

Energy density functional description of nuclear low-lying spectrum and shape transition

Zhipan Li (李志攀)

*School of Physical Science and Technology,
Southwest University, Chongqing 400715, China*

Physics Department, Faculty of Science, University of Zagreb, Croatia



Outline

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Introduction

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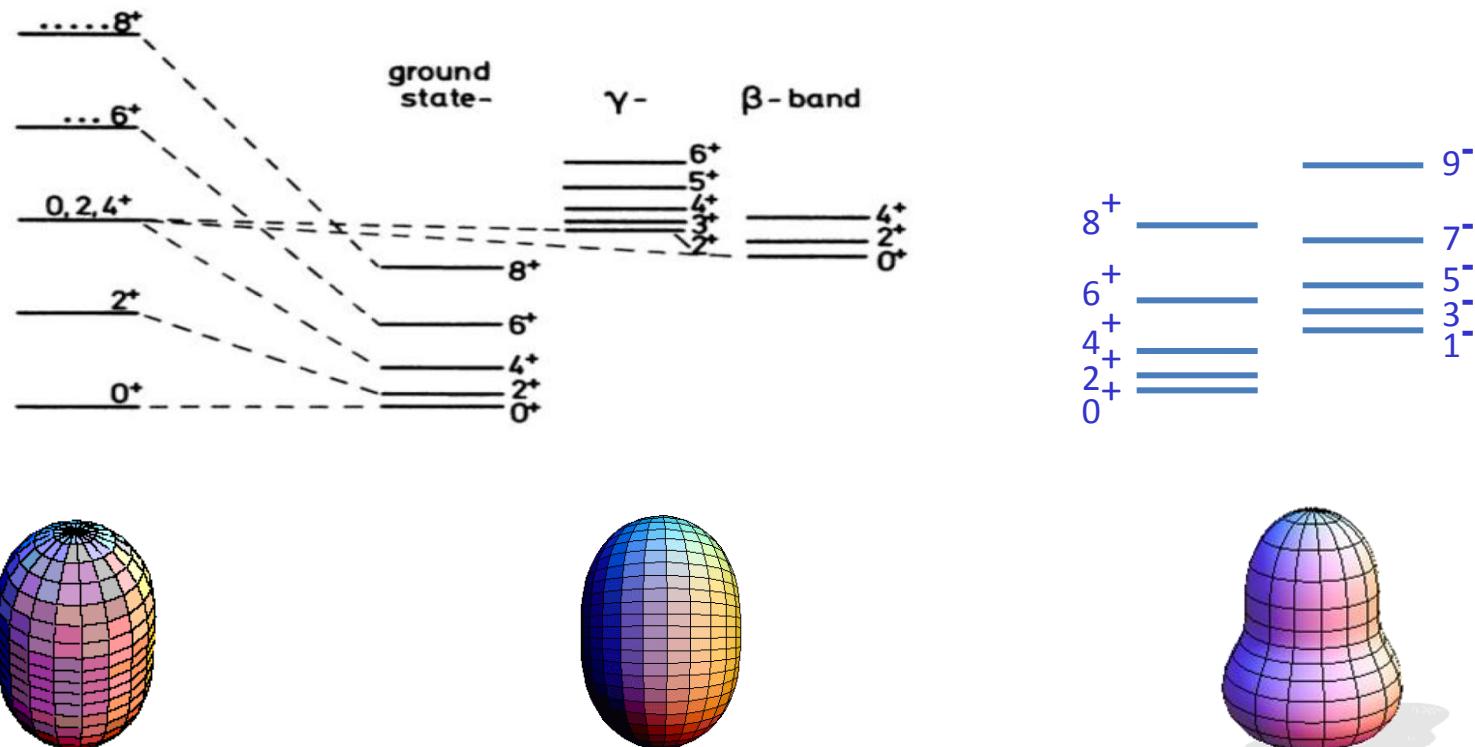
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Summary

Nuclear Low-lying Spectrum

- Nuclear low-lying spectroscopy can reveal rich structure information of atomic nuclei

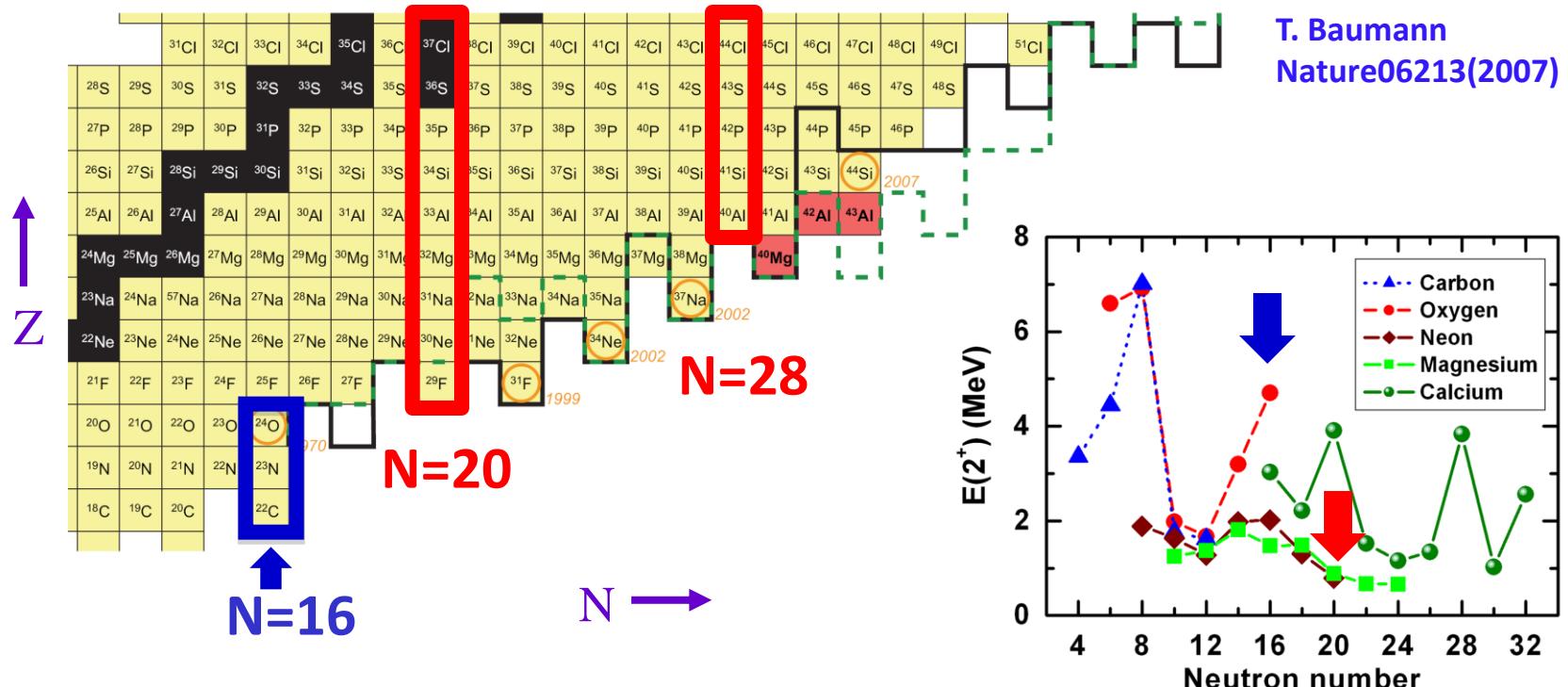
- ◆ Shape and shape transition



Nuclear Low-lying Spectrum

- Nuclear low-lying spectroscopy can reveal rich structure information of atomic nuclei

- ◆ Shape and shape transition
- ◆ Evolution of the shell structure



Nuclear Low-lying Spectrum

➤ CEDF: nuclear structure over almost the whole nuclide chart

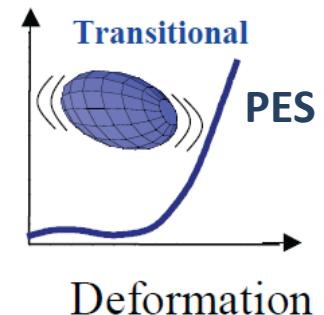
Ring96, Vretenar2005, Meng2006

- ◆ Spin-orbit splitting
- ◆ Scalar and vector fields: nuclear saturation properties
- ◆ Origin of the pseudo-spin symmetry
- ◆ Spin symmetry in anti-nucleon spectrum
- ◆

➤ Spectrum: beyond the mean-field approximation

- ◆ Restoration of broken symmetry, e.g. rotational
- ◆ Mixing of different shape configurations

AMP+GCM: Niksic2006, Yao2010



5D Collective Hamiltonian based on CEDF

Theoretical Framework

Coll. Potential

Moments of inertia

Mass parameters

Triaxial
Covariant
Density
Functional
ph + pp

Construct
5-dimensional
Hamiltonian

(*vib + rot*)

Diagonalize:
Nuclear spectroscopy

$E(J^\pi)$, BE2 ...
Cal. \leftrightarrow Exp.

