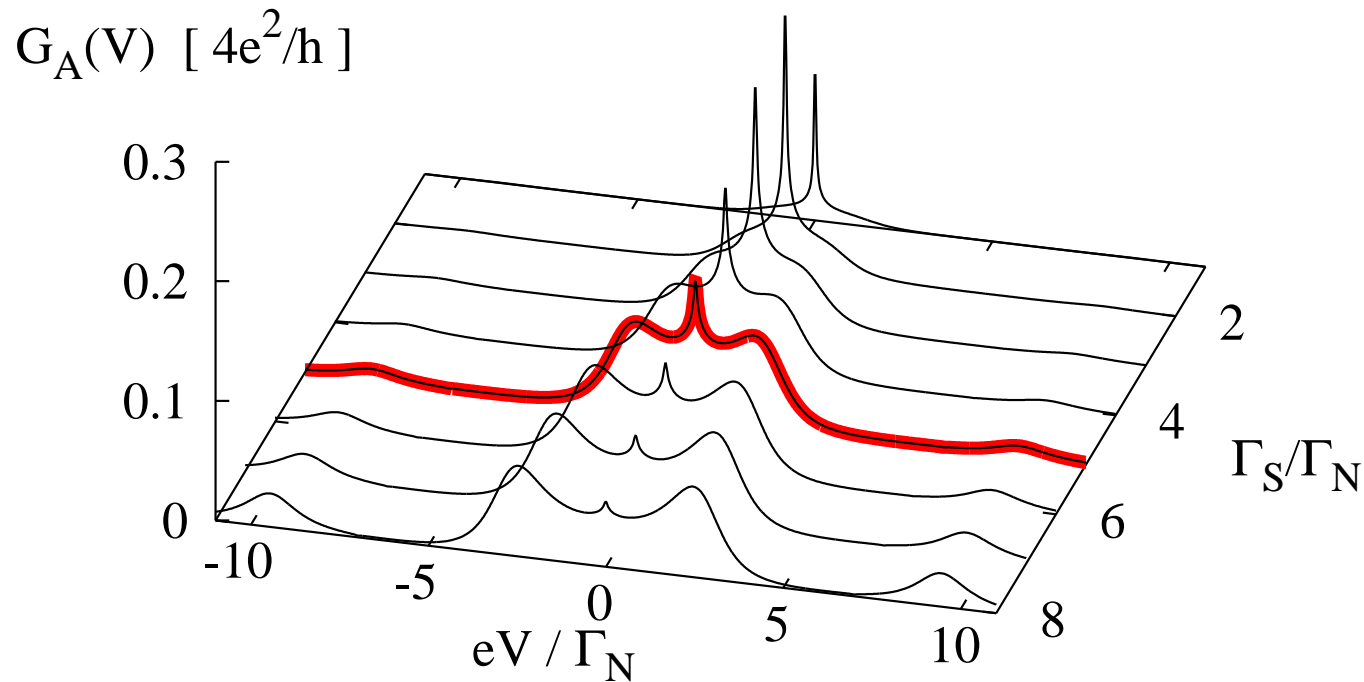
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 5$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

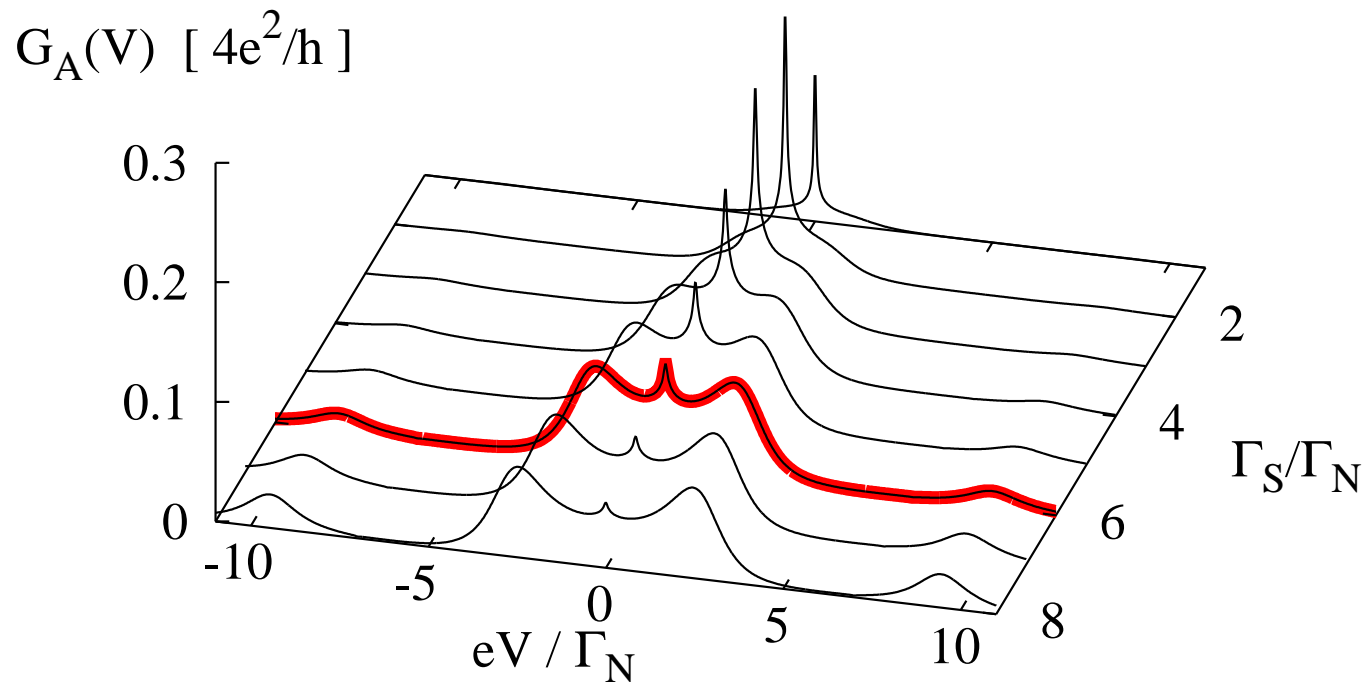
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 6$$



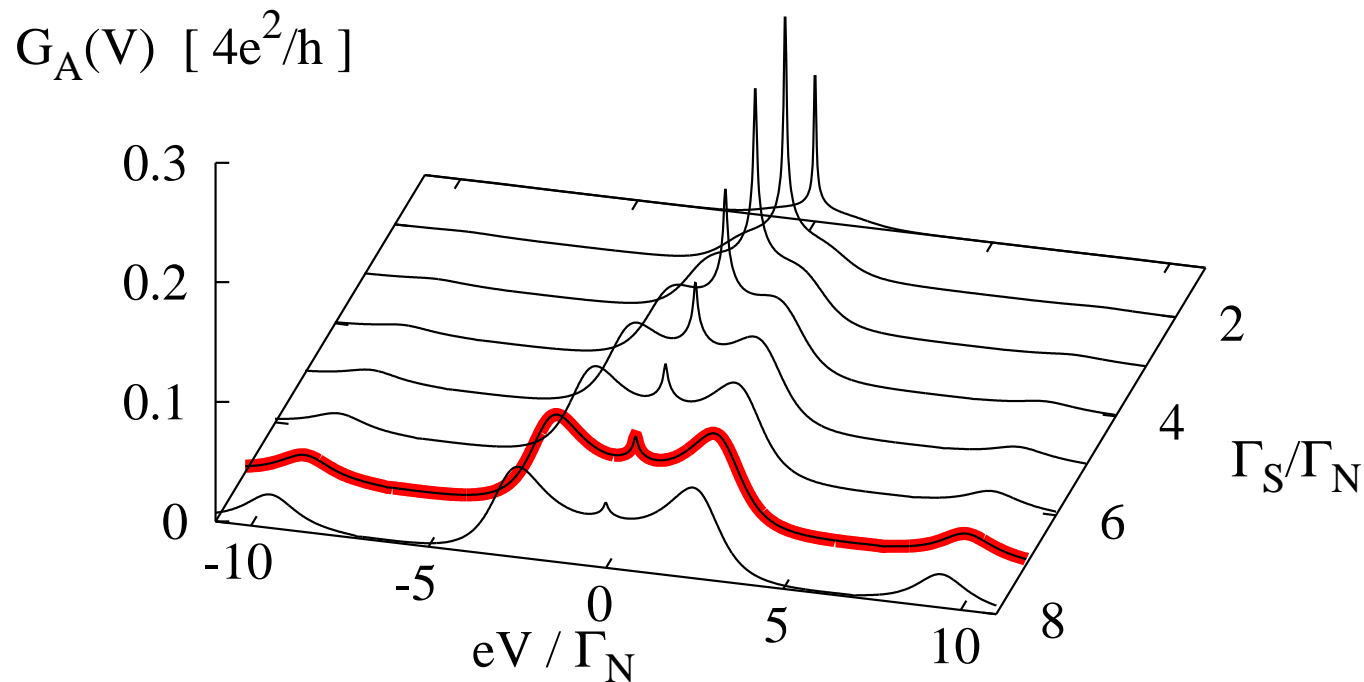
T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 7$$



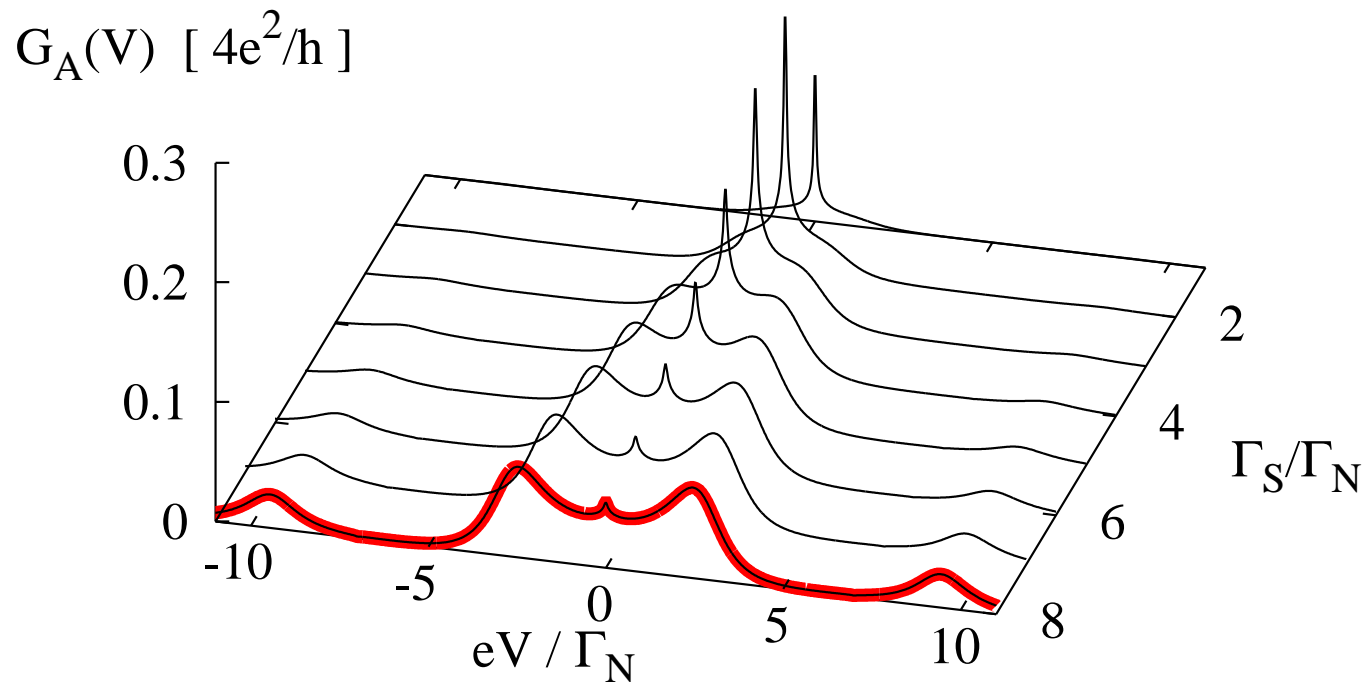
T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

**Correlated QD** – effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 8$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).



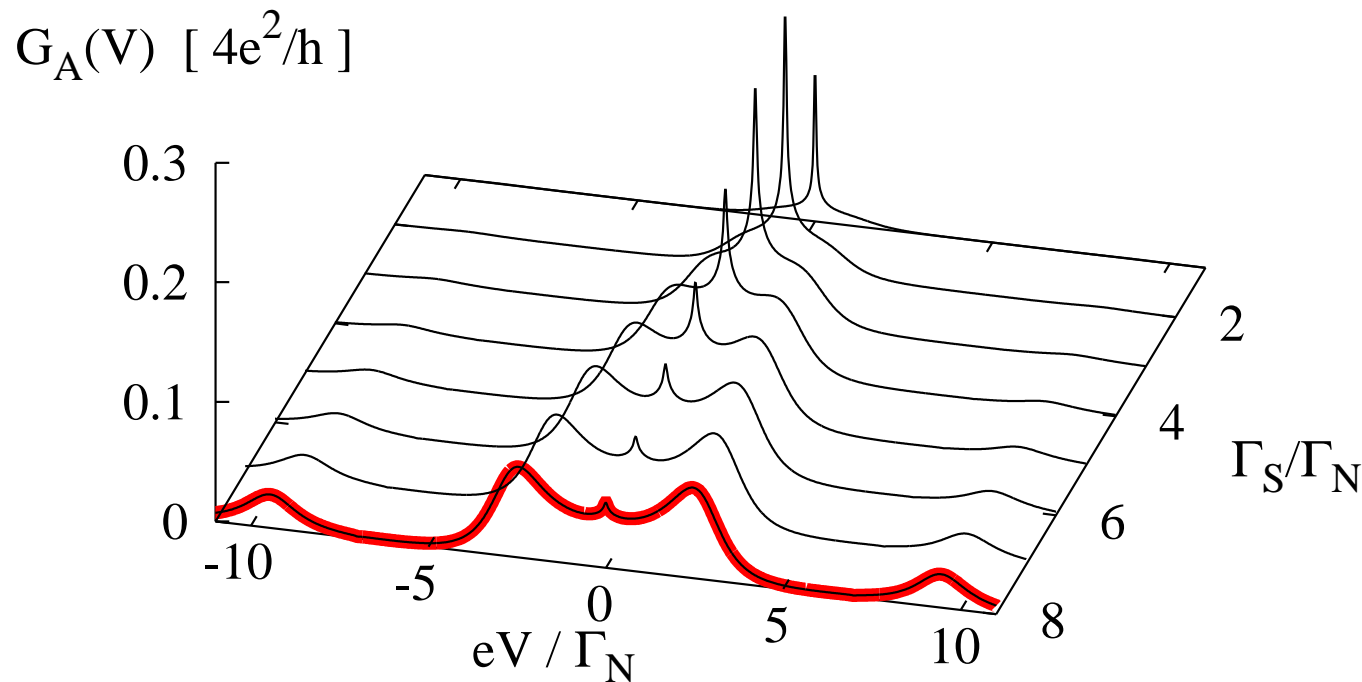
## Correlated QD

– effect of the asymmetry  $\Gamma_S/\Gamma_N$

Andreev conductance  $G_A(V)$  for:

