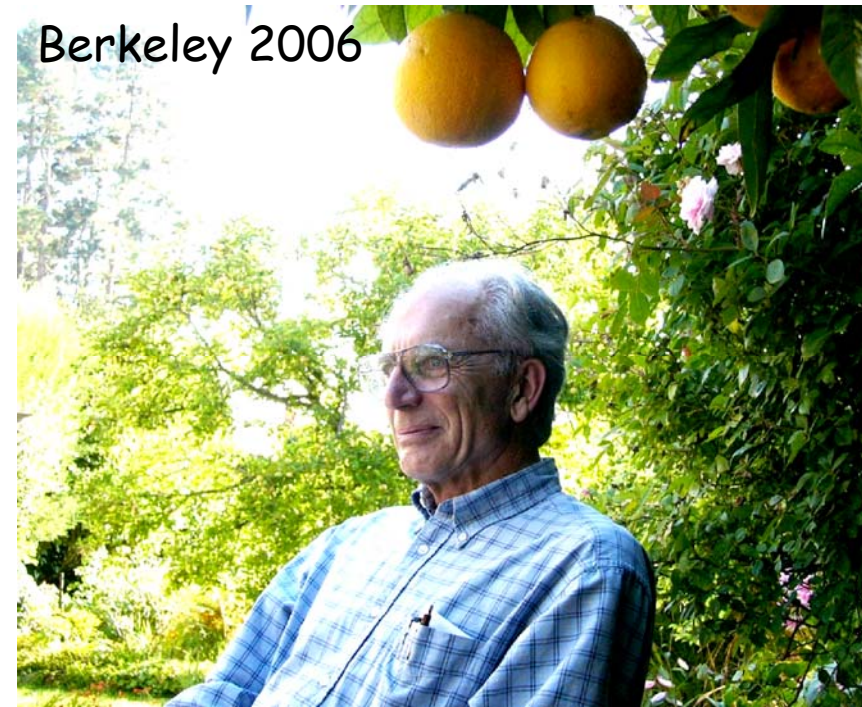


# The „Fusion by Diffusion” model revisited



Dedicated to the memory of  
Władek our Mentor and  
Friend

Kazimierz, September 2010

2003 - formulation of the „Fusion by Diffusion model”  
- FBD

W.J. Świątecki, K. Siwek-Wilczyńska, J. Wilczyński  
Acta Physica Polonica 34, 2049 (2003) (53 citations)

2005 - version II of the FBD

W.J. Świątecki, K. Siwek-Wilczyńska, J. Wilczyński  
Physical Review C 71, 014602 (2005) (61 citations)

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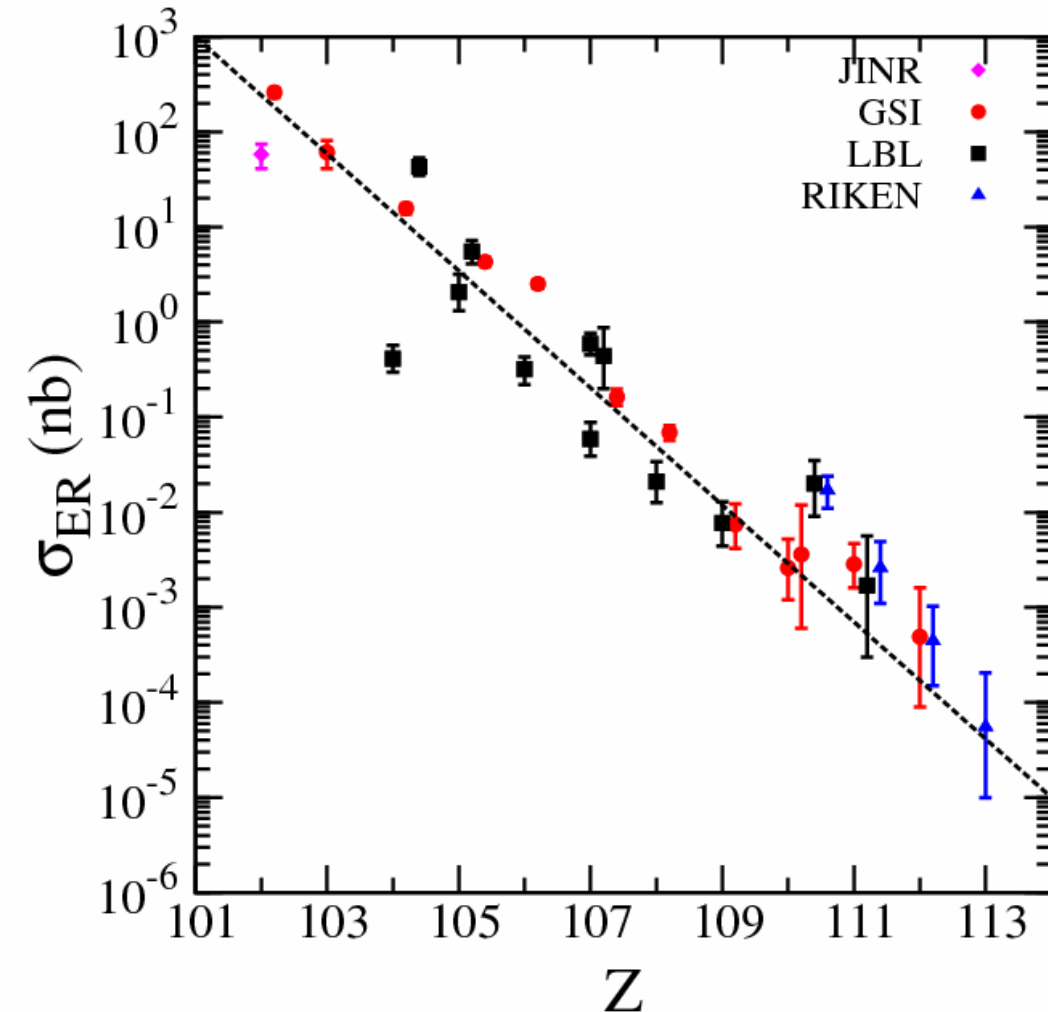
last joint paper - January 2010

W.J. Świątecki, K. Siwek-Wilczyńska, J. Wilczyński  
Physical Review C81, 019804 (2010)

Władek's legacy:

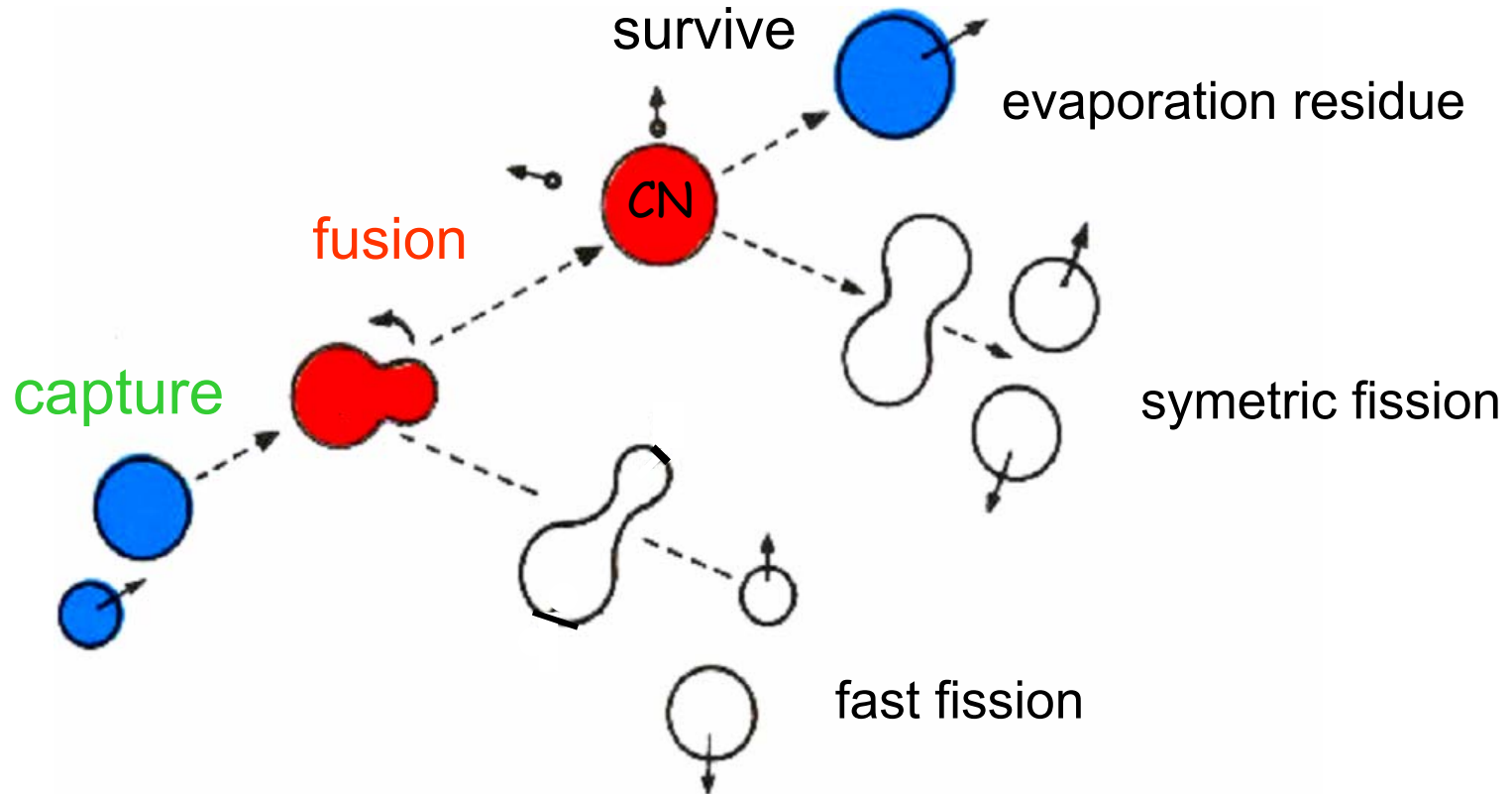
- to improve the model,
  - to involve young people into this project
- Tomek Cap**, J. Wilczyński & KSW

