

**TRANSPORT COEFFICIENTS FOR A SLOW
FERMI-PARTICLE MOTION**

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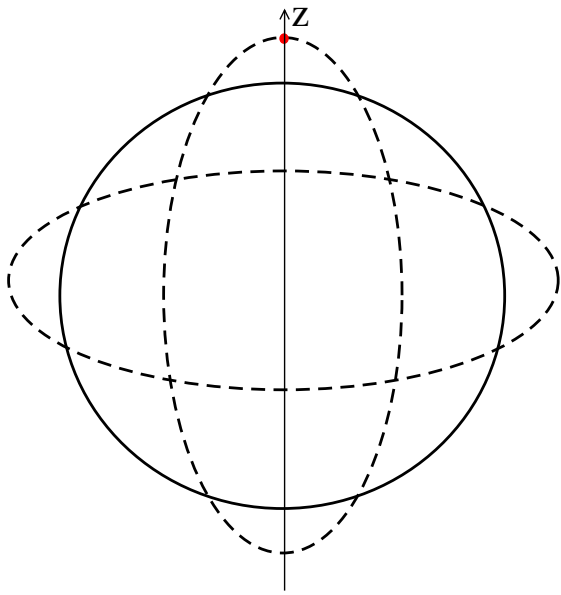
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Plan of talk

1. **INTRODUCTION.** Quantum-classical correspondence in the nuclear collective dynamics
2. **RESPONSE THEORY AND TRANSPORT COEFFICIENTS**
3. **SEMICLASSICAL FRICTION AND INERTIA**
4. **DISCUSSIONS OF THE RESULTS**
5. **CONCLUSIONS**



$$1) \quad \frac{dE}{dt} = \rho \bar{v} \oint \dot{n}^2 d\sigma$$

WALL FORMULA, Swiatecki et al. 1977,78

2) **ORDER-CHAOS, POINCARÉ SECTIONS
LYAPUNOV EXPONENTS**, Swiatecki,
Blocki et al. 1995-1999

3) **QUANTUM AND CLASSICAL RESULTS
FOR EXCITATION ENERGY**, Swiatecki,
Blocki, Skalski et al. 1995-2007;

Blocki, Yatsyshyn, Magner, 2010-11, **FOR 10-20 OSC. PERIODS:**

$$dE/dt = M\dot{\alpha}\ddot{\alpha} + \gamma\dot{\alpha}^2 \quad \Rightarrow \quad \langle dE/dt \rangle = \gamma\omega^2\alpha^2/2$$

4) **SEMICLASSICAL FRICTION AND INERTIA OF ONE-BODY
COLLECTIVE DYNAMICS**, Koonin, Randrup
Gutzwiller POT: Gzhebinsky, Magner, Fedotkin, 2007
SMOOTH FRICTION: LOCAL WALL FORMULA)

5) **QUANTUM AND SEMICLASSICAL FRICTION AND
INERTIA, NONLOCAL TRAJECTORY CORRECTIONS
WITHIN POT**

2. RESPONSE THEORY AND TRANSPORT COEFFICIENTS

COLLECTIVE RESPONSE to $V_{\text{ext}}(t) = \hat{F}\alpha_{\text{ext}}^\omega e^{-i\omega t}$

$$\delta\langle\hat{F}\rangle_\omega = -\chi_{\text{coll}}(\omega)\alpha_{\text{ext}}^\omega, \quad \hat{F}(\mathbf{r}) = \left(\frac{\partial V}{\partial \alpha}\right)_{\alpha=0} = -V_0 R \delta(r - R) Y_{L0}(\theta)$$

DYNAMICAL NUCLEAR SURFACE VIBRATIONS

$$\delta R(\theta, \alpha) \propto R\alpha(t)Y_{L0}(\theta), \quad \alpha(t) = \alpha_\omega e^{-i\omega t}$$

SELFCONSISTENCY CONDITION, Bohr & Mottelson (1975)

$$\delta\langle\hat{F}\rangle_\omega = \kappa \delta\alpha_\omega \quad \text{with the coupling constant } \kappa$$

$$\chi_{\text{coll}}(\omega) = \kappa \frac{\chi(\omega)}{\chi(\omega) + \kappa}, \quad \left[\delta\langle\hat{F}\rangle_\omega = -\chi(\omega) (\delta\alpha_\omega + \alpha_{\text{ext}}^\omega) \right]$$

STRENGTH FUNCTION $S(\omega) = \frac{1}{\pi} \text{Im}\chi_{\text{coll}}(\omega)$

TRANSPORT COEFFICIENTS, Hofmann et al. (1997,2008)

$$\frac{1}{\chi_{\text{coll}}(\omega)} = \frac{1}{\chi_{\alpha\alpha}^{\text{osc}}(\omega)} = \frac{1}{\chi_{\alpha\alpha}(\omega)} + \kappa = -M\omega^2 - i\gamma\omega + C$$

INTRINSIC TRANSPORT COEFFICIENTS in the limit $\omega \rightarrow 0$

$$\chi(\omega) = \chi(0) - i\gamma(0)\omega - M(0)\omega^2 + \dots, \quad \chi(0) = -\kappa - C(0)$$

$$C = [1 + C(0)/\chi(0)] C(0), \quad \gamma = [1 + C(0)/\chi(0)]^2 \gamma(0)$$

$$M = [1 + C(0)/\chi(0)]^2 [M(0) + \gamma^2(0)/\chi(0)]$$

INTRINSIC RESPONSE \Leftrightarrow GREEN'S FUNCTION

$$\chi(\omega) \propto \int_0^\infty d\varepsilon n(\varepsilon) \int dr_1 \int dr_2 \hat{F}(r_1) \hat{F}(r_2) \text{Im}G(r_1, r_2; \varepsilon) \\ \times [G^*(r_1, r_2; \varepsilon - \hbar\omega) + G(r_1, r_2; \varepsilon + \hbar\omega)], \quad G = \sum_i \frac{\psi_i(r_1) \bar{\psi}_i(r_2)}{\varepsilon - \varepsilon_i + i\epsilon}$$

$$\gamma(0) = (\partial\chi(\omega)/\partial\omega)_{\omega=0}, \quad M(0) = (1/2) (\partial^2\chi(\omega)/\partial\omega^2)_{\omega=0}$$