$$U = 10\Gamma_N$$

$$\Gamma_S / \Gamma_N = 8$$



T. Domański and A. Donabidowicz, PRB **78**, 073105 (2008).

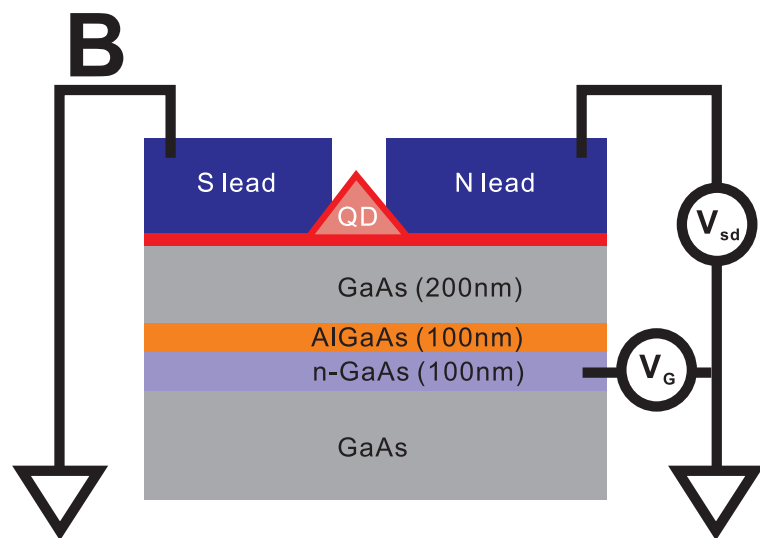
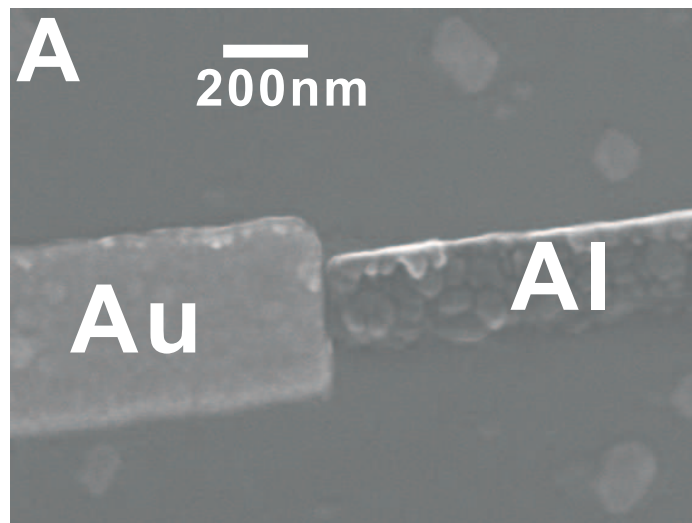
**Kondo resonance slightly enhances the zero-bias  
Andreev conductance, especially for  $\Gamma_S \sim \Gamma_N$  !**

**Experimental setup**

**/ University of Tokyo /**

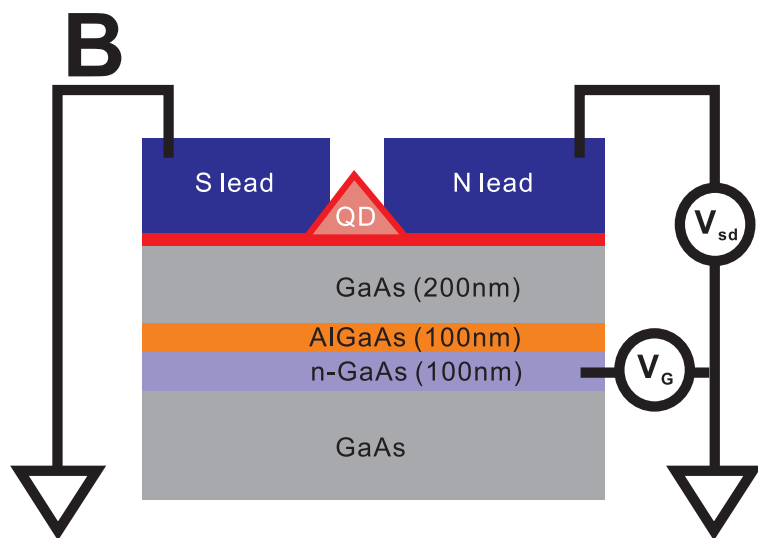
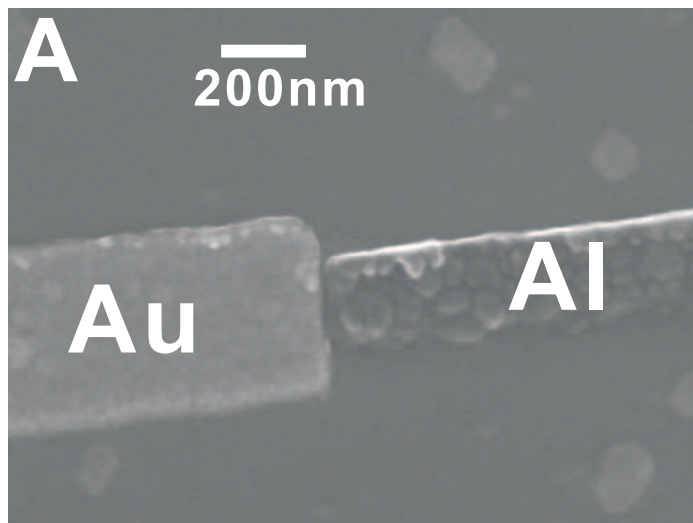
# Experimental setup

/ University of Tokyo /



# Experimental setup

/ University of Tokyo /



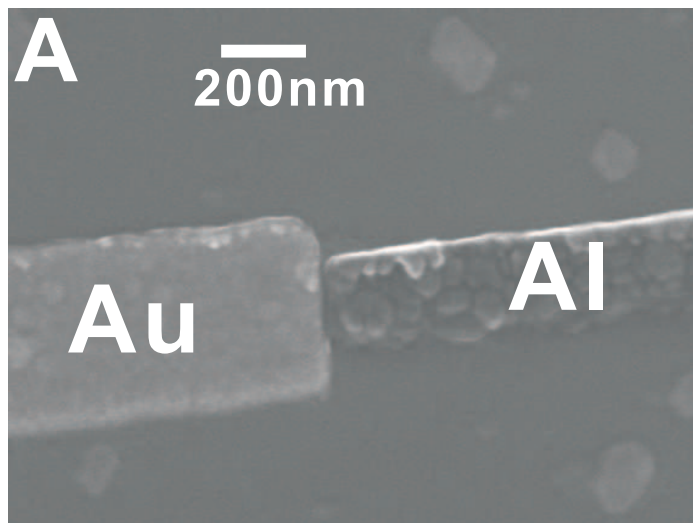
**QD** : self-assembled InAs

**diameter**  $\sim$  100 nm

**backgate** : Si-doped GaAs

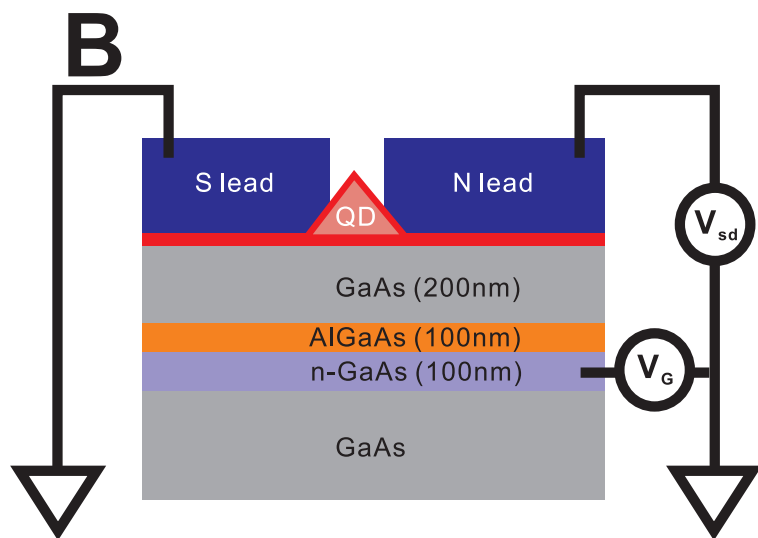
## Experimental setup

/ University of Tokyo /



$$T_c \simeq 1\text{K}$$

$$\Delta \simeq 152\mu\text{eV}$$



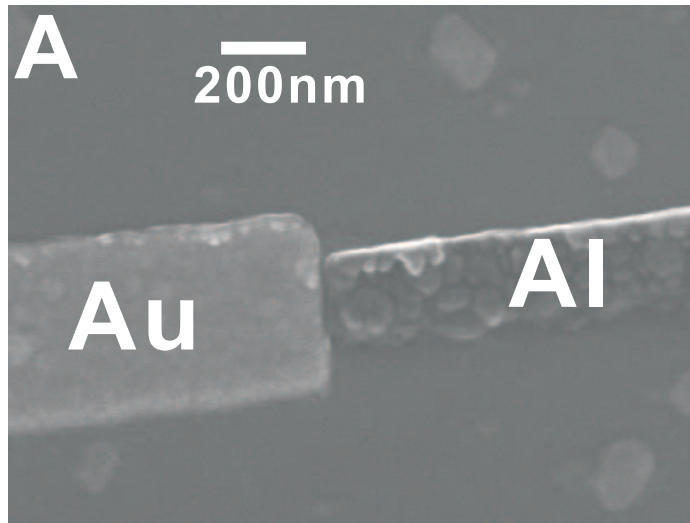
**QD** : self-assembled InAs

**diameter**  $\sim 100$  nm

**backgate** : Si-doped GaAs

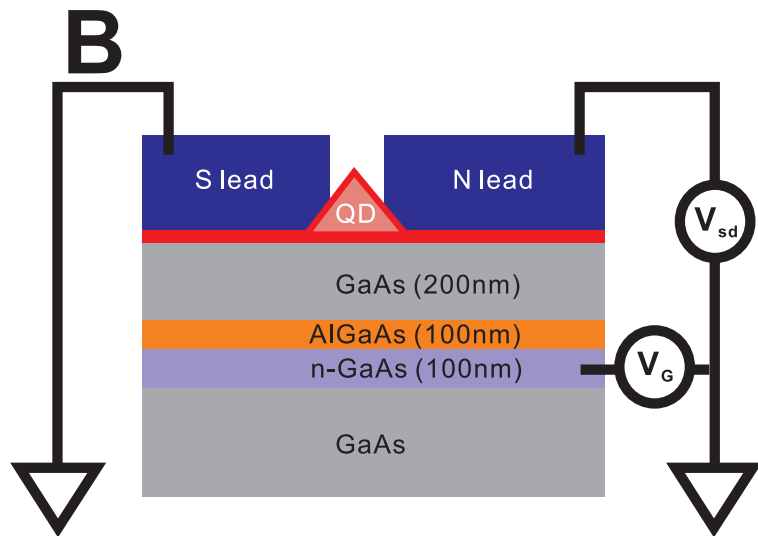
## Experimental setup

/ University of Tokyo /



$$T_c \simeq 1\text{K}$$

$$\Delta \simeq 152\mu\text{eV}$$



**QD** : self-assembled InAs

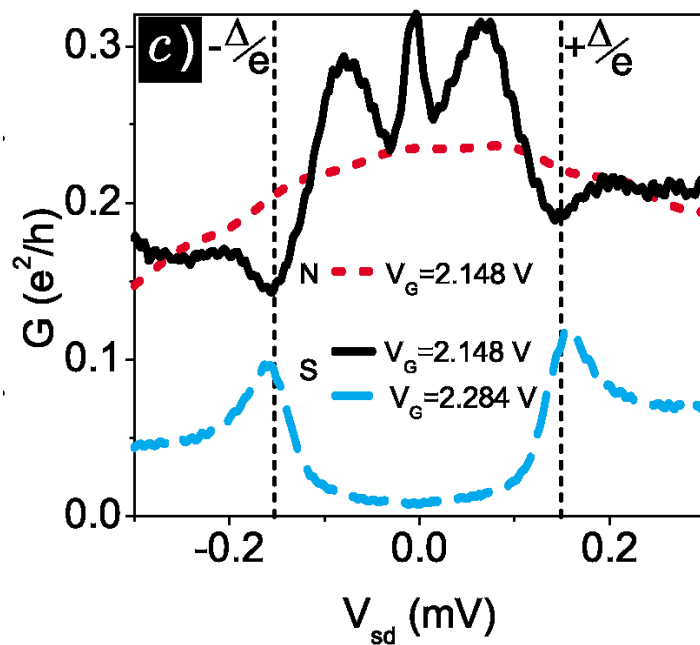
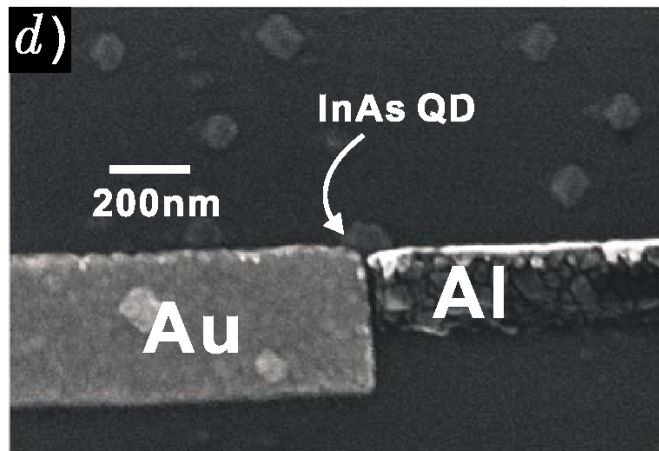
**diameter**  $\sim 100$  nm