# Cold fusion reactions



- $102 < Z < 113$
- target:  $^{208}\text{Pb}$  or  $^{209}\text{Bi}$
- projectile:  $^{48}\text{Ca} \dots ^{70}\text{Zn}$
- $E^* = 10 - 20 \text{ MeV}$
- $10^2 < \sigma_{ER} < 10^{-5} \text{ nb}$

# A collision of two nuclei leading to the formation of a super-heavy nucleus



$$\sigma(\text{synthesis}) = \sigma(\text{capture}) \times P(\text{fusion}) \times P(\text{survive})$$

W. J. Świątecki, K. Siwek-Wilczyńska, and J. Wilczyński, *Acta Phys. Pol.* 34, 2049 (2003).

W. J. Świątecki, K. Siwek-Wilczyńska, and J. Wilczyński, *Phys. Rev. C* 71, 014602 (2005).

Original version of FBD - for central collisions only

Extended FBD - L-dependence included

$$\sigma(\text{synthesis}) = \pi \hat{\lambda}^2 \sum_{l=0}^{l_{\max}} (2l+1) P_l(\text{fusion}) P_l(\text{survive})$$

$l_{\max}$  - calculated from the capture cross section

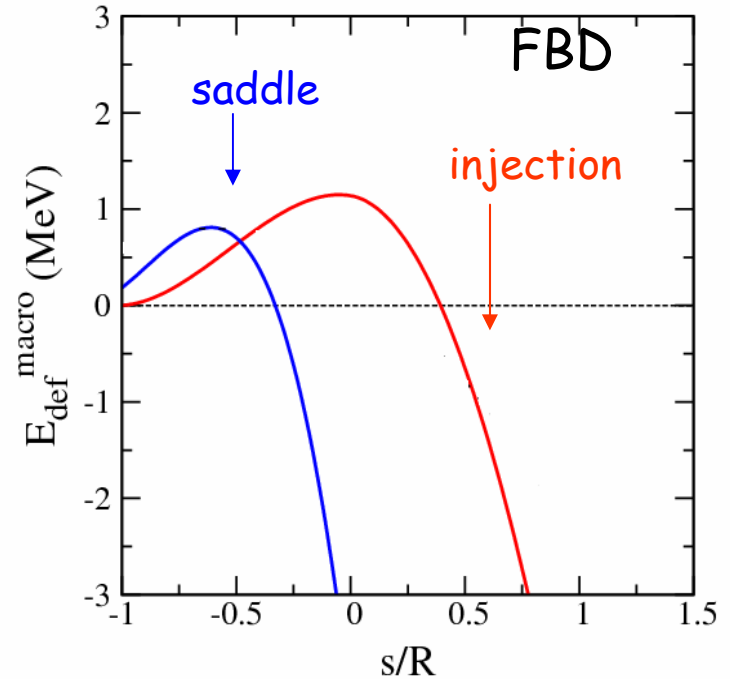
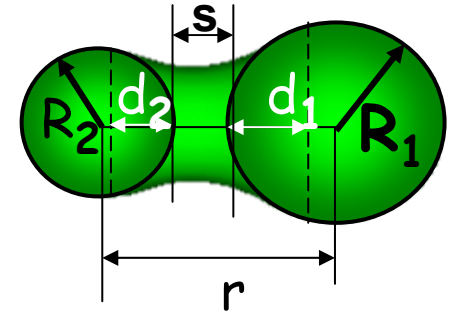
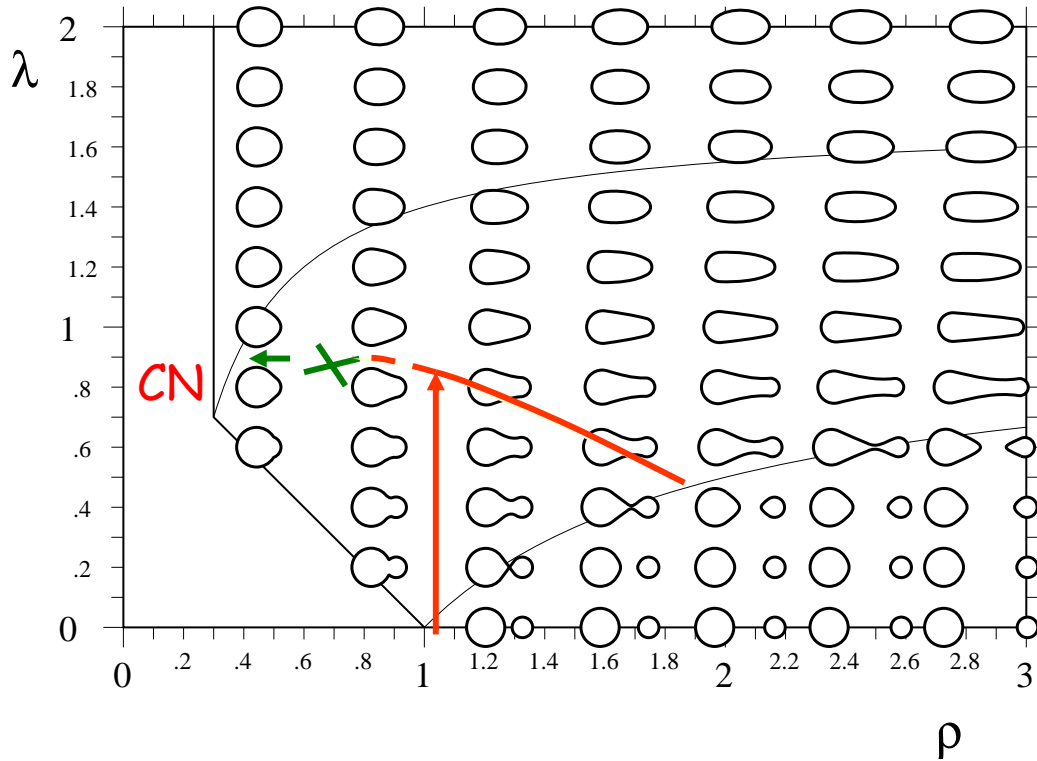
$$\sigma_{cap} = \pi \hat{\lambda}^2 (l_{\max} + 1)^2$$

$\sigma(\text{capture})$  from the „diffused barrier formula“  
(K. Siwek-Wilczyńska, J. Wilczyński Phys. Rev. C 69 (2004) 024611)  
(fusion cross section parametrized assuming the  
Gaussian shape of the barrier distribution).