4. SEMICLASSICAL FRICTION AND INERTIA (sph. box)

$$\gamma = \int_0^\pi \sin\psi d\psi P_n(\cos\psi) \gamma(\psi), \quad \gamma(\psi) \propto (\partial^2 \text{Im}G / \partial r_1 \partial r_2)_{S,F}^2$$

$$M = \int_0^\pi \sin\psi d\psi P_n(\cos\psi) M(\psi), \quad M(\psi) \propto \int_0^{\epsilon_F} d\epsilon \left(\frac{\partial^2 \text{Im}G}{\partial r_1 \partial r_2} \frac{\partial^2}{\partial \epsilon^2} \frac{\partial^2 \text{Re}G}{\partial r_1 \partial r_2} \right)_S$$

GREEN'S FUNCTION (Gutzwiller, 1971) for $k_F R_0 \sim A^{1/3} \gg 1$:

$$G(r_1, r_2; \epsilon) = \sum_{CT} \mathcal{A}_{CT}(r_1, r_2; \epsilon) \exp \left[\frac{i}{\hbar} S_{CT}(r_1, r_2; \epsilon) - \frac{i\pi}{2} \mu_{CT} - \phi_d \right]$$

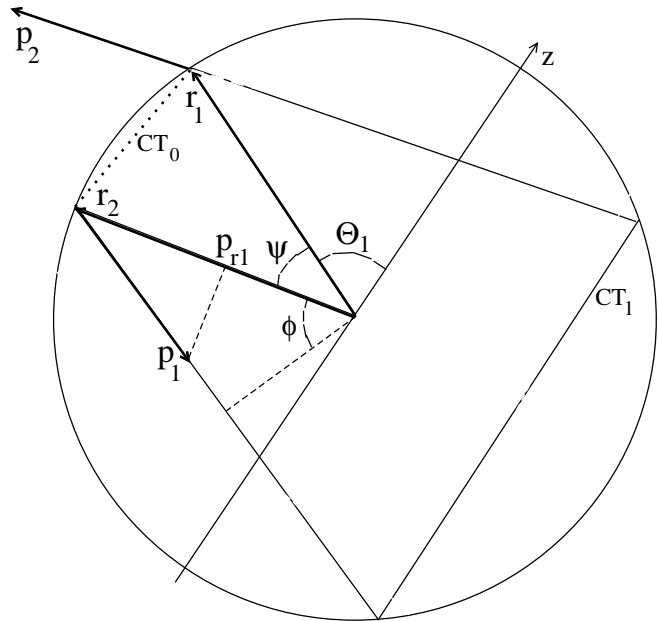


Fig. 4. The trajectories CT_0 and CT_1

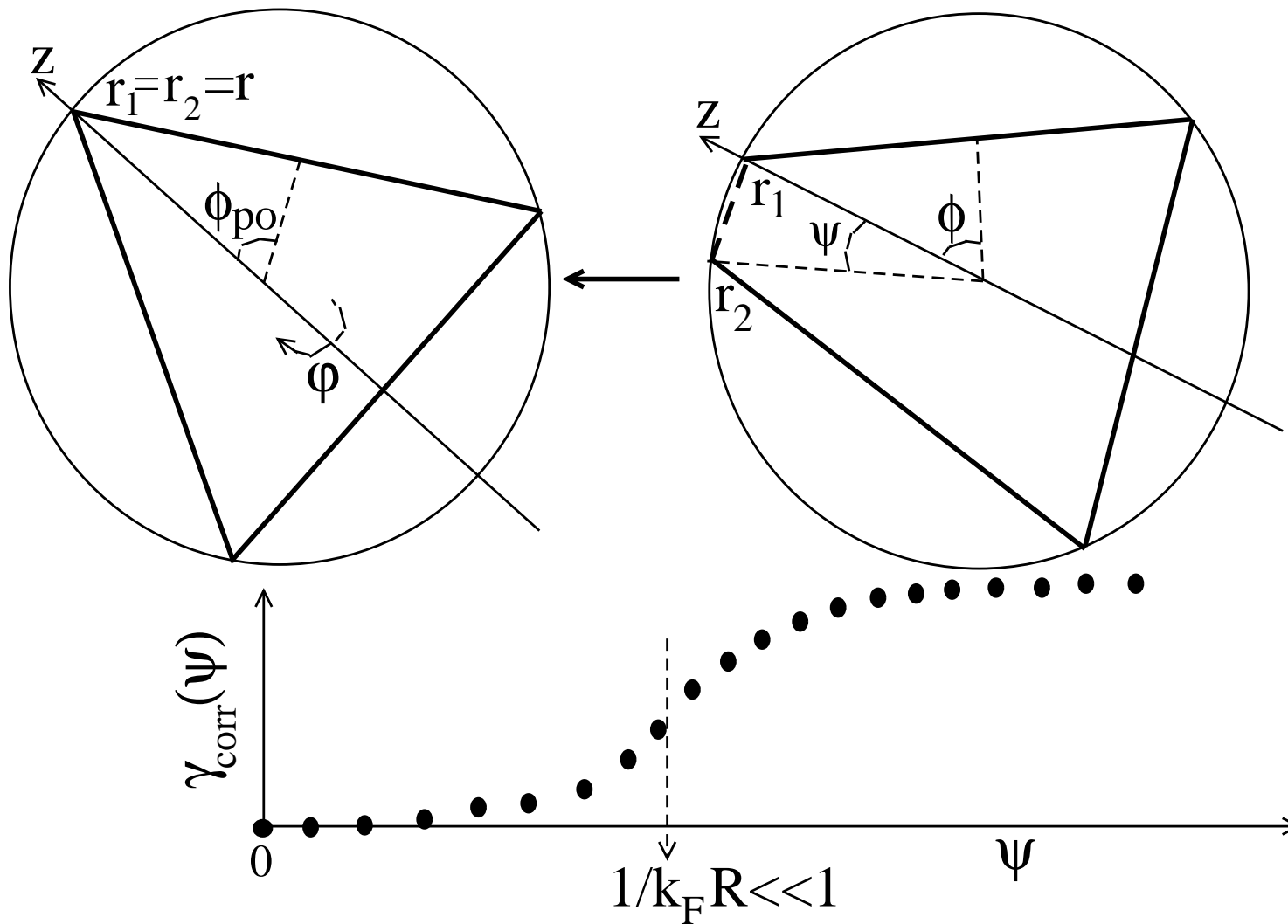
$$r_1 \{r_1, \theta_1, \varphi_1\} \Rightarrow r_2 \{r_2, \theta_2, \varphi_2\}$$

