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# Chiral symmetry in atomic nuclei

Jie Meng 孟杰



北京大学物理学院  
School of Physics  
Peking University (PKU)



北京航空航天大学物理与核能工程学院  
School of Physics and Nuclear Energy Engineering  
Beihang University (BUAA)

# Outline

- Introduction
- Chirality in atomic nuclei
- Experimental progress
- Theoretical progress
- Summary and perspectives

# Chiral symmetries exist commonly in nature

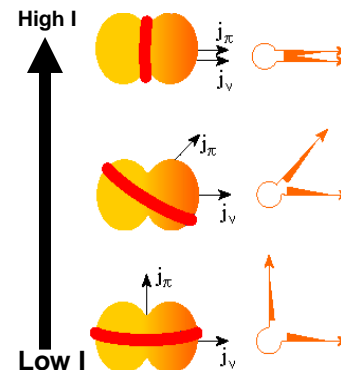
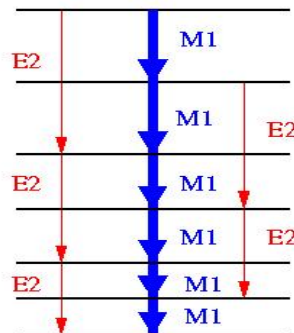
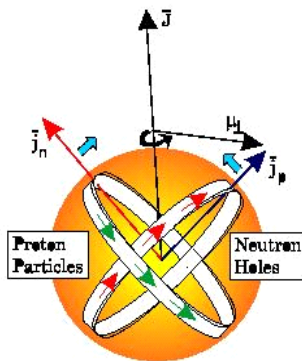
- ❑ Macroscopic spirals of snail shells and the Human hands...
- ❑ In geometry, a figure is chiral if it is not identical to its mirror image, or it cannot be mapped onto its mirror image by rotations and translations alone
- ❑ Particle physics, chirality is a dynamic property distinguishing between the parallel and anti-parallel orientations of the intrinsic spin with respect to the momentum of the massless particle.
- ❑ Chemistry, the study of chirality is a very active in inorganic, organic, physical, biochemistry and supramolecular chemistry.
- ❑ Nobel Prize in Chemistry for 2001: the development of catalytic asymmetric synthesis
- ❑ ...



# Chirality in atomic nucleus: short history

Magnetic Rotation: Hot topics of high spin physics in 1990s

- Near spherical nucleus
- $\Delta I=1$  bands with strong BM1
- Orientation of valence p and n provides spin