# $P_f(\text{fusion})$

J. Błocki, W. J. Świątecki, Nuclear Deformation Energies, Report LBL 12811 (1982)

$\Delta=0.3$



$\lambda = (d_1 + d_2) / (R_1 + R_2)$  - neck variable

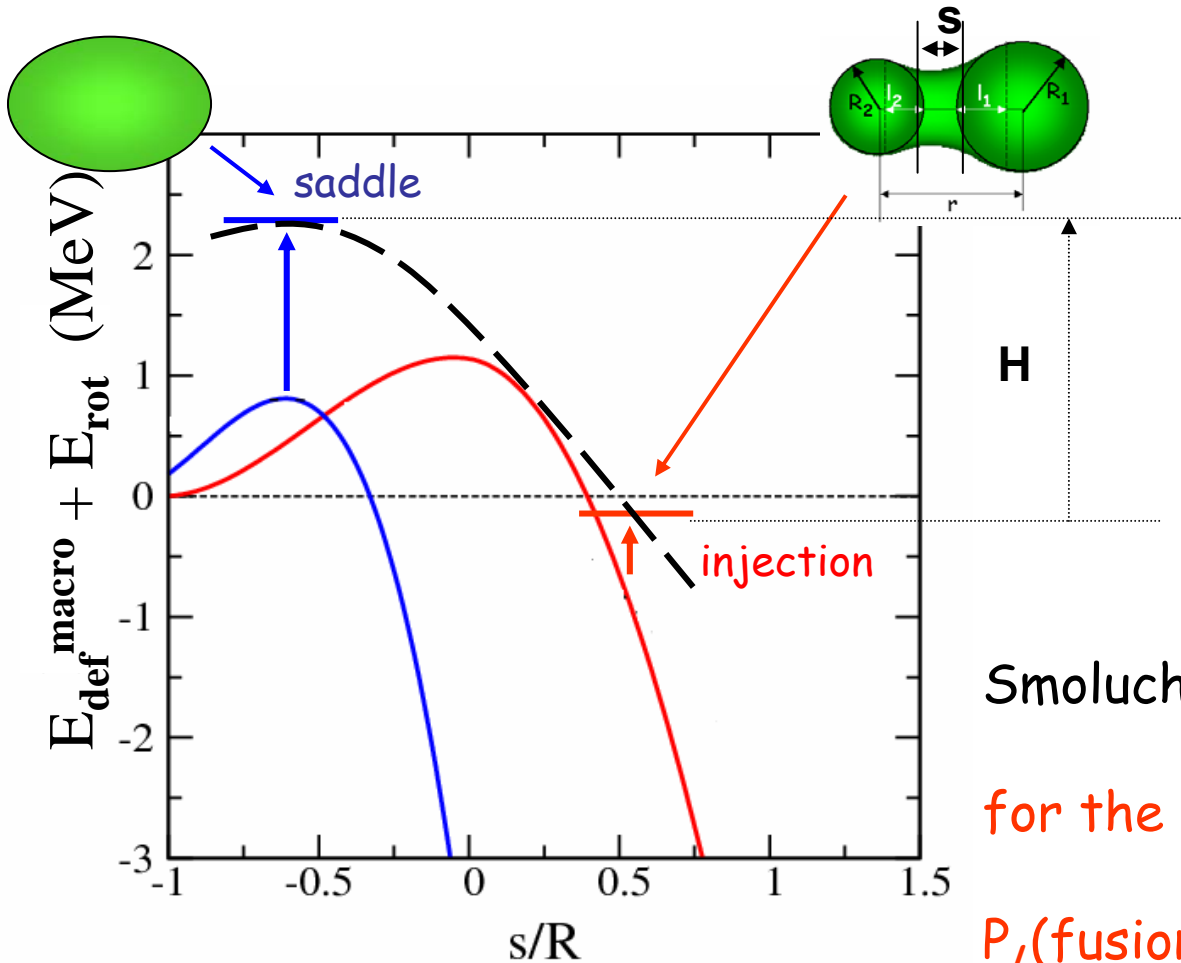
$\rho = r / (R_1 + R_2)$  - relative distance between centers of mass

$\Delta = (R_1 - R_2) / (R_1 + R_2)$  - asymmetry

Deformation energy for **symmetric** and **asymmetric** fission

$P_l$  (fusion)

EFBD



Smoluchowski Diffusion equation

for the parabolic potential

$$P_l(\text{fusion}) = \frac{1}{2}(1 - \text{erf} \sqrt{H(l)/T})$$

$P_l$  (fusion)  $\ll 1$  for heavy systems



$$P_l (\text{survival}) = \frac{\Gamma_n}{\Gamma_n + \Gamma_f} \times P_{<}$$

Probability that after 1 n emission the excitation energy is less than the threshold for second chance fission or 2n emission.

Partial widths for **emission of light particles** - Weisskopf formula

$$\Gamma_i = \frac{m_i}{\pi^2 \hbar^2} (2s_i + 1) \int_0^{E_i^{\max}} \varepsilon_i \sigma_i \frac{\rho_i(E_i^{\max} - \varepsilon_i)}{\rho(E^*)} d\varepsilon_i$$

where:  $E_i^{\max} = E^* - E_{rot}^i - B_i - V_i^C - P$  Upper limit of the final-state excitation energy after emission of a particle  $i$

$\sigma_i$  - cross section for the production of the compound nucleus in the inverse process

$m_i, s_i, \varepsilon_i$  - mass, spin and kinetic energy of the emitted particle

$\rho, \rho_i$  - level densities of the parent and daughter nuclei

The **fission** width (transition state method),  $E^* < 40 \text{ MeV}$

$$\Gamma_{fiss} = \frac{1}{2\pi} \int_0^{E_f^{\max}} \frac{\rho_{fiss}(E_f^{\max} - K)}{\rho(E^*)} dK$$

$E_f^{\max} = E^*(\text{saddle}) - E_{rot}(\text{saddle}) - P$  Upper limit of the thermal excitation energy at the saddle



The level density is calculated using the Fermi-gas-model formula

$$\rho(E) \propto \exp(2\sqrt{aE})$$

- Shell effects

included as proposed by Ignatyuk

(A.V. Ignatyuk et al., Sov. J. Nucl. Phys. 29 (1975) 255)

$$a = a_{macro} \left[ 1 + \frac{\delta_{shell}}{U} \left( 1 - e^{-U/E_d} \right) \right]$$

where:  $U$  - excitation energy,  $E_d$  - damping parameter

$\delta_{shell}$  - shell correction energy,  $\delta_{shell}$  (g.s.) (Möller et al., At. Data Nucl. Data Tables 59 (1995) 185),  
 $\delta_{shell}$  (saddle)  $\approx 0$

$$a_{macro} = 0.04543 r_0^3 A + 0.1355 r_0^2 A^{2/3} B_s^j + 0.1426 r_0 A^{1/3} B_k^j$$

$$r_0 = 1.153 \text{ fm}$$

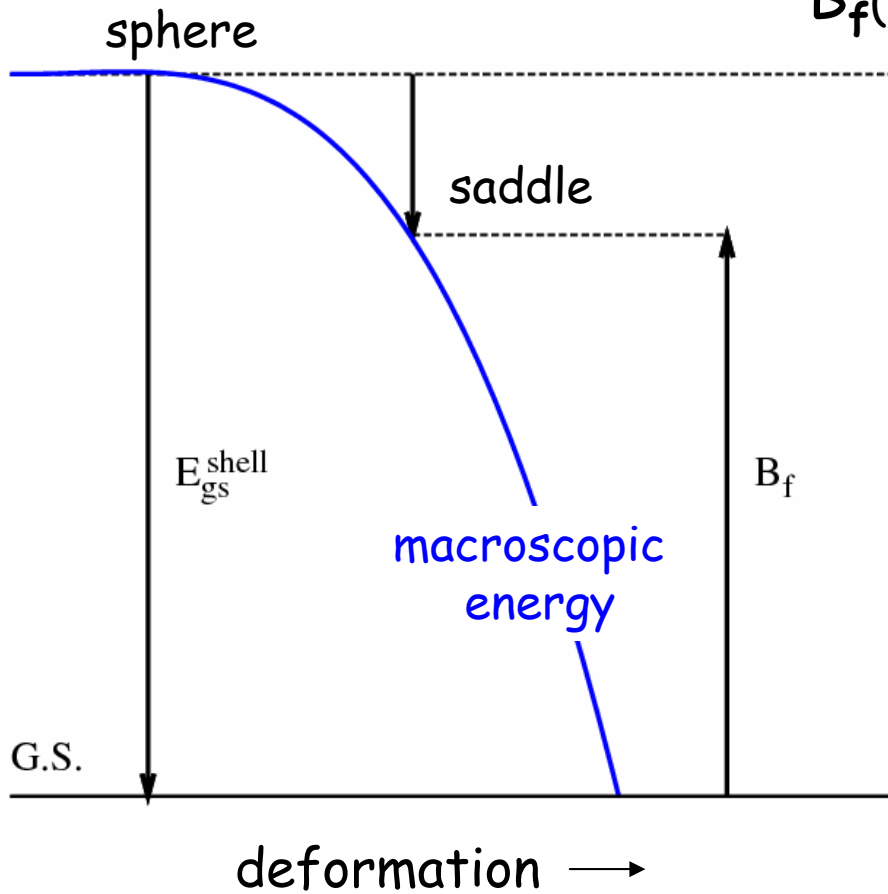
$$E_d = 18.5 \text{ MeV} \quad (\text{W. Reisdorf, Z. Phys. A. - Atoms and Nuclei 300 (1981) 227})$$