$$CT: \phi = \phi_{PO} - \frac{\psi}{2v}, \quad \phi_{PO} = \frac{\pi w}{v}$$

v is the **vertex** number and w is the **winding** number, $\psi = \theta_2 - \theta_1$

$$G = G_{CT_0} + G_{CT_1} \Rightarrow \gamma(\psi) = \gamma_{wf} + \gamma_{corr}$$

$$G_{CT_0} \approx - (m/2\pi\hbar^2 r_{12}) \exp \left(\frac{i}{\hbar} r_{12} p \right), \quad G_1 = \sum_{CT \neq CT_0} G_{CT}$$



$$\gamma_{corr}(\psi) = \gamma_{wf} \sum_{v,w} \frac{\sin^3 \phi \cos \phi}{2v \sin \psi} \left[1 - J_0 \left(\sqrt{\frac{\Delta S}{\hbar}} \right) \right], \quad \phi = \phi_{PO} - \frac{\psi}{2v}$$

$\Delta S/\hbar = \pi k_F R_0 \cos \phi_{PO} \sin \psi$, $\phi_{PO} = \pi w/v$, uniform approximation
 by M.Brack, R.K.Bhaduri, Semiclassical physics, 2003

$k_F R \setminus n$	2	3	4	5	6	7	8	9	10
10.0	0.73	1.15	0.92	1.05	0.97	1.02	0.99	1.01	0.99
20.0	0.69	1.20	0.82	1.12	0.82	1.07	0.95	1.04	0.97
30.0	0.84	1.16	0.80	1.14	0.87	1.10	0.91	1.07	0.95
40.0	0.95	1.10	0.86	1.13	0.86	1.11	0.89	1.08	0.93
50.0	0.97	1.06	0.93	1.09	0.88	1.10	0.89	1.09	0.91
∞	1.11	0.92	1.06	0.95	1.08	0.96	1.03	0.97	1.03
K&R	0.00	0.85	0.45	0.90	0.62	0.93	0.71	0.94	0.76

Table: Friction coefficients γ/γ_{wf} vs multipolarity n for small slow vibrations; K&R is Koonin&Randrup,1977, Abrosimov et al., 1999.

INERTIA: $M(\psi) = M_{ETF}(\psi) + M_{corr}(\psi), \quad \phi = \phi_{PO} - \frac{\psi}{2v}$

$$M_{corr}(\psi) \propto - \sum_{v,w} \frac{\sin^3 \phi \cos \phi |\sin \phi|}{\sin \psi} \left[1 - J_0 \left(\sqrt{\frac{\Delta S}{\hbar}} \right) \right]$$

$k_F R \setminus n$	2	3	4	5	6	7	8	9	10
5	1.4	2.0	2.5	3.6	3.9	5.0	5.3	6.4	6.6
6	1.1	1.5	1.9	2.8	2.9	3.8	4.0	4.8	5.0
7	0.9	1.2	1.5	2.2	2.4	3.1	3.2	3.9	4.0
8	0.8	1.0	1.3	1.9	2.0	2.6	2.7	3.3	3.4
9	0.7	0.9	1.1	1.7	1.7	2.3	2.3	2.9	2.9
10	0.4	0.8	1.0	1.5	1.5	2.0	2.0	2.6	2.6
∞	-1.6	1.9	2.2	-2.2	2.3	-2.3	11	8.8	2.7
K&R		0.05		0.025		0.02		0.015	

Table 1: **Inertia** M in units of the irrotational flow inertia M_{irr} vs multipolarity n and $k_F R_0$.

CONCLUSIONS

- Within the response function theory, we derived **semiclassically by using the POT the smooth friction and inertia coefficients for slow small-amplitude vibrations near the spherical shape accounting for the nonlocal trajectory corrections.**
- It was important to solve the fundamental problem of the **symmetry breaking** related to an appearance of the **periodic-orbit families** at a spherical equilibrium shape.
- The transport coefficients depend essentially on the particle numbers, especially for smaller multipolarities, in contrast to earlier results.
- **As perspectives,** our **semiclassical** results might be helpful for better understanding the **one-body friction and inertia** of the nuclear collective dynamics within the **Microscopic - Macroscopic Model** like in fission and heavy ion collisions.

THANKS!