Libert, Girod & Delaroche, PRC60, 054301 (99)

Prochniak & Rohozinski, JPG36, 123101 (09)

Niksic, Li, Vretenar, Prochniak, Meng & Ring, PRC79, 034303 (09)

Collective Hamiltonian

The five-dimensional Hamiltonian: [Bohr1952]

$$\hat{H}(\beta, \gamma, \Omega) = \hat{T}_{\text{rot}}(\beta, \gamma, \Omega) + \hat{T}_{\text{vib}}(\beta, \gamma) + V_{\text{coll}}(\beta, \gamma). \quad (1)$$

- rotational kinetic energy:

$$\hat{T}_{\text{rot}} = \frac{1}{2} \sum_{k=1}^3 \hat{j}_k^2 / \mathcal{I}_k, \quad (2)$$

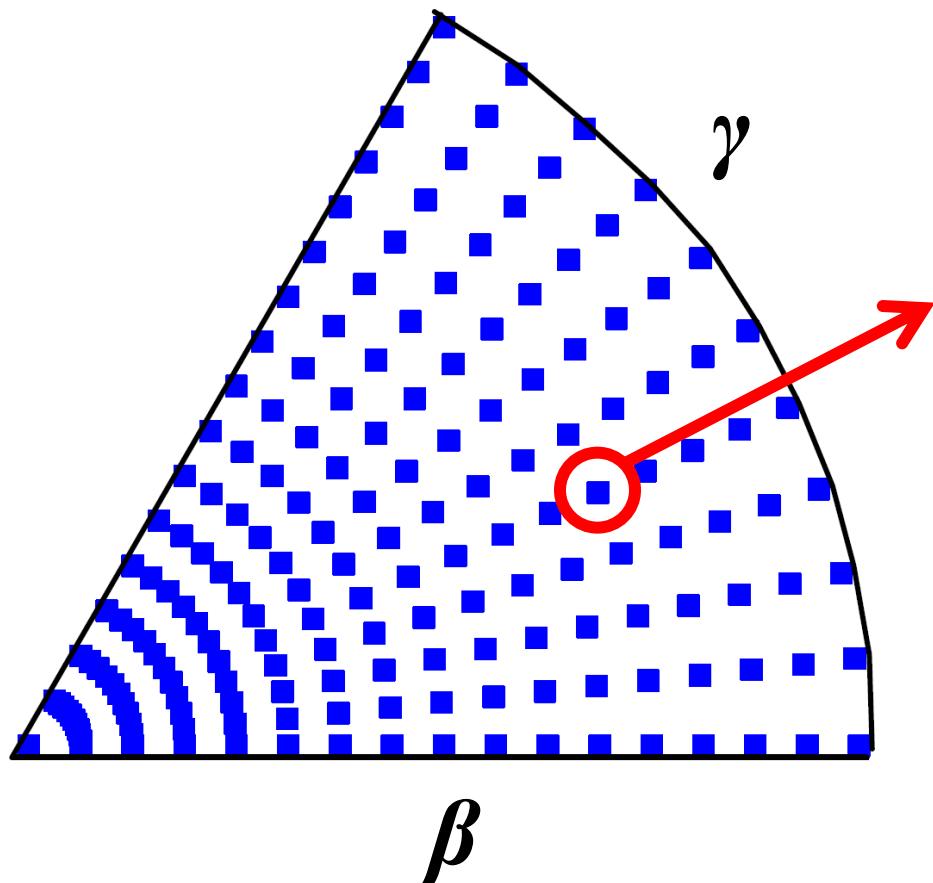
- vibrational kinetic energy:

$$\begin{aligned} \hat{T}_{\text{vib}} = & -\frac{\hbar^2}{2\sqrt{wr}} \left\{ \frac{1}{\beta^4} \left[\partial_\beta \left(\beta^4 \sqrt{\frac{r}{w}} B_{\gamma\gamma} \partial_\beta \right) - \partial_\beta \left(\beta^3 \sqrt{\frac{r}{w}} B_{\beta\gamma} \partial_\gamma \right) \right] \right. \\ & \left. + \frac{1}{\beta \sin 3\gamma} \left[-\partial_\gamma \left(\sqrt{\frac{r}{w}} \sin 3\gamma B_{\beta\gamma} \partial_\beta \right) + \frac{1}{\beta} \partial_\gamma \left(\sqrt{\frac{r}{w}} \sin 3\gamma B_{\beta\beta} \partial_\gamma \right) \right] \right\}. \end{aligned} \quad (3)$$

- V_{col} is the collective potential.

7 collective parameters (β, γ) : $\mathcal{I}_1, \mathcal{I}_2, \mathcal{I}_3, B_{\beta\beta}, B_{\beta\gamma}, B_{\gamma\gamma}$, and V_{coll}

Collective Hamiltonian



$I_1(\beta, \gamma)$	$I_2(\beta, \gamma)$	$I_3(\beta, \gamma)$
$B_{\beta\beta}(\beta, \gamma)$	$B_{\beta\gamma}(\beta, \gamma)$	$B_{\gamma\gamma}(\beta, \gamma)$
$V_{\text{coll}}(\beta, \gamma)$		

How to determine
these 7 collective
parameters in a
microscopic way?

Covariant Energy Density Functional (CEDF)

→ The CEDF (*contact interaction or Point-Coupling*):

$$\begin{aligned} E_{\text{RMF}} &= \int d\mathbf{r} \mathcal{E}_{\text{RMF}}(\mathbf{r}) \\ &= \sum_k \int d\mathbf{r} v_k^2 \bar{\psi}_k(\mathbf{r}) (-i\gamma\boldsymbol{\nabla} + m) \psi_k(\mathbf{r}) \\ &\quad + \int d\mathbf{r} \left(\frac{\alpha_S}{2} \rho_S^2 + \frac{\beta_S}{3} \rho_S^3 + \frac{\gamma_S}{4} \rho_S^4 + \frac{\delta_S}{2} \rho_S \Delta \rho_S + \frac{\alpha_V}{2} j_\mu j^\mu + \frac{\gamma_V}{4} (j_\mu j^\mu)^2 + \frac{\delta_V}{2} j_\mu \Delta j^\mu \right. \\ &\quad \left. + \frac{\alpha_{TV}}{2} j_{TV}^\mu (j_{TV})_\mu + \frac{\delta_{TV}}{2} j_{TV}^\mu \Delta (j_{TV})_\mu + \frac{\alpha_{TS}}{2} \rho_{TS}^2 + \frac{\delta_{TS}}{2} \rho_{TS} \Delta \rho_{TS} + \frac{e}{2} \rho_p A^0 \right), \end{aligned}$$

Burvenich02, Niksic08, Zhao10

→ Collective parameters: Cranking approximation

$$\mathcal{I}_k = \sum_{i,j} \frac{(u_i v_j - v_i u_j)^2}{E_i + E_j} |\langle i | \hat{j}_i | j \rangle|^2$$

$$B_{\mu\nu} = \frac{1}{2} \left[\mathcal{M}_{(1)}^{-1} \mathcal{M}_{(3)} \mathcal{M}_{(1)}^{-1} \right]_{\mu\nu},$$

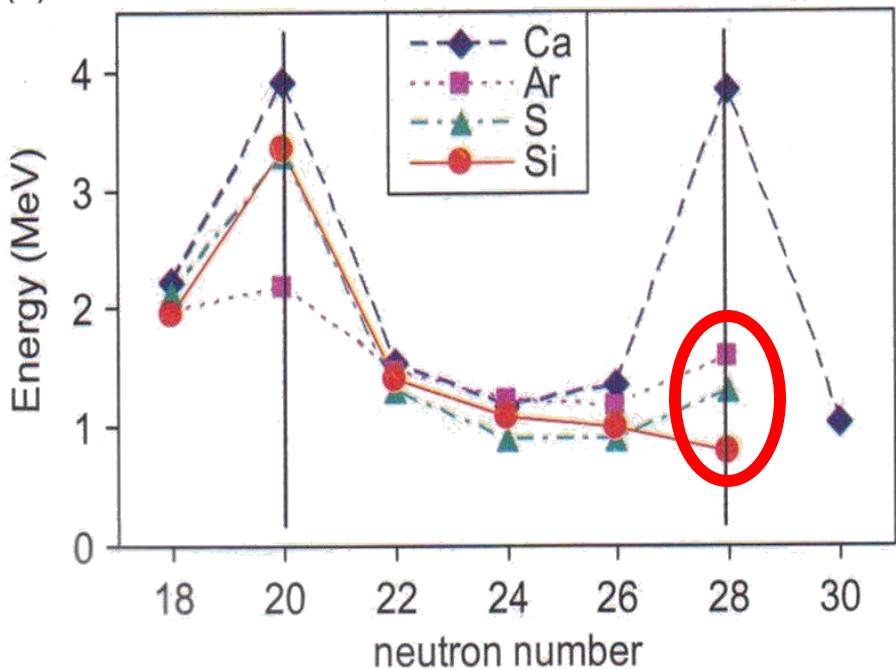
$$V_{\text{coll}} = E_{\text{tot}} - \Delta V_{\text{vib}} - \Delta V_{\text{rot}},$$

$$\mathcal{M}_{(n),\mu\nu} = \sum_{i,j} \frac{(u_i v_j + v_i u_j)^2}{(E_i + E_j)^n} \langle i | \hat{Q}_{2\mu} | j \rangle \langle j | \hat{Q}_{2\nu} | i \rangle$$

Results and discussion

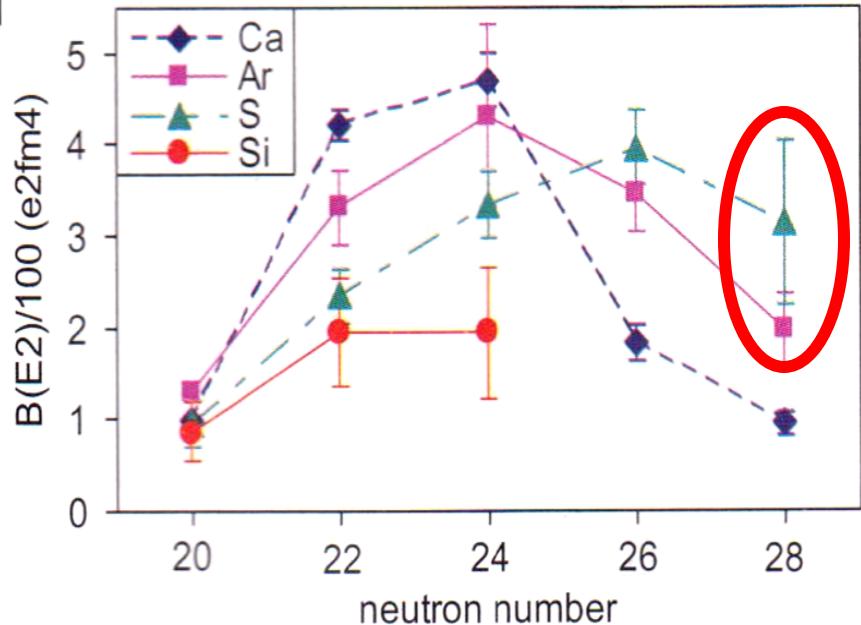
- ◆ **Rapid shape evolution in $N=28$ isotones**
- ◆ **Enhanced collectivity in Sn isotopes**
- ◆ **Effect of time-odd mean field:**
Expansion method for TV inertia parameter