$$B_s, B_k \quad (\text{W.D. Myers and W.J. Świątecki, Ann. Phys. 84 (1974) 186})$$

# Fission barrier

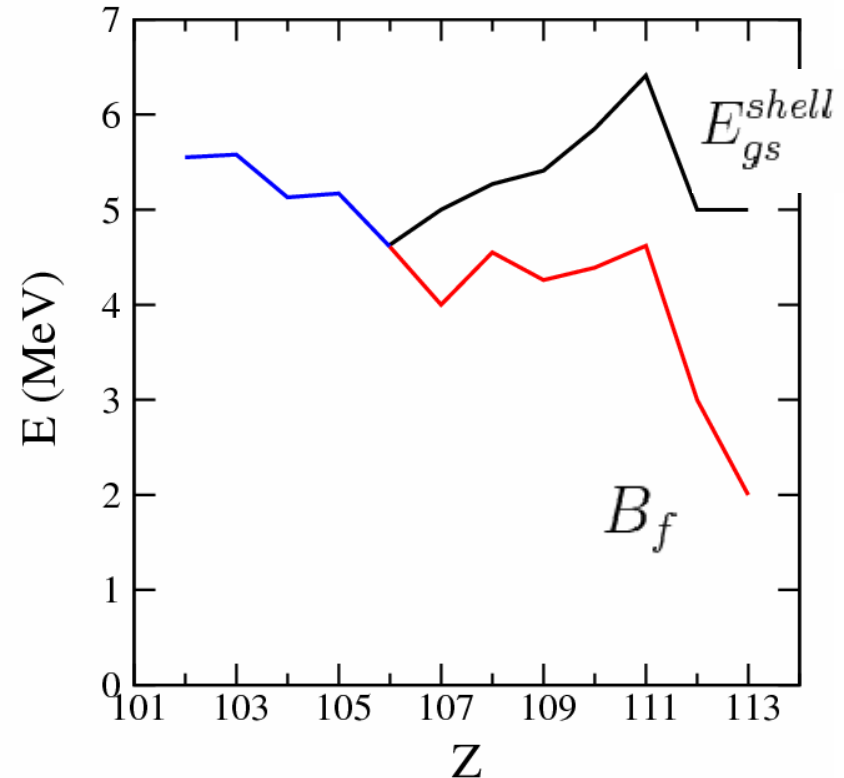
$$B_f = B_f(\text{macro}) - E_{gs}^{\text{shell}}$$

$$B_f(\text{macro}) = E_{\text{saddle}}(\text{macro}) - E_{\text{sphere}}(\text{macro})$$



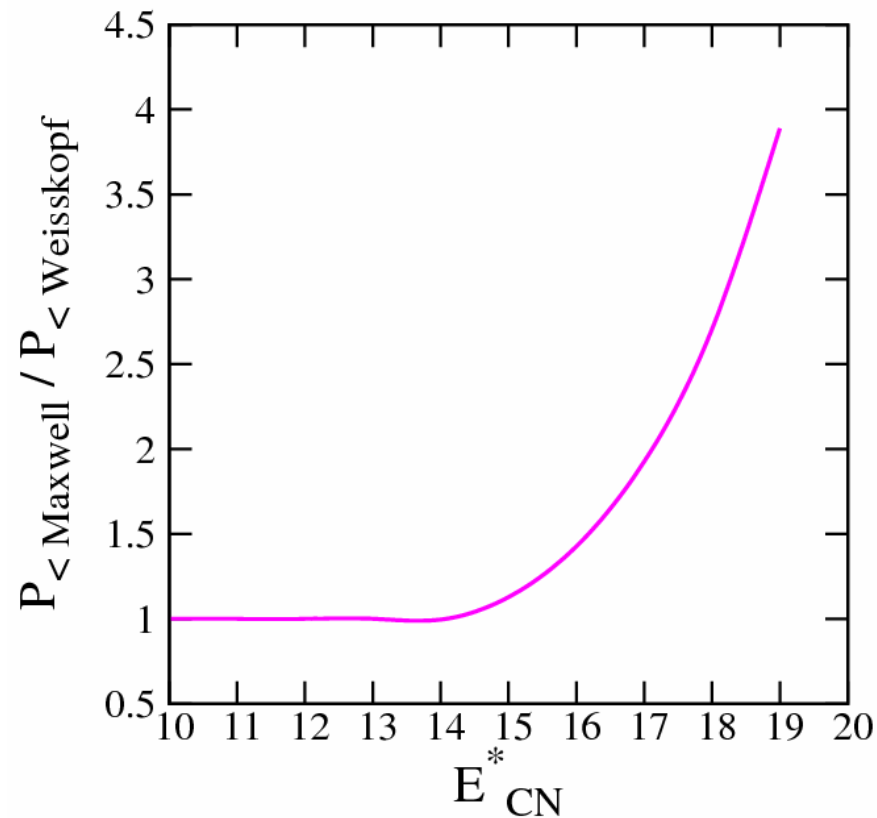
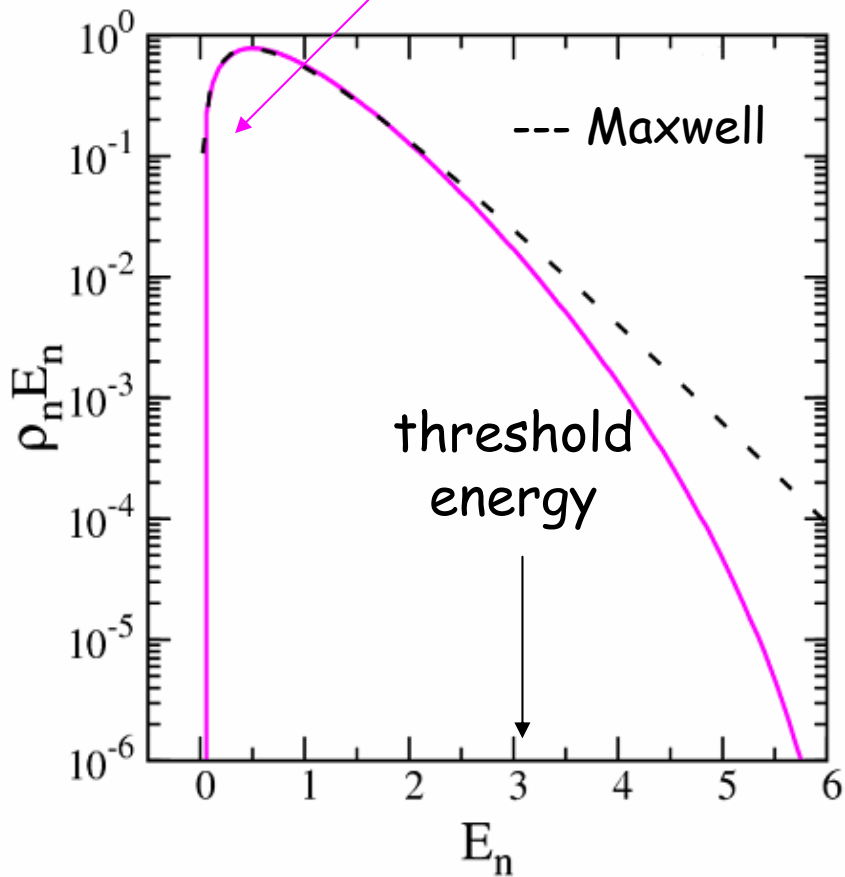
the macroscopic energy in saddle configuration  $\rightarrow$  deformation at the saddle point (A. Sobiczewski et al. private communication)

