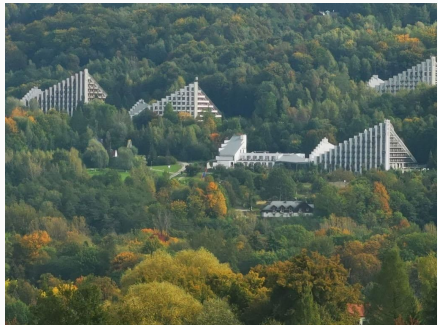


Correlations in superconducting nanostructures under nonequilibrium cases

Tadeusz Domański

**M. Curie-Skłodowska University
Lublin, Poland**



MAIN ISSUES

1. Pairing and correlations

[in nanoscopic scale]

2. Dynamical phenomena

[affecting in-gap states]

3. Out-of-equilibrium transport

[via Andreev scattering]

Superconducting nanostructures

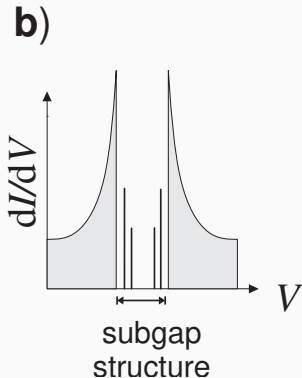
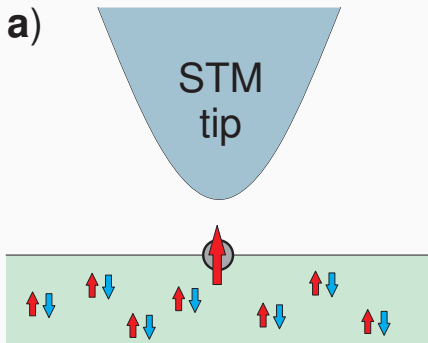
Superconducting nanostructures

some examples ...

SUPERCONDUCTING NANOSTRUCTURES

1. Local spectroscopy:

quantum impurity on a surface of superconductor + STM tip

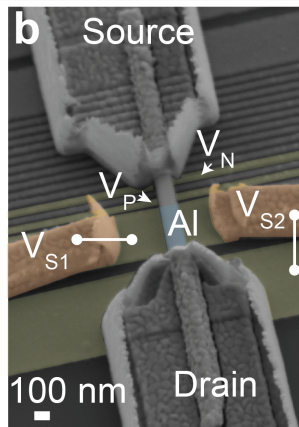
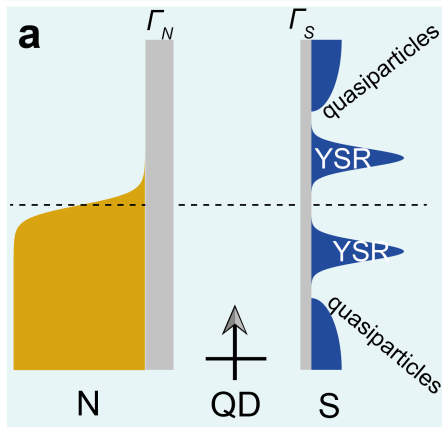


Differential conductance probes the effective spectrum of impurity.

SUPERCONDUCTING NANOSTRUCTURES

2. Andreev junctions:

normal metal (N) - quantum dot (QD) - superconductor (S)

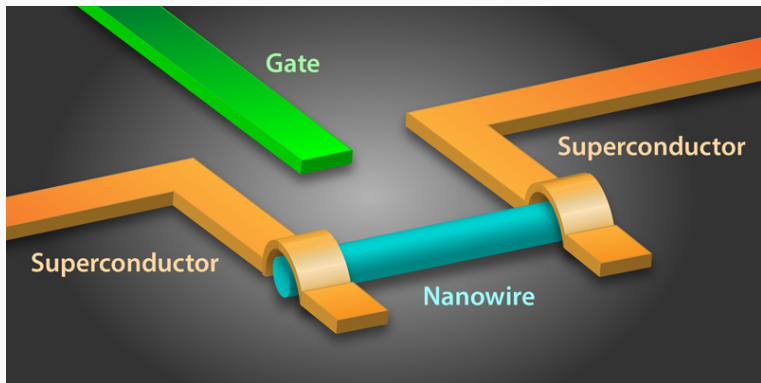


Subgap tunneling via the electron-to-hole (Andreev) scattering.

SUPERCONDUCTING NANOSTRUCTURES

3. Josephson junctions:

superconductor (S) - quantum dot (QD) - superconductor (S)



Tunneling of Cooper pairs via bound states in Josephson junction.

SUPERCONDUCTING PROXIMITY EFFECT

Quantum dot (QD) coupled to bulk superconductor (SC) experiences:

⇒ **on-dot pairing**

SUPERCONDUCTING PROXIMITY EFFECT

Quantum dot (QD) coupled to bulk superconductor (SC) experiences:

⇒ **on-dot pairing**

which is spectroscopically manifested by:

⇒ **in-gap bound states**

SUPERCONDUCTING PROXIMITY EFFECT

Quantum dot (QD) coupled to bulk superconductor (SC) experiences:

⇒ **on-dot pairing**

which is spectroscopically manifested by:

⇒ **in-gap bound states**

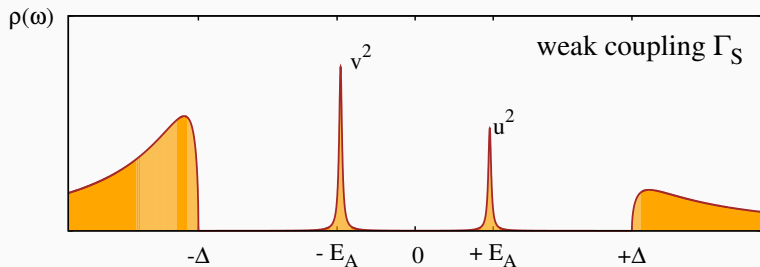
originating from:

⇒ **leakage of the Cooper pairs onto QD** (Andreev)

⇒ **interaction of QD spin with SC** (Yu-Shiba-Rusinov)

IN-GAP STATES

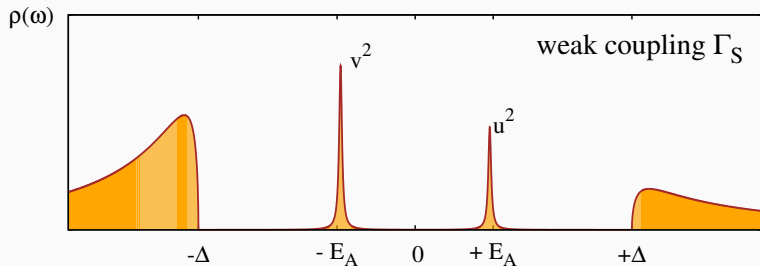
Spectrum of the quantum impurity coupled to superconductor



Bound states appear at $\pm E_A$ in the subgap region $E \in \langle -\Delta, \Delta \rangle$

IN-GAP STATES

Spectrum of the quantum impurity coupled to superconductor

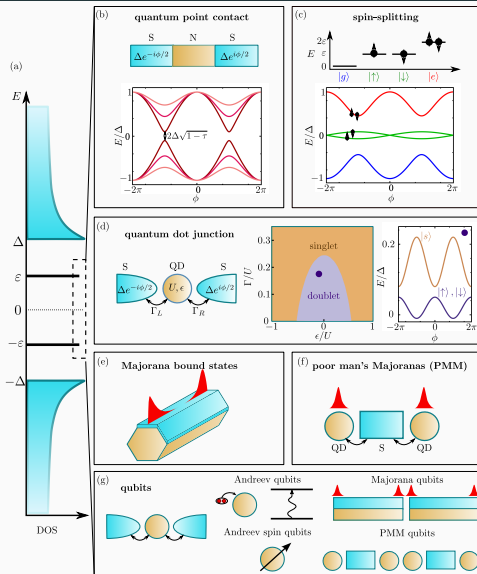


Bound states appear at $\pm E_A$ in the subgap region $E \in \langle -\Delta, \Delta \rangle$

Let's focus on such in-gap bound states ...

Why ?

VARIETY OF SUPERCONDUCTING QUBITS



Examples ...

SUPERCONDUCTING PROCESSOR: WILLOW

In December 2024 Google demonstrated 105-qubit processor based on superconducting qubits (transmons).

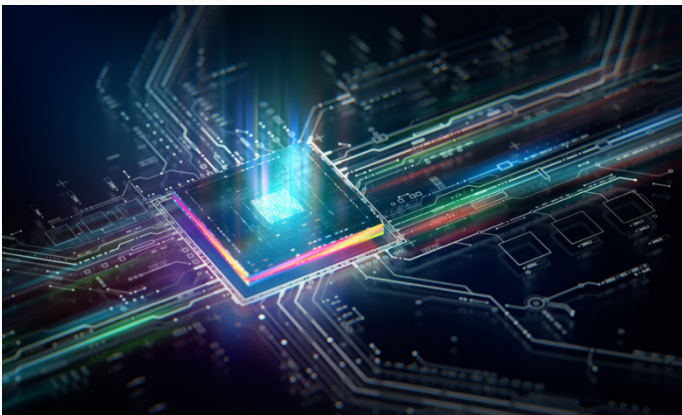


Google Quantum AI and collaborators, Nature 638, 920 (2024).

SUPERCONDUCTING PROCESSOR: WILLOW

Simulation of the probability distribution obtained in 5 minutes by processor

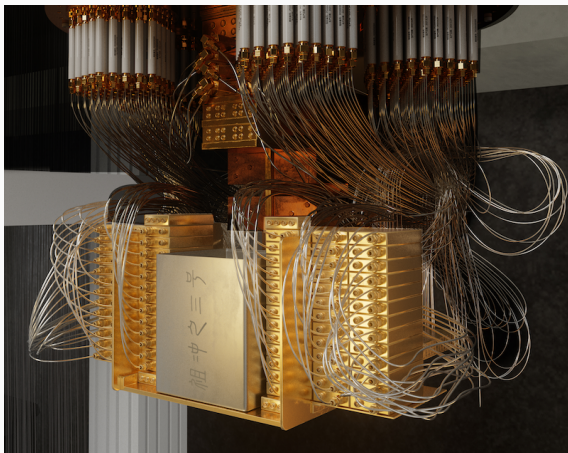
Willow would take about 10^{25} years by the fastest classical computer.