Shape evolution & coexistence in $N=28$ isotones



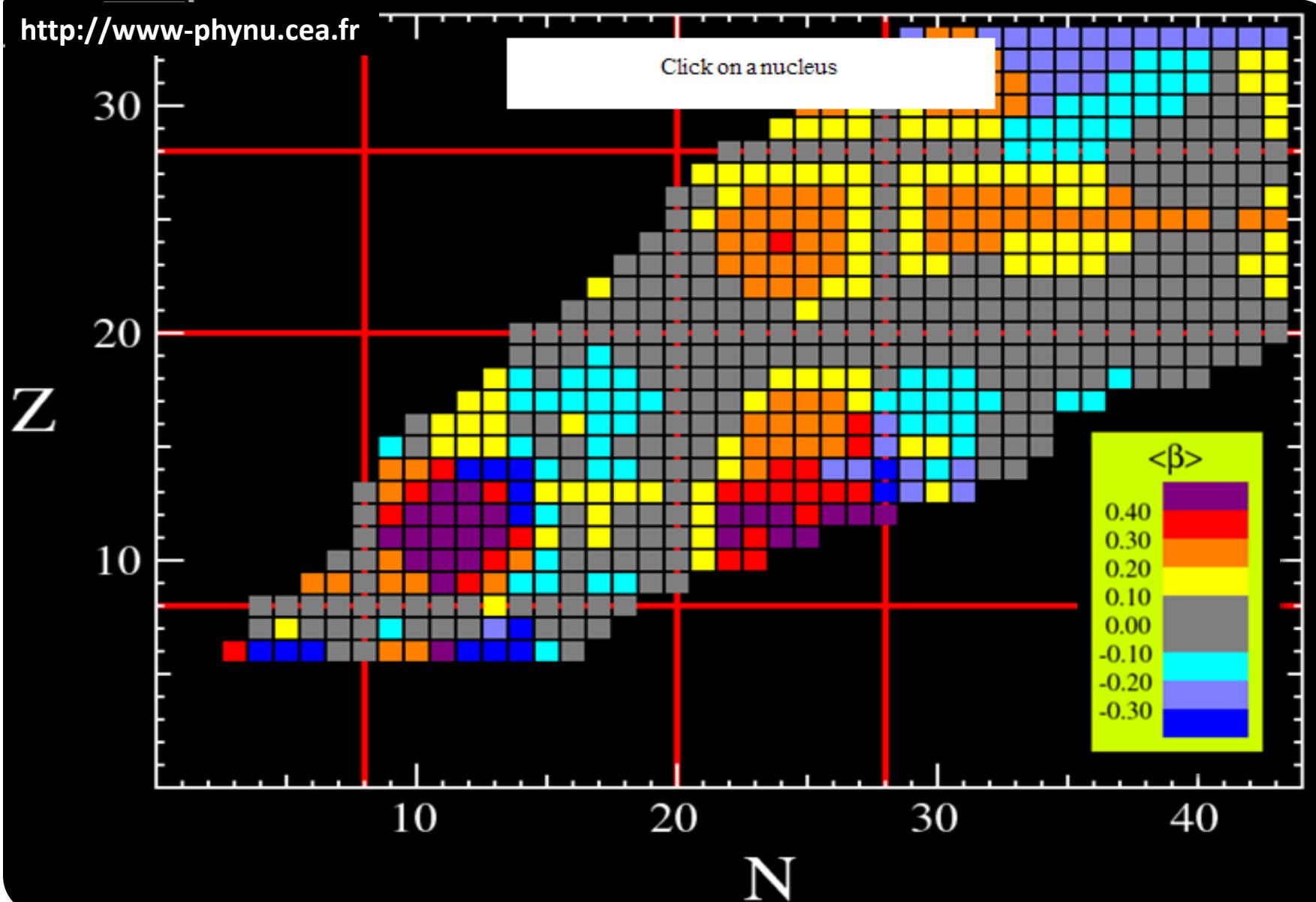
$E(2_1^+)$

$B(E2)$



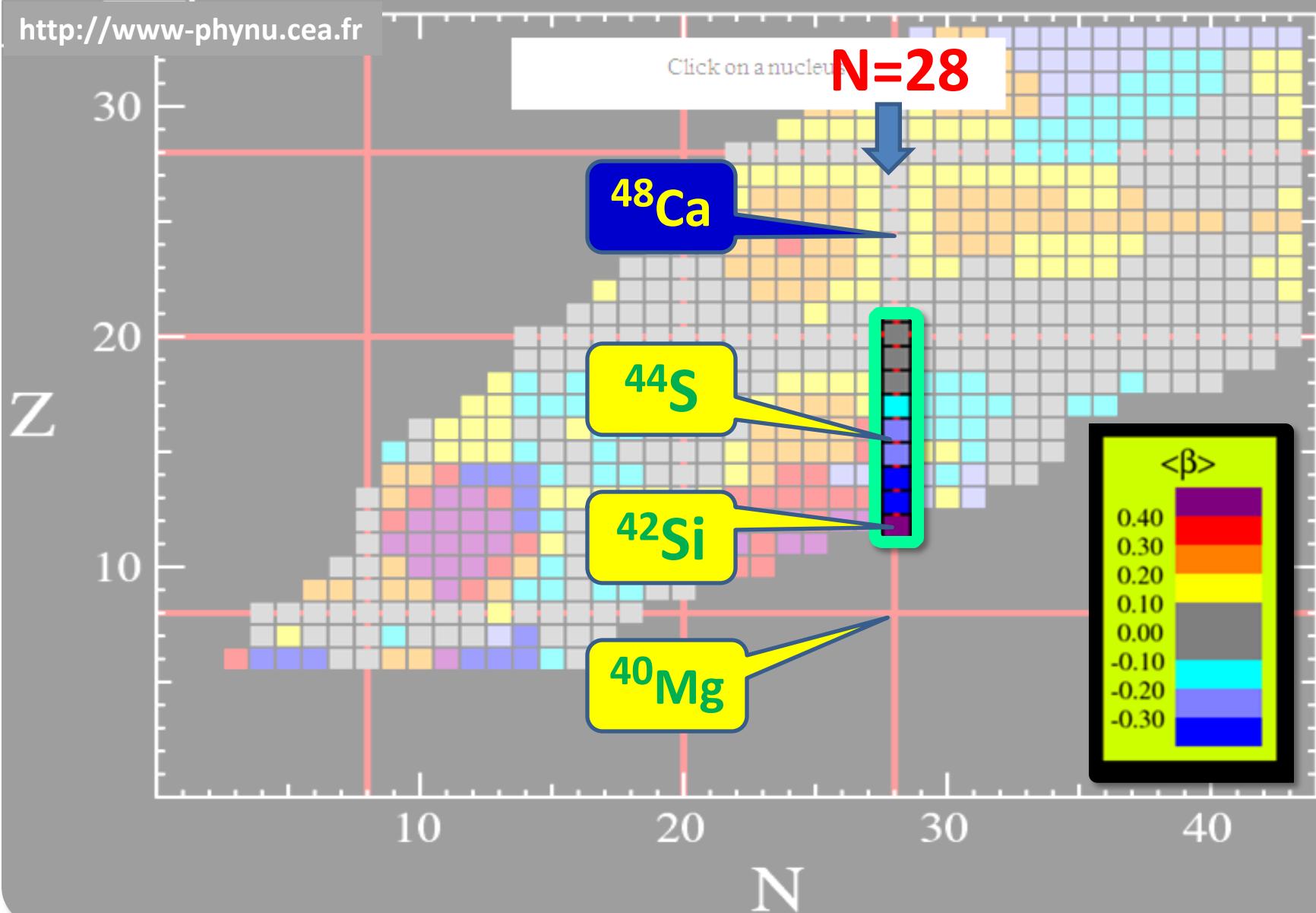
Shape evolution & coexistence in $N=28$ isotones

