Dynamical effects in quantum dots coupled to superconductors

Tadeusz Domański

M. Curie-Skłodowska University Lublin, Poland



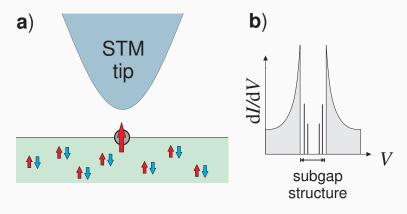


PRELIMINARIES

Superconducting nanostructures

NANOSTRUCTURES WITH SUPERCONDUCTOR(S)

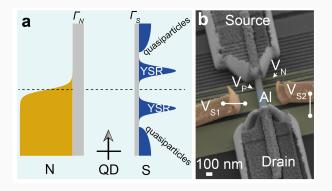
Quantum impurity on a surface of superconductor + STM tip



STM is a tool to probe the spectra of impurities

NANOSTRUCTURES WITH SUPERCONDUCTOR(S)

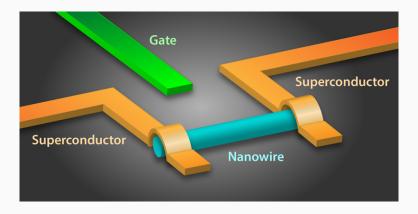
normal metal (N) - quantum dot (QD) - superconductor (S)



Tunneling by the electron-to-hole (Andreev) scattering.

NANOSTRUCTURES WITH SUPERCONDUCTOR(S)

superconductor (S) - quantum dot (QD) - superconductor (S)



Tunneling of Cooper pairs via bound states in Josephson junction.

SUPERCONDUCTING PROXIMITY EFFECT

Coupling of the localized (QD) to itinerant (SC) electrons induces:

⇒ on-dot pairing

SUPERCONDUCTING PROXIMITY EFFECT

- Coupling of the localized (QD) to itinerant (SC) electrons induces:
- ⇒ on-dot pairing
- manifested spectroscopically by:
- ⇒ bound states inside the gap of SC

SUPERCONDUCTING PROXIMITY EFFECT

- Coupling of the localized (QD) to itinerant (SC) electrons induces:
- ⇒ on-dot pairing
- manifested spectroscopically by:
- ⇒ bound states inside the gap of SC
- originating from:
- ⇒ leakage of Cooper pairs on QD (Andreev)
- ⇒ exchange int. of QD with SC (Yu-Shiba-Rusinov)

IMPURITY + CONVENTIONAL SUPERCONDUCTOR

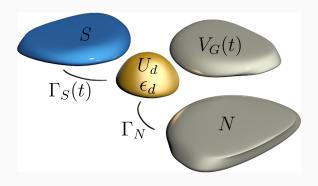
Characteristic time-scales

IMPURITY + CONVENTIONAL SUPERCONDUCTOR

Characteristic time-scales

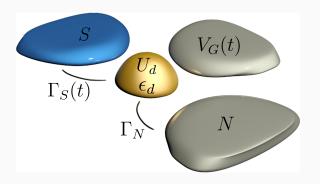
[emergence of in-gap states]

TIME-RESOLVED BOUND STATES



Protocols of non-equilibrium conditions:

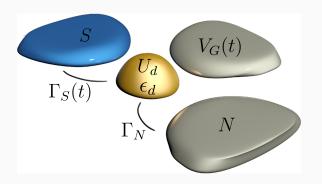
TIME-RESOLVED BOUND STATES



Protocols of non-equilibrium conditions:

 \Rightarrow variation of the coupling Γ_S to superconductor

TIME-RESOLVED BOUND STATES

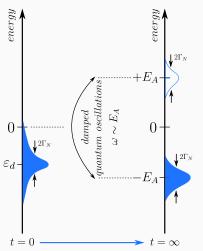


Protocols of non-equilibrium conditions:

- \Rightarrow variation of the coupling Γ_S to superconductor
- \Rightarrow abrupt change of the gate potential V_G

BUILDUP OF IN-GAP STATES

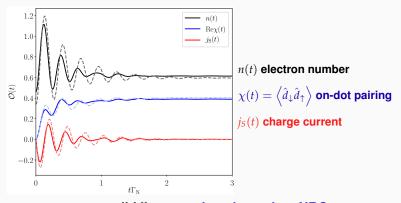
Emergence of Andreev states due to the sudden coupling $0 o \Gamma_S$