# $\Delta I = 1$ Enhanced magnetic dipole transition

508

M. Neffgen et al. / Nuclear Physics A 595 (1995) 499-512

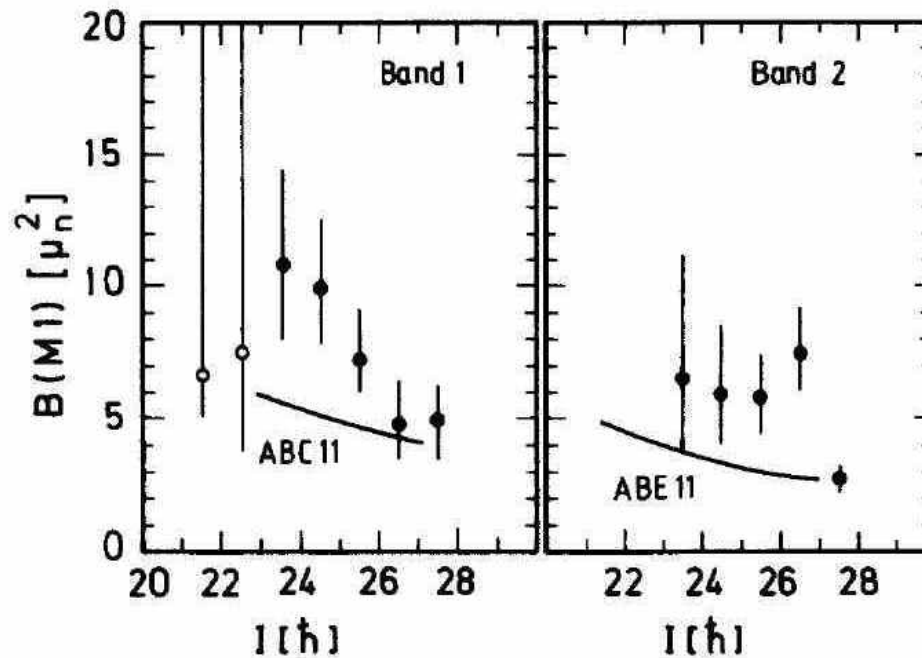


Fig. 5. Experimental (points) and calculated (lines) reduced magnetic dipole transition probabilities for bands 1 and 2 in  $^{199}\text{Pb}$  as a function of spin. Open circles: transitions in the band-crossing region.

How does  $B(M1)$  change with spin  $I$  ?