<http://www-phynu.cea.fr>

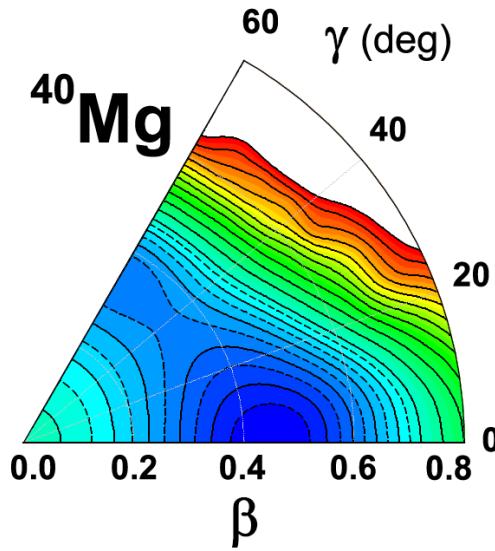
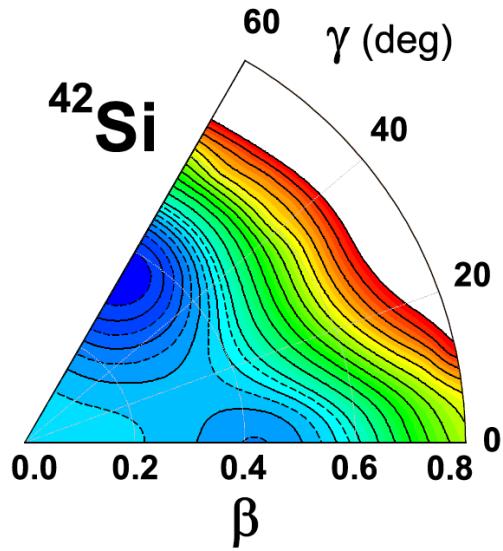
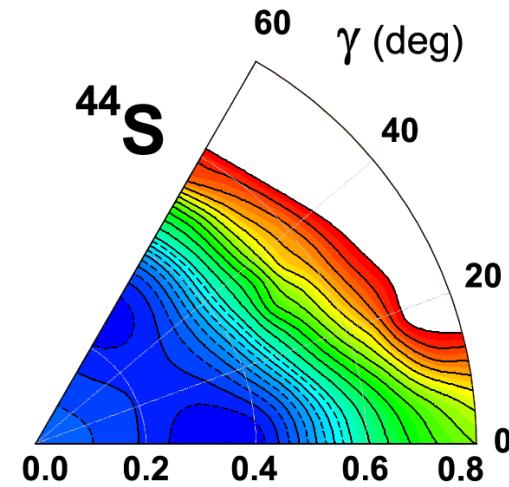
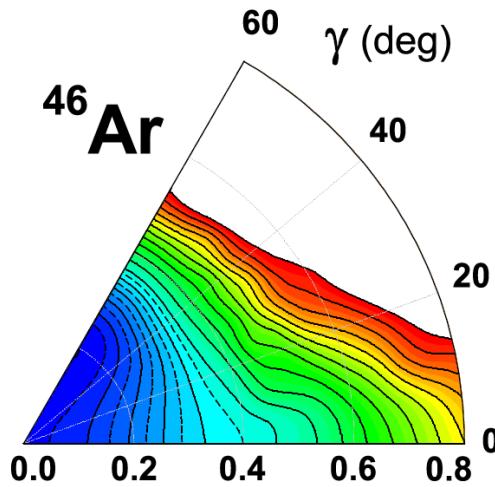
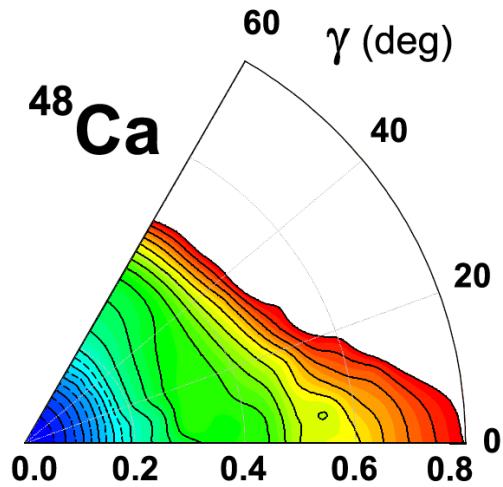


Shape evolution & coexistence in $N=28$ isotones

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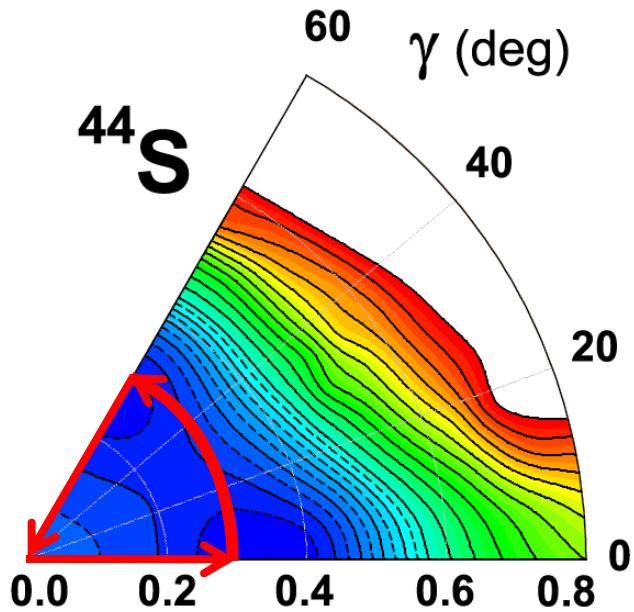


Shape evolution & coexistence in $N=28$ isotones



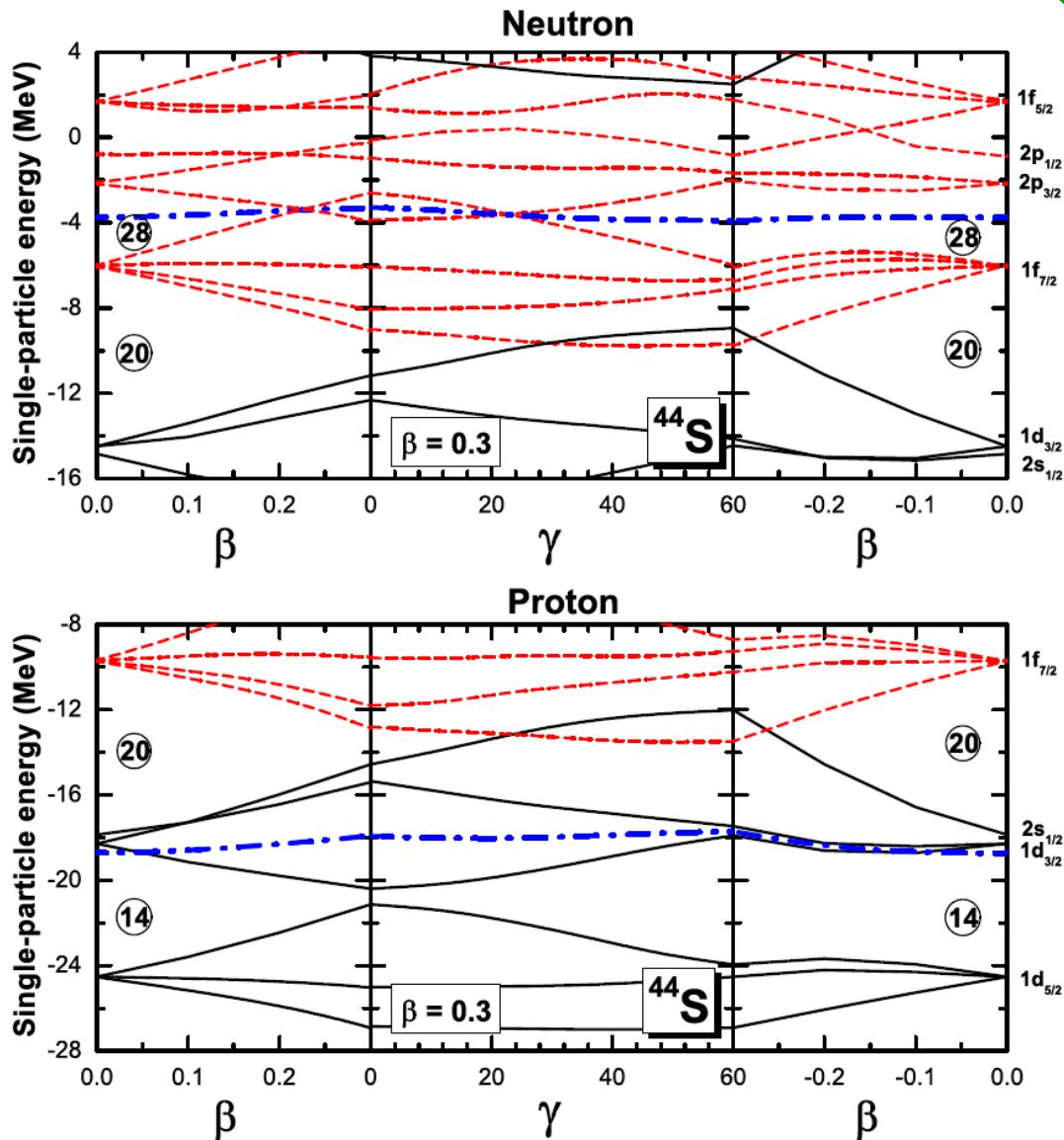
	$\Delta_{N=28}^{\text{sph.}}$	β_{\min}
^{48}Ca	4.73	0.00
^{46}Ar	4.48	-0.19
^{44}S	3.86	0.34
^{42}Si	3.13	-0.35
^{40}Mg	2.03	0.45