K. Wrześniewski, B. Baran, R. Taranko, T. Domański & I. Weymann, PRB 103, 155420 (2021).

BUILDUP OF IN-GAP STATES

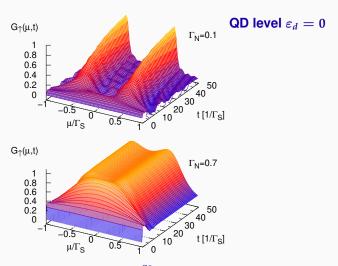
Time-dependent observables driven by the quantum quench $0 ightarrow \Gamma_S$



solid lines - time dependent NRG dashed lines - Hartree-Fock-Bogolubov

K. Wrześniewski, B. Baran, R. Taranko, T. Domański & I. Weymann, PRB 103, 155420 (2021).

TIME-DEPENDENT TUNNELING CONDUCTANCE



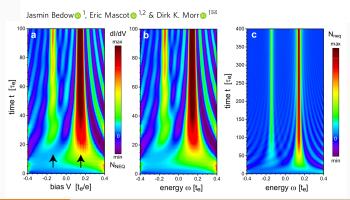
Subgap tunneling conductance $G_{\sigma}=\frac{\partial I_{\sigma}}{\partial t}$ vs time (t) and voltage (μ)

BOUND STATES OF CLASSICAL IMPURITY

communications physics

ARTICLE
https://doi.org/10.1038/s42005-022-01050-7
OPEN

Emergence and manipulation of non-equilibrium Yu-Shiba-Rusinov states



Correlation effects

Correlation effects

[Experts: T. Novotný, M. Žonda, V. Pokorny, P. Zalom]

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can described by

$$\hat{H}_{QD} = \sum \epsilon_d \; \hat{d}_{\sigma}^{\dagger} \; \hat{d}_{\sigma} \; + \; U_d \; \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_S \; \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow}^{\dagger} + \text{h.c.} \right)$$

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can described by

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \; \hat{d}^{\dagger}_{\sigma} \; \hat{d}_{\sigma} \; + \; U_d \; \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_S \; \hat{d}^{\dagger}_{\uparrow} \hat{d}^{\dagger}_{\downarrow} + \text{h.c.} \right)$$

Eigen-states of this problem are represented by:

$$\begin{array}{lll} |\!\!\uparrow\rangle & \text{and} & |\!\!\downarrow\rangle & \Leftarrow & \text{doublet states (spin $\frac{1}{2}$)} \\ u \, |0\rangle - v \, |\!\!\uparrow\downarrow\rangle & \\ v \, |0\rangle + u \, |\!\!\uparrow\downarrow\rangle & \Leftrightarrow & \text{singlet states (spin 0)} \end{array}$$

SINGLY OCCUPIED VS BCS-TYPE CONFIGURATIONS

Quantum dot proximitized to superconductor can described by

$$\hat{H}_{QD} = \sum_{\sigma} \epsilon_d \; \hat{d}_{\sigma}^{\dagger} \; \hat{d}_{\sigma} \; + \; U_d \; \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} - \left(\Gamma_S \; \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow}^{\dagger} + \text{h.c.} \right)$$

Eigen-states of this problem are represented by:

$$\begin{array}{lll} |\!\!\uparrow\rangle & \text{and} & |\!\!\downarrow\rangle & \Leftarrow & \text{doublet states (spin $\frac{1}{2}$)} \\ u \, |0\rangle - v \, |\!\!\uparrow\downarrow\rangle \\ v \, |0\rangle + u \, |\!\!\uparrow\downarrow\rangle \end{array} \right\} & \Leftarrow & \text{singlet states (spin 0)} \end{array}$$

Upon varrying the ratio ε_d/U_d or Γ_S/U_d the doublet-singlet transition can be induced between these ground states.