For nuclei with  $Z > 106$  there is no macroscopic barrier



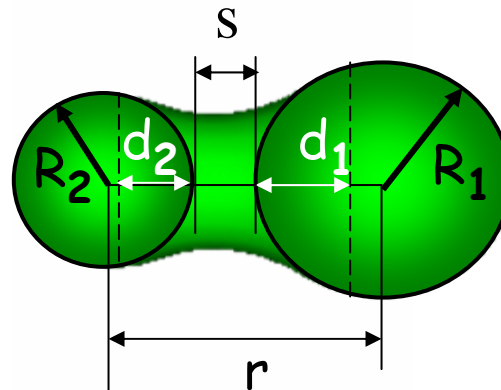
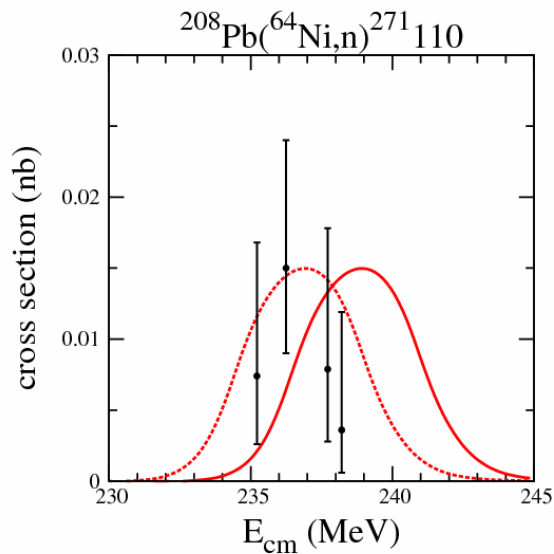
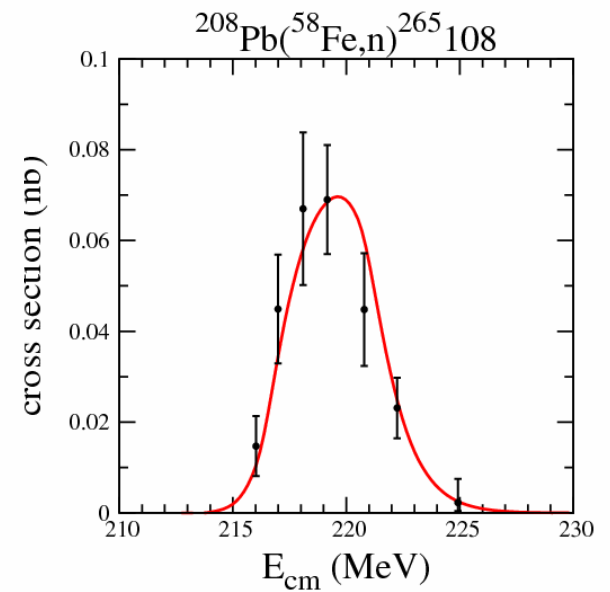
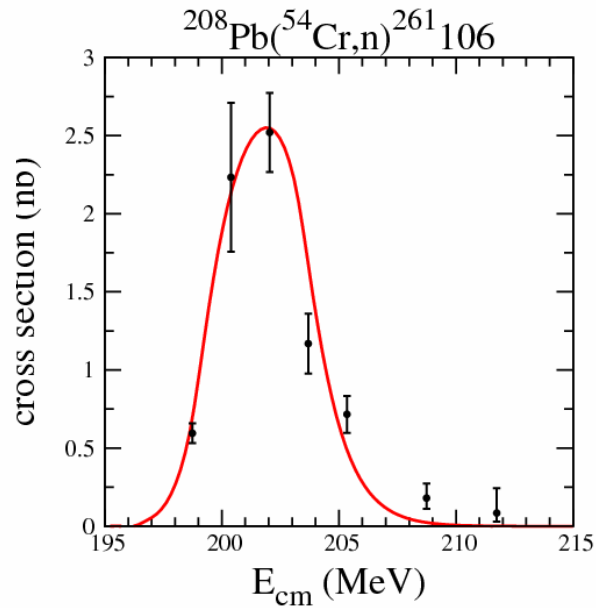
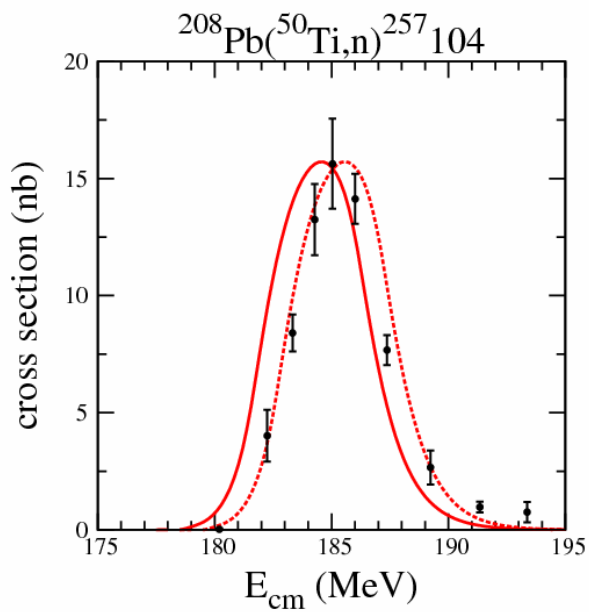
# $P_{<}$ Shape of the kinetic energy spectrum

$$\int_0^{E_n^{max}} \rho_n(A-1, \epsilon^*) (E_n^{max} - \epsilon^*) d\epsilon^*$$



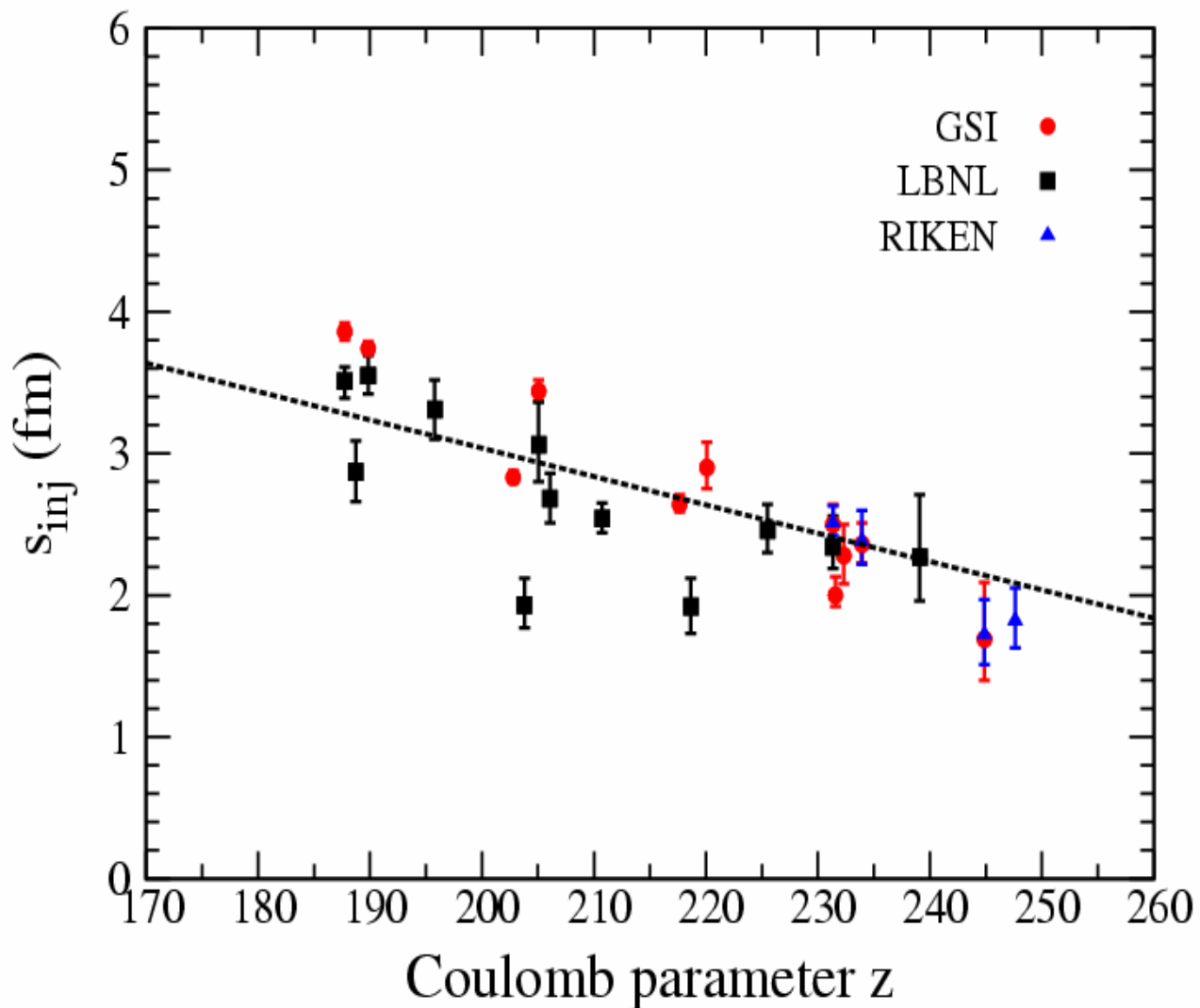
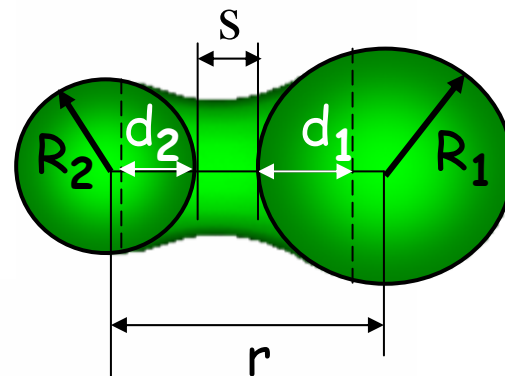
Probability  $P_{<}$