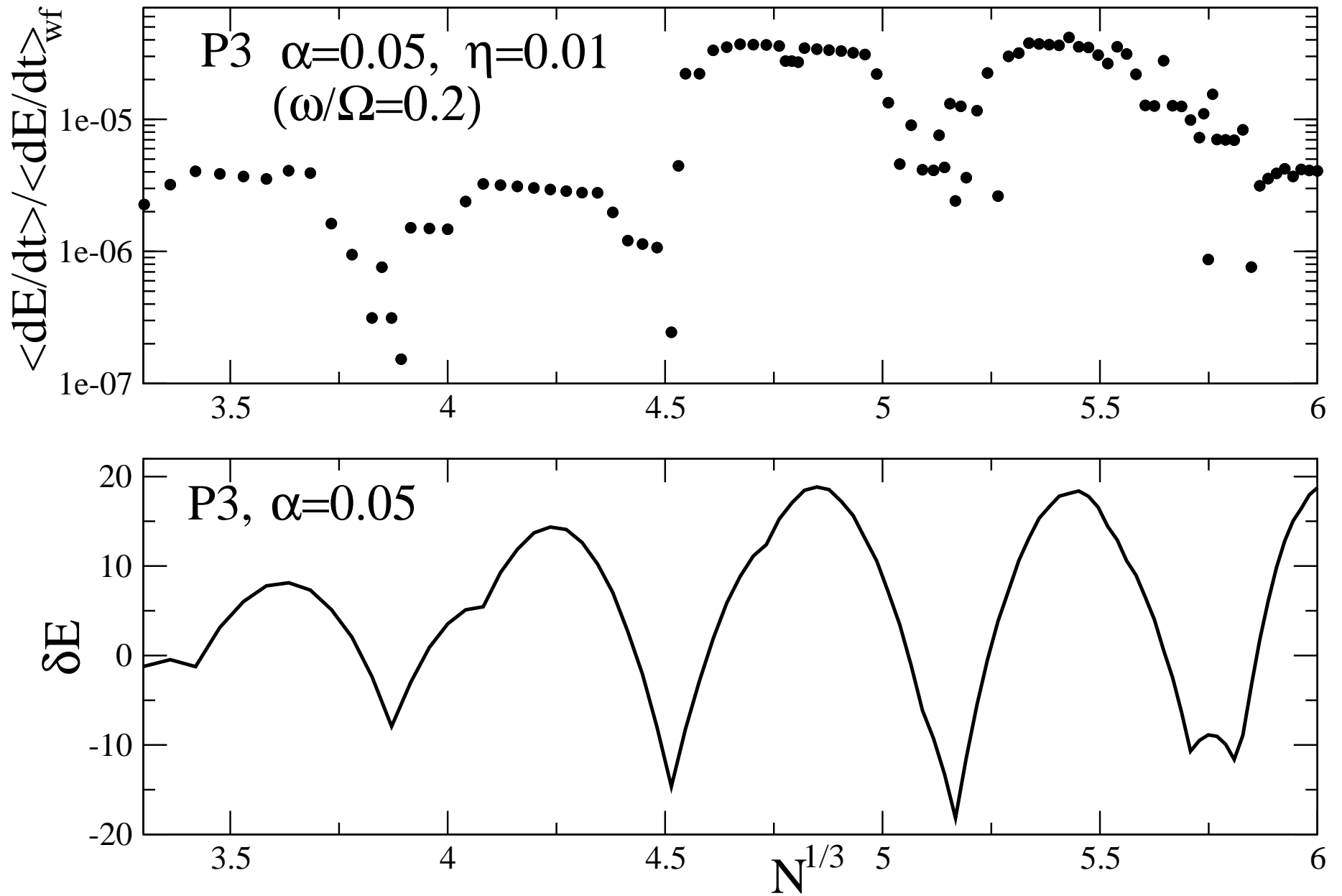


Fig. 11. Mean time derivatives $\langle dE/dt \rangle$ in the w.f. units (top) and shell corrections δE (bottom) vs. particle numbers, $N^{1/3}$.

SHELL CORRECTIONS

$$\delta E = \sum_{PO} \left(\frac{\hbar}{t_{PO}} \right)^2 \mathcal{B}_{PO} \cos \left[\frac{1}{\hbar} S_{PO}(\varepsilon_F) - \frac{\pi}{2} \mu_{PO} \right], \quad (\text{Strutinsky \& M., 1975, 76})$$

$$\delta \gamma \propto \sum_{CT, CT'} \Upsilon_{CT, CT'} \cos \left\{ \frac{1}{\hbar} [S_{CT}(\varepsilon_F) \pm S_{CT'}(\varepsilon_F)] \right\}, \quad \hbar \Omega = 2\pi \hbar / t_{CT}$$

$k_F R \setminus n$	0	1	2	3	4	5	6	7	8	9	10
5.0	0.00	1.31	0.92	1.07	0.98	1.02	0.99	1.01	1.00	1.00	1.00
10.0	0.20	1.35	0.73	1.15	0.92	1.05	0.97	1.02	0.99	1.01	0.99
15.0	0.60	1.29	0.67	1.19	0.86	1.09	0.88	1.04	0.97	1.02	0.99
20.0	0.86	1.22	0.69	1.20	0.82	1.12	0.82	1.07	0.95	1.04	0.97
25.0	0.96	1.16	0.76	1.19	0.80	1.14	0.77	1.08	0.94	1.05	0.97
30.0	0.93	1.19	0.84	1.16	0.80	1.14	0.87	1.10	0.91	1.07	0.95
35.0	0.86	1.10	0.91	1.13	0.83	1.14	0.86	1.11	0.90	1.07	0.94
40.0	0.80	1.09	0.95	1.10	0.86	1.13	0.86	1.11	0.89	1.08	0.93
45.0	0.75	1.09	0.97	1.08	0.89	1.11	0.87	1.11	0.89	1.09	0.92
50.0	0.75	1.09	0.97	1.06	0.93	1.09	0.88	1.10	0.89	1.09	0.91
∞	1.59	0.83	1.11	0.92	1.06	0.95	1.08	0.96	1.03	0.97	1.03
K&R	0.00	0.00	0.00	0.85	0.45	0.90	0.62	0.93	0.71	0.94	0.76

Table: Friction coefficients γ/γ_{wf} vs multipolarity n for small slow.

vibrations; K&R is Koonin&Randrup,1977, Abrosimov et al., 1999.