Dynamical phase transition

Dynamical phase transition

[transition in time-domain]

For t < 0 we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0\ket{\Psi_0}=E_0\ket{\Psi_0}$$

For t < 0 we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0\ket{\Psi_0}=E_0\ket{\Psi_0}$$

Next, at time t=0, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For t < 0 we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0\ket{\Psi_0}=E_0\ket{\Psi_0}$$

Next, at time t = 0, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For t>0 the Schrödinger eqn $i\frac{d}{dt}\ket{\Psi(t)}=\hat{H}\ket{\Psi(t)}$ implies:

$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

For t < 0 we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0\ket{\Psi_0}=E_0\ket{\Psi_0}$$

Next, at time t = 0, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For t>0 the Schrödinger eqn $i\frac{d}{dt}\ket{\Psi(t)}=\hat{H}\ket{\Psi(t)}$ implies:

$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

Fidelity (similarity) of these states is:

$$\langle \Psi_0 | \Psi(t)
angle = \left\langle \Psi_0 | e^{-it\hat{H}} | \Psi_0
ight
angle$$

For t < 0 we assume the system \hat{H}_0 to be in its ground state:

$$\hat{H}_0\ket{\Psi_0}=E_0\ket{\Psi_0}$$

Next, at time t = 0, we impose an abrupt change (quench):

$$\hat{H}_0 \longrightarrow \hat{H}$$

For t>0 the Schrödinger eqn $i\frac{d}{dt}\ket{\Psi(t)}=\hat{H}\ket{\Psi(t)}$ implies:

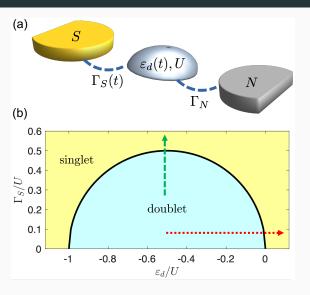
$$|\Psi(t)\rangle = e^{-it\hat{H}} |\Psi_0\rangle$$

Fidelity (similarity) of these states is:

$$\langle \Psi_0 | \Psi(t) \rangle = \left\langle \Psi_0 | e^{-it\hat{H}} | \Psi_0 \right\rangle$$

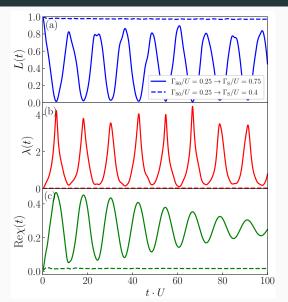
Loschmidt amplitude

QUENCH ACROSS STATIC QPT BOUNDARY



K. Wrześniewski, I. Weymann, N. Sedlmayr & T. Domański, Phys. Rev. B 105, 094514 (2022).

tnrg results: Abrupt change of Γ_S



$$\varepsilon_d = -U/2$$
 $\Gamma_N = U/100$

K. Wrześniewski, I. Weymann, N. Sedlmayr & T. Domański, Phys. Rev. B 105, 094514 (2022).

Quantum impurity/dot embedded into bulk superconductor:

induces the Rabi-type oscillations

(due to particle-hole mixing)

Quantum impurity/dot embedded into bulk superconductor:

- induces the Rabi-type oscillations
 (due to particle-hole mixing)
- leading to the buildup (re-arrangement) of in-gap states

Quantum impurity/dot embedded into bulk superconductor:

- induces the Rabi-type oscillations
 (due to particle-hole mixing)
- leading to the buildup (re-arrangement) of in-gap states
- which can undergo dynamical transitions
 - (qualitative changeover of the ground states)

Quantum impurity/dot embedded into bulk superconductor:

- induces the Rabi-type oscillations
 (due to particle-hole mixing)
- leading to the buildup (re-arrangement) of in-gap states
- which can undergo dynamical transitions
 - (qualitative changeover of the ground states)

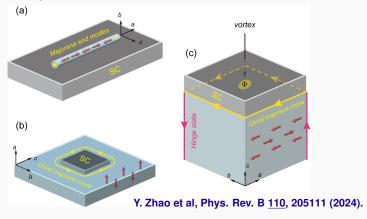
These phenomena could be detected by the charge transport and evidenced in time-resolved Andreev/Josephsnon conductance.

Part 2. topological superconductors

(Majorana-type quasiparticles)

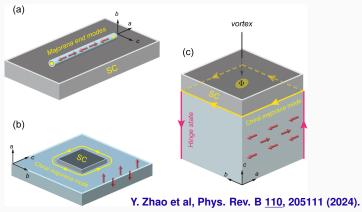
MOTIVATION

The boundary modes (localized, chiral or Hinge states) of topological superconductors realized in different dimensions



MOTIVATION

The boundary modes (localized, chiral or Hinge states) of topological superconductors realized in different dimensions



can be detected, using the charge tunneling spectroscopies (with attachment of external electrodes) in nonequilibrium conditions.