# Fit to 27 excitation functions



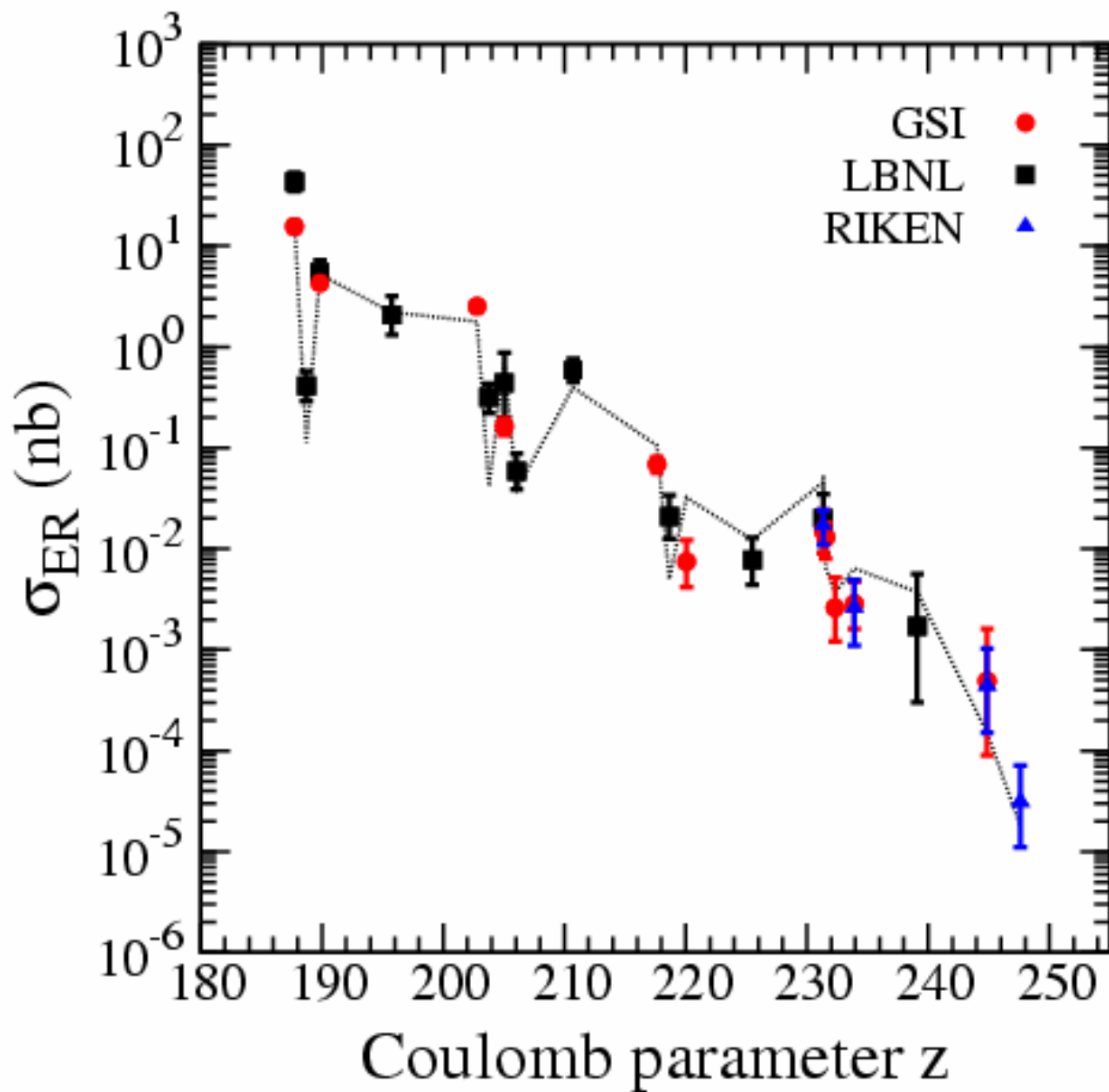
$S_{\text{inj}}$

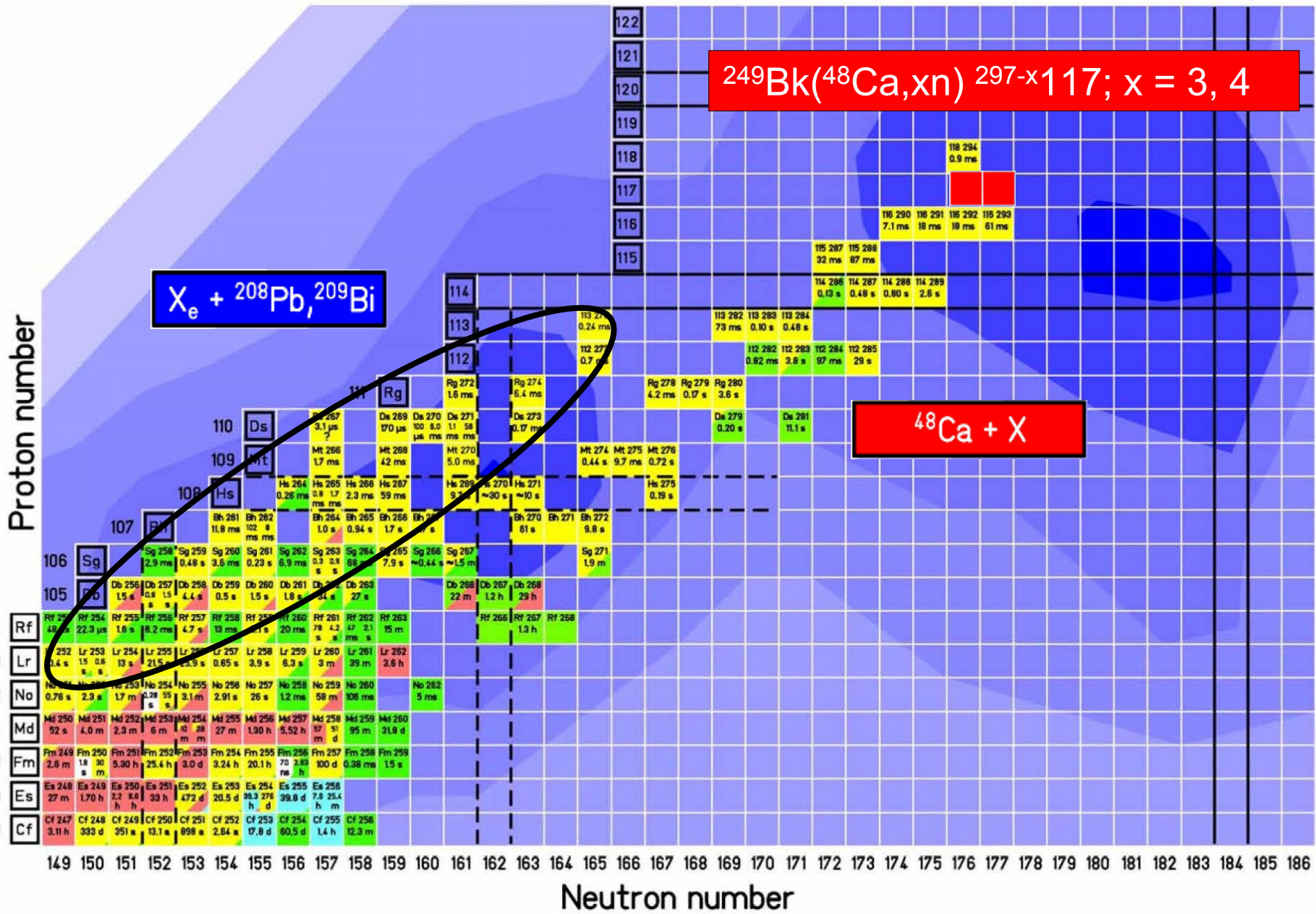
# Parametrization of $s_{inj}$



$$z = \frac{Z_p Z_T}{A_P^{1/3} + A_T^{1/3}}$$

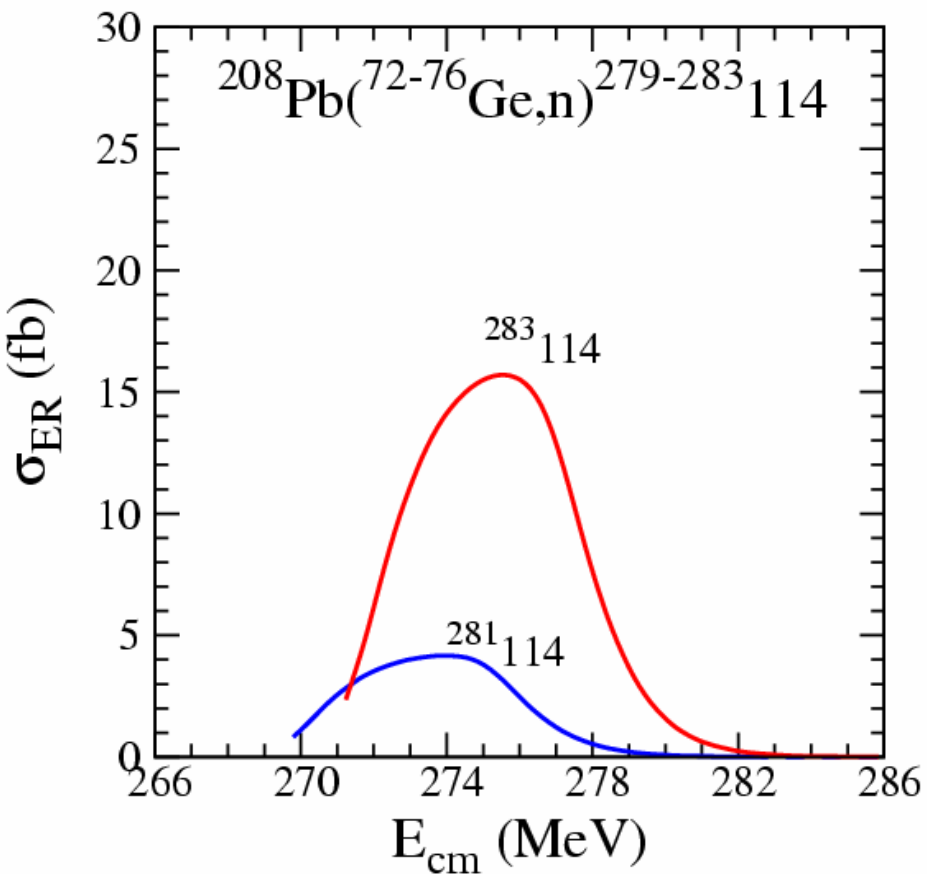
# results





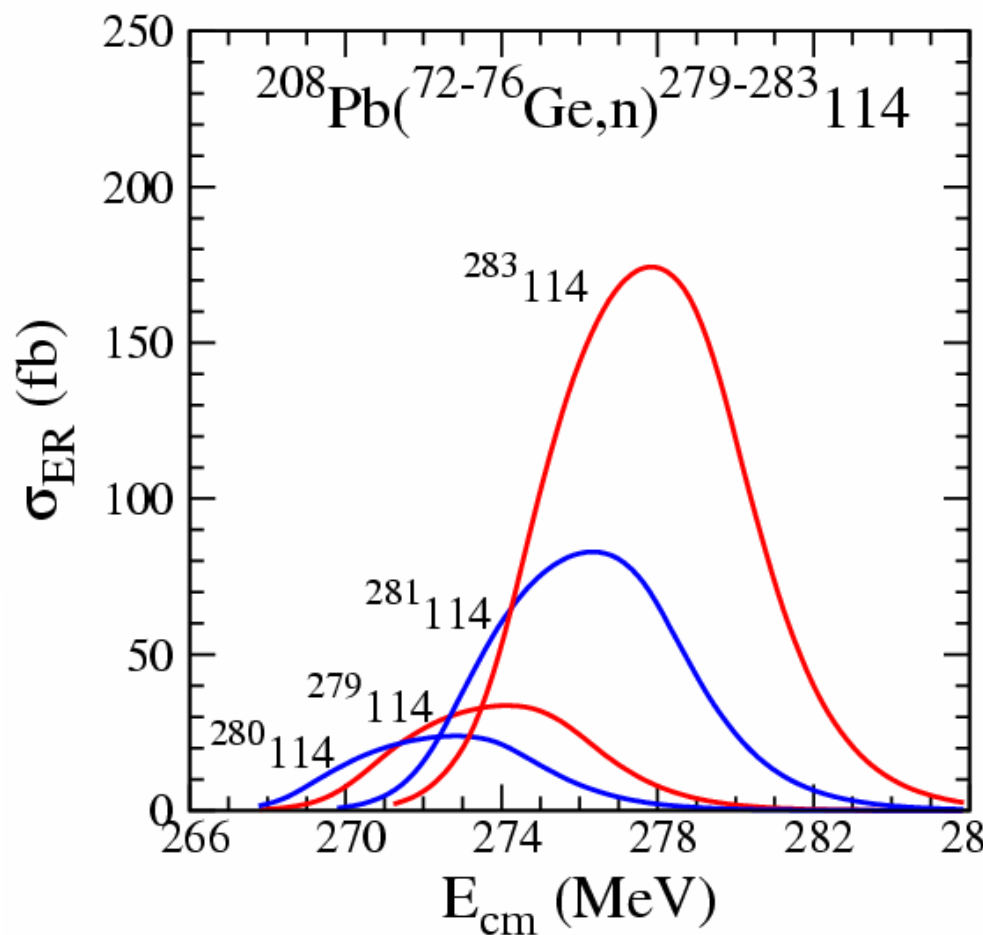