**backgate** : Si-doped GaAs

*R.S. Deacon et al, Phys. Rev. Lett. **104**, 076805 (2010).*

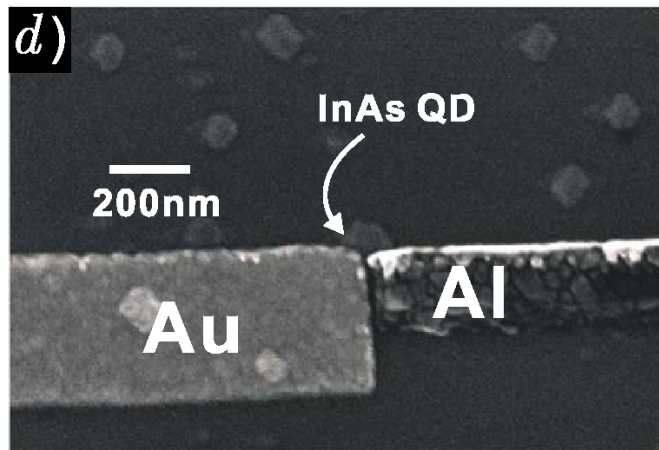
## Interplay with the Kondo effect

## Interplay with the Kondo effect

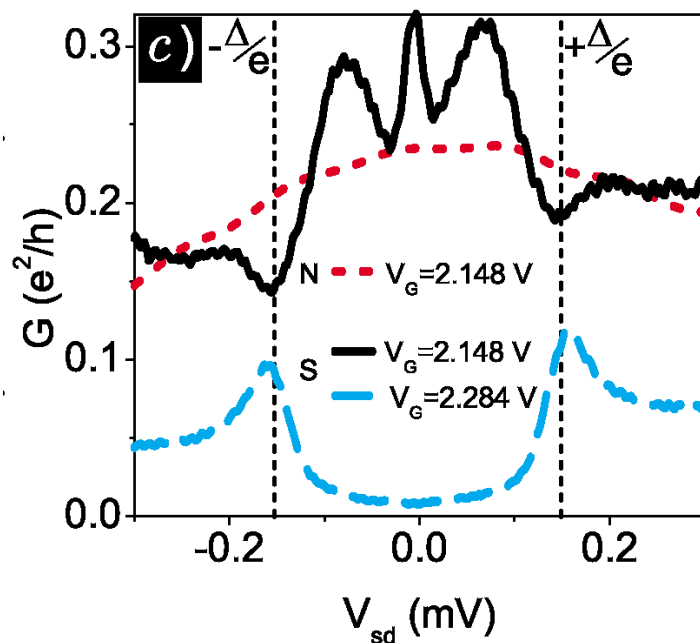




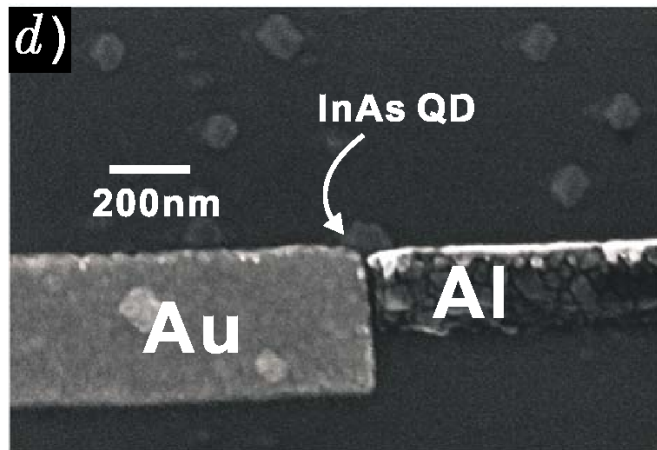
## Interplay with the Kondo effect



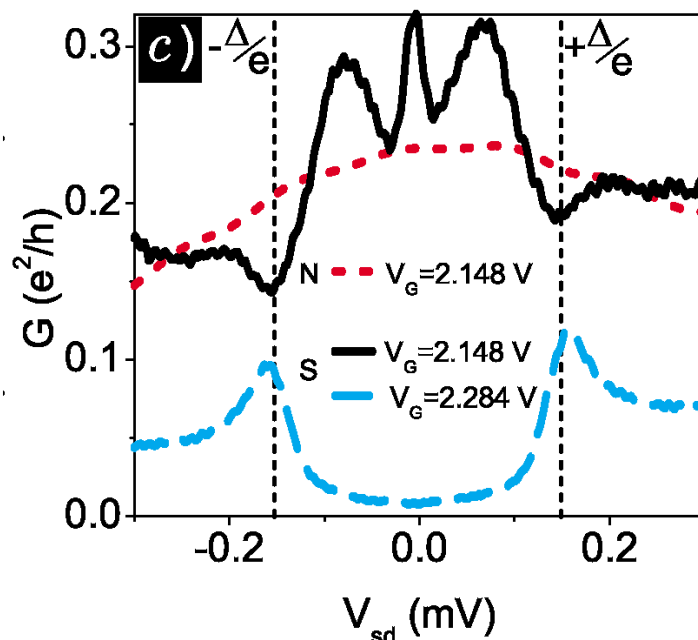
*"The zero-bias conductance peak is consistent with Andreev transport enhanced by the Kondo singlet state"*



## Interplay with the Kondo effect



*"The zero-bias conductance peak is consistent with Andreev transport enhanced by the Kondo singlet state"*



*"We note that the feature exhibits excellent qualitative agreement with a recent theoretical treatment by Domanski et al"*

**Summary**

/ for part 2 /

# Summary

/ for part 2 /

**QD coupled between N and S electrodes:**

## Summary

/ for part 2 /

QD coupled between N and S electrodes:

⇒ absorbs the superconducting order / proximity effect /

## Summary

/ for part 2 /

**QD coupled between N and S electrodes:**

- ⇒ **absorbs the superconducting order** / proximity effect /
- ⇒ **is affected by the correlations** / Kondo & charging effects /

## **Summary**

/ for part 2 /

**QD coupled between N and S electrodes:**

- ⇒ **absorbs the superconducting order** / proximity effect /
- ⇒ **is affected by the correlations** / Kondo & charging effects /

**Interplay between the proximity and correlation effects  
is manifested in a subgap Andreev transport by:**

## Summary

/ for part 2 /

QD coupled between N and S electrodes:

- ⇒ absorbs the superconducting order / proximity effect /
- ⇒ is affected by the correlations / Kondo & charging effects /

Interplay between the proximity and correlation effects is manifested in a subgap Andreev transport by:

- ⇒ particle-hole splitting / when  $\varepsilon_d \sim \mu_S$  /



## Summary

/ for part 2 /

QD coupled between N and S electrodes:

- ⇒ absorbs the superconducting order / proximity effect /
- ⇒ is affected by the correlations / Kondo & charging effects /

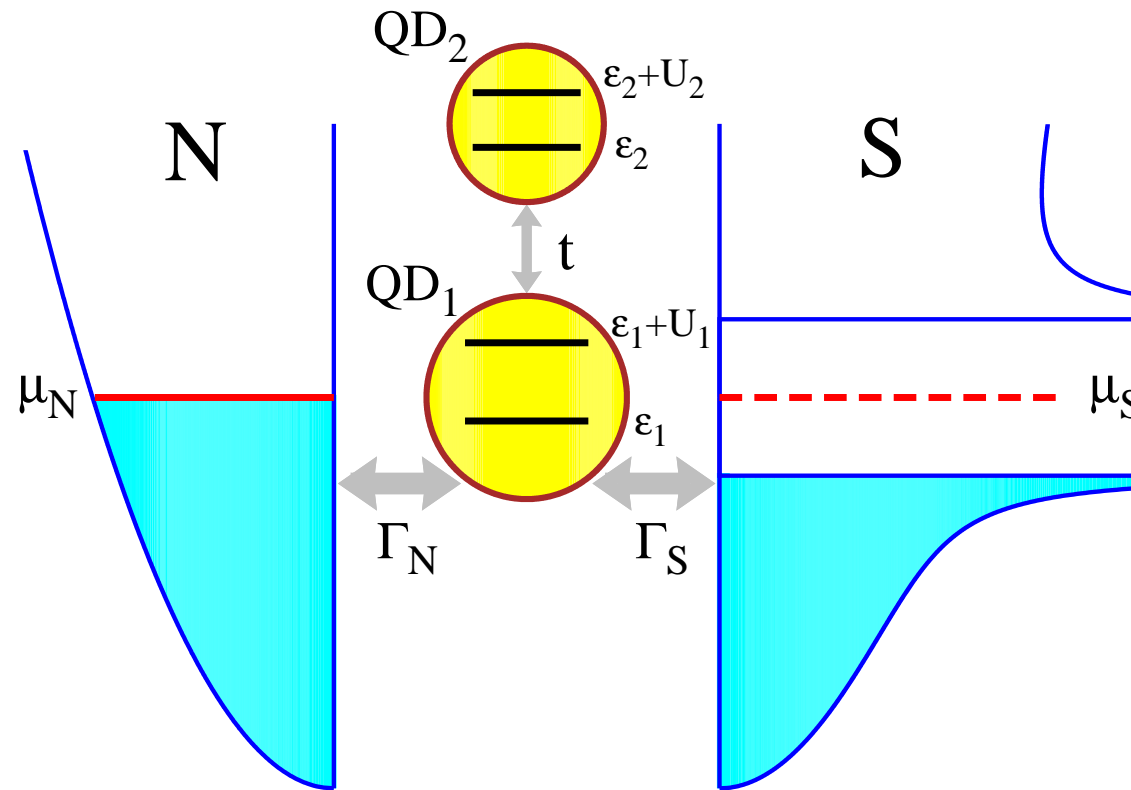
Interplay between the proximity and correlation effects is manifested in a subgap Andreev transport by:

- ⇒ particle-hole splitting / when  $\varepsilon_d \sim \mu_S$  /
- ⇒ zero-bias enhancement / below  $T_K$  /

### **3. Further extensions**

## Double QD

– between a metal and superconductor



( T-shape configuration)

### Relevant issues:

- ⇒ induced on-dot pairing ..... (due to  $\Gamma_S$ )
- ⇒ Coulomb blockade & Kondo effect ..... (via  $U_1$  and  $\Gamma_N$ )
- ⇒ quantum interference ..... (because of  $t$ )