HYBRID TOPOLOGICAL STRUCTURES

Topological superconductors can hybridize with other (topologically trivial) objects:

 \Rightarrow through some interface,

 \Rightarrow forming boundary modes.

HYBRID TOPOLOGICAL STRUCTURES

Topological superconductors can hybridize with other (topologically trivial) objects:

 \Rightarrow through some interface,

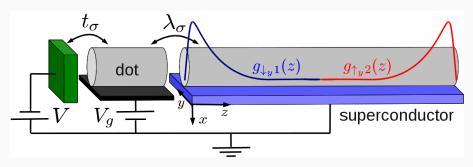
 \Rightarrow forming boundary modes.

The simplest case:

⇒ single-level impurity + Majorana mode(s).

LEAKAGE OF MAJORANA MODE ON QUANTUM DOT

Hybrid structure: quantum dot + topological superconductor

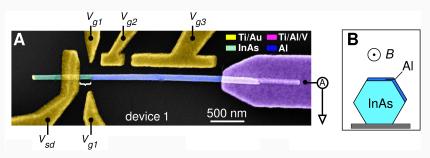


Idea: Majorana mode is partly transferred onto quantum dot where it can be detected by tunneling spectroscopy

M. Leijnse and K. Flensberg, Phys. Rev. B 84, 140501(R) (2011).

FIRST EXPERIMENTAL REALIZATION

Setup: Epitaxial AI shell (blue) grown on two facets of the hexagonal InAs core (cyan), with a thickness of \sim 10 nm.



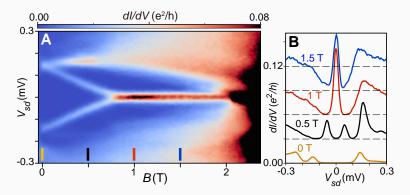
Data: Transport measurements have been collected, varying the magnetic field oriented parallelly to the nanowire.

M.T. Deng et al, Science 354, 1557 (2016).

EVIDENCE FOR MAJORANA LEAKAGE

Panel (A): Tunneling spectrum for resonant dot-wire coupling obtained at $V_{bg}=-8.5$ V, $V_{g1}=22$ V, and $V_{g2}=V_{g3}=-10$ V.

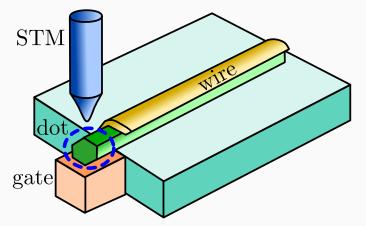
Panel (B): Differential conductance at various values of the magnetic field.



M.T. Deng et al, Science 354, 1557 (2016).

GATE-CONTROLLED BOUND STATES

Hybrid structure: trivial + topological segments of nanowire

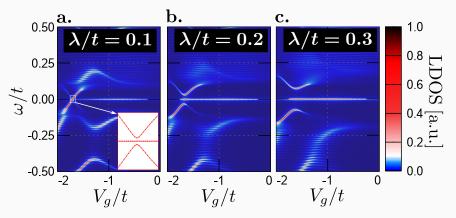


Issue: bound states of trivial segment attached to topological sc

A. Ptok, A. Kobiałka, T. Domański, Phys. Rev. B <u>96</u>, 195430 (2017).

GATE-CONTROLLED BOUND STATES

Hybrid structure: trivial + topological segments of nanowire



Variation the trivial (Andreev) & topological (Majorana) states vs the gate potential V_g for several spin-orbit couplings λ .

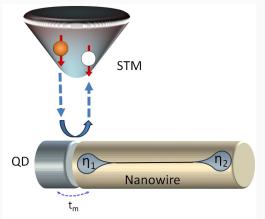
A. Ptok, A. Kobiałka, T. Domański, Phys. Rev. B 96, 195430 (2017).

/ induced by Coulomb repulsion /

What about correlations?

CORRELATIONS VS LEAKAGE

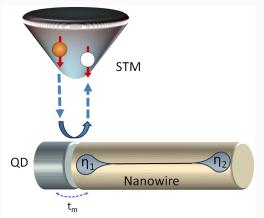
Hybrid structure: Anderson impurity + topological superconductor



Question: Does the Coulomb repulsion affect the Majorana mode(s) leakage?

CORRELATIONS VS LEAKAGE

Hybrid structure: Anderson impurity + topological superconductor



Question: Does the Coulomb repulsion affect the Majorana mode(s) leakage? Is there any competition?