Shape evolution & coexistence in $N=28$ isotones

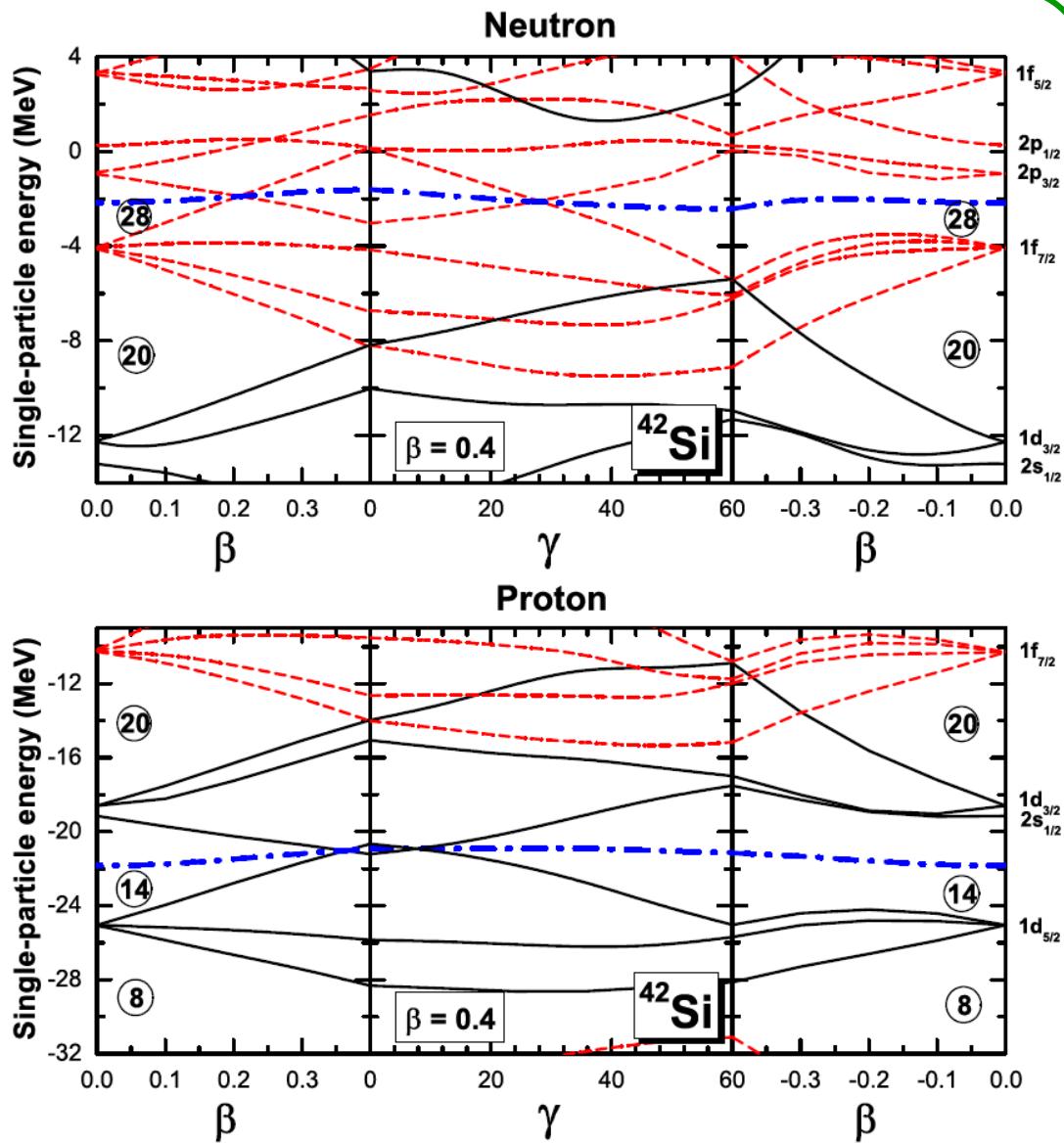
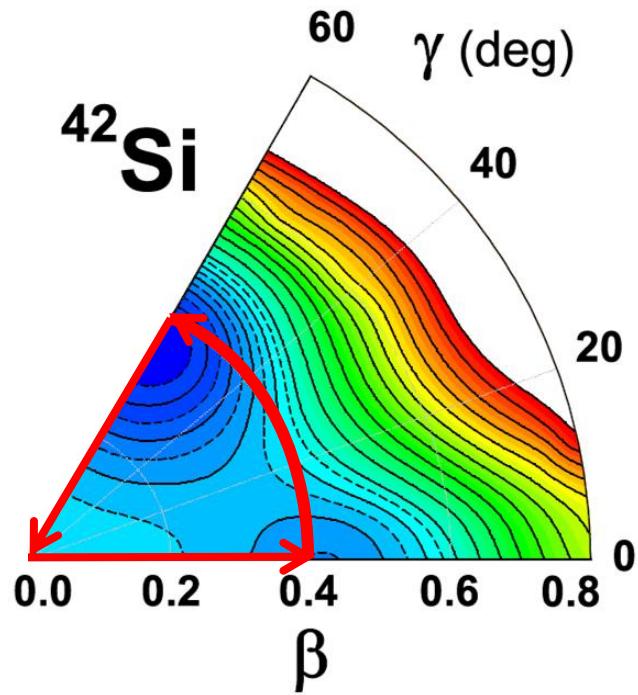


Oblate: $\Delta n = 3.91$ MeV

Prolate: $\Delta p = 5.05$ MeV



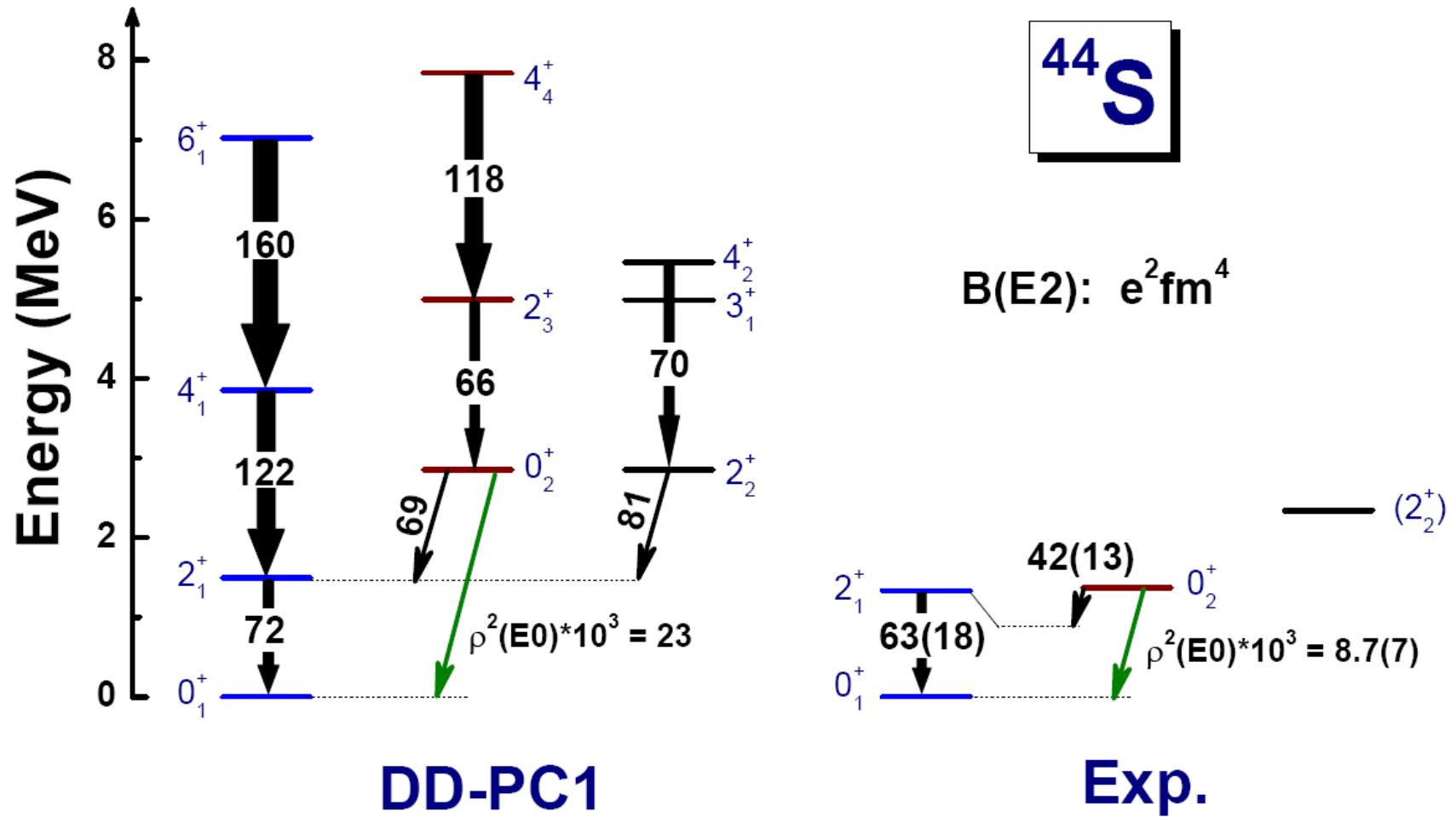
Shape evolution & coexistence in $N=28$ isotones