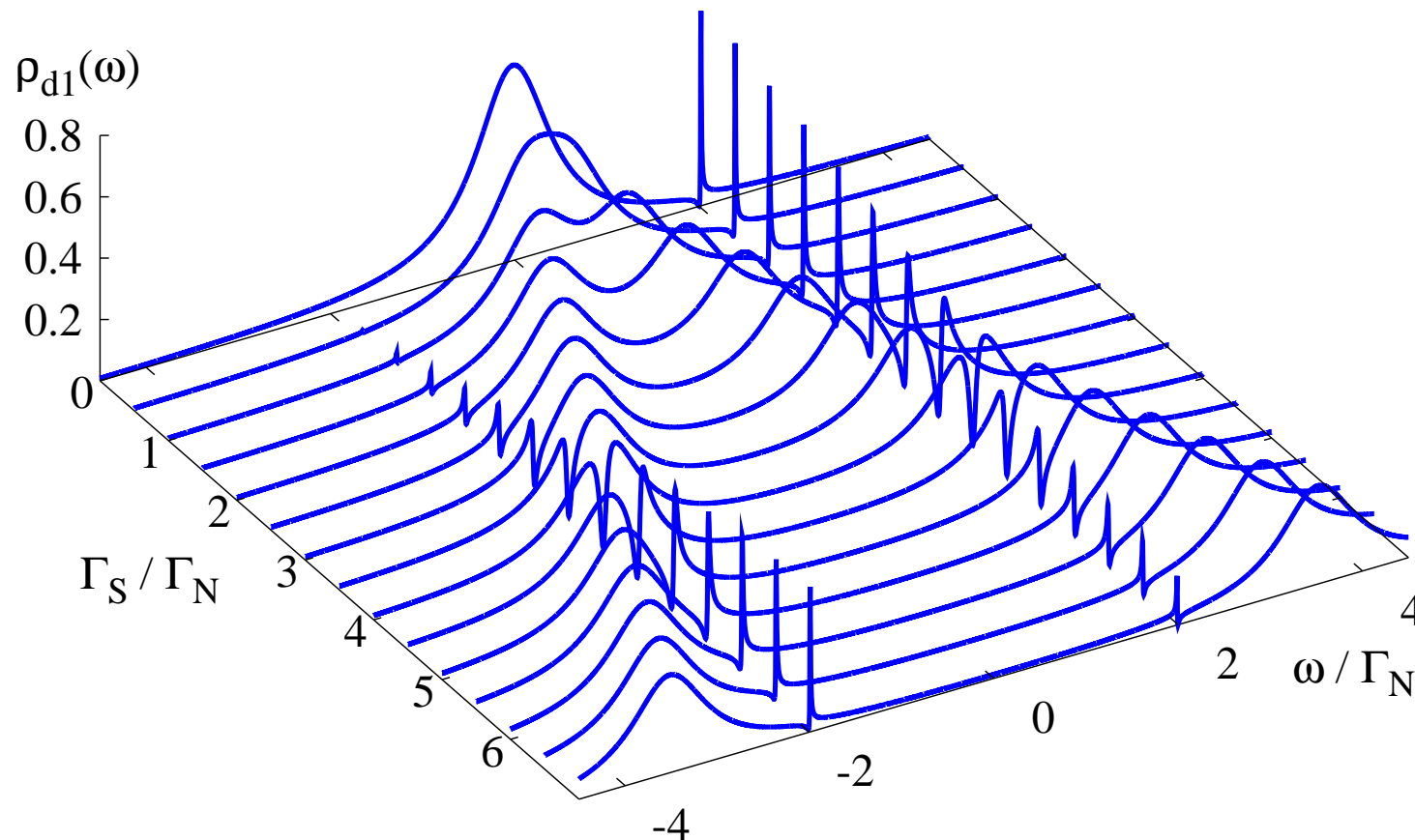
## Quantum interference

– in the particle and hole channels

## Quantum interference

– in the particle and hole channels

**Fano-type lineshapes appear simultaneously at  $\pm\varepsilon_2$**

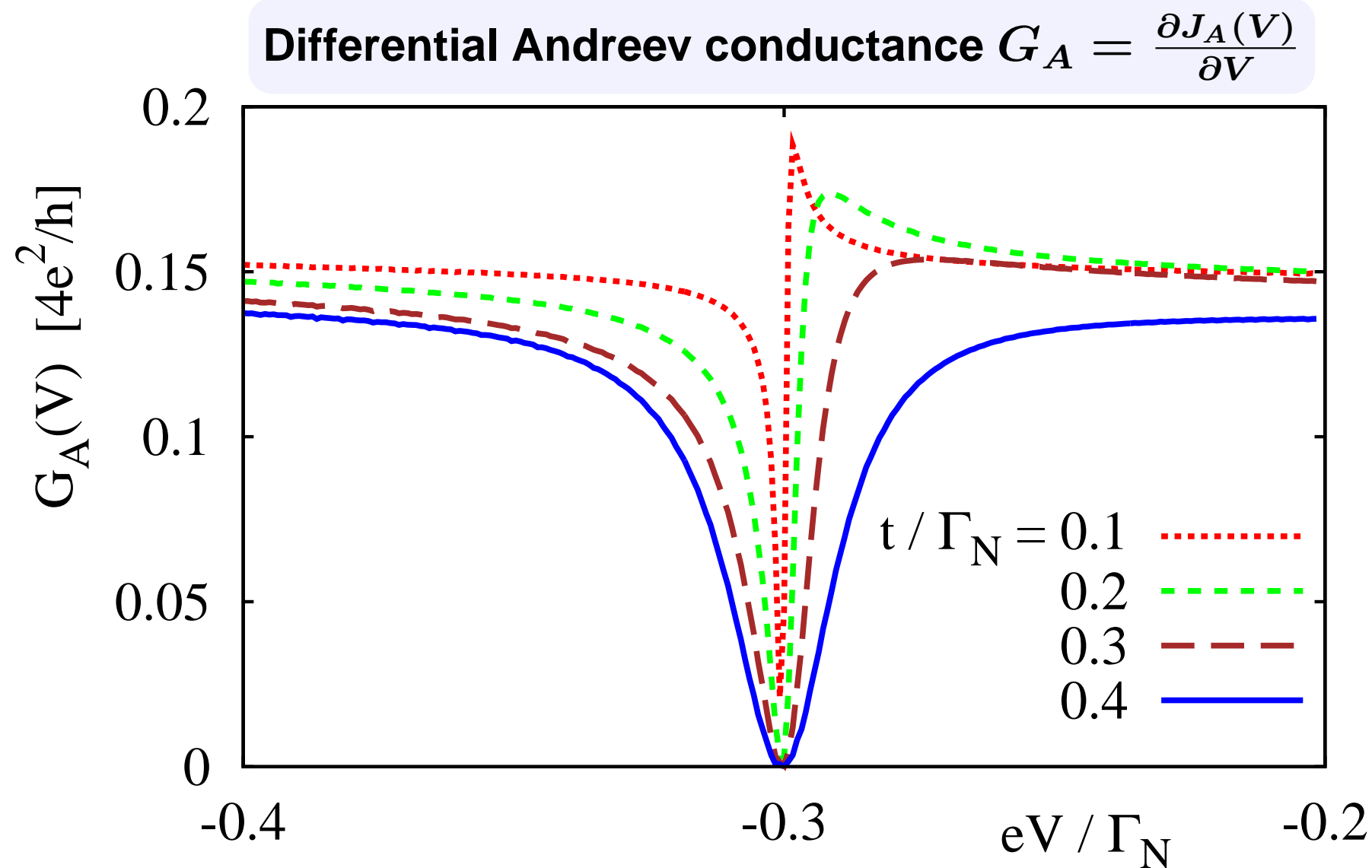


/ The case:  $U_1 = 0$  and  $U_2 = 0$  /

J. Barański and T. Domański, Phys. Rev. B (2012).

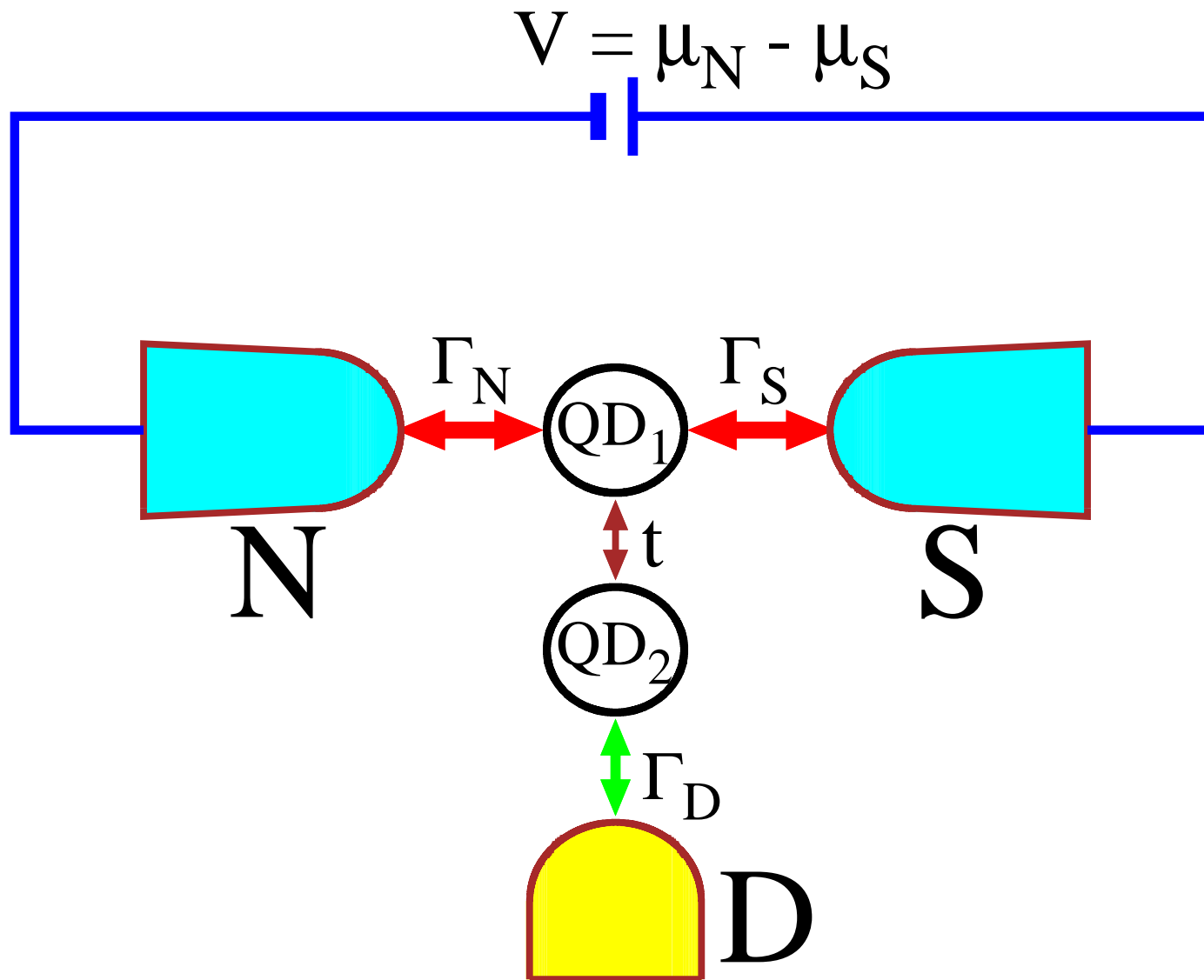
## Quantum interference

– in the particle and hole channels



## Double QD

– decoherence effects



In this setup the floating lead (D) is responsible for a dephasing.

## Quantum interference

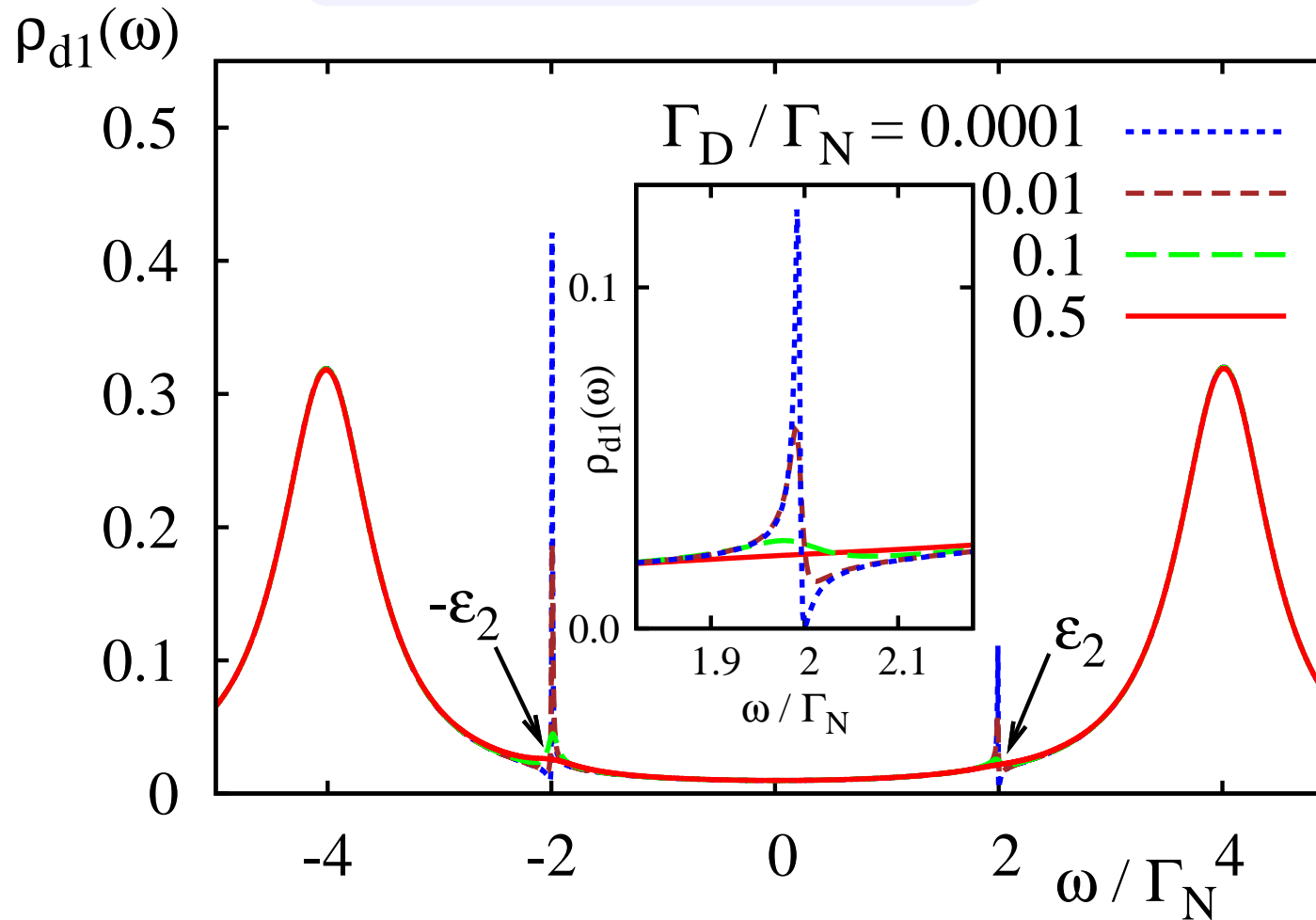
– influence of the decoherence



# Quantum interference

– influence of the decoherence

## Density of states $\rho_{d1}(\omega)$

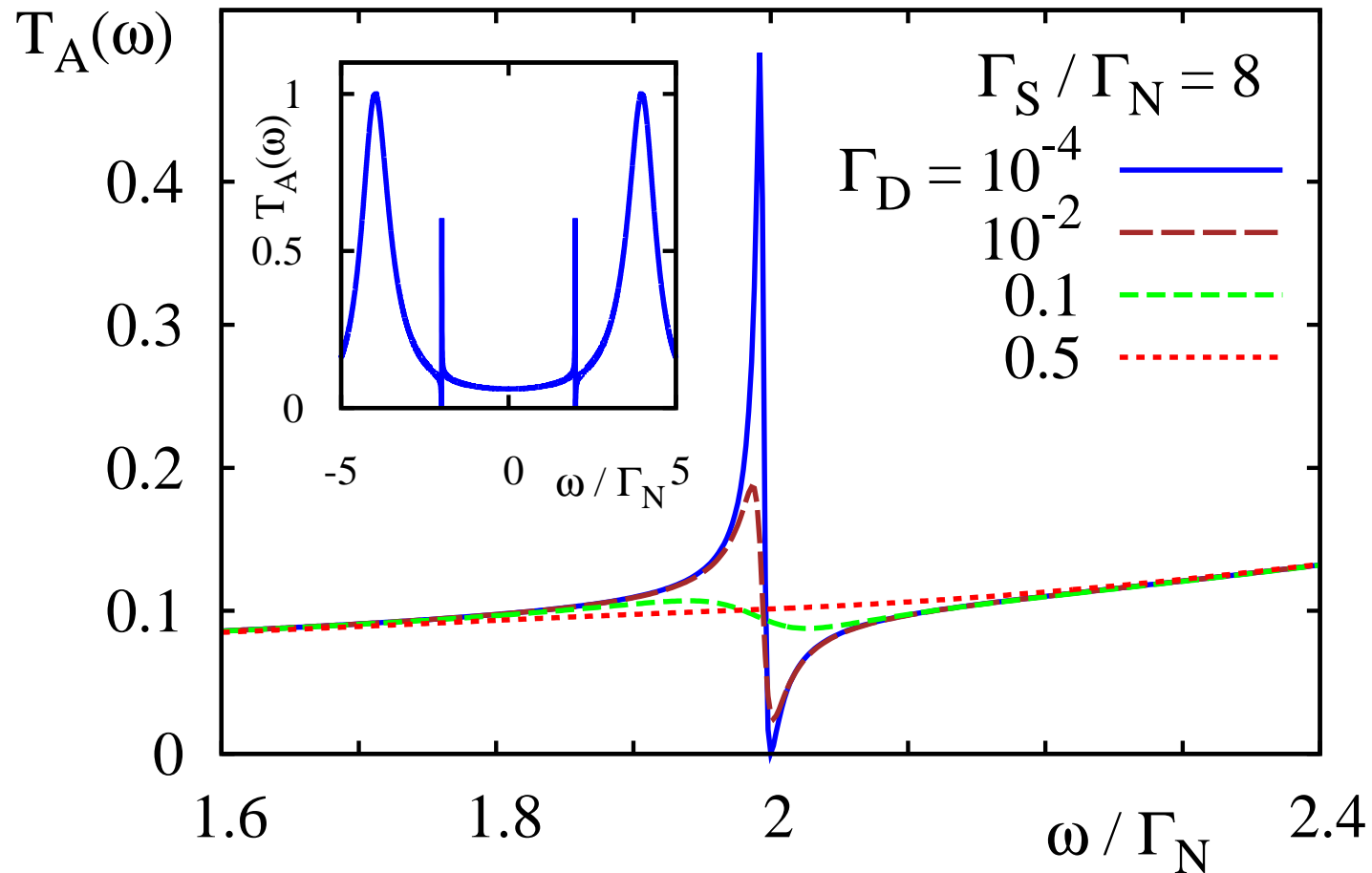


J. Barański and T. Domański, Phys. Rev. B (2012).

# Quantum interference

— influence of the decoherence

## Andreev transmittance $T_A(\omega)$



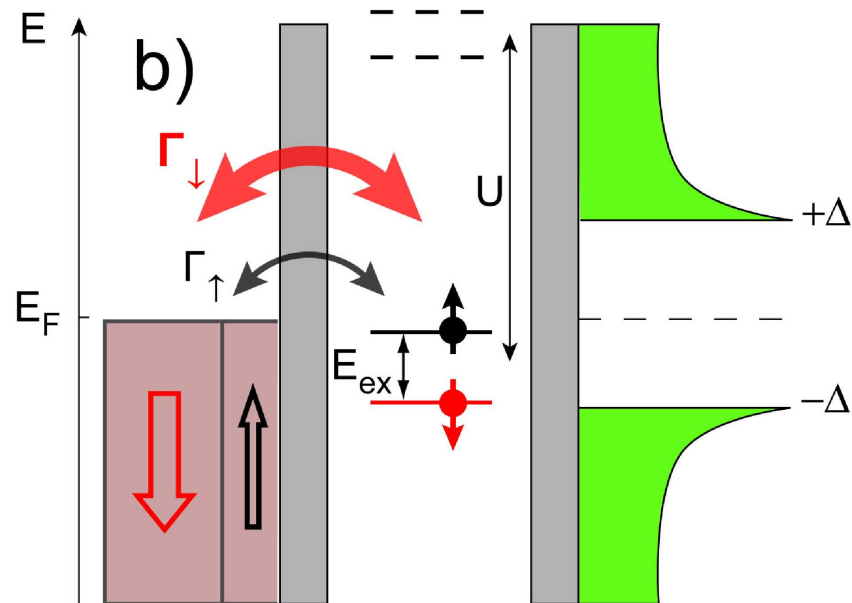
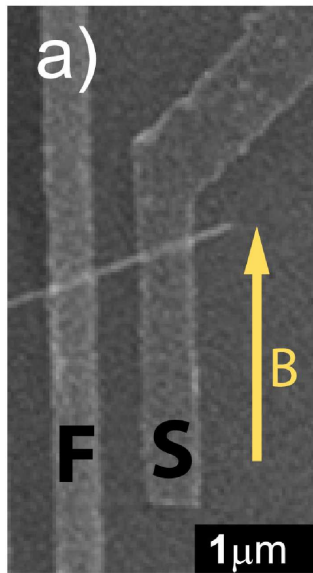
J. Barański and T. Domański, Phys. Rev. B (2012).