J. Barański, M. Barańska, T. Zienkiewicz & T. Domański, J. Phys.: Cond. Matt. 37, 055302 (2025).

For microscopic considerations we used the Anderson-type model

$$\hat{H} = \hat{H}_{QD} + \lambda (\hat{d}_{\downarrow}^{\dagger}\hat{\eta}_{1} + \hat{\eta}_{1}\hat{d}_{\downarrow}) + i\epsilon_{m}\hat{\eta}_{1}\hat{\eta}_{2}$$

For microscopic considerations we used the Anderson-type model

$$\hat{H} = \hat{H}_{QD} + \lambda (\hat{d}_{\downarrow}^{\dagger}\hat{\eta}_{1} + \hat{\eta}_{1}\hat{d}_{\downarrow}) + i\epsilon_{m}\hat{\eta}_{1}\hat{\eta}_{2}$$

where the correlated quantum dot is described by

$$\hat{H}_{\mathrm{QD}} = \sum_{\sigma} arepsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{\uparrow} \hat{n}_{\downarrow}$$

For microscopic considerations we used the Anderson-type model

$$\hat{H}=\hat{H}_{QD}+\lambda(\hat{d}_{\downarrow}^{\dagger}\hat{\eta}_{1}+\hat{\eta}_{1}\hat{d}_{\downarrow})+i\epsilon_{m}\hat{\eta}_{1}\hat{\eta}_{2}$$

where the correlated quantum dot is described by

$$\hat{H}_{QD} = \sum_{\sigma} arepsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{\uparrow} \hat{n}_{\downarrow}$$

recasting the Majorana operators in terms of conventional fermions

$$\hat{\eta}_1 = rac{1}{\sqrt{2}}(\hat{f}^\dagger + \hat{f})$$

For microscopic considerations we used the Anderson-type model

$$\hat{H} = \hat{H}_{QD} + \lambda (\hat{d}_{ot}^{\dagger}\hat{\eta}_1 + \hat{\eta}_1\hat{d}_{ot}) + i\epsilon_m\hat{\eta}_1\hat{\eta}_2$$

where the correlated quantum dot is described by

$$\hat{H}_{ ext{QD}} = \sum_{\sigma} arepsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{\uparrow} \hat{n}_{\downarrow}$$

recasting the Majorana operators in terms of conventional fermions

$$\hat{\eta}_1=rac{1}{\sqrt{2}}(\hat{f}^\dagger+\hat{f})$$
 and $\hat{\eta}_2=rac{i}{\sqrt{2}}(\hat{f}^\dagger-\hat{f})$

For microscopic considerations we used the Anderson-type model

$$\hat{H} = \hat{H}_{QD} + \lambda (\hat{d}_{\downarrow}^{\dagger}\hat{\eta}_{1} + \hat{\eta}_{1}\hat{d}_{\downarrow}) + i\epsilon_{m}\hat{\eta}_{1}\hat{\eta}_{2}$$

where the correlated quantum dot is described by

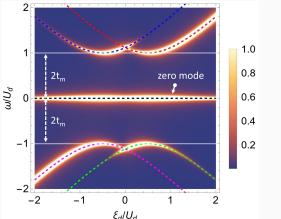
$$\hat{H}_{ ext{QD}} = \sum_{\sigma} arepsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U_d \hat{n}_{\uparrow} \hat{n}_{\downarrow}$$

recasting the Majorana operators in terms of conventional fermions

$$\hat{\eta}_1=rac{1}{\sqrt{2}}(\hat{f}^\dagger+\hat{f})$$
 and $\hat{\eta}_2=rac{i}{\sqrt{2}}(\hat{f}^\dagger-\hat{f})$

Quasiparticle states of the quantum dot can be determined analytically.

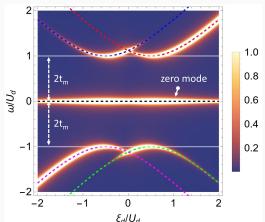
Hybrid structure: Anderson impurity + topological nanowire



J. Barański et al, (2025).

Spectrum of spin-√electrons which are directly coupled to the Majorana mode.

Hybrid structure: Anderson impurity + topological nanowire

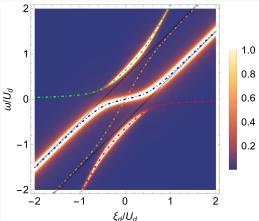


J. Barański et al, (2025).

Spectrum of spin- \downarrow electrons which are directly coupled to the Majorana mode. Zero-energy mode appears near ϵ_d and $\epsilon_d + U_d$.