# Interpretation and quality of the tilted axis cranking approximation

Stefan Frauendorf<sup>1</sup>, Jie Meng<sup>1,2,\*</sup>

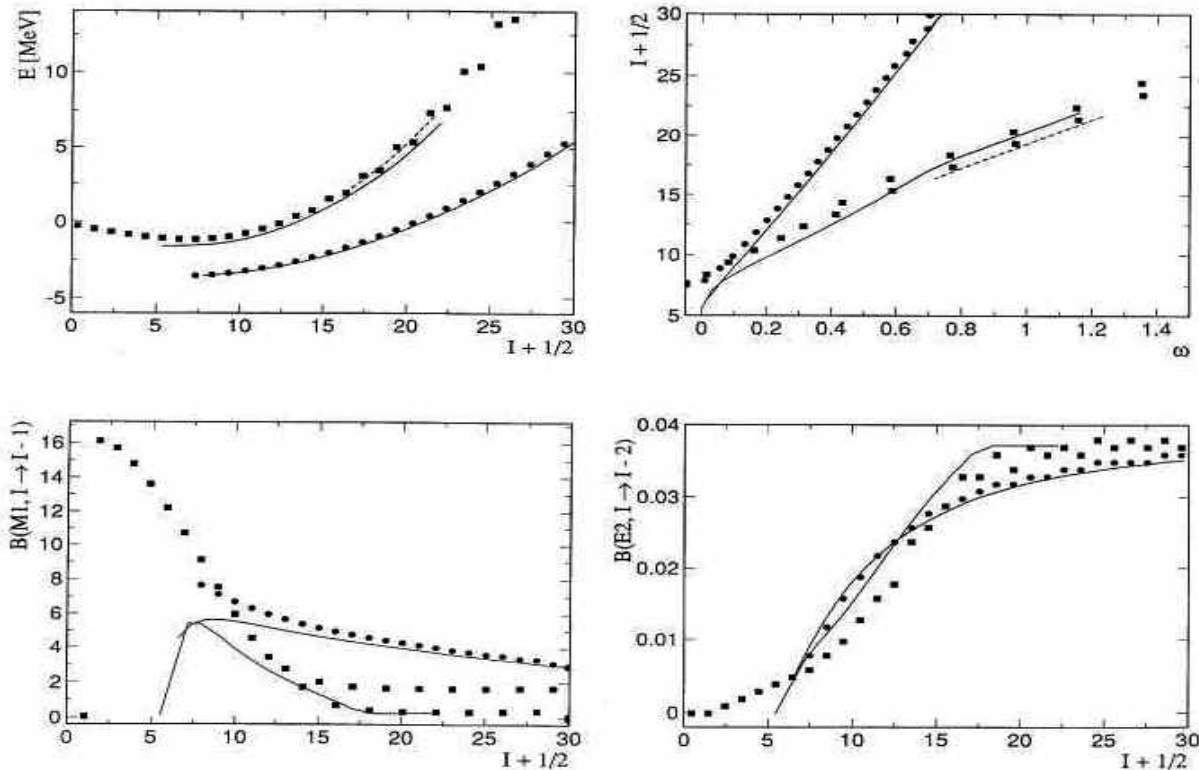
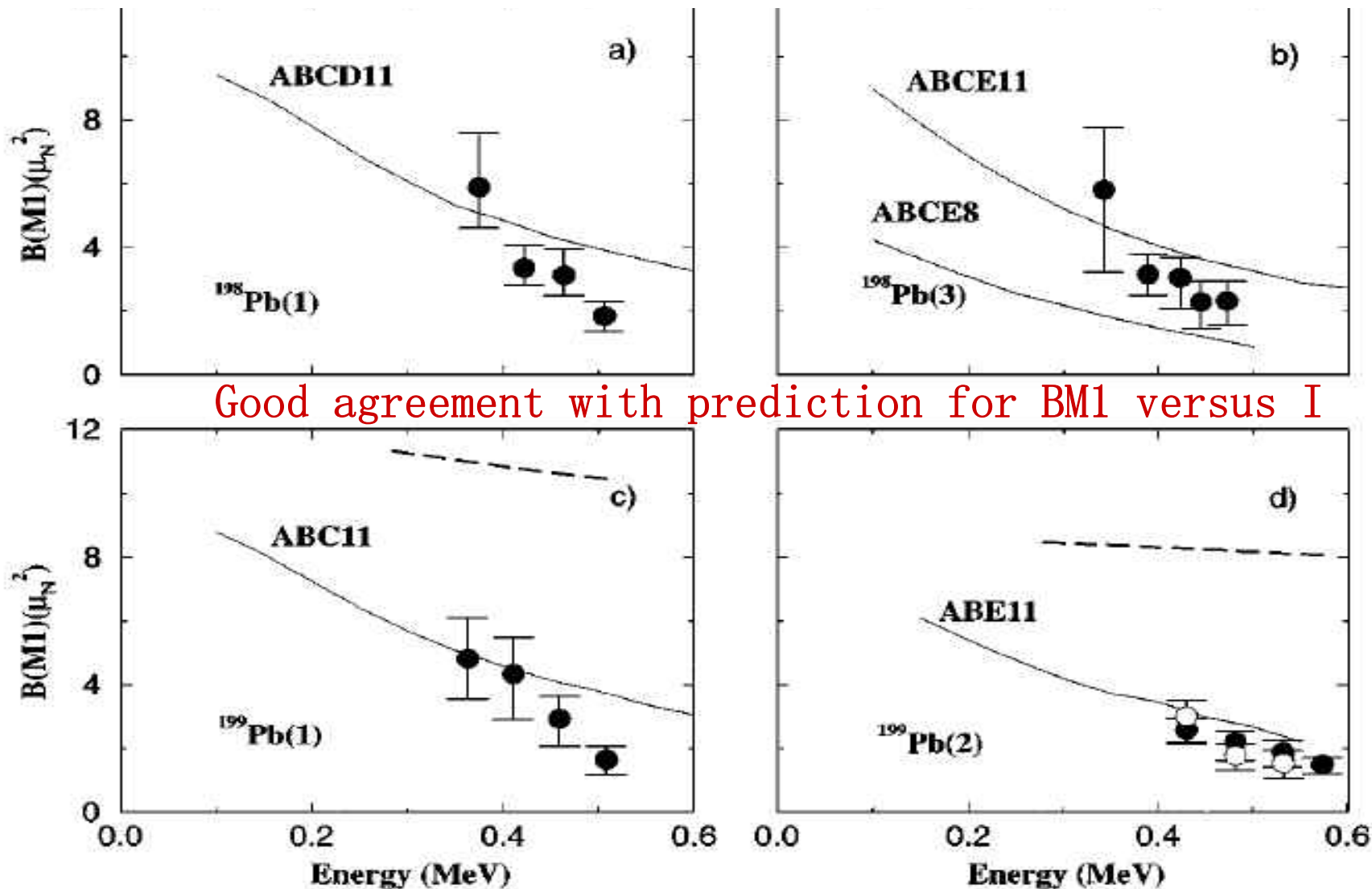


Fig. 8. Energy, angular momentum,  $B(M1)$  and  $B(E2)$  values for lowest band of the combination of a proton RAL hole with a neutron DAL hole. Circles: PRM  $C = 0.25$  MeV, squares: PRM  $C = 0.10$  MeV, full lines: TAC, dashed lines: PAC signature.