# Interplay with ferromagnetism

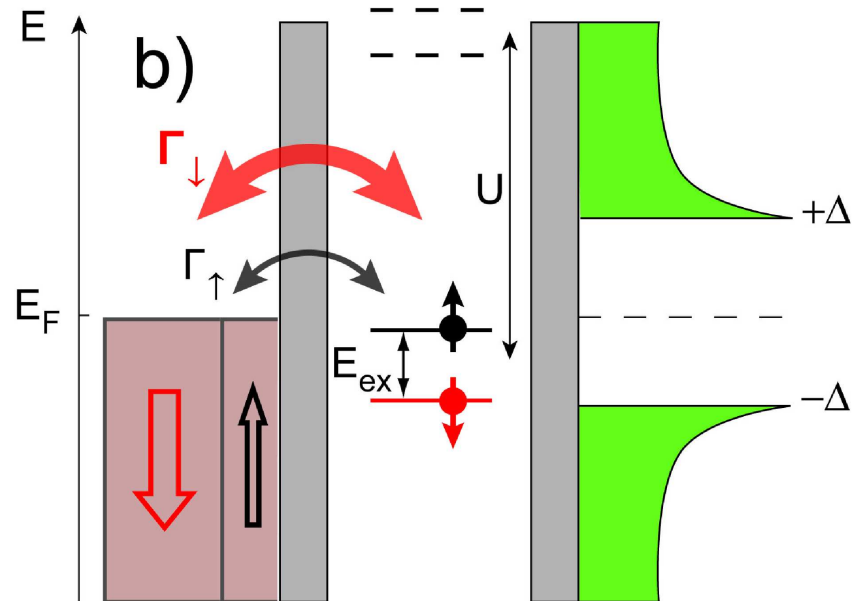
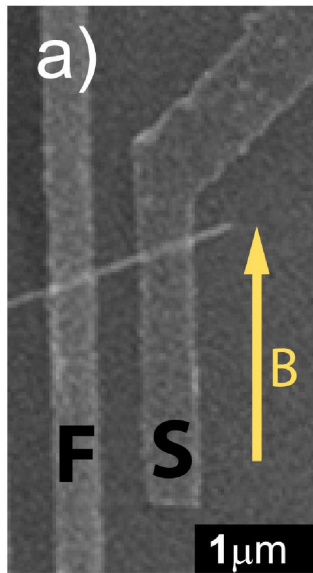
– Univ. of Basel group

# Interplay with ferromagnetism

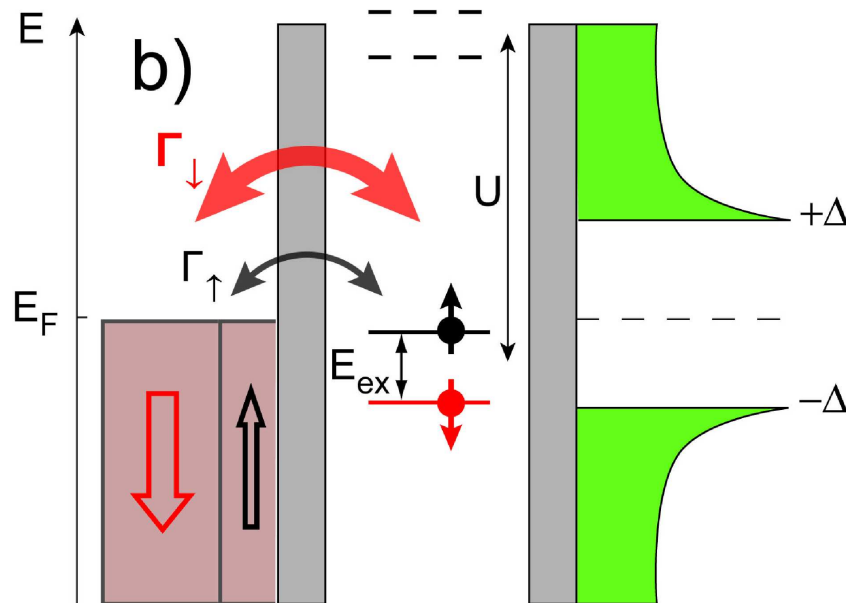
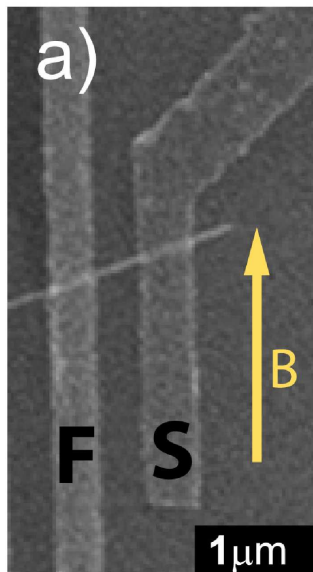
– Univ. of Basel group



L. Hofstetter et al, *Phys. Rev. Lett.* **104**, 246804 (2010).

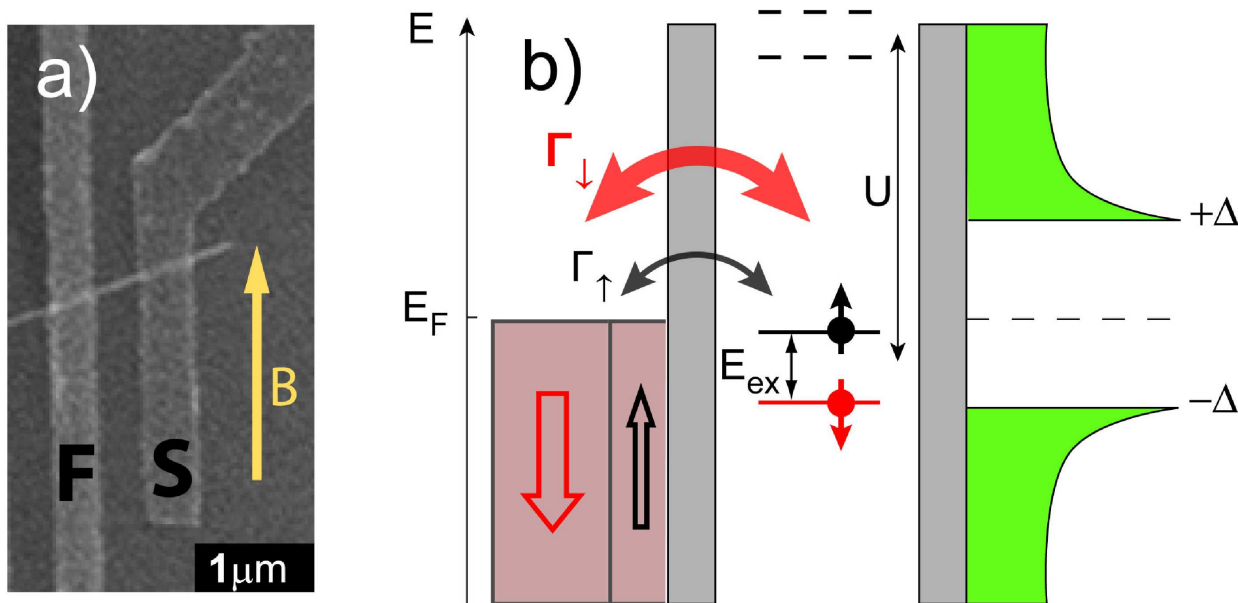


*L. Hofstetter et al, Phys. Rev. Lett. 104, 246804 (2010).*



Effects of ferromagnetism and superconductivity

*L. Hofstetter et al, Phys. Rev. Lett. 104, 246804 (2010).*



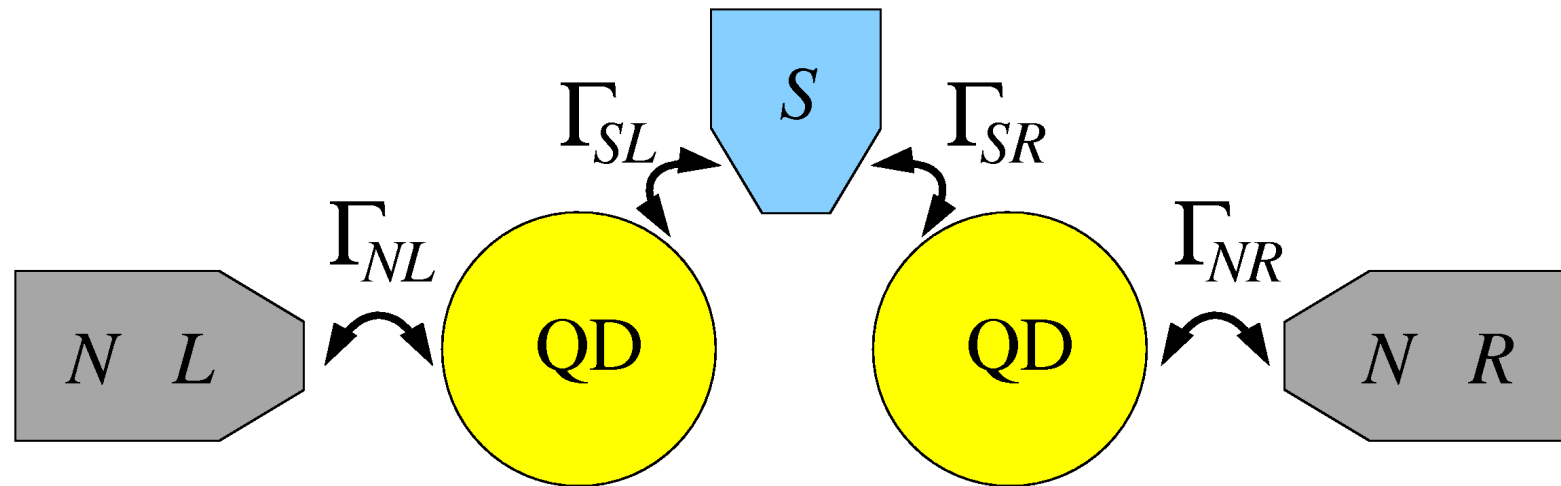
## Effects of ferromagnetism and superconductivity

- left** – Ni/Co/Pd trilayer ferromagnet
- QD** – InAs nanowire
- right** – Ti/Al bilayer superconductor