Notation: $\xi_d \equiv \epsilon_d + U_d/2$.

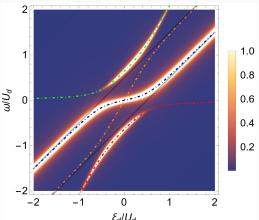
Hybrid structure: Anderson impurity + topological nanowire



T. Barański et al, (2025).

Spectrum of spin- \uparrow electrons which are not directly coupled to the Majorana mode.

Hybrid structure: Anderson impurity + topological nanowire



T. Barański et al, (2025).

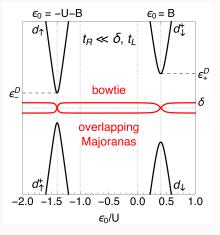
Spectrum of spin- \uparrow electrons which are not directly coupled to the Majorana mode. Majorana features are missing.

/ overlapping Majorana modes /

Short topological nanowire

OVERLAPPING MAJORANA MODES, $\epsilon_M \neq 0$

Hybrid structure: quantum impurity + short topological nanowire

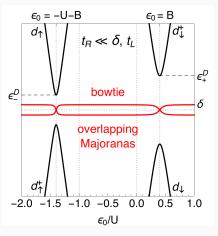


E. Prada et al, PRB <u>96</u>, 085418 (2017).

Quasiparticle spectrum of the quantum dot obtained for $\epsilon_M \neq 0$.

OVERLAPPING MAJORANA MODES, $\epsilon_M eq 0$

Hybrid structure: quantum impurity + short topological nanowire



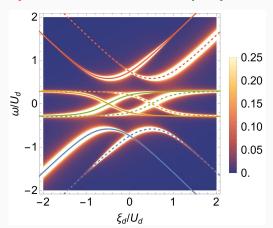
E. Prada et al, PRB <u>96</u>, 085418 (2017).

Quasiparticle spectrum of the quantum dot obtained for $\epsilon_M \neq 0$.

Notice: bowtie features near the crossing points.

OVERLAPPING MAJORANA MODES, $\epsilon_M eq 0$

Hybrid structure: Anderson impurity + short topological nanowire

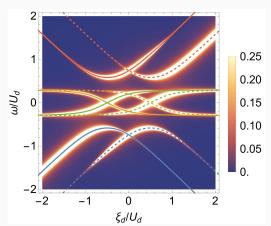


J. Barański et al, (2025).

Quasiparticle spectrum of spin- \downarrow electrons obtained for $\epsilon_M
eq 0$.

OVERLAPPING MAJORANA MODES, $\epsilon_M \neq 0$

Hybrid structure: Anderson impurity + short topological nanowire



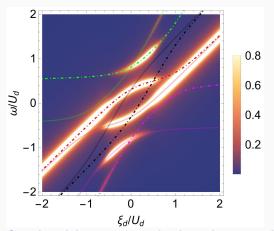
J. Barański et al, (2025).

Quasiparticle spectrum of spin- \downarrow electrons obtained for $\epsilon_M
eq 0$.

Appearance of <u>two bowtie features</u> inside the topological gap.

OVERLAPPING MAJORANA MODES, $\epsilon_M \neq 0$

Hybrid structure: Anderson impurity + short topological nanowire

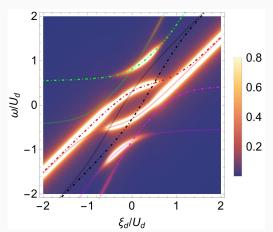


J. Barański et al, (2025).

Quasiparticle spectrum of spin- \uparrow electrons obtained for $\epsilon_M
eq 0$.

OVERLAPPING MAJORANA MODES, $\epsilon_M eq 0$

Hybrid structure: Anderson impurity + short topological nanowire



J. Barański et al, (2025).

Quasiparticle spectrum of spin- \uparrow electrons obtained for $\epsilon_M
eq 0$.

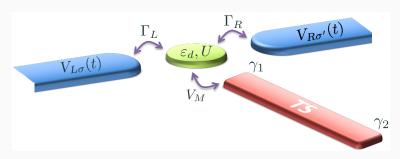
Majorana quasiparticles are completely absent.