## Deformed nuclei

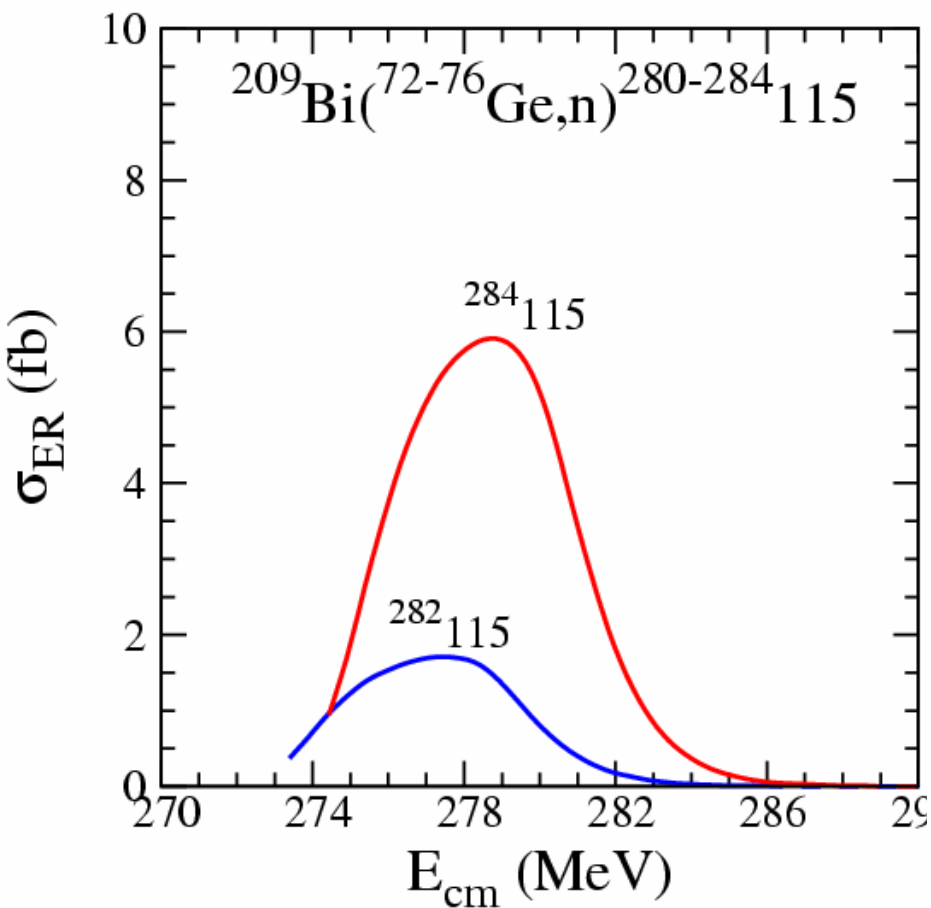


## Predictions for $Z = 114$

### Spherical nuclei

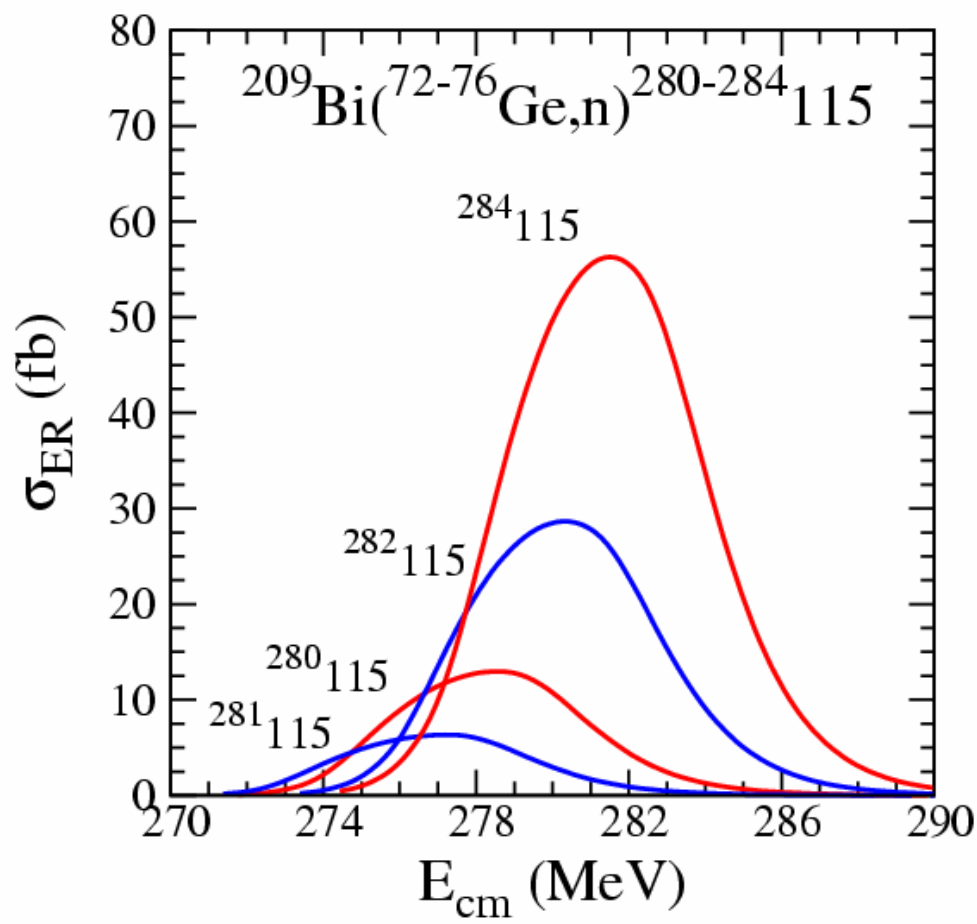


## Deformed nuclei



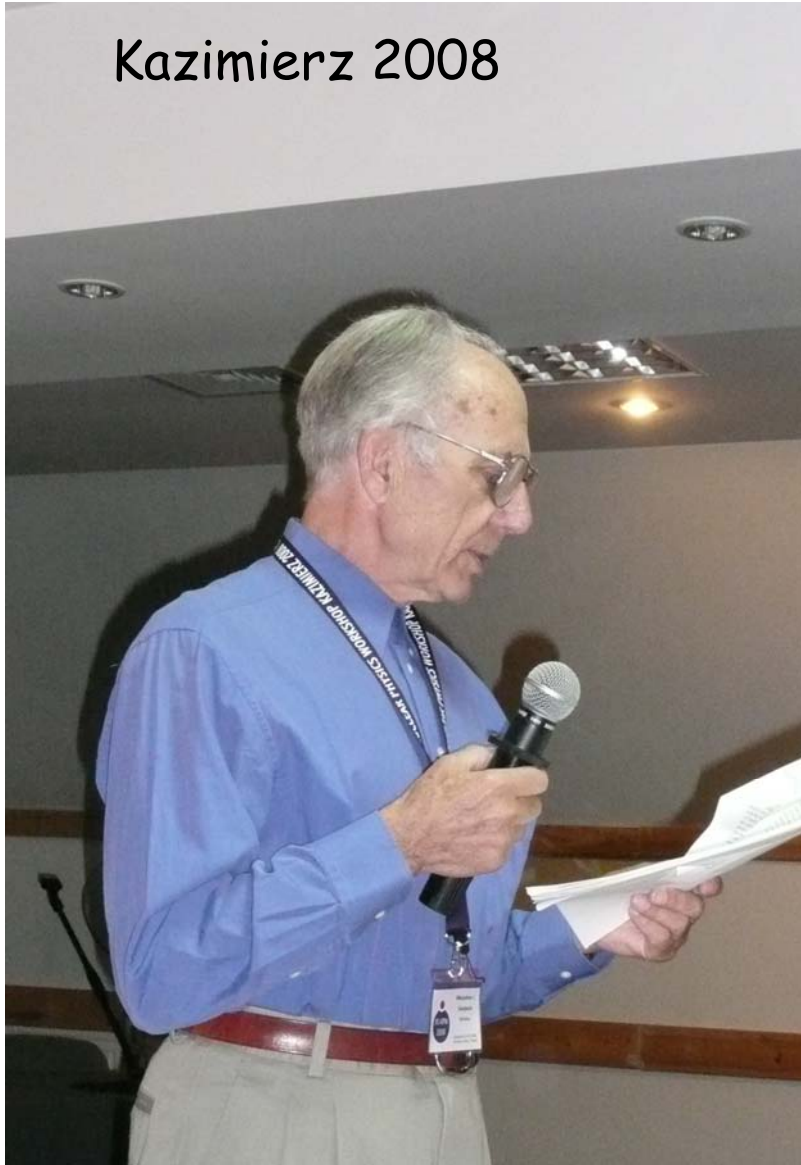
## Predictions for $Z = 115$

### Spherical nuclei



# Summary

Kazimierz 2008



- Extended version of the FBD model was presented.
- Good agreement with existing experimental data (entrance channel effects reproduced).
- Predictions for the production of new isotopes of elements 114 and 115 were presented.

Kazimierz 2010



WŁADEK ŚWIĄTECKI  
NUCLEAR THEORY LOUNGE



September 10, 2010