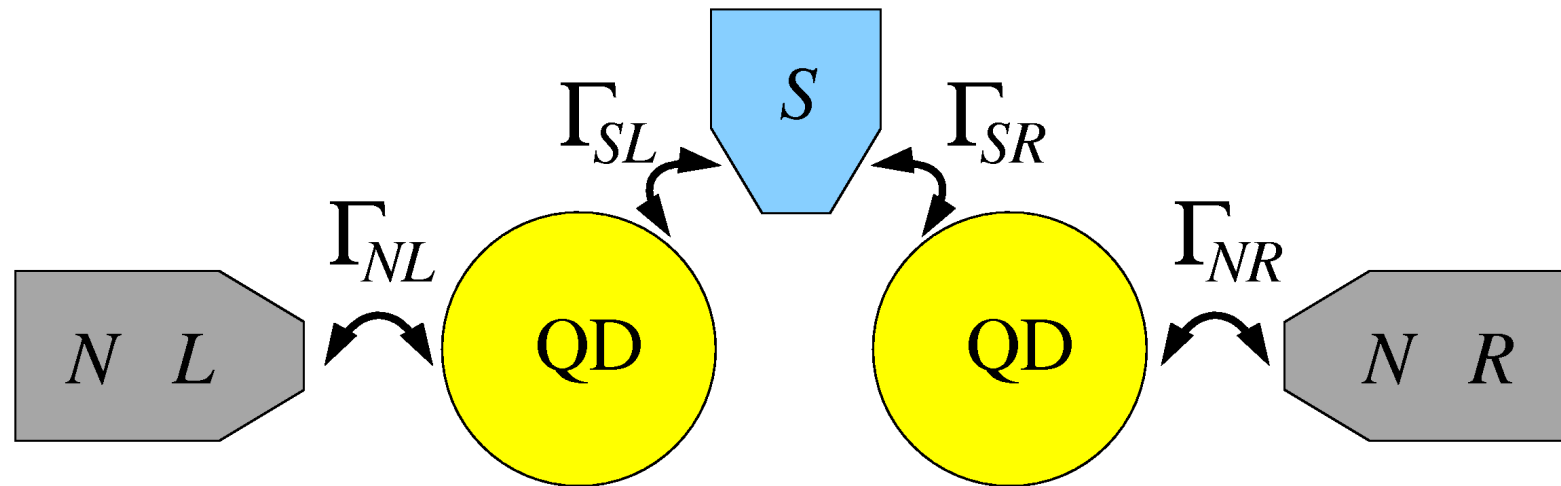
## Three terminal junctions



## Three terminal junctions

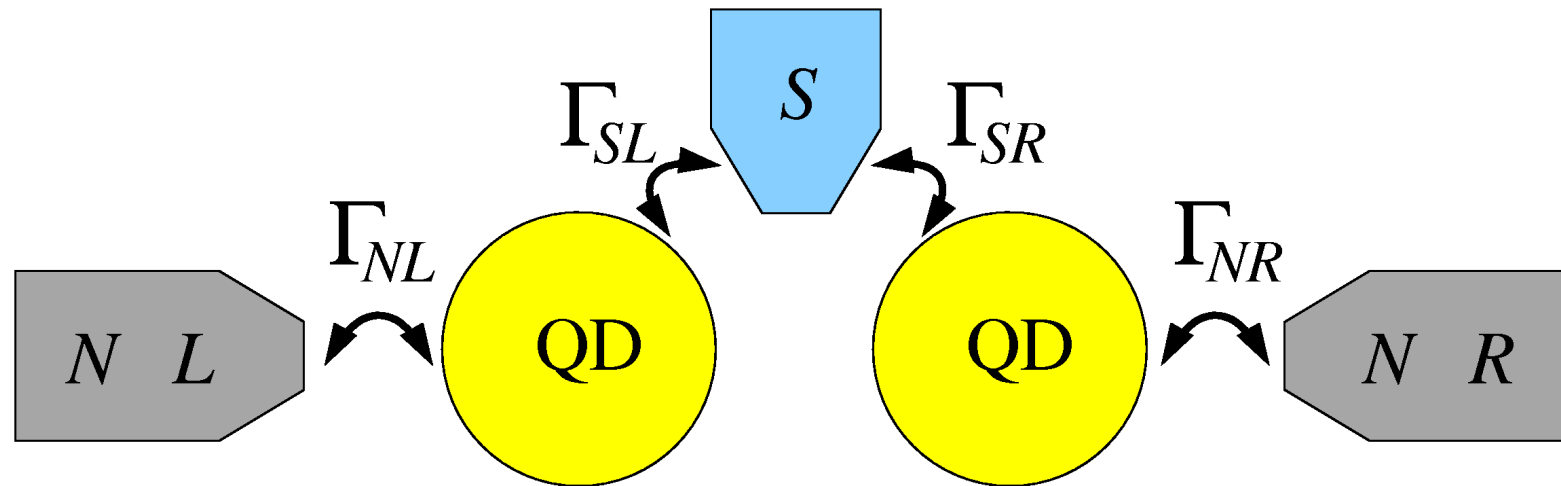


## Three terminal junctions



**Crossed Andreev reflections tunable via gate voltages**

## Three terminal junctions

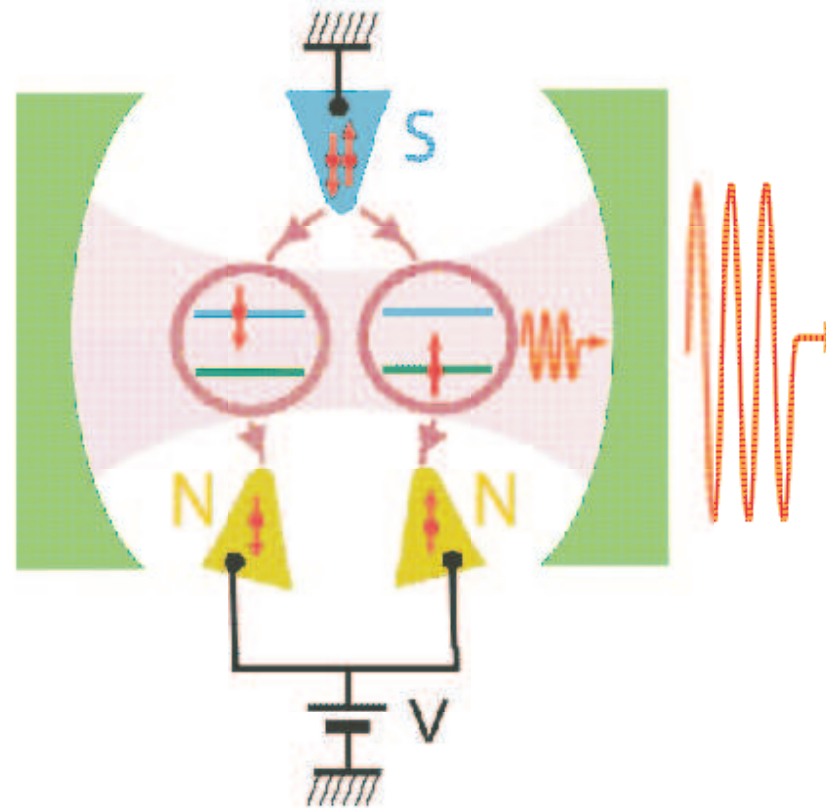


**Crossed Andreev reflections tunable via gate voltages**

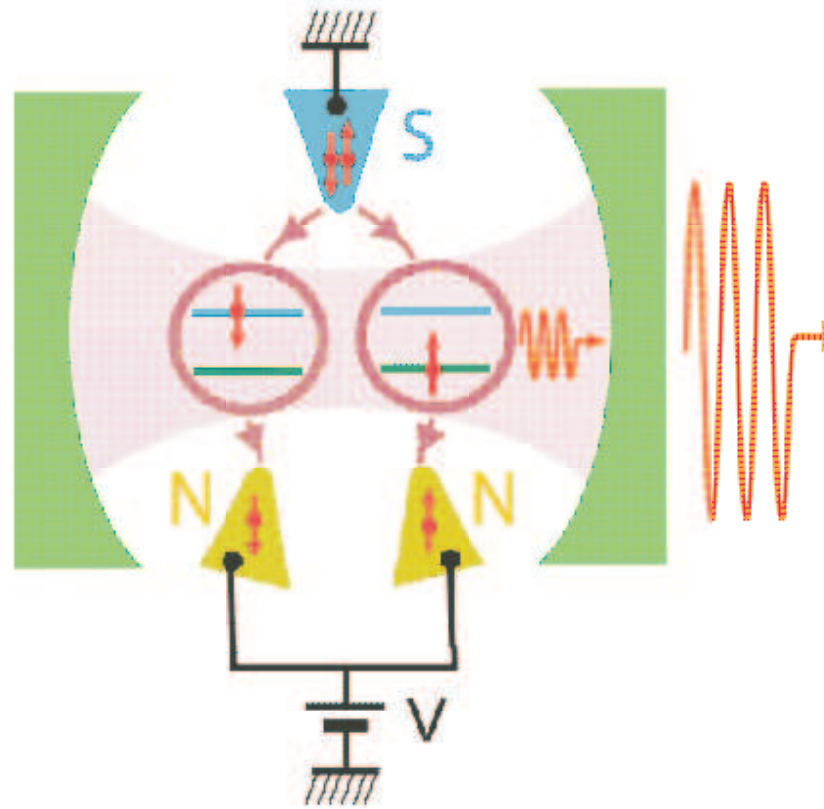
*J. Eldridge, M.G. Pala, M. Governale, J. König, Phys. Rev. B **82**, 184507 (2010)*

# Cooper pair splitter

# Cooper pair splitter

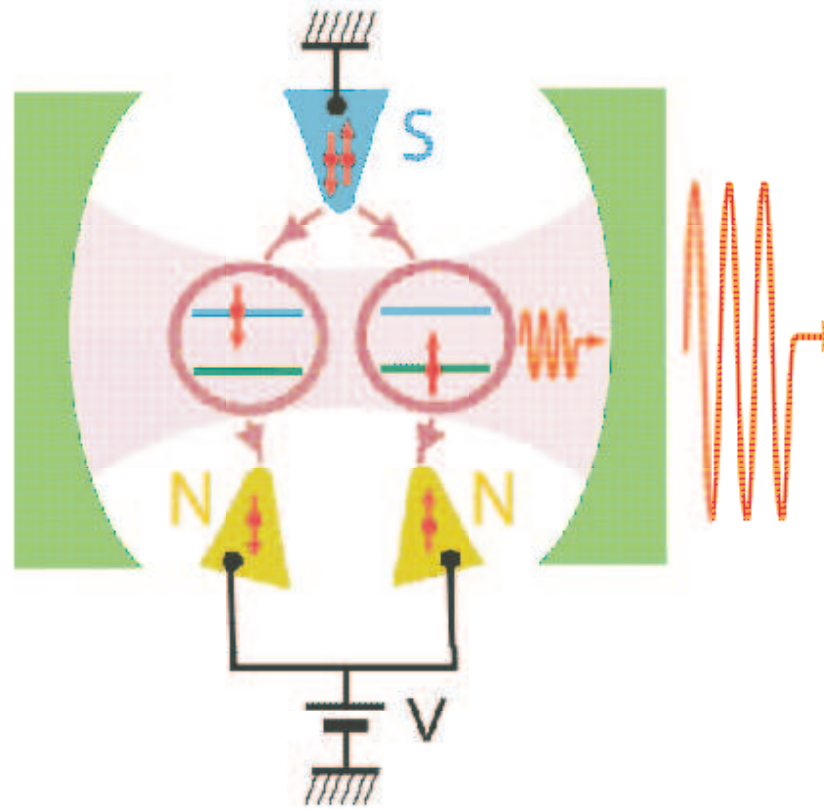


## Cooper pair splitter



Realization of the Cooper pair splitting in a microwave cavity

## Cooper pair splitter



Realization of the Cooper pair splitting in a microwave cavity

*A. Cottet, T. Kontos, and A. Levy Yeyati, Phys. Rev. Lett. **108**, 166803 (2012)*

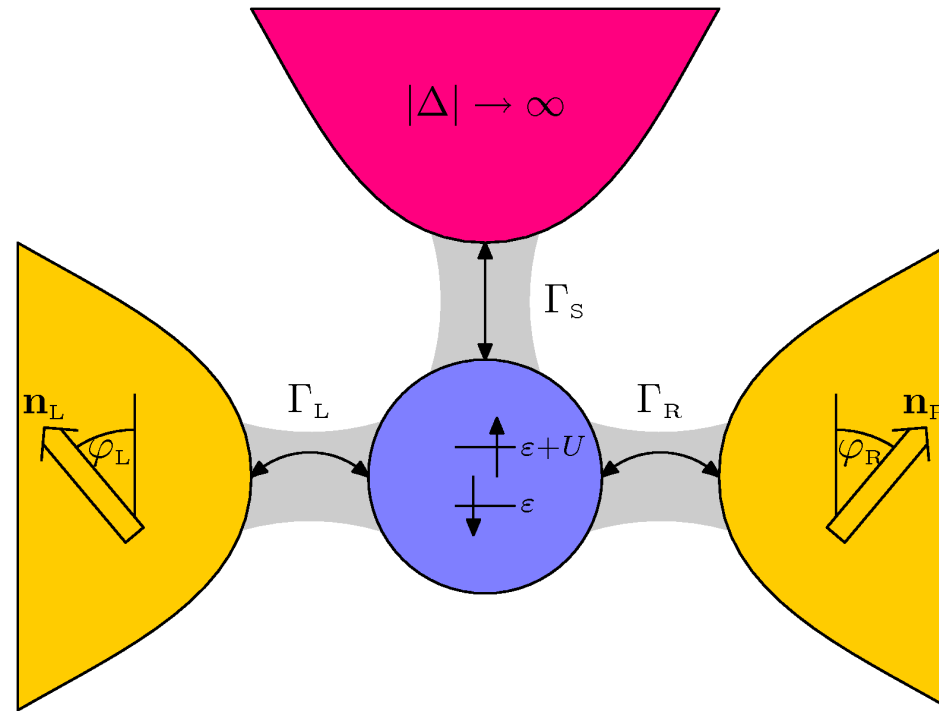
**QD spin-valve**

– in three terminal junctions



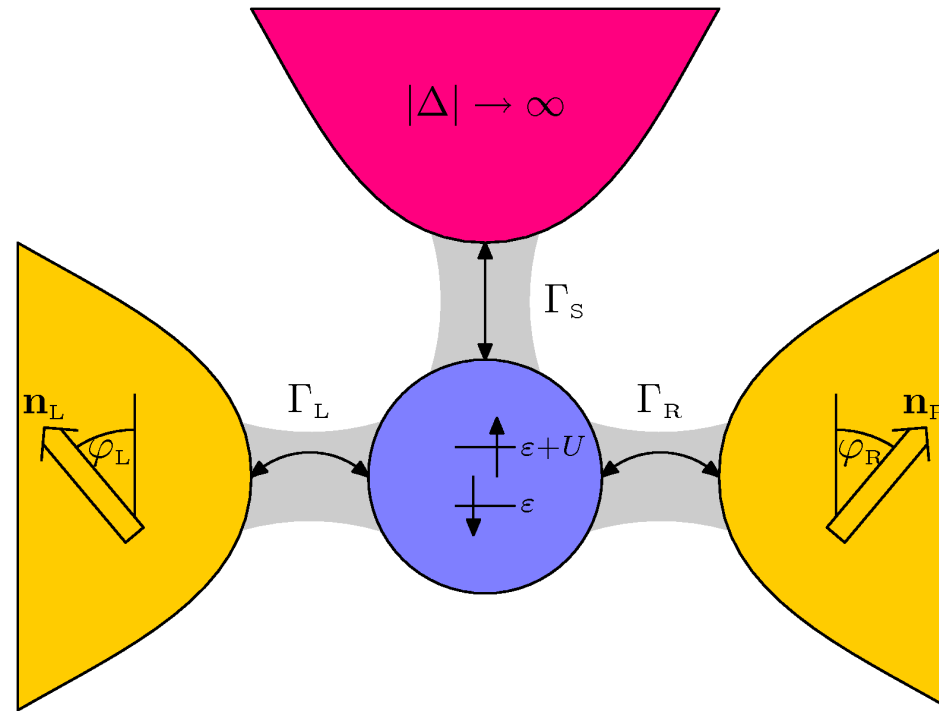
# QD spin-valve

– in three terminal junctions



## QD spin-valve

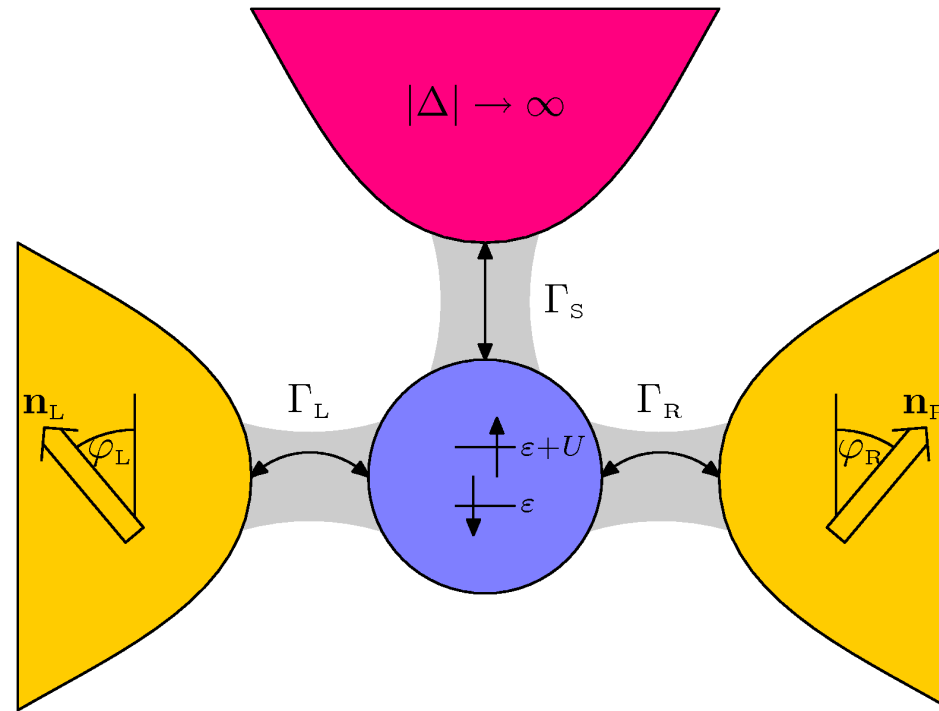
– in three terminal junctions



Idea of the spin valves using the Andreev reflections

## QD spin-valve

– in three terminal junctions



Idea of the spin valves using the Andreev reflections

*B. Sothmann, D. Futterer, M. Governale, J. König, Phys. Rev. B* **82**, 094514 (2010).

## **4. Bulk superconductors**

**Andreev spectroscopy**

– **for bulk superconductors**

## **Andreev spectroscopy**

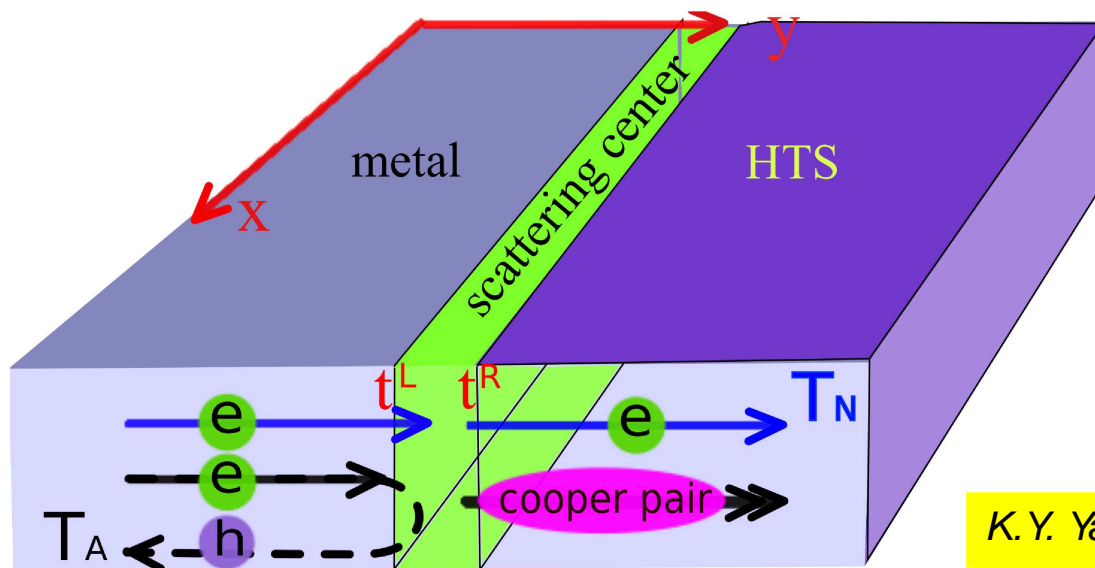
– **for bulk superconductors**

The subgap Andreev spectroscopy is also a valuable tool for studying various superconducting compounds.

## Andreev spectroscopy

— for bulk superconductors

The subgap Andreev spectroscopy is also a valuable tool for studying various superconducting compounds.