Kondo vs Majorana

(means to distinguish them)

MAJORANA SIGNATURES IN AC-CONDUCTANCE

Quantum dot coupled to the topological nanowire under ac-voltage

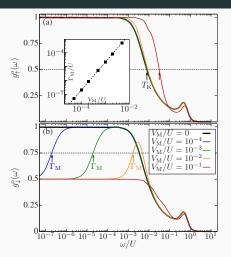


Question:

Can we resolve Majorana and Kondo states in ac-response?

K.P. Wójcik, T. Domański, I. Weymann, Phys. Rev. B 109, 075432 (2024).

CONDUCTANCE OF ac-DRIVEN JUNCTION



Spin-resolved conductance: Signatures of the Coulomb peak and the Kondo effect can be clearly distinguished at finite-frequencies.

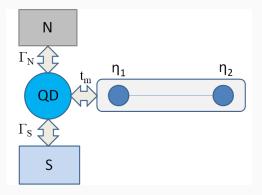
K.P. Wójcik, T. Domański, I. Weymann, Phys. Rev. B 109, 075432 (2024).

Time - resolved effects

(with Majorana modes)

TIME-RESOLVED LEAKAGE OF MAJORANA MODE

Hybrid structure: quantum dot attached to topological nanowire



Question:

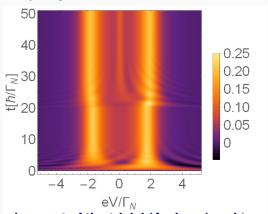
How much time does it take to transfer the Majorana mode on QD?

J. Barański, M. Barańska, T. Zienkiewicz, R. Taranko, T.Domański, PRB 103, 235416 (2021).

TIME-RESOLVED LEAKAGE OF MAJORANA MODE

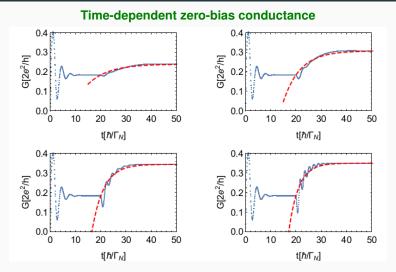
Transient effects:

- \Rightarrow at t = 0 QD is coupled to the external N and S electrodes,
- \Rightarrow at t = 20 topological nanowire is attached to N-QD-S setup.



Gradual development of the trivial (Andreev) and topological (Majorana) states manifested in the differential conductance.

TIME-RESOLVED LEAKAGE OF MAJORANA MODE



Majorana zero-bias feature establishes in about nanoseconds.

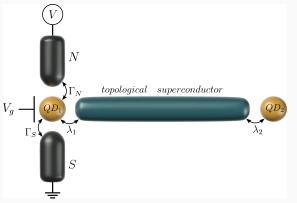
J. Barański, M. Barańska, T. Zienkiewicz, R. Taranko, T.Domański, PRB 103, 235416 (2021).

Distant cross-correlations

/ transmitted via Majorana modes /

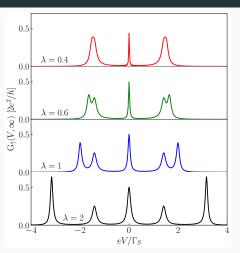
DYNAMICAL CROSS-CORRELATIONS

Two quantum dots interconnected via topological superconductor



Question: Is any nonlocal communication transmitted between ${\bf QD_1}$ and ${\bf QD_2}$ through the Majorana boundary modes ?

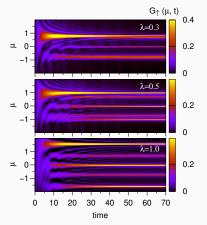
STEADY-LIMIT CONDUCTANCE



Differential conductance $G(V, t \to \infty)$ versus bias V for several couplings λ between QD_{1,2} and topological superconductor.

TIME-RESOLVED CONDUCTANCE

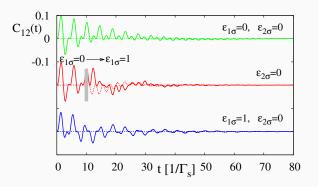
Time-dependent conductance of the biased N-QD₁-S junction



Signatures of the (trivial) molecular bound states and (topological) Majorana mode obtained for $\varepsilon_1=0,\,\varepsilon_2=2.$

NONLOCAL CROSS-CORRELATIONS

Evolution of the interdot electron pairing $C_{12}(t) = \left<\hat{d}_{1\downarrow}\hat{d}_{2\uparrow}\right>$

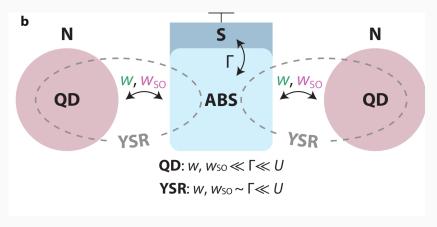


The nonlocal electron pairing persists only over a short transient time-scale. It could be detected by crossed Andreev refelections.