Shape evolution & coexistence in $N=28$ isotones

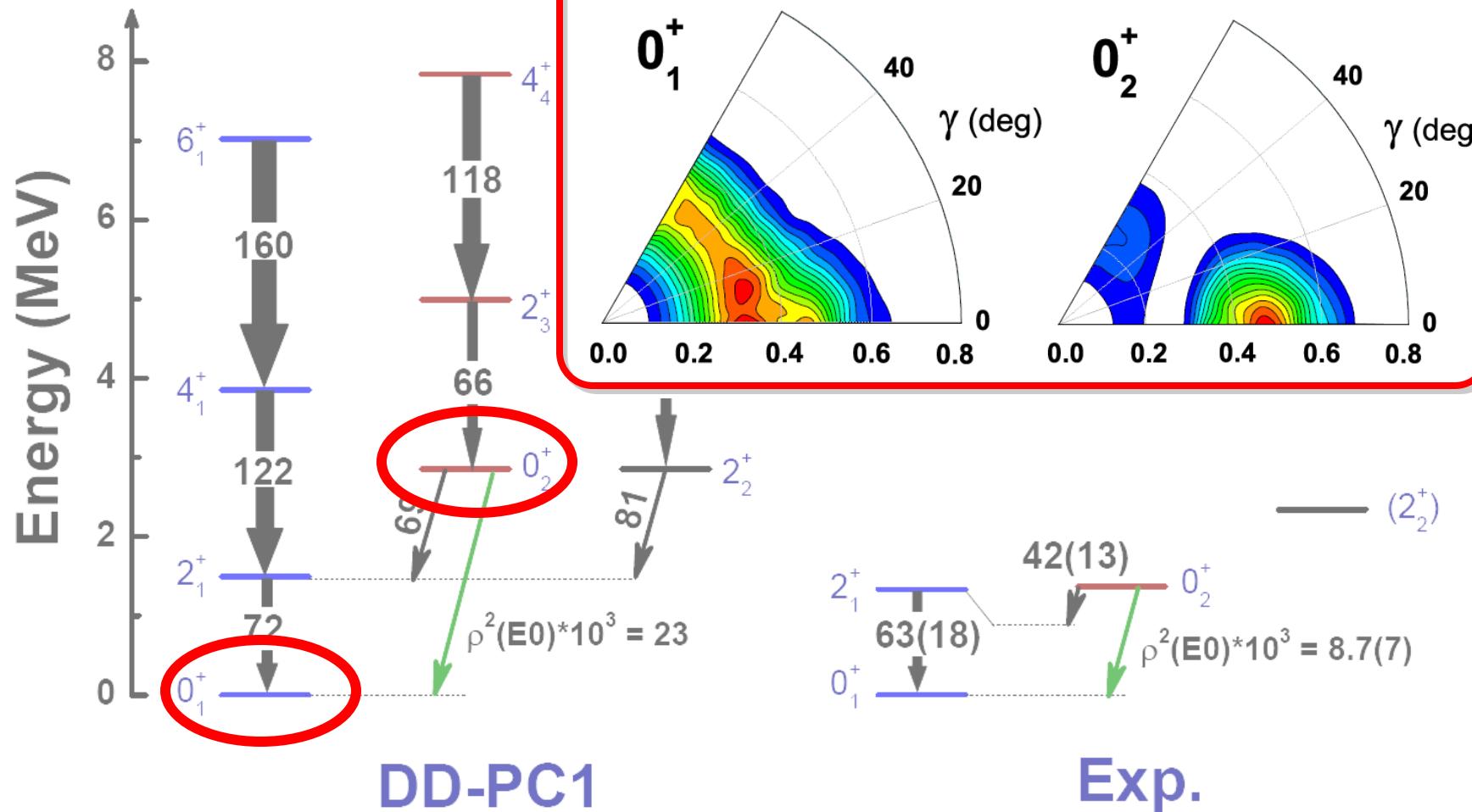
➤ Low-lying spectrum

C. Force et al., PRL105

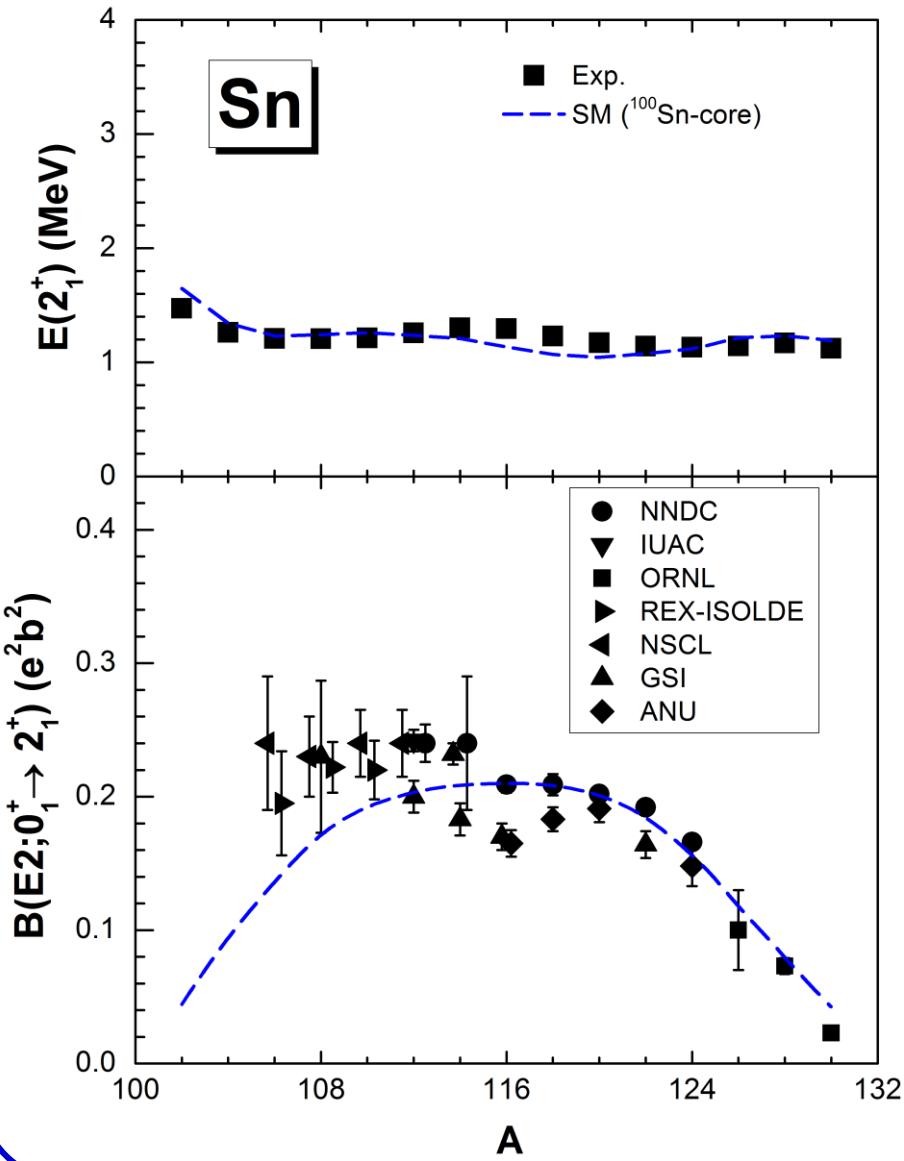


Shape evolution & coexistence in $N=28$ isotones

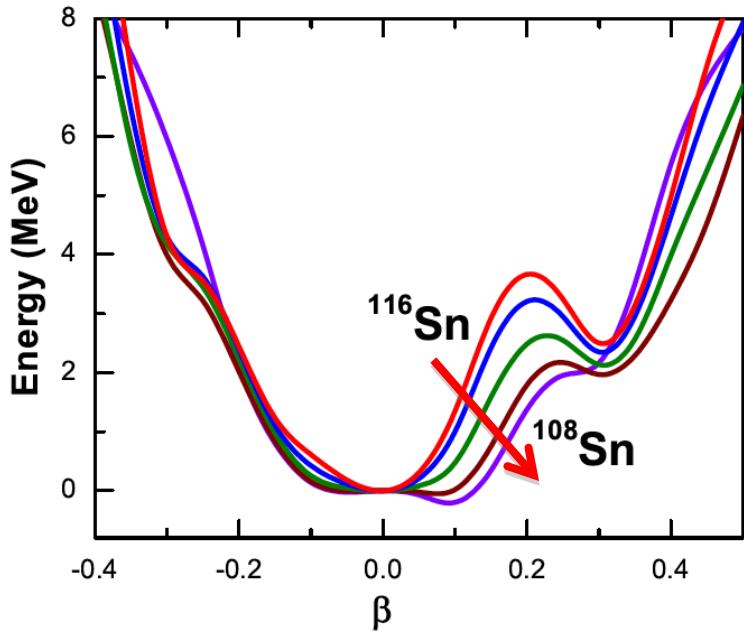
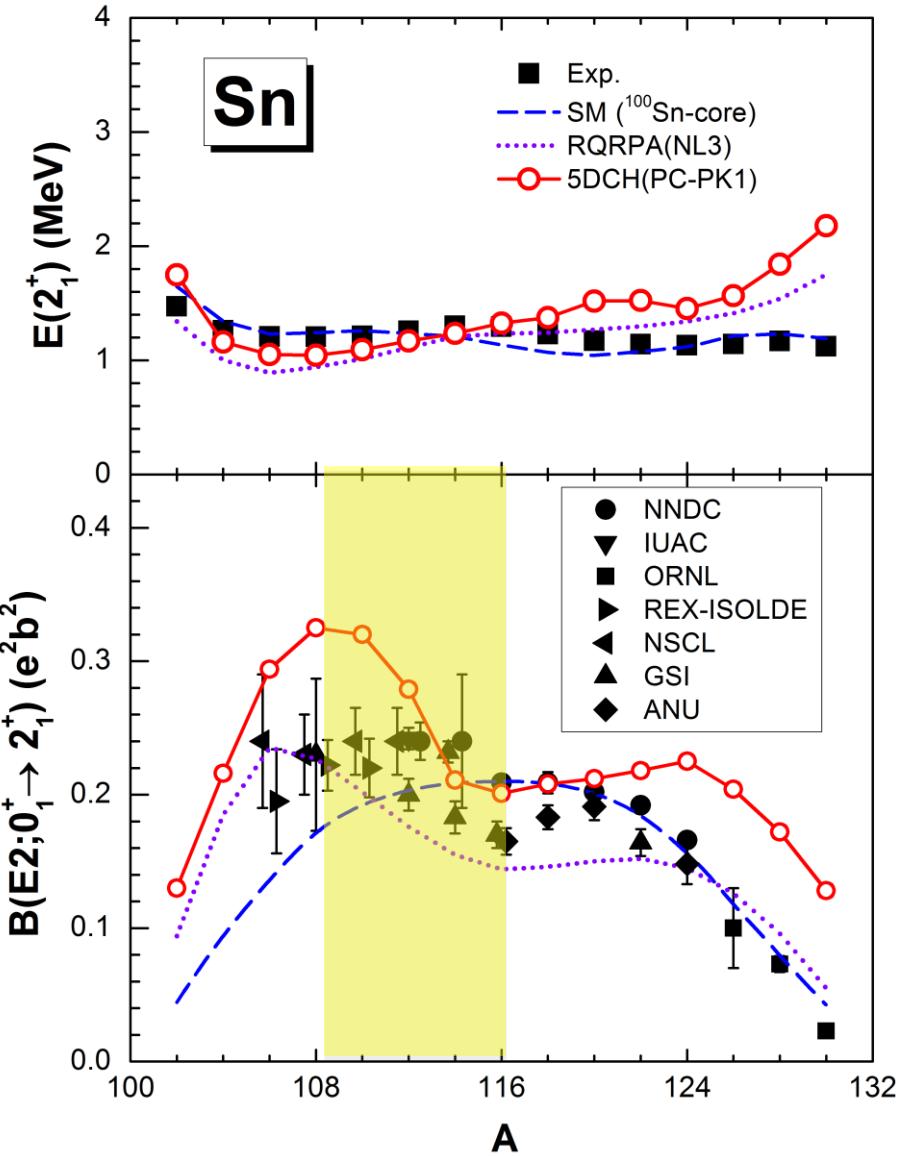
➤ Low-lying spectrum