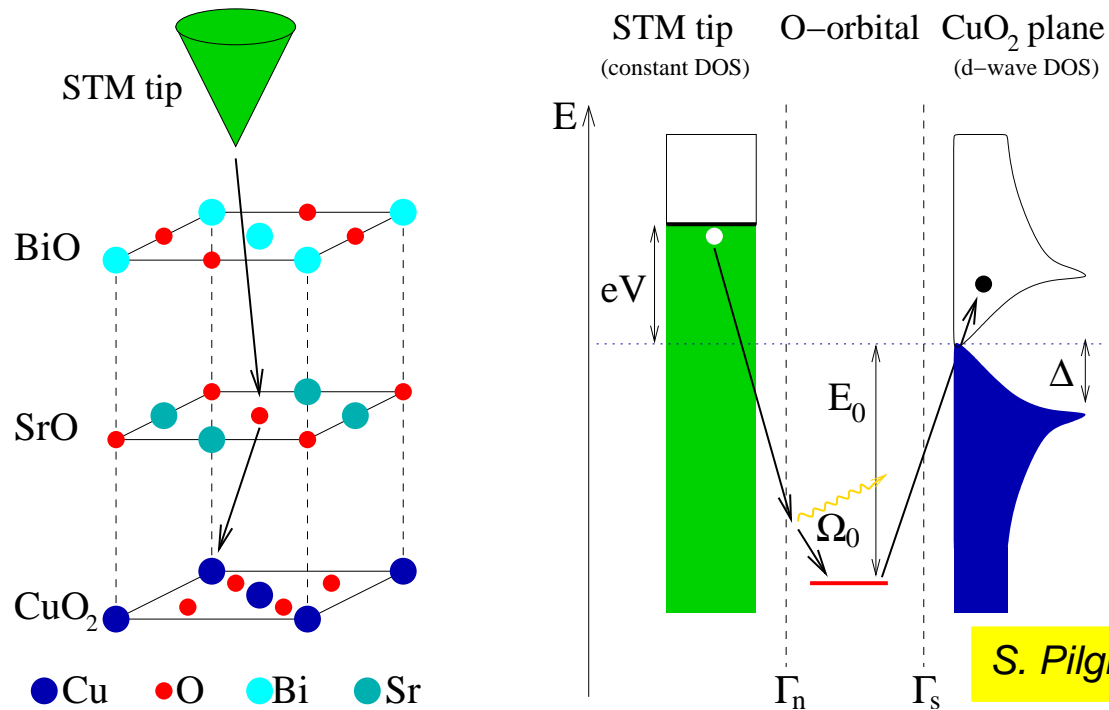
*K.Y. Yang et al, Phys. Rev. Lett. 105, 167004 (2010).*

For practical experimental realizations one can e.g. use an insulating barrier sandwiched between the conducting (N) and the probed superconductor (S).

# Andreev spectroscopy

— for bulk superconductors

The subgap Andreev spectroscopy is also a valuable tool for studying various superconducting compounds.



*S. Pilgram et al, Phys. Rev. Lett. 97, 117003 (2006).*

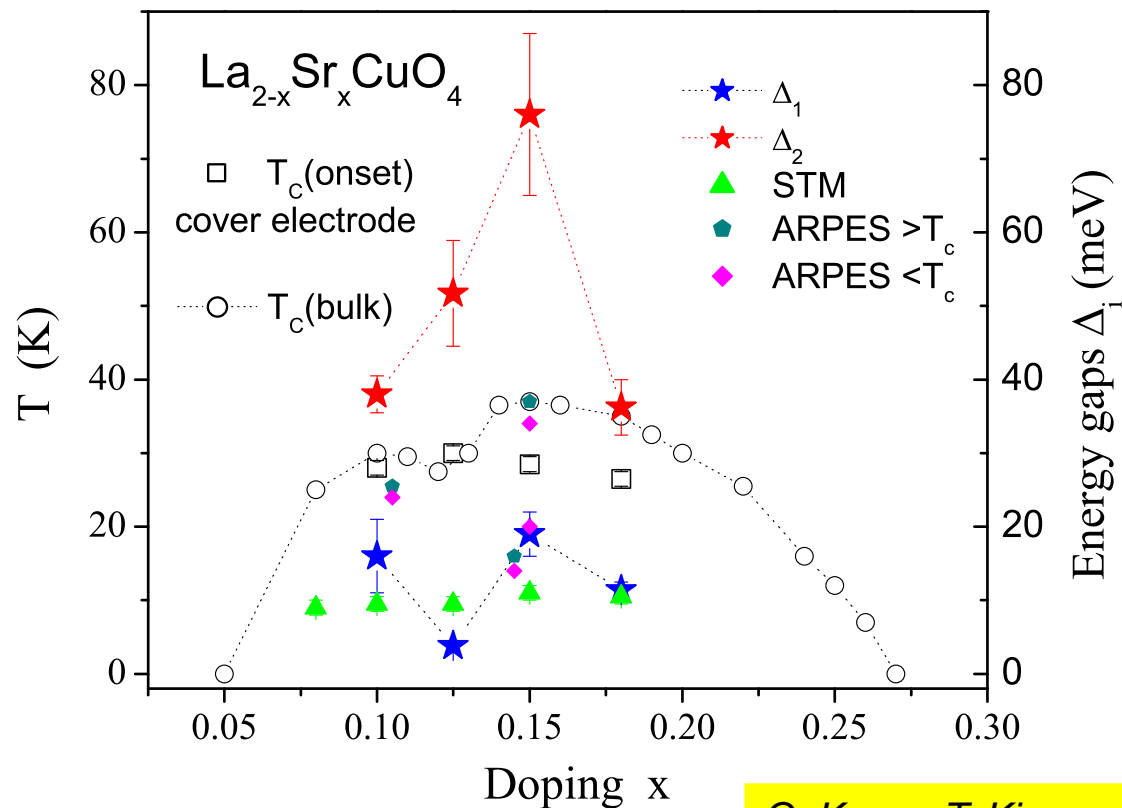
Other experimental realizations are also possible in the STM configuration, where the apex oxygen atoms play a role similar to QD in the N-QD-S setup.



## Andreev spectroscopy

— for bulk superconductors

The subgap Andreev spectroscopy is also a valuable tool for studying various superconducting compounds.



G. Koren, T. Kirzner, *Phys. Rev. Lett.* **106**, 017002 (2011).

Such Andreev spectroscopy has revealed the intriguing two-gap feature.

## **5. Ultracold gasses**

**Andreev spectroscopy**

– **for ultracold atoms**

## **Andreev spectroscopy**

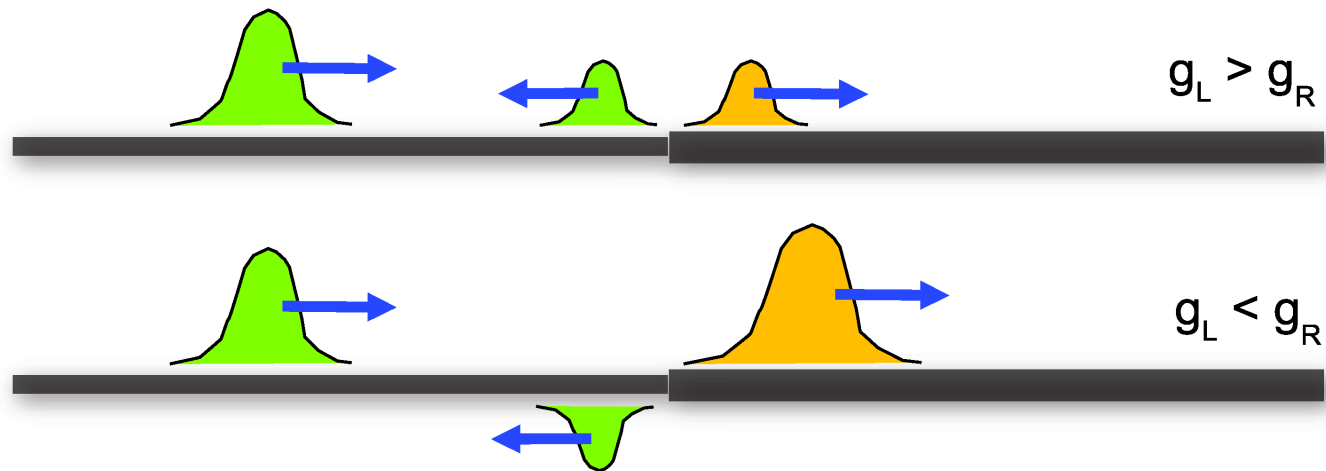
– **for ultracold atoms**

**Proposal for the Andreev-type spectroscopy has been discussed also in a context of the superfluid ultracold fermion atom systems.**

## Andreev spectroscopy

– for ultracold atoms

Proposal for the Andreev-type spectroscopy has been discussed also in a context of the superfluid ultracold fermion atom systems.



*A.J. Daley, P. Zoller, and B. Trauzettel, Phys. Rev. Lett. **100**, 110404 (2008).*

The wave packet propagating along the 1-dimensional optical lattice can be scattered at an interaction boundary in the Andreev-type fashion.

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

describes:

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

describes:

$\hat{c}_{\sigma}^{(\dagger)}(\mathbf{r})$  fermion atoms ..... (*open channel*)



## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

describes:

$\hat{c}_{\sigma}^{(\dagger)}(\mathbf{r})$  fermion atoms ..... (*open channel*)

$\hat{b}^{(\dagger)}(\mathbf{r})$  molecules ..... (*closed channel*)

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

**describes:**

$\hat{c}_{\sigma}^{(\dagger)}(\mathbf{r})$  fermion atoms ..... (*open channel*)

$\hat{b}^{(\dagger)}(\mathbf{r})$  molecules ..... (*closed channel*)

**resonantly interacting via:**

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

**describes:**

$\hat{c}_{\sigma}^{(\dagger)}(\mathbf{r})$  fermion atoms ..... (*open channel*)

$\hat{b}^{(\dagger)}(\mathbf{r})$  molecules ..... (*closed channel*)

**resonantly interacting via:**

$\hat{b}^{\dagger} \hat{c}_{\downarrow} \hat{c}_{\uparrow} + h.c.$  ..... (*Feshbach resonance*)

## Feshbach resonance

[ local problem ]

$$\begin{aligned}\hat{H}_{loc}(\mathbf{r}) = & \sum_{\sigma} \varepsilon(\mathbf{r}) \hat{c}_{\sigma}^{\dagger}(\mathbf{r}) \hat{c}_{\sigma}(\mathbf{r}) + E(\mathbf{r}) \hat{b}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \\ & + g \left( \hat{b}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}(\mathbf{r}) \hat{c}_{\uparrow}(\mathbf{r}) + \hat{c}_{\uparrow}^{\dagger}(\mathbf{r}) \hat{c}_{\downarrow}^{\dagger}(\mathbf{r}) \hat{b}(\mathbf{r}) \right)\end{aligned}$$

describes:

$\hat{c}_{\sigma}^{(\dagger)}(\mathbf{r})$  fermion atoms ..... (*open channel*)

$\hat{b}^{(\dagger)}(\mathbf{r})$  molecules ..... (*closed channel*)

resonantly interacting via:

$\hat{b}^{\dagger} \hat{c}_{\downarrow} \hat{c}_{\uparrow} + h.c.$  ..... (*Feshbach resonance*)

M.L. Chiofalo, S.J.J.M.F. Kokkelmans, J.N. Milstein, and M.J. Holland, *Phys. Rev. Lett.* **88**, 090402 (2002).

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

$\varepsilon$

..... energy of non-bonding state

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

$\varepsilon$  ..... energy of non-bonding state

$Z(T)$  ..... the spectral weight



## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

$\varepsilon$  ..... energy of non-bonding state

$Z(T)$  ..... the spectral weight

$\varepsilon_{\pm} = E/2 \pm \sqrt{(\varepsilon - E/2)^2 + g^2}$  ..... BCS-like excitation energies

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

$\varepsilon$  ..... energy of non-bonding state

$Z(T)$  ..... the spectral weight

$\varepsilon_{\pm} = E/2 \pm \sqrt{(\varepsilon - E/2)^2 + g^2}$  ..... BCS-like excitation energies

$u^2, v^2 = \frac{1}{2} \left[ 1 \pm (\varepsilon - E/2) / \sqrt{(\varepsilon - E/2)^2 + g^2} \right]$  ..... BCS-like coefficients

## Local solution

[ exact ]

$$\mathcal{G}_{loc}(i\omega_n) = [1 - Z(T)] \left( \frac{u^2}{i\omega_n - \varepsilon_+} + \frac{v^2}{i\omega_n - \varepsilon_-} \right) + \frac{Z(T)}{i\omega_n - \varepsilon}$$

where

$\varepsilon$  ..... energy of non-bonding state

$Z(T)$  ..... the spectral weight

$\varepsilon_{\pm} = E/2 \pm \sqrt{(\varepsilon - E/2)^2 + g^2}$  ..... BCS-like excitation energies

$u^2, v^2 = \frac{1}{2} \left[ 1 \pm (\varepsilon - E/2) / \sqrt{(\varepsilon - E/2)^2 + g^2} \right]$  ..... BCS-like coefficients

*T. Domański, Eur. Phys. J. B* **33**, 41 (2003);

*T. Domański et al, Sol. State Commun.* **105**, 473 (1998).

## Approximate general solution

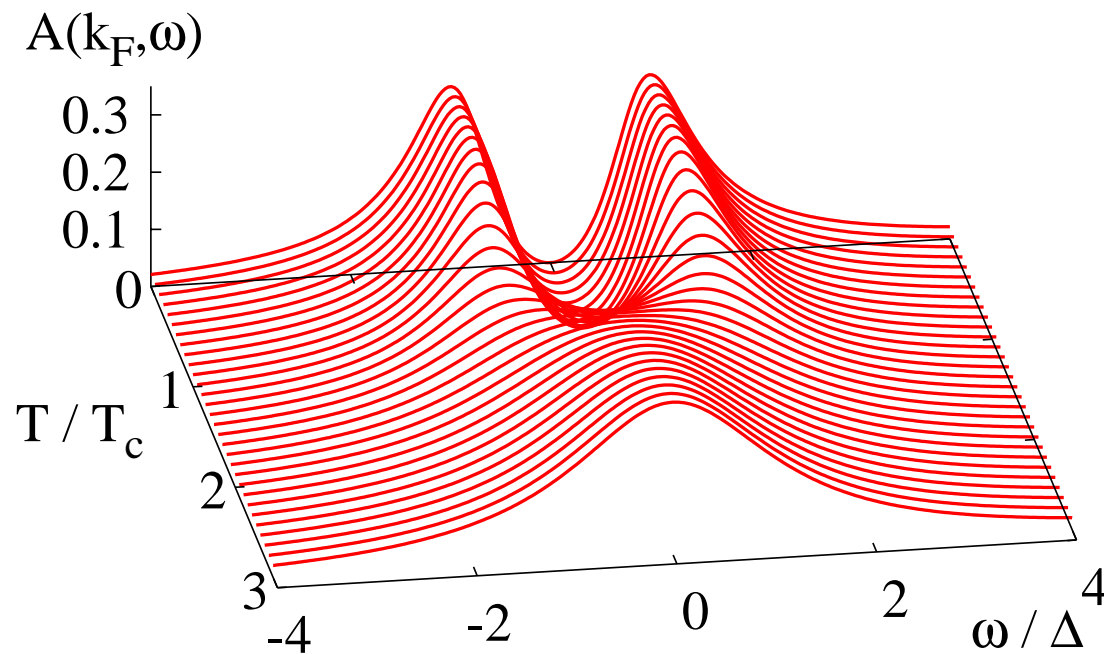
[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$

## Approximate general solution

[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$

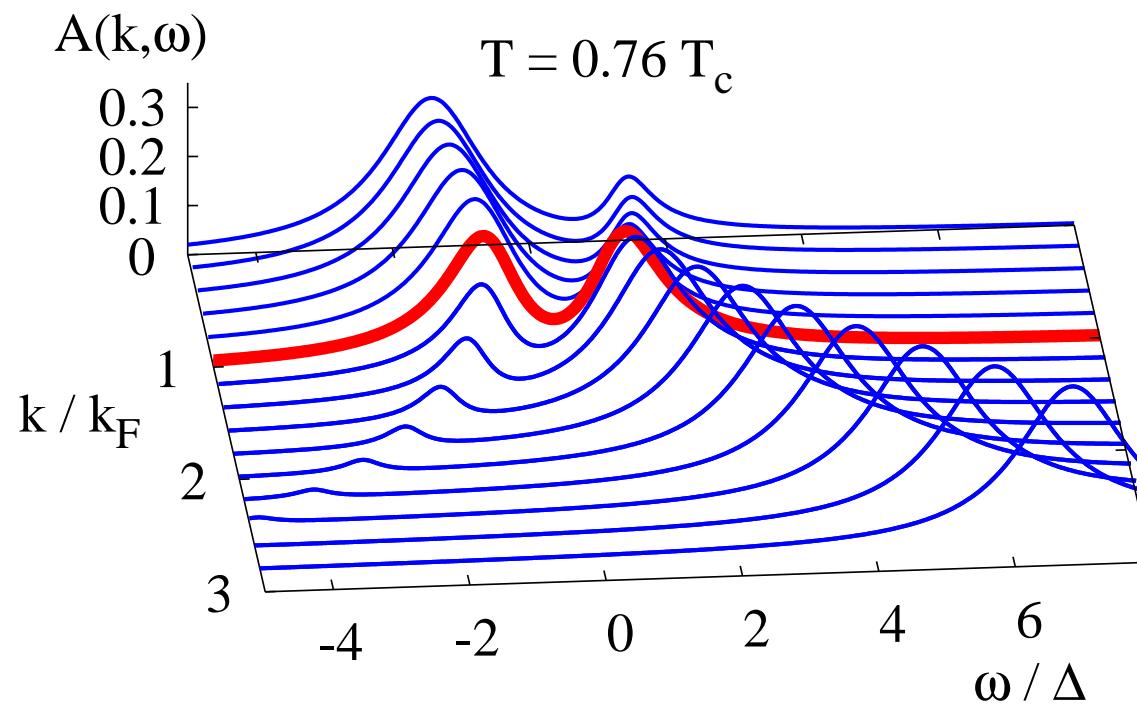


*T. Domański, Phys. Rev. A* **84**, 023634 (2011).