Recent activities

MINIMAL KITAEV CHAIN

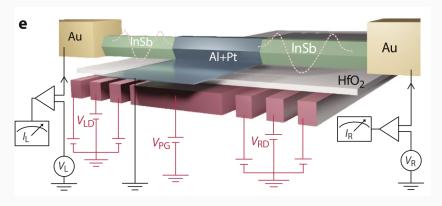
Effective triplet pairing can be realized using two quantum dots interconnected by superconductor (Poor Man's Majorana states)



T. Dvir, ... & L.P. Kouwenhoven, Nature 614, 445 (2023).

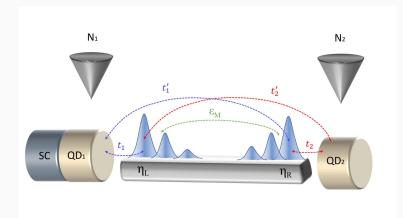
MINIMAL KITAEV CHAIN

Two spin-polarized quantum dots in an InSb nanowire strongly coupled by elastic co-tunneling and crossed Andreev reflection



T. Dvir, ... & L.P. Kouwenhoven, Nature 614, 445 (2023).

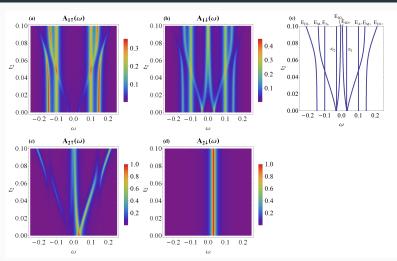
QUASIPARTICLE SPECTRUM OF QUANTUM DOTS



Issue: Molecular quasiparticle spectrum of the quantum dots connected via the overlapping Majorana modes

G. Górski, K.P. Wójcik, J. Barański, I. Weymann & T. Domański, Sci. Rep. 14, 13848 (2024).

QUASIPARTICLE SPECTRUM OF QUANTUM DOTS



Quasiparticles are shared by both quantum dots, however, appearing with different spectral weights.

G. Górski, K.P. Wójcik, J. Barański, I. Weymann & T. Domański, Sci. Rep. 14, 13848 (2024).

QUANTUM ENTANGLEMENT OF DOUBLE DOTS

Setup: Quantum dots interconnected via short topological nanowire

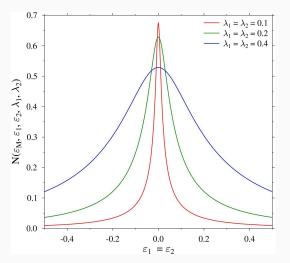


Issue: Quantum dots acquire entanglement via Majorana qps

C. Jasiukiewicz, A. Sinner, I. Weymann, T. Domański & L. Chotorlishvili, Phys. Rev. B <u>111</u>, 075415 (2025).

QUANTUM ENTANGLEMENT OF DOUBLE DOTS

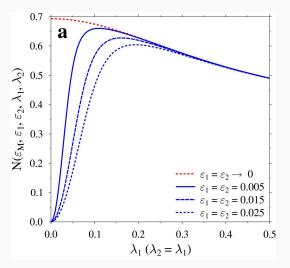
Setup: Quantum dots interconnected via short topological nanowire



Logarithmic negativity versus the energy levels QD's obtained for $\varepsilon_M \neq 0$.

QUANTUM ENTANGLEMENT OF DOUBLE DOTS

Setup: Quantum dots interconnected via short topological nanowire



Logarithmic negativity versus the coupling $\lambda_1 = \lambda_2$ obtained for $\varepsilon_M \neq 0$.

SUMMARY (PART 2)

Quantum dots attached to topological superconductors:

- allow for leakage of Majorana mode(s) (no competition with on-site repulsion),
- ⇒ become nonlocally cross-correlated (but only under non-equilibrium conditions),
- ⇒ acquire quantum entanglement (solely via short topological superconductors).