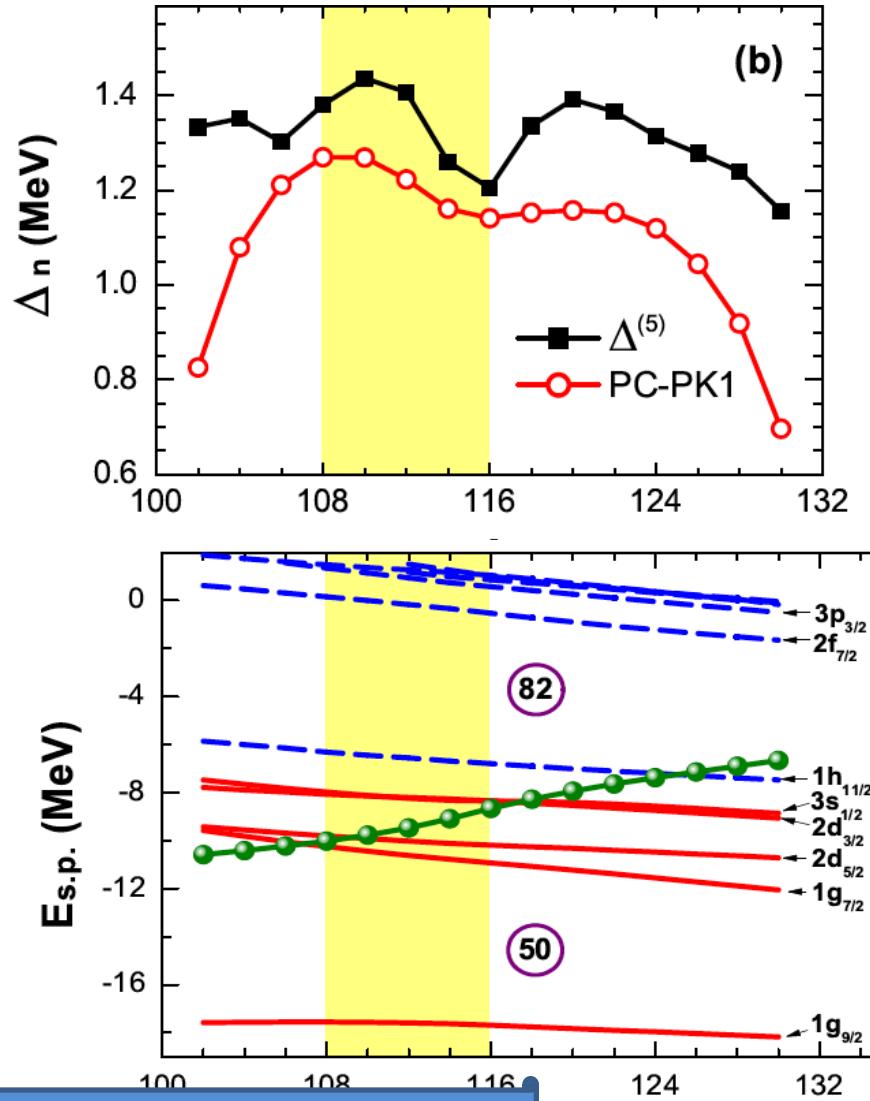
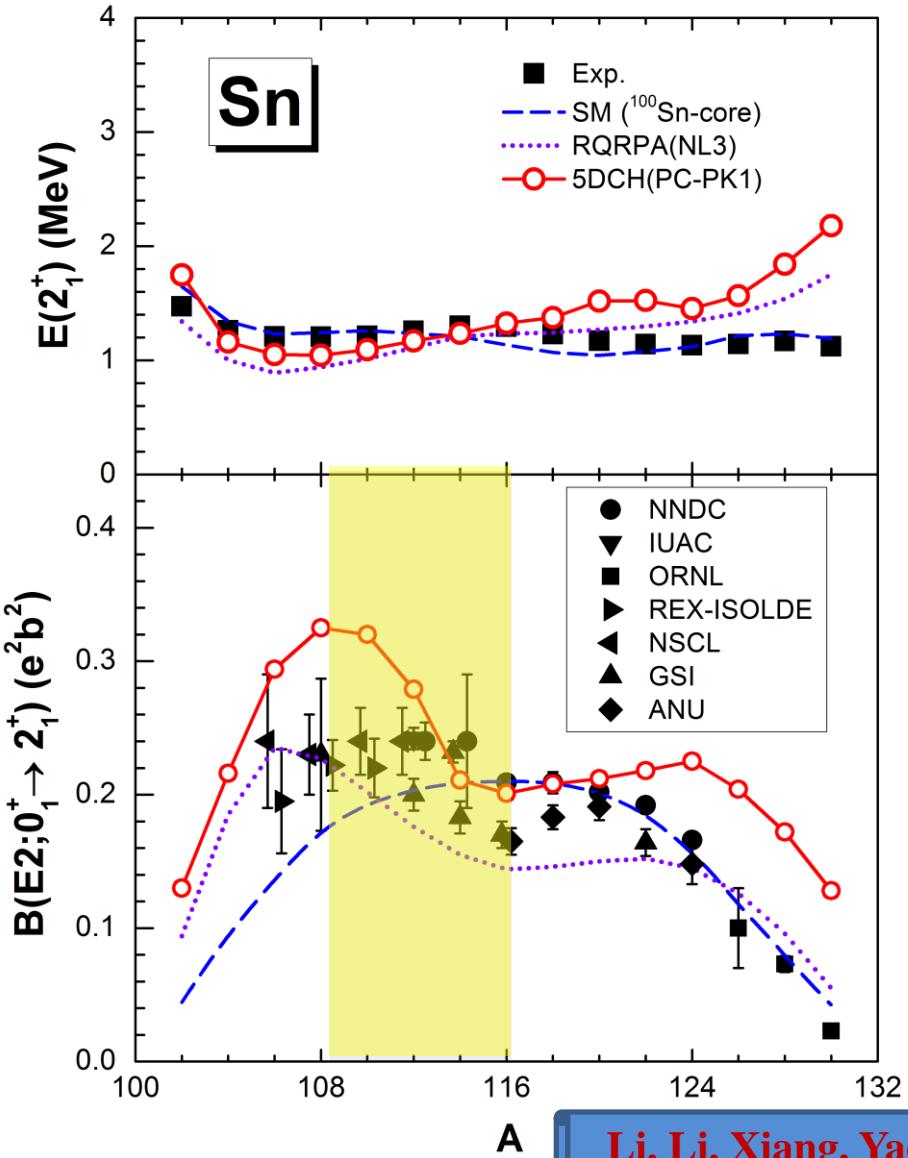
Enhanced collectivity in neutron-deficient Sn isotope



Enhanced collectivity in neutron-deficient Sn isotope



Enhanced collectivity in neutron-deficient Sn isotope



Effect of time-odd mean-field

➤ Expansion method for TV inertia parameter

$$M_{\mu\mu'}(\mathbf{q}) = \begin{pmatrix} P^* & -P \end{pmatrix}_\mu \begin{pmatrix} A & -B \\ -B^* & A^* \end{pmatrix}^{-1} \begin{pmatrix} P \\ -P^* \end{pmatrix}_{\mu'}$$

P. Ring and
P. Schuck,
*"The Nuclear
Many-Body
Problem"*

$$A_{php'h'} = (\epsilon_p - \epsilon_h) \delta_{pp'} \delta_{hh'} + V_{phh'p'}, \quad B_{php'h'} = V_{php'h'}.$$

$$\mathcal{M}^{-1} = \begin{pmatrix} A & -B \\ -B^* & A^* \end{pmatrix} = \mathcal{M}_0^{-1} + \mathcal{V}$$

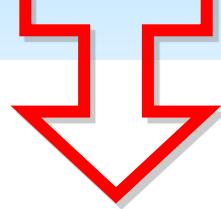
$$V_{adbc} = -\alpha_V \int [\psi_a^\dagger \boldsymbol{\alpha} \psi_b] [\psi_d^\dagger \boldsymbol{\alpha} \psi_c] d^3r$$

Effect of time-odd mean-field

➤ Expansion method for TV inertia parameter

$$\mathcal{M} = [\mathcal{M}_0^{-1} + \mathcal{V}]^{-1} = \mathcal{M}_0[1 + \mathcal{V}\mathcal{M}_0]^{-1}$$

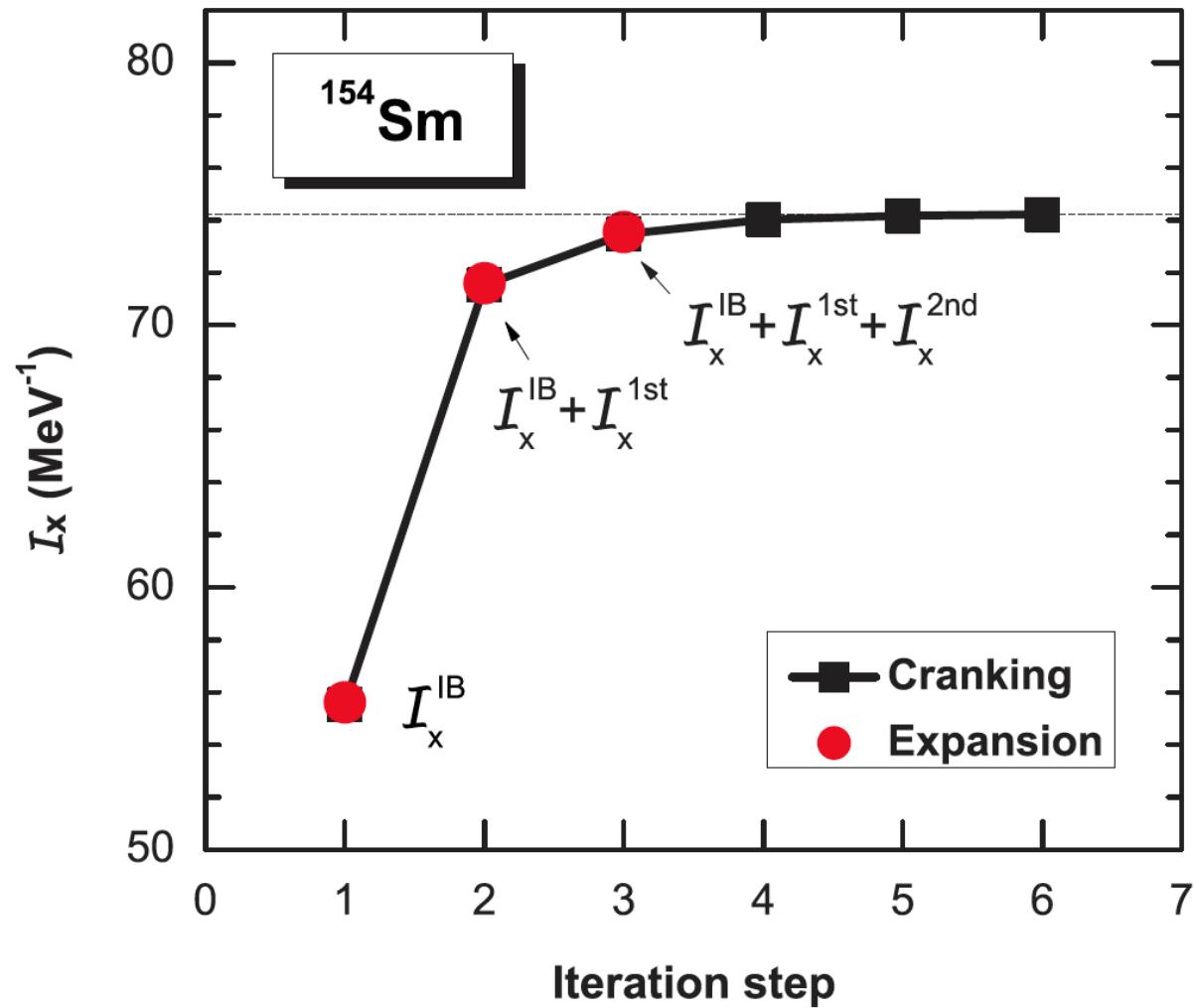
$$\mathcal{M} = \boxed{\mathcal{M}_0} - \boxed{\mathcal{M}_0\mathcal{V}\mathcal{M}_0 + \mathcal{M}_0\mathcal{V}\mathcal{M}_0\mathcal{V}\mathcal{M}_0 + \dots}$$