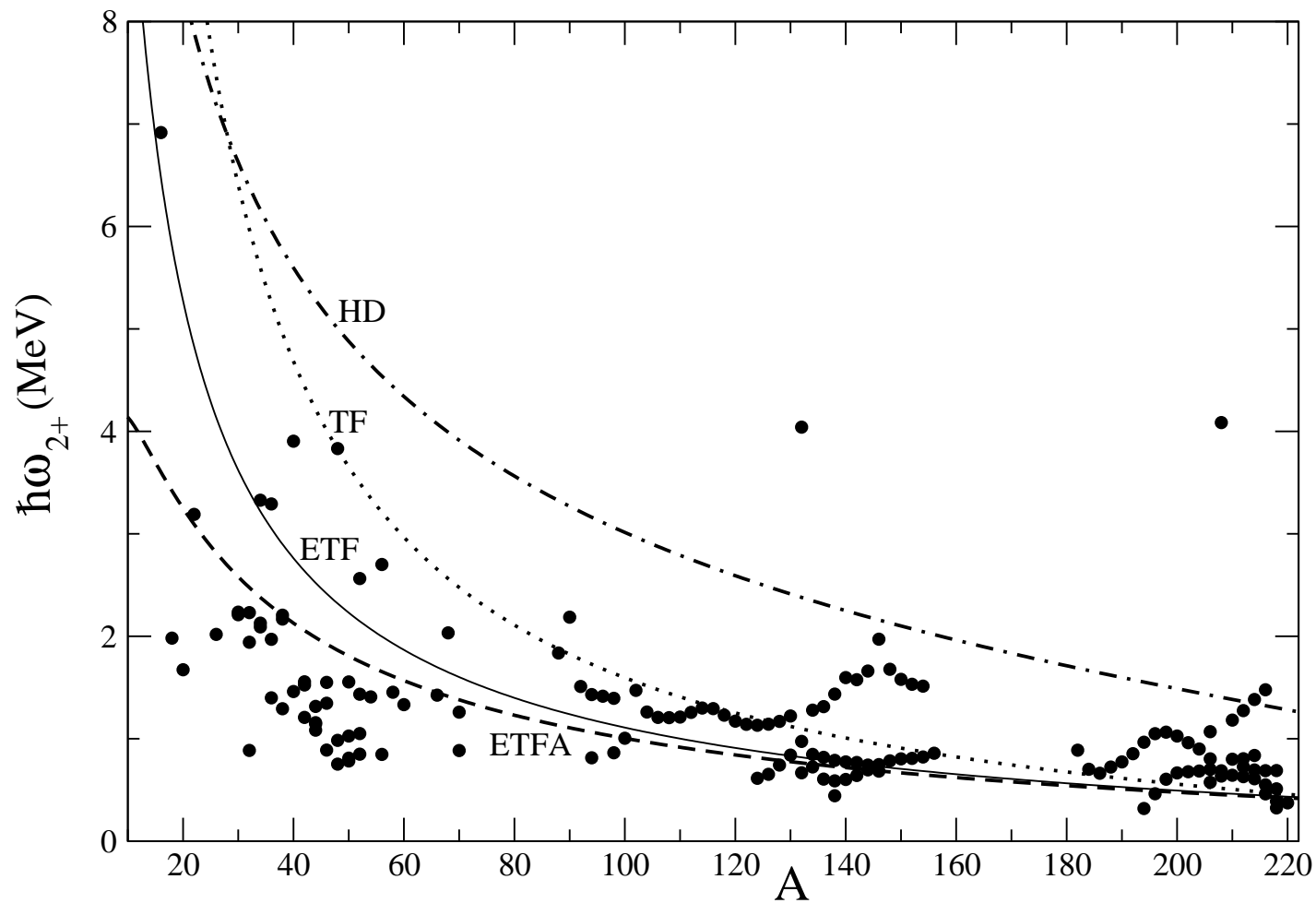


Fig. 2. Low-lying quadrupole vibration energies $\hbar\omega_{2+}$ vs particle number A ; frequent dots are TF; thin solid is ETF; thick dashed is ETFA; dash-dotted shows HD. Heavy full dots are the experimental data [Raman, 2001] for nearly spherical nuclei with $q_2 < 0.05$; $\rho = 0.16 \text{ fm}^{-3}$, $b_v = 16 \text{ MeV}$, $b_s = 18 \text{ MeV}$, $K = 220 \text{ MeV}$, $b_{\text{sym}} = 60 \text{ MeV}$.

$k_F R \setminus n$	0	1	2	3	4	5	6	7	8	9	10
5		2.7	1.4	2.0	2.5	3.6	3.9	5.0	5.3	6.4	6.6
6		2.5	1.1	1.5	1.9	2.8	2.9	3.8	4.0	4.8	5.0
7		2.4	0.9	1.2	1.5	2.2	2.4	3.1	3.2	3.9	4.0
8		2.3	0.8	1.0	1.3	1.9	2.0	2.6	2.7	3.3	3.4
9		2.3	0.7	0.9	1.1	1.7	1.7	2.3	2.3	2.9	2.9
10		2.3	0.4	0.8	1.0	1.5	1.5	2.0	2.0	2.6	2.6
∞		2.0	-1.6	1.9	2.2	-2.2	2.3	-2.3	11	8.8	2.7
K&R		1.00		0.05		0.025		0.02		0.015	

Table 2: Inertia M in units of the irrotational flow inertia M_{irr} vs multipolarity n and $k_F R_0$.