H. Neven (Google blog, 9 December 2024).

SC PROCESSOR: ZUCHONGZHI 3.0

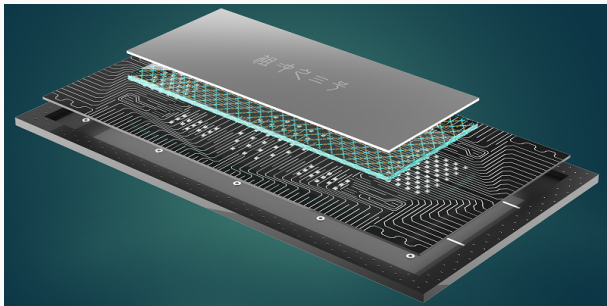
105-qubit processor constructed by the group of prof. Jian-Wei Pan
(University of Science and Technology, China)



D. Gao et al, Phys. Rev. Lett. 134, 090601 (2025).

SC PROCESSOR: ZUCHONGZHI 3.0

Simulation of the probability distribution obtained in 100 seconds by processor Zuchongzhi 3.0 would take at least several 10^6 years by the fastest classical computer.



Zuchongzhi 3.0 processor consists of 105 qubits: 15 qubits in 7 arrays.

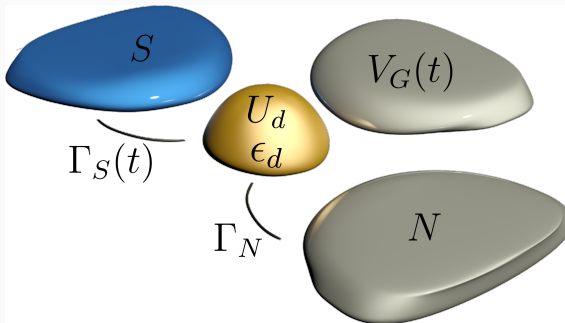
D. Gao et al, Phys. Rev. Lett. 134, 090601 (2025).

Characteristic time-scales

Characteristic time-scales

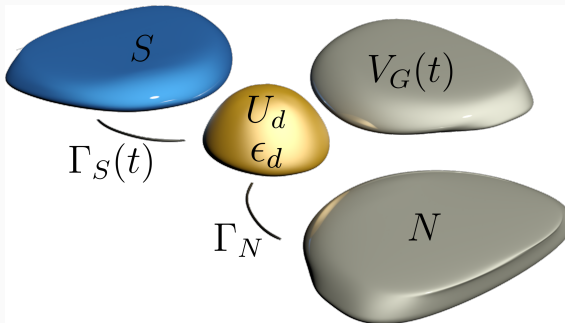
relevant to operations on bound states

DYNAMICS OF IN-GAP STATES



Empirical protocols for time-resolved phenomena:

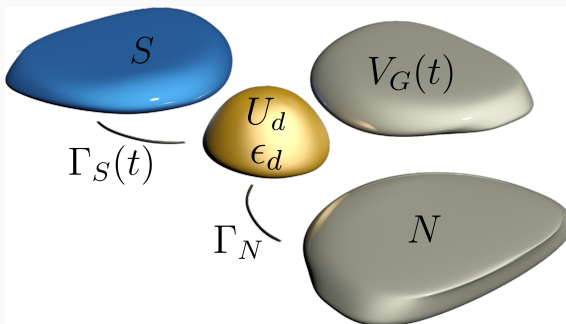
DYNAMICS OF IN-GAP STATES



Empirical protocols for time-resolved phenomena:

\Rightarrow variation of the coupling Γ_S to superconductor

DYNAMICS OF IN-GAP STATES



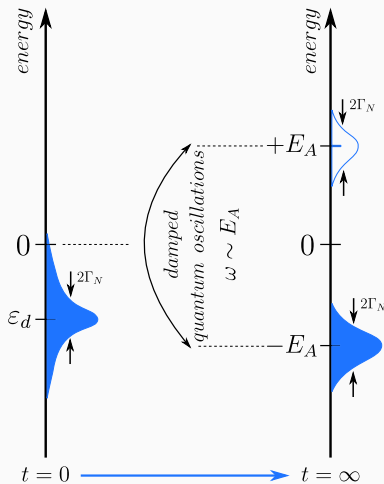
Empirical protocols for time-resolved phenomena:

\Rightarrow variation of the coupling Γ_S to superconductor

\Rightarrow change of the gate potential V_G

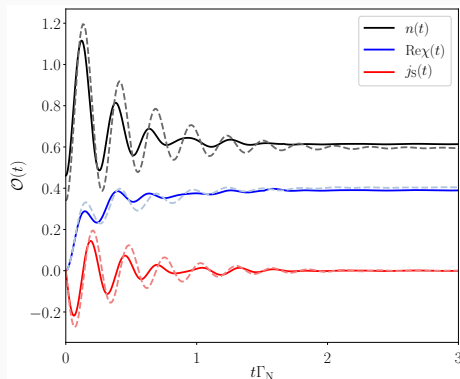
BUILDUP OF IN-GAP STATES

Sudden coupling of QD to superconductor $0 \rightarrow \Gamma_S$



BUILDUP OF IN-GAP STATES

Time-dependent observables driven by the quantum quench $0 \rightarrow \Gamma_S$



$n(t)$ electron number

$\chi(t) = \langle \hat{d}_\downarrow \hat{d}_\uparrow \rangle$ on-dot pairing

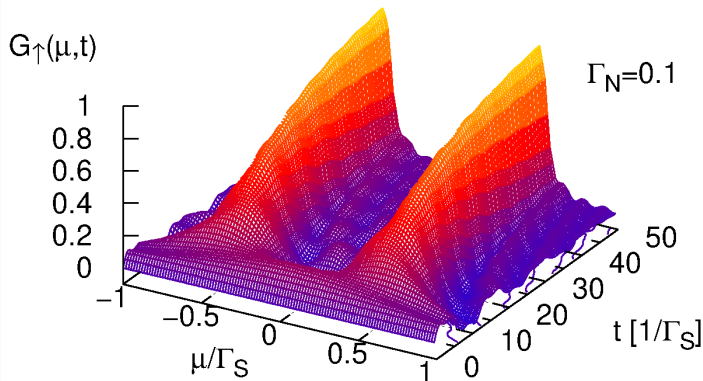
$j_S(t)$ charge current

solid lines - time dependent NRG

dashed lines - Hartree-Fock-Bogolubov

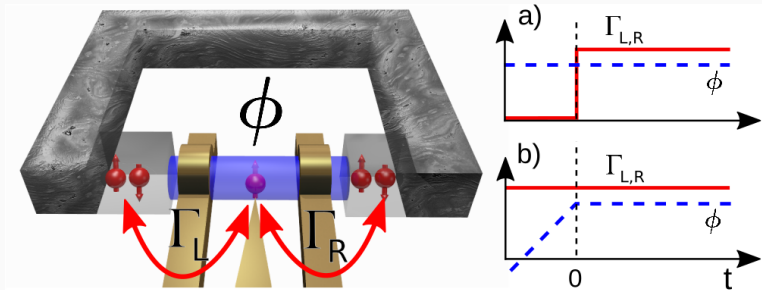
TIME-DEPENDENT TUNNELING CONDUCTANCE

Subgap tunneling conductance $G_{\sigma} = \frac{\partial I_{\sigma}(t)}{\partial \mu}$ vs time (t) and voltage (μ)



BOUND STATES IN JOSEPHSON JUNCTION

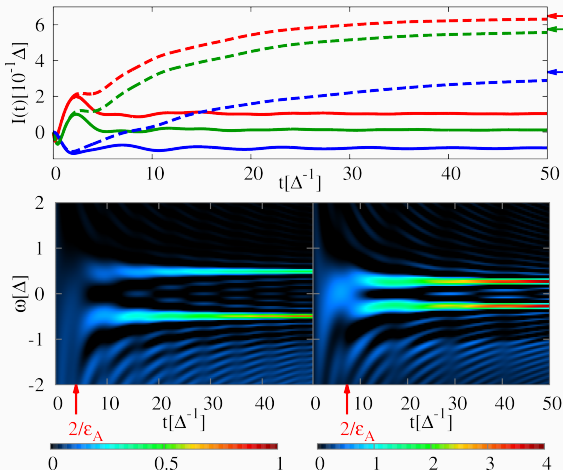
Quantum quench imposed on QD in Josephson junction geometry



R. Seoane Souto, A. Martín-Rodero, A. Levy Yeyati, *Phys. Rev. Lett.* **117**, 267701 (2016).

BOUND STATES IN JOSEPHSON JUNCTION

Transient current and quasiparticle spectrum obtained for different ratios of Γ/Δ (from top to bottom: 10, 5 and 1).



YU-SHIBA-RUSINOV STATES OF CLASSICAL IMPURITY

communications physics

ARTICLE

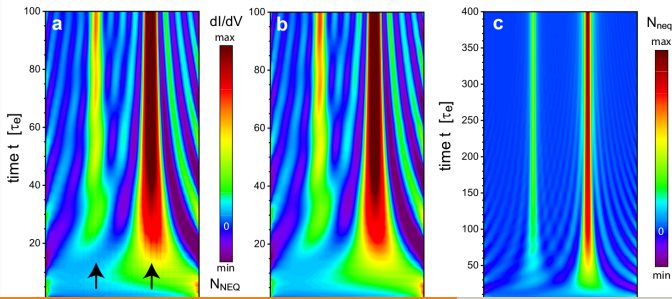
<https://doi.org/10.1038/s42005-022-01050-7>

OPEN



Emergence and manipulation of non-equilibrium Yu-Shiba-Rusinov states

Jasmin Bedow ¹, Eric Mascot ^{1,2} & Dirk K. Morr ^{1✉}



Correlation effects

Correlation effects

[singlet-doublet (quantum phase) transition]

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can be described by

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_s \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow}^{\dagger} + \text{h.c.} \right)$$

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can be described by

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_s \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow}^{\dagger} + \text{h.c.} \right)$$

Eigen-states of this problem are represented by:

$|\uparrow\rangle$ **and** $|\downarrow\rangle$ \Leftarrow **doublet states (spin $\frac{1}{2}$)**

$\left. \begin{array}{l} u |0\rangle - v |\uparrow\downarrow\rangle \\ v |0\rangle + u |\uparrow\downarrow\rangle \end{array} \right\}$ \Leftarrow **singlet states (spin 0)**

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can be described by

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_S \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow}^{\dagger} + \text{h.c.} \right)$$

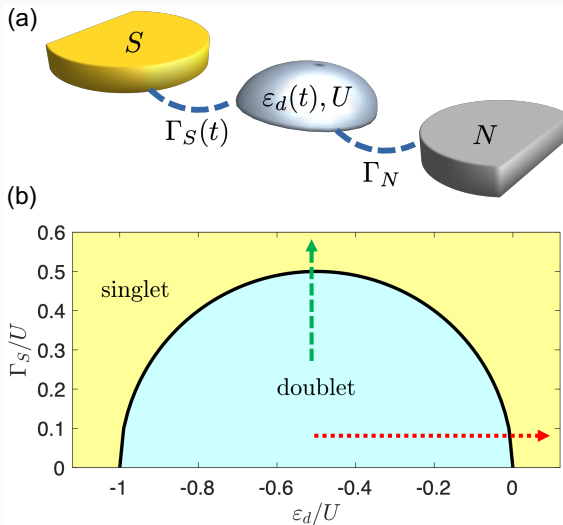
Eigen-states of this problem are represented by:

$|\uparrow\rangle$ **and** $|\downarrow\rangle$ \Leftarrow **doublet states (spin $\frac{1}{2}$)**

$\left. \begin{array}{l} u |0\rangle - v |\uparrow\downarrow\rangle \\ v |0\rangle + u |\uparrow\downarrow\rangle \end{array} \right\}$ \Leftarrow **singlet states (spin 0)**

Upon varying the ratio ϵ_d/U_d or Γ_S/U_d the doublet-singlet **transition** can be induced between these ground states.

QUENCH ACROSS STATIC QPT BOUNDARY



GENERAL OUTLINE

For $t < 0$ we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0 |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

GENERAL OUTLINE

For $t < 0$ we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0 |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Next, at time $t = 0$, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

GENERAL OUTLINE

For $t < 0$ we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0 |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Next, at time $t = 0$, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For $t > 0$ the Schrödinger eqn $i \frac{d}{dt} |\Psi(t)\rangle = \hat{H} |\Psi(t)\rangle$ implies:

$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

GENERAL OUTLINE

For $t < 0$ we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0 |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Next, at time $t = 0$, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For $t > 0$ the Schrödinger eqn $i \frac{d}{dt} |\Psi(t)\rangle = \hat{H} |\Psi(t)\rangle$ implies:

$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

Fidelity (similarity) of these states is:

$$\langle \Psi_0 | \Psi(t) \rangle = \langle \Psi_0 | e^{-it\hat{H}} | \Psi_0 \rangle$$

GENERAL OUTLINE

For $t < 0$ we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0 |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Next, at time $t = 0$, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For $t > 0$ the Schrödinger eqn $i \frac{d}{dt} |\Psi(t)\rangle = \hat{H} |\Psi(t)\rangle$ implies:

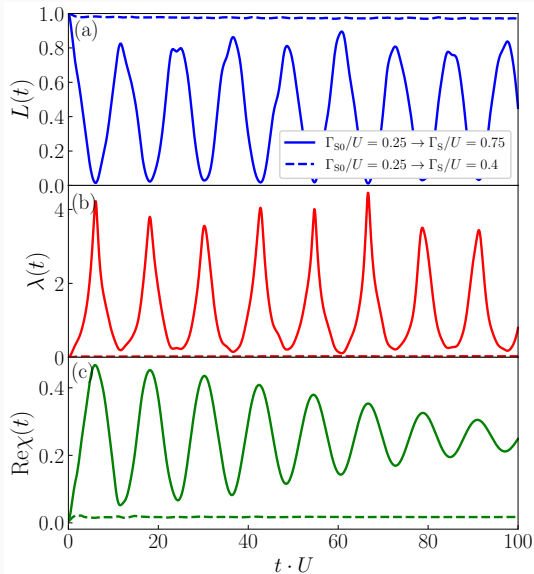
$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

Fidelity (similarity) of these states is:

$$\langle \Psi_0 | \Psi(t) \rangle = \langle \Psi_0 | e^{-it\hat{H}} | \Psi_0 \rangle$$

Loschmidt amplitude

t NRG RESULTS: ABRUPT CHANGE OF Γ_S



$$\varepsilon_d = -U/2$$

$$\Gamma_N = U/100$$

Triplet blockade

in junctions with two quantum dots

Theory of Andreev blockade in a double quantum dot with a superconducting lead

David Pekker, Po Zhang and Sergey M. Frolov

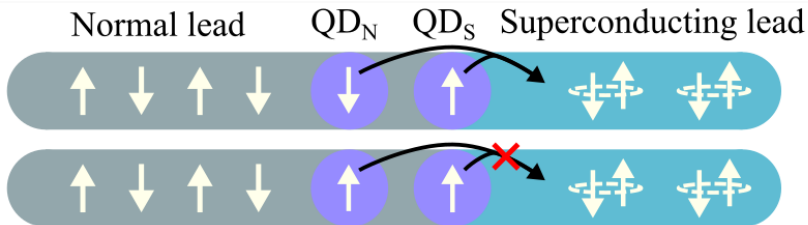
Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, 15260

ANDREEV BLOCKADE: CONCEPT

Theory of Andreev blockade in a double quantum dot with a superconducting lead

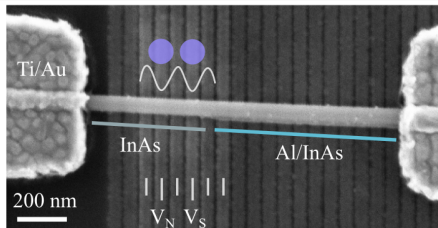
David Pekker, Po Zhang and Sergey M. Frolov

Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, 15260

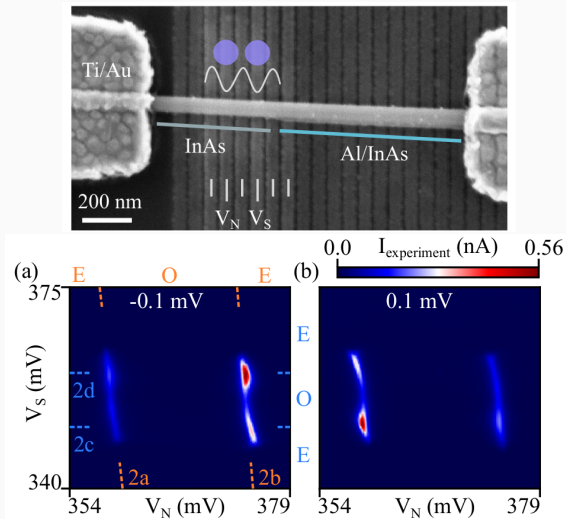


Superconducting proximity effect would be blocked by triplet configuration of the quantum dots (Andreev current forbidden).