$$M_{\mu\mu'}^0(\mathbf{q}) = \frac{1}{\hbar^2} (P^* - P)_\mu \mathcal{M}_0 \begin{pmatrix} P \\ -P^* \end{pmatrix}_{\mu'} = \frac{2}{\hbar^2} \sum_{ph} \frac{|\hat{P}_{ph}|^2}{\epsilon_p - \epsilon_h}$$

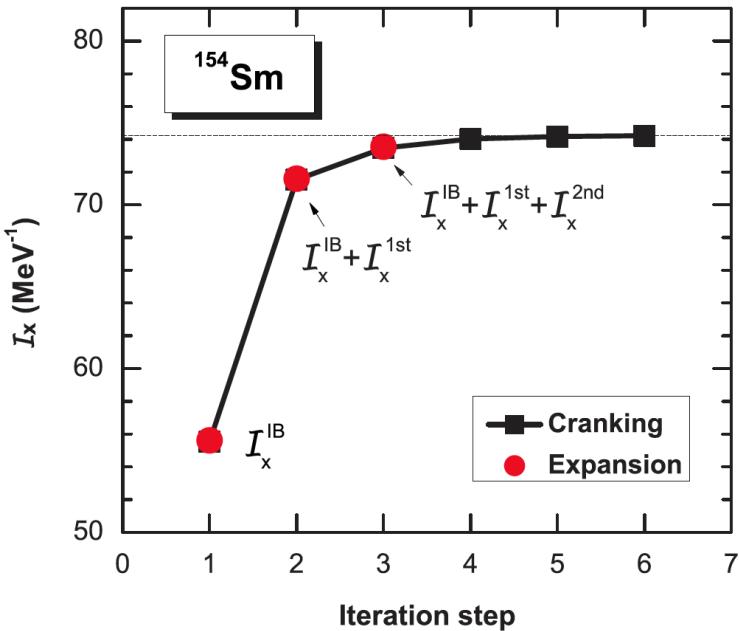
Effect of time-odd mean-field

➤ Numerical check

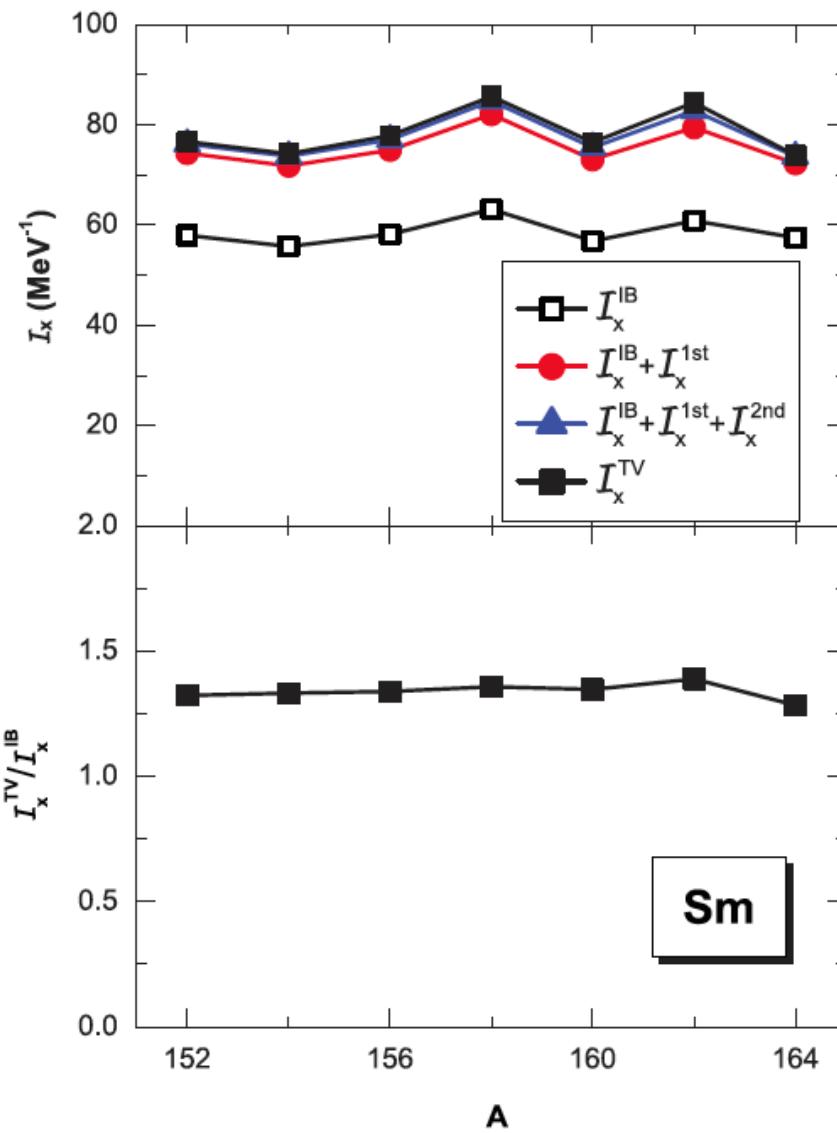


Effect of time-odd mean-field

➤ The ATDHF mass tensor



The ratio ranges ~1.
3- ~1.4



Summary

- ◆ Microscopic Collective Hamiltonian based on CEDF has been developed
- ◆ Application to the interesting topics
 - Rapid shape evolution in $N=28$ isotones
 - Enhanced collectivity in Sn isotopes
- ◆ Expansion method for Thouless-Valatin inertia parameter was introduced



Thank You !

J. Meng & ICNP group



D. Vretenar & T. Niksic



P. Ring



L. Prochniak



J. M. Yao





Shape fluctuation

$$\Delta\beta = \sqrt{\langle\beta^4\rangle - \langle\beta^2\rangle^2}/2\langle\beta\rangle,$$

$$\Delta\gamma = \sqrt{\frac{\langle\beta^6 \cos^2 3\gamma\rangle}{\langle\beta^6\rangle} - \frac{\langle\beta^3 \cos 3\gamma\rangle^2}{\langle\beta^4\rangle\langle\beta^2\rangle}} / 3 \sin(3\langle\gamma\rangle),$$

where the average values of β and γ ,

$$\langle\beta\rangle = \sqrt{\langle\beta^2\rangle},$$

$$\langle\gamma\rangle = \arccos(\langle\beta^3 \cos 3\gamma\rangle / \sqrt{\langle\beta^4\rangle\langle\beta^2\rangle})/3,$$

Observables

➤ Observables and expected values

For $E_I\alpha$, the wave function:

$$\Psi_{\alpha}^{IM}(\beta, \gamma, \Omega) = \sum_{K \in \Delta I} \psi_{\alpha K}^I(\beta, \gamma) \Phi_{MK}^I(\Omega).$$

The reduced $E2$ transition

$$B(E2; \alpha I \rightarrow \alpha' I') = \frac{1}{2I+1} |\langle \alpha' I' | \hat{\mathcal{M}}(E2) | \alpha I \rangle|^2$$

$$\hat{\mathcal{M}}(E2; \mu) = \sqrt{\frac{5}{16\pi}} \left[D_{\mu 0}^2 Q_0(\beta, \gamma) + \frac{1}{\sqrt{2}} (D_{\mu 2}^2 + D_{\mu -2}^2) Q_2(\beta, \gamma) \right]$$

Spectroscopic quadrupole moment

$$Q_{\text{spec}, \alpha I} = \frac{1}{\sqrt{2I+1}} C_{II20}^{II} \langle \alpha I | \hat{\mathcal{M}}(E2) | \alpha I \rangle.$$

The average value of β & γ

$$\begin{aligned} \langle \beta \rangle_{I\alpha} &= \sqrt{\langle \beta^2 \rangle_{I\alpha}}, \\ \langle \gamma \rangle_{I\alpha} &= \frac{1}{3} \arccos \frac{\langle \beta^3 \cos 3\gamma \rangle_{I\alpha}}{\sqrt{\langle \beta^2 \rangle_{I\alpha} \langle \beta^4 \rangle_{I\alpha}}}; \end{aligned}$$

Distribution of K component

$$N_K = 6 \int_0^{\pi/3} \int_0^\infty |\psi_{\alpha, K}^I(\beta, \gamma)|^2 \beta^4 |\sin 3\gamma| d\beta d\gamma.$$