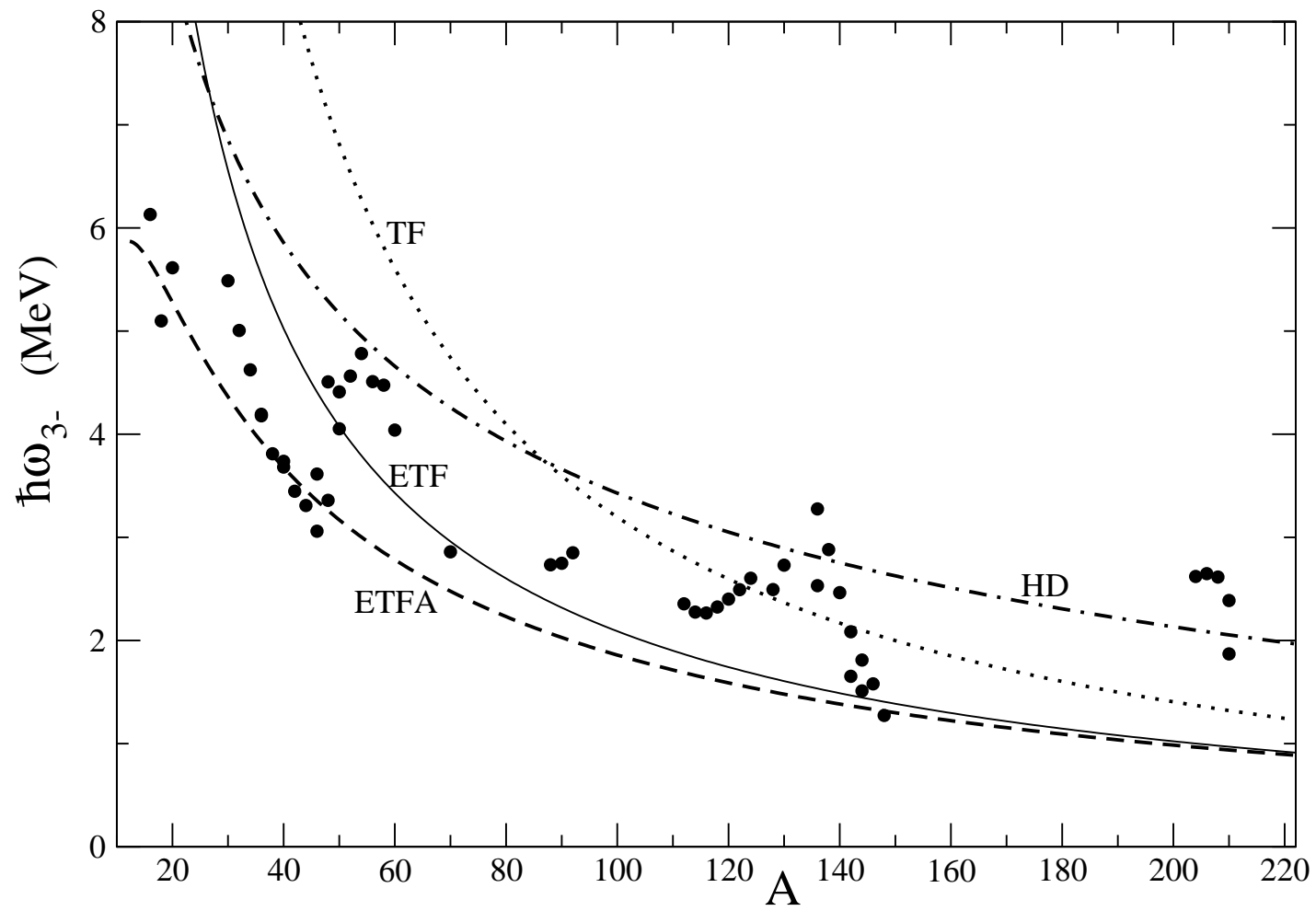


Fig. 3. The same for the low-lying octupole vibration energies $\hbar\omega_{3-}$; the experimental data are taken from [Kibedi, 2005].