ANDREEV BLOCKADE: REALIZATION



ANDREEV BLOCKADE: REALIZATION







P. Zhang, H. Wu, J. Chen, S.A. Khan, P. Krogstrup, D. Pekker, and S.M. Frolov, *Phys. Rev. Lett.* **128**, 046801 (2022).

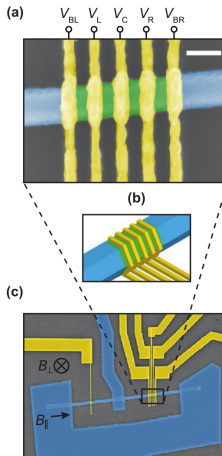
BLOCKADE IN JOSEPHSON JUNCTION

PHYSICAL REVIEW B **102**, 220505(R) (2020)

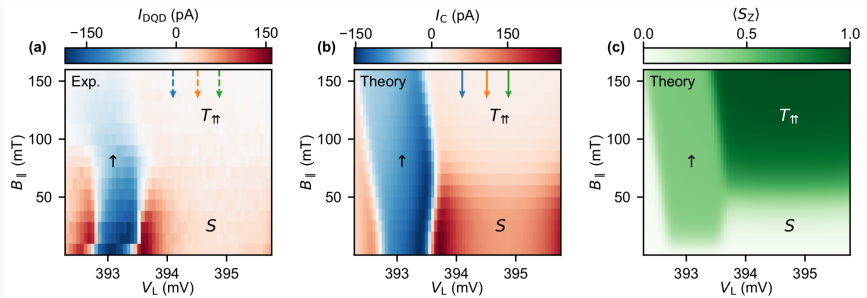
Rapid Communications

Triplet-blockaded Josephson supercurrent in double quantum dots

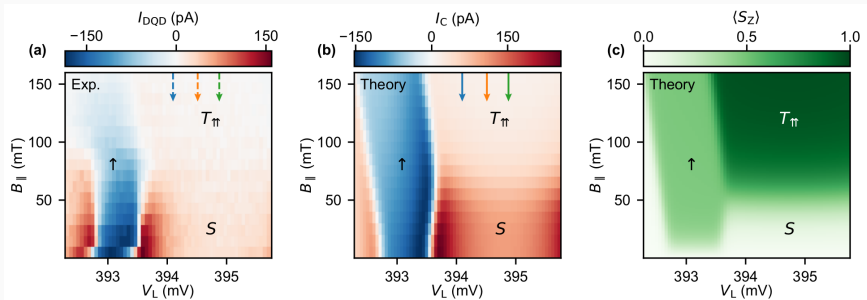
Daniël Bouman ¹, Ruben J. J. van Gulik,¹ Gorm Steffensen,² Dávid Pataki ³, Péter Boross,⁴ Peter Krogstrup,² Jesper Nygård ², Jens Paaske,² András Pályi,³ and Attila Geresdi ^{1,5,*}



BLOCKADE IN JOSEPHSON JUNCTION



BLOCKADE IN JOSEPHSON JUNCTION



Experimental observation:

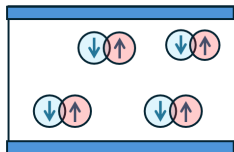
“magnetic field dependence of the supercurrent amplitude in the even occupied state reveals the presence of supercurrent blockade in the spin-triplet ground state”

Andreev blockade

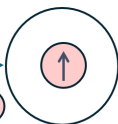
[dynamical realizations]

DYNAMICAL ANDREEV BLOCKADE

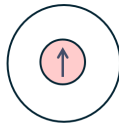
Superconducting lead



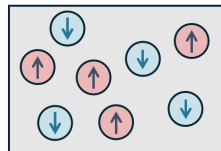
QD₁



QD₂



Normal lead

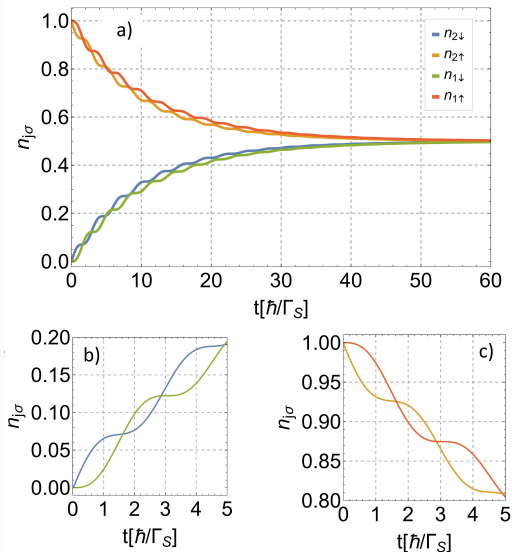


Superconducting proximity effect is blocked:

⇒ **when both quantum dots are singly occupied**

⇒ **by the same spin (for example \uparrow) electrons**

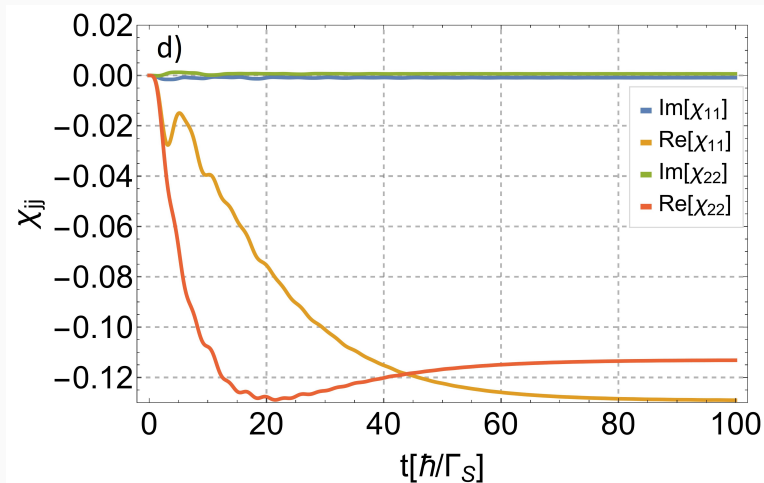
TRANSIENT BLOCKADE



Occupancy of the quantum dots initially occupied by \uparrow electrons.

TRANSIENT BLOCKADE

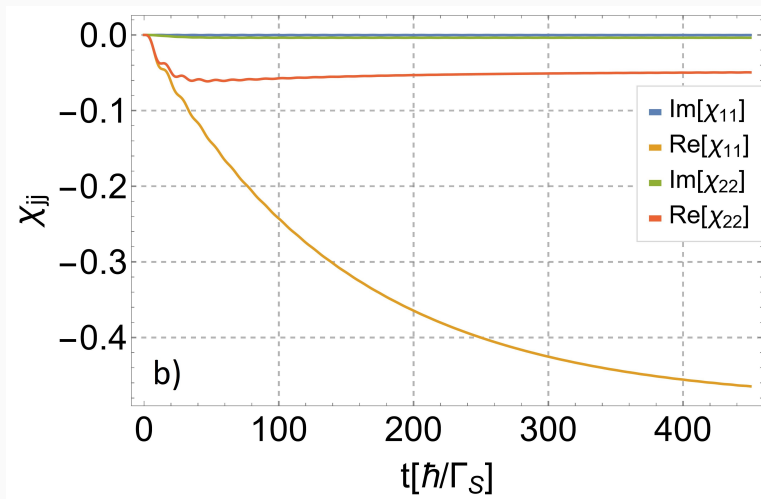
Results for $\chi_{jj}(t) \equiv \langle \hat{a}_{j\downarrow} \hat{a}_{j\uparrow} \rangle$ in the strong inter-dot coupling $V_{12} = \Gamma_S$.



Pairing in the quantum dots initially occupied by \uparrow electrons

TRANSIENT BLOCKADE

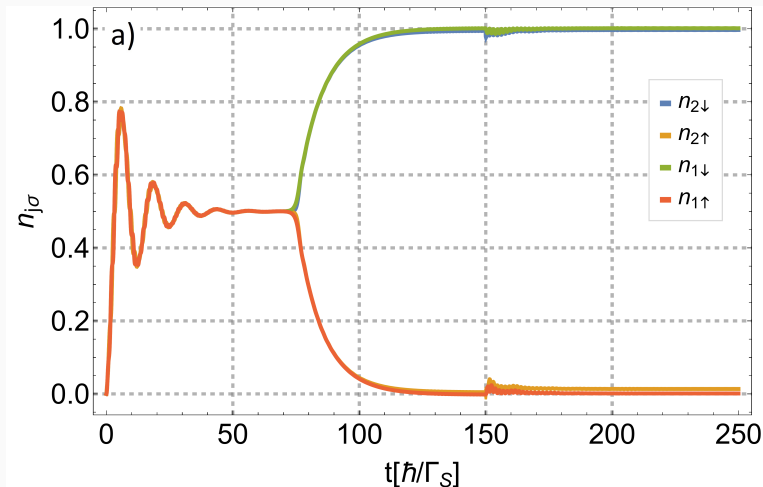
Results for $\chi_{jj}(t) \equiv \langle \hat{d}_{j\downarrow} \hat{d}_{j\uparrow} \rangle$ in the weak inter-dot coupling $V_{12} = 0.1\Gamma_S$.



Pairing in the quantum dots initially occupied by \uparrow electrons

ZEEMAN INDUCED BLOCKADE

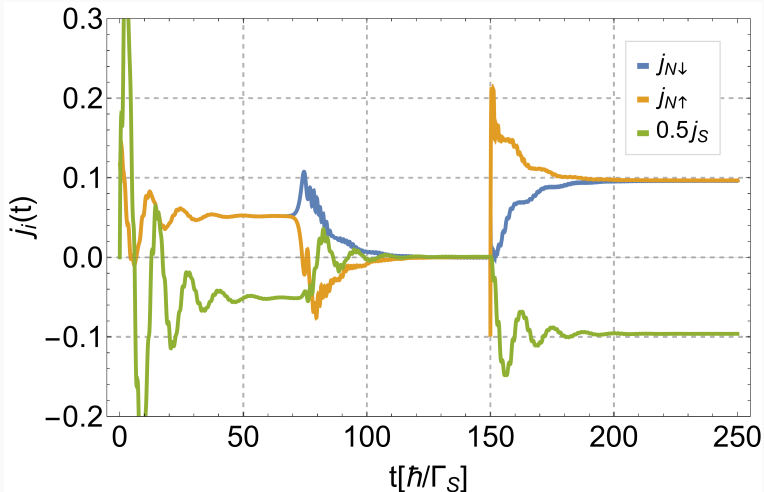
Magnetic field $B = 10\Gamma_S/\mu_B$ is switched on at $t = 75\Gamma_S/\hbar$



Occupancy of the initially empty quantum dots.