## Approximate general solution

[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$

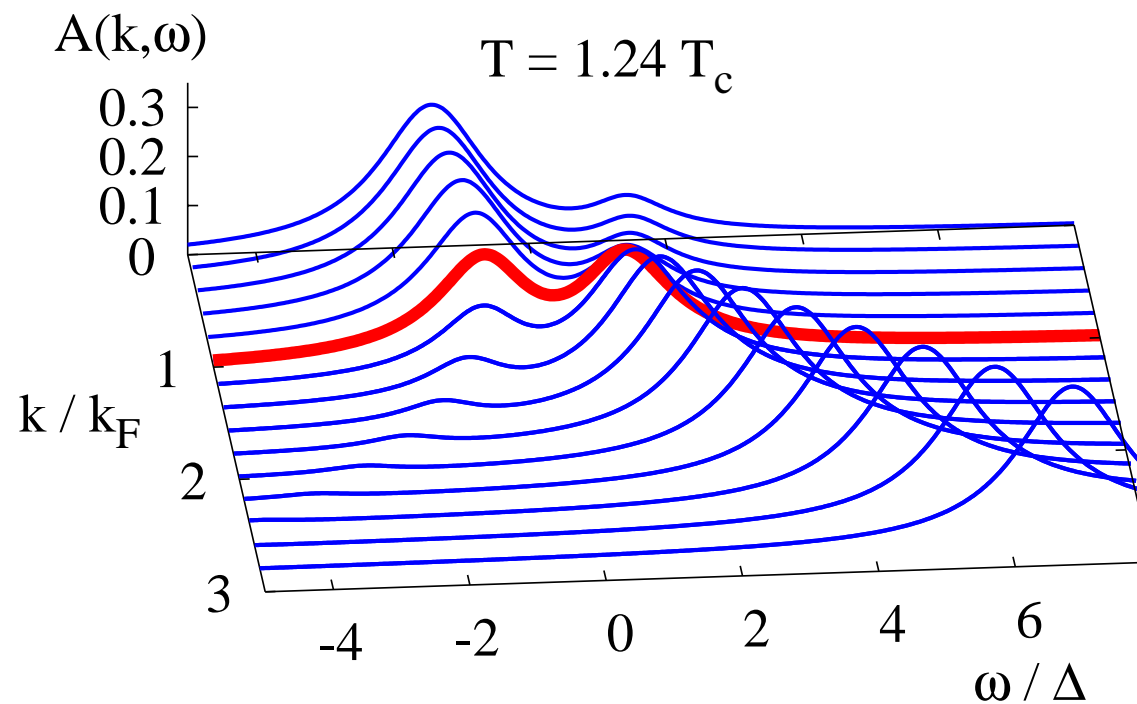


*T. Domański, Phys. Rev. A* **84**, 023634 (2011).

## Approximate general solution

[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$

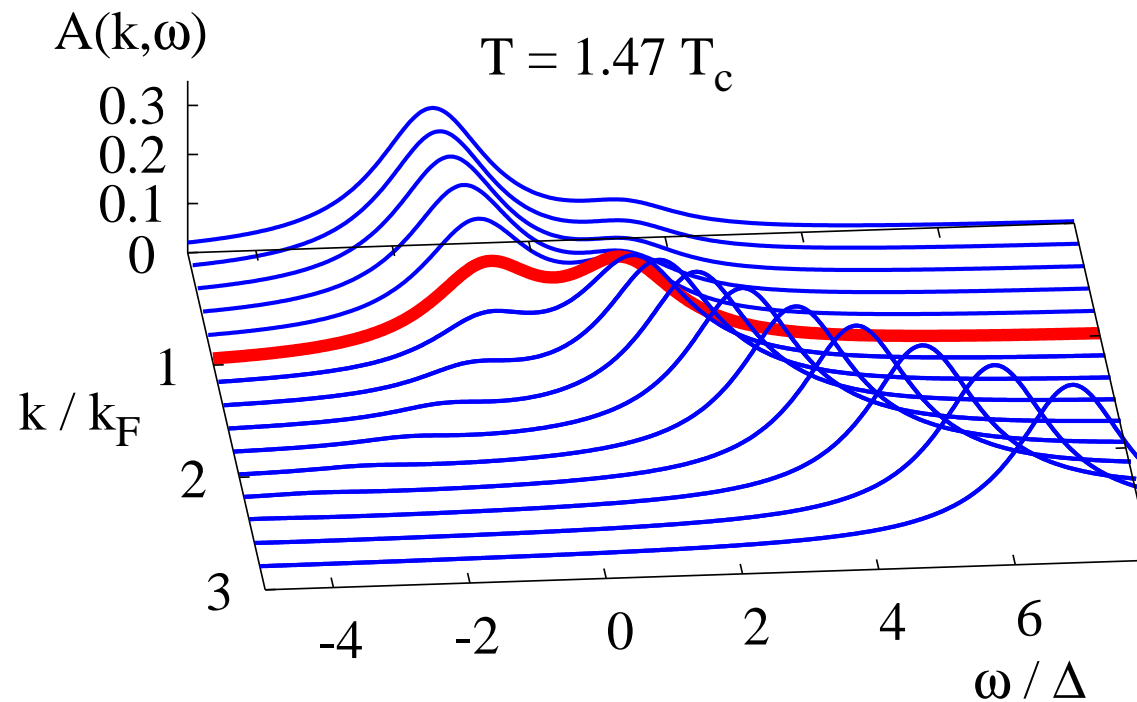


*T. Domański, Phys. Rev. A* **84**, 023634 (2011).

## Approximate general solution

[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$



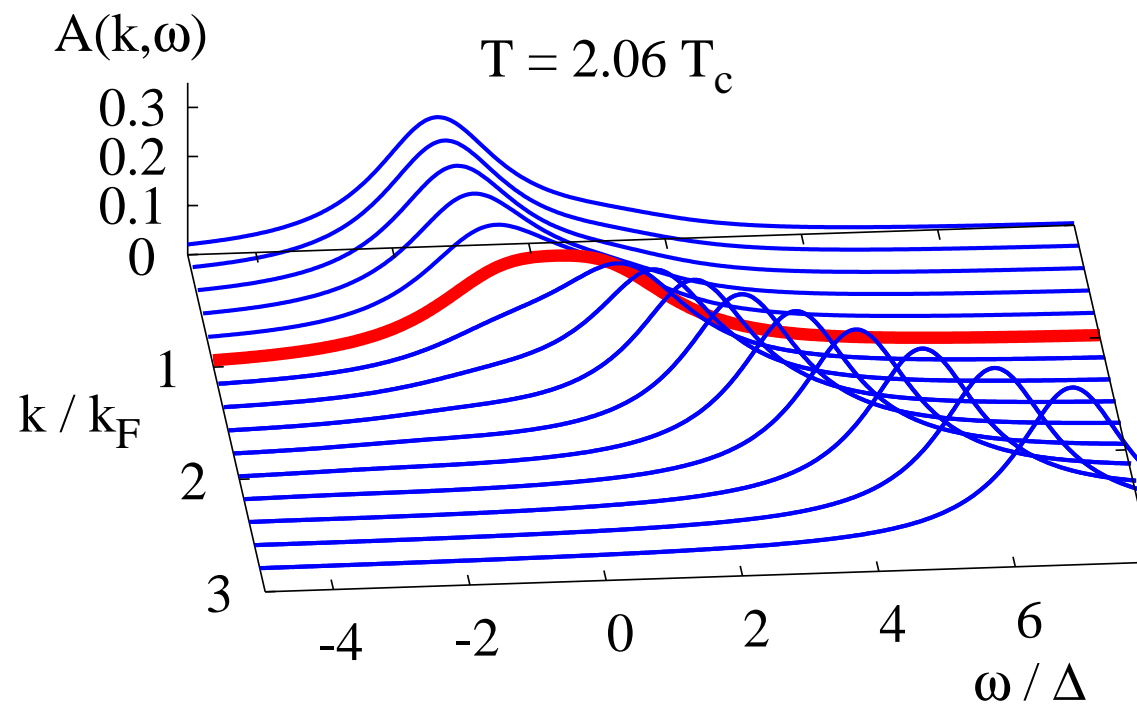
*T. Domański, Phys. Rev. A* **84**, 023634 (2011).



## Approximate general solution

[ near the unitary limit ]

$$\hat{H} = \int dr \left( \hat{T}_{kin}(r) + \hat{H}_{loc}(r) \right)$$

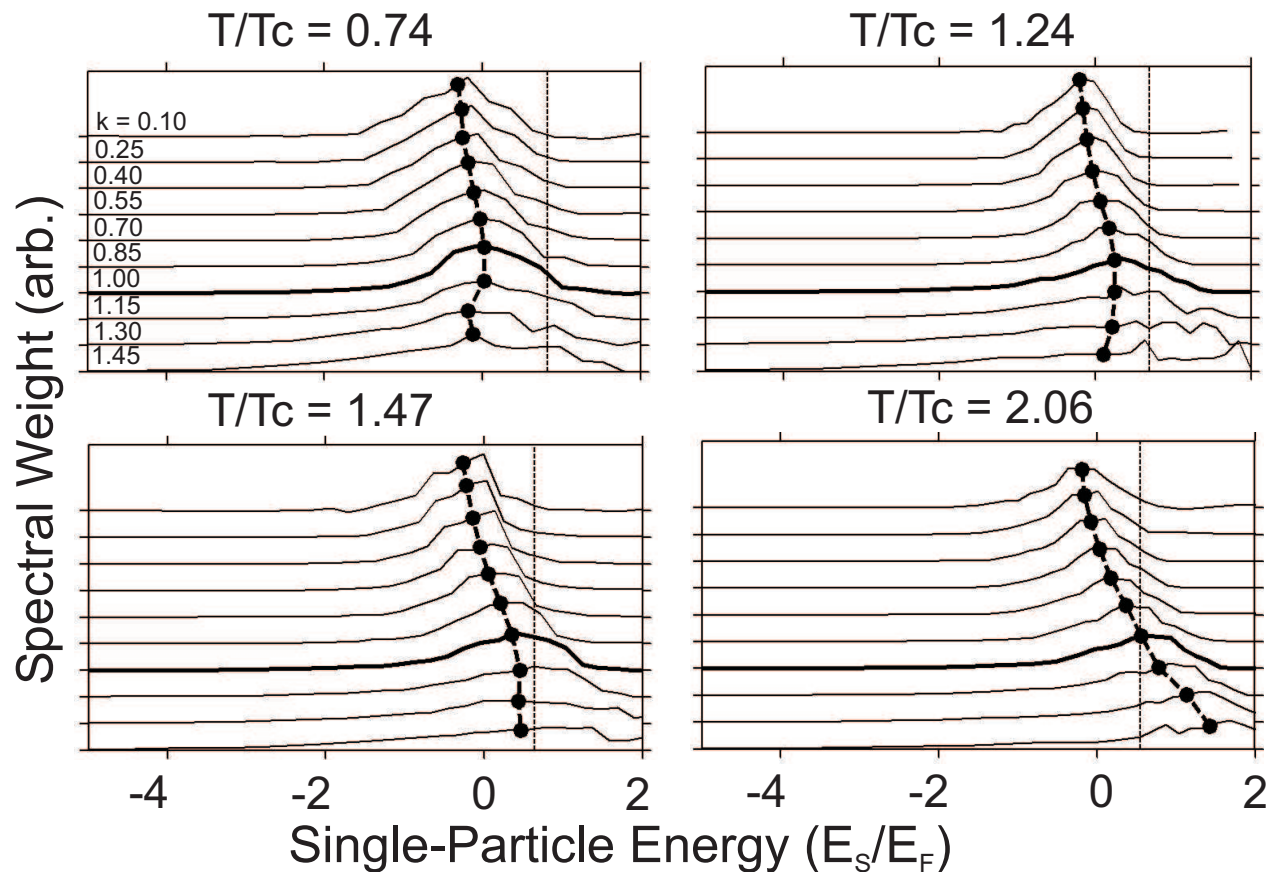


*T. Domański, Phys. Rev. A* **84**, 023634 (2011).

## Evidence for Bogoliubov QPs above $T_c$

# Evidence for Bogoliubov QPs above $T_c$

## D. Jin group (Boulder, USA)



Results for the ultracold  $^{40}\text{K}$  atoms

*J.P. Gaebler et al, Nature Phys. 6, 569 (2010).*

# Conclusions

# Conclusions

**Andreev spectroscopy :**

# Conclusions

**Andreev spectroscopy :**

⇒ **is a suitable tool for probing the pair-coherence**

# Conclusions

**Andreev spectroscopy :**

- ⇒ **is a suitable tool for probing the pair-coherence**
- ⇒ **simultaneously exploring the particle and hole states**

# Conclusions

**Andreev spectroscopy :**

- ⇒ **is a suitable tool for probing the pair-coherence**
- ⇒ **simultaneously exploring the particle and hole states**

**It can be applied to:**



# Conclusions

**Andreev spectroscopy :**

- ⇒ **is a suitable tool for probing the pair-coherence**
- ⇒ **simultaneously exploring the particle and hole states**

**It can be applied to:**

- **nanoscopic objects**  
**/ coupled to superconducting electrodes /**

# Conclusions

**Andreev spectroscopy :**

- ⇒ **is a suitable tool for probing the pair-coherence**
- ⇒ **simultaneously exploring the particle and hole states**

**It can be applied to:**

- **nanoscopic objects**  
/ coupled to superconducting electrodes /
- **bulk materials**  
/ superconductors, atomic superfluids, black holes (?) /