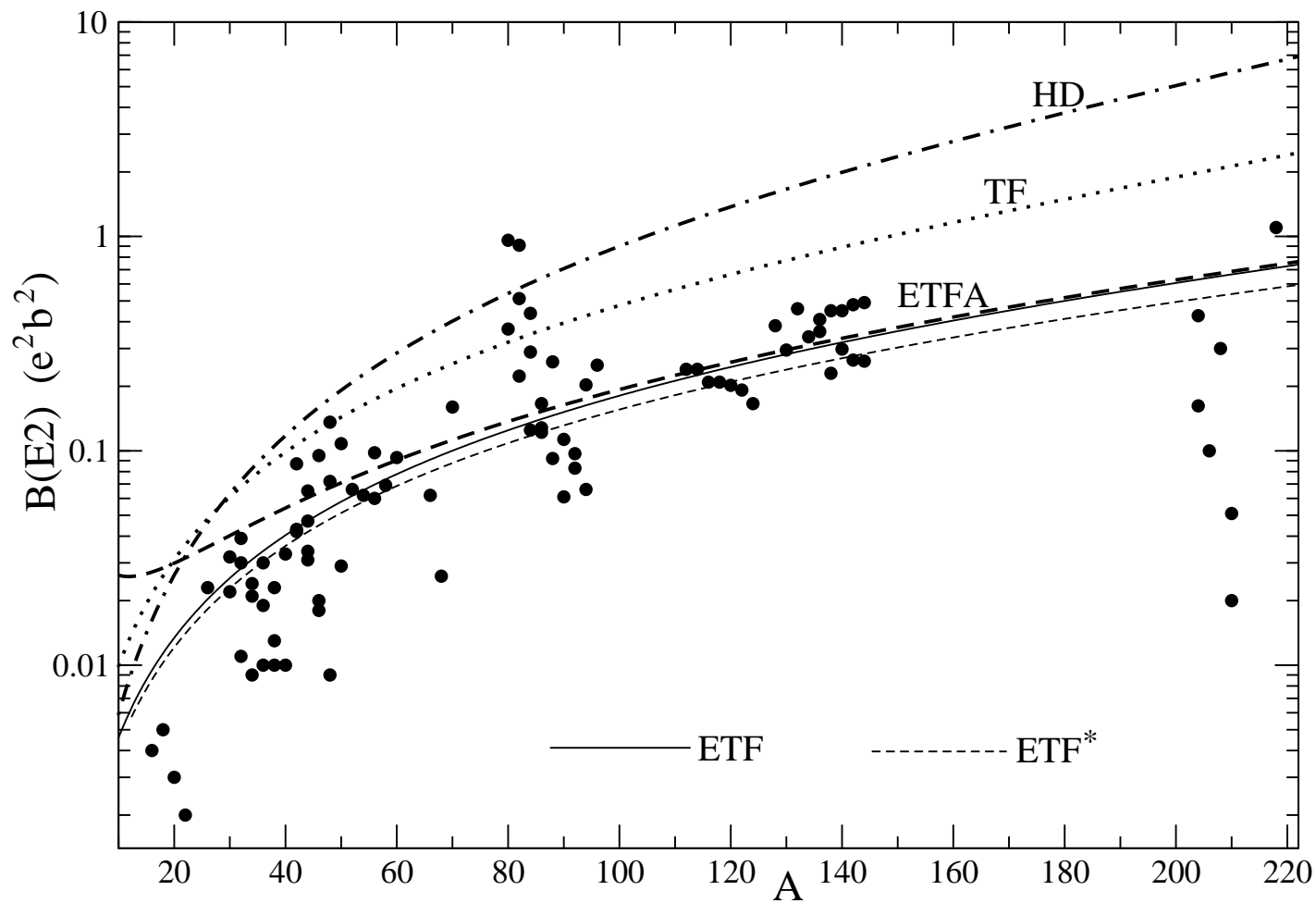


Fig. 4. The reduced probabilities $B(E2)$ in units of $e^2 b^2$; full heavy points are experimental data [Raman, 2001]; frequent dashed is semiclassical $B(E2)$ with (solid) and without (dashed) account for the η dependence; other notations are the same as in Fig. 2.

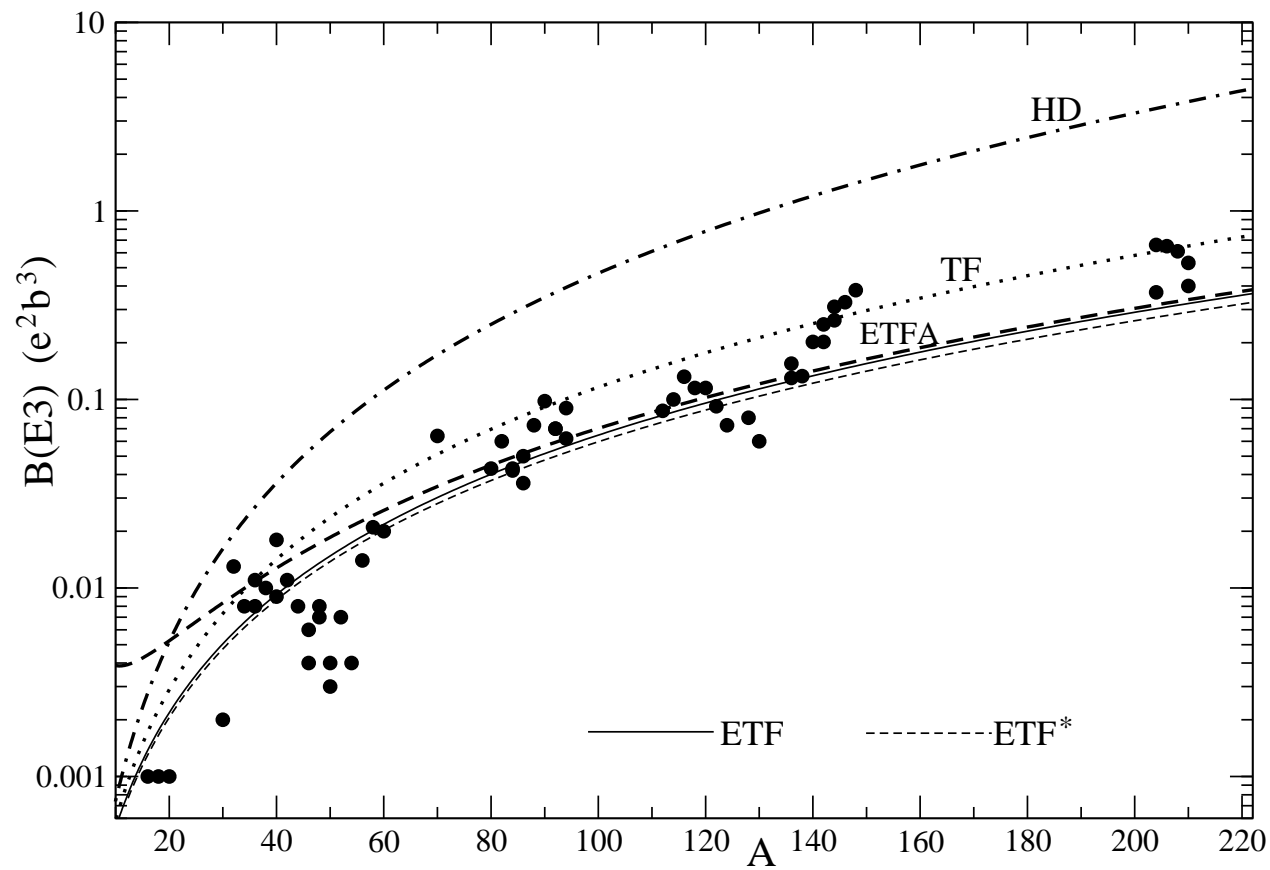


Fig. 5. The reduced transition probabilities $B(E3)$ in units of $e^2 b^3$; full heavy points are experimental data [Kibedi, 2005]; other notations are the same as in Fig. 4.