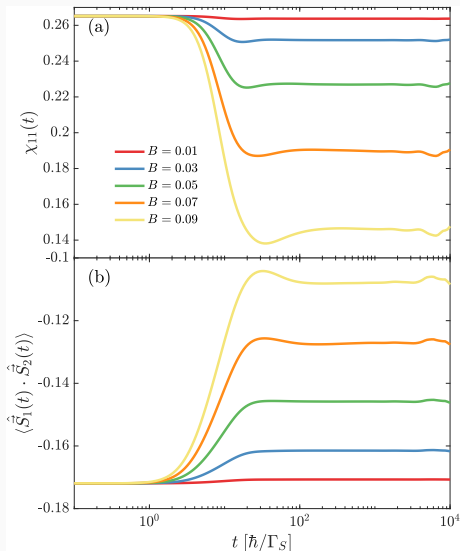
NON-EQUILIBRIUM CHARGE TRANSPORT

Magnetic field $B = 10\Gamma_S/\mu_B$ is switched on at $t = 75\Gamma_S/\hbar$ and bias voltage is strongly amplified at $t = 150\Gamma_S/\hbar$



CORRELATED SYSTEM

Results obtained by time-dependent NRG calculations



Model parameters:

$$U_1 = 0.9\Gamma_S, \quad U_2 = 0,$$

$$\varepsilon_1 = -U_1/2, \quad \varepsilon_2 = 0,$$

$$V_{12} = 0.225\Gamma_S$$

SUMMARY

**By attaching the quantum impurity to bulk superconductor
(or when its energy level / coupling strength is varied):**

- **Rabi-type oscillations are induced** (due to particle-hole mixing)
- **leading to buildup (re-arrangement) of the in-gap states**
- **dynamical phase transition can occur** (changeover of ground state)

SUMMARY

**By attaching the quantum impurity to bulk superconductor
(or when its energy level / coupling strength is varied):**

- **Rabi-type oscillations are induced** (due to particle-hole mixing)
- **leading to buildup (re-arrangement) of the in-gap states**
- **dynamical phase transition can occur** (changeover of ground state)

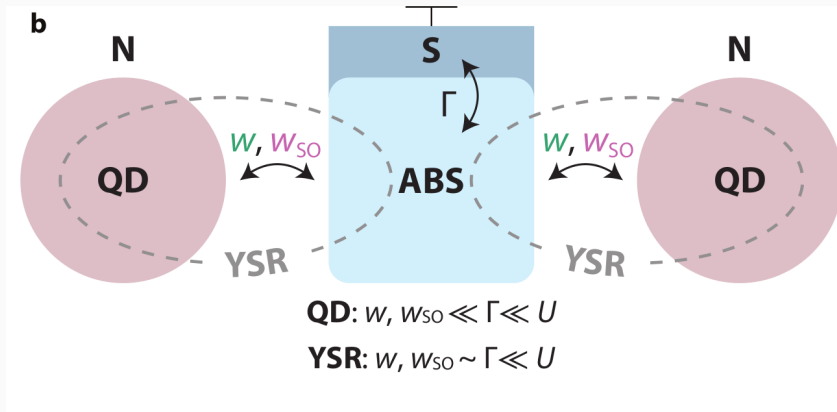
These phenomena could be detected in the charge transport measurements, using time-resolved Andreev spectroscopy.

Outlook

[triplet conf. in topological superconductors]

MINIMAL KITAEV CHAIN

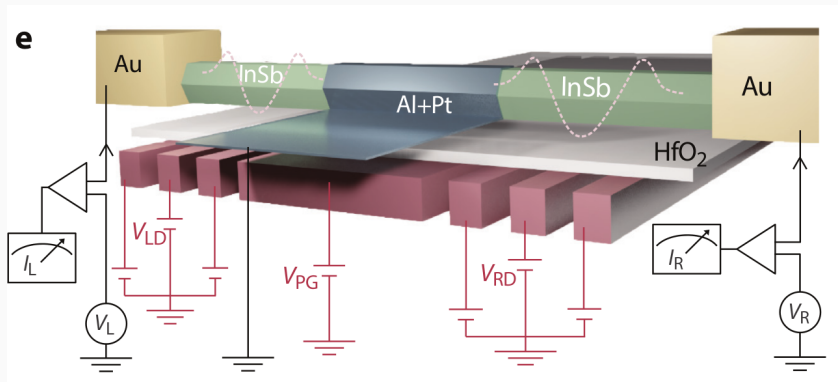
Effective triplet pairing can be realized using two quantum dots interconnected by superconductor (**Poor Man's Majorana states**)



T. Dvir, ... & L.P. Kouwenhoven, Nature 614, 445 (2023).

MINIMAL KITAEV CHAIN

Two spin-polarized quantum dots in an InSb nanowire strongly coupled by elastic co-tunneling and crossed Andreev reflection



T. Dvir, ... & L.P. Kouwenhoven, Nature 614, 445 (2023).