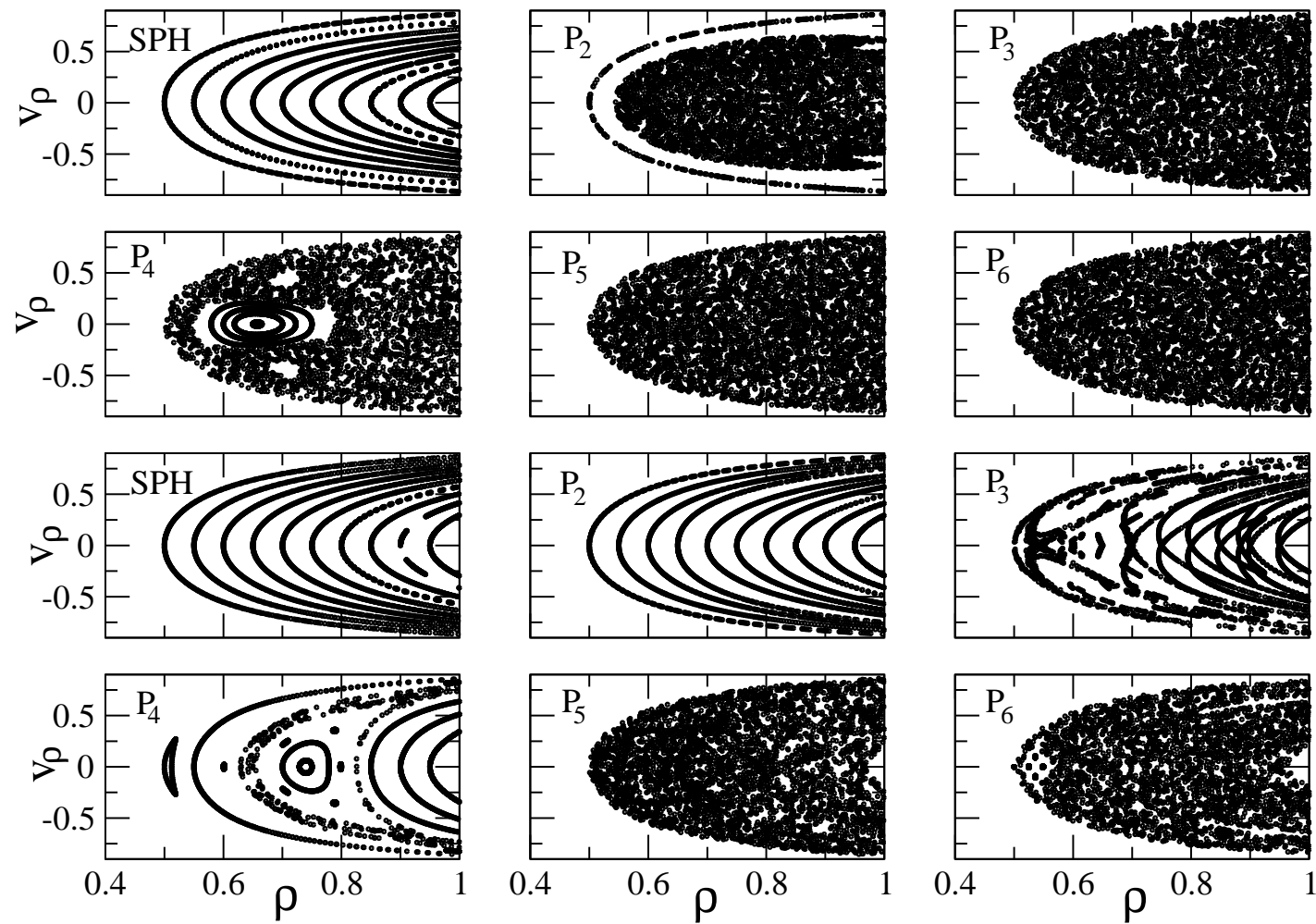


Fig.4. Poincare sections for six shapes at the deformation $\alpha = 0.05$ (two lower rows) and deformation $\alpha = 0.4$ (two upper rows) for the projections of the angular momentum on the symmetry axis $K = 0.5$ with respect to the maximal value

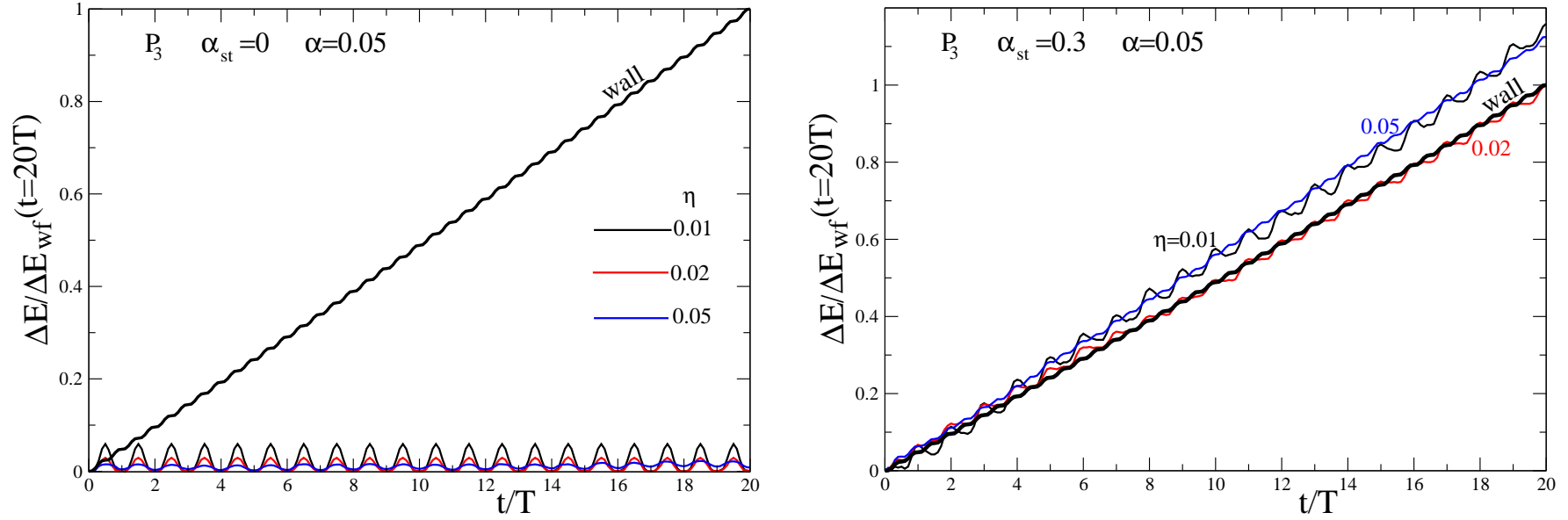


Fig.6. The classical dynamical excitations ΔE in units of the wall formula value $\Delta E_{wf}(t = 20T)$ as functions of the time t in units of the period $T = 2\pi\alpha/\eta$ for vibrations near the spherical ($\alpha_{st} = 0$, left) and deformed ($\alpha_{st} = 0.3$, right) shapes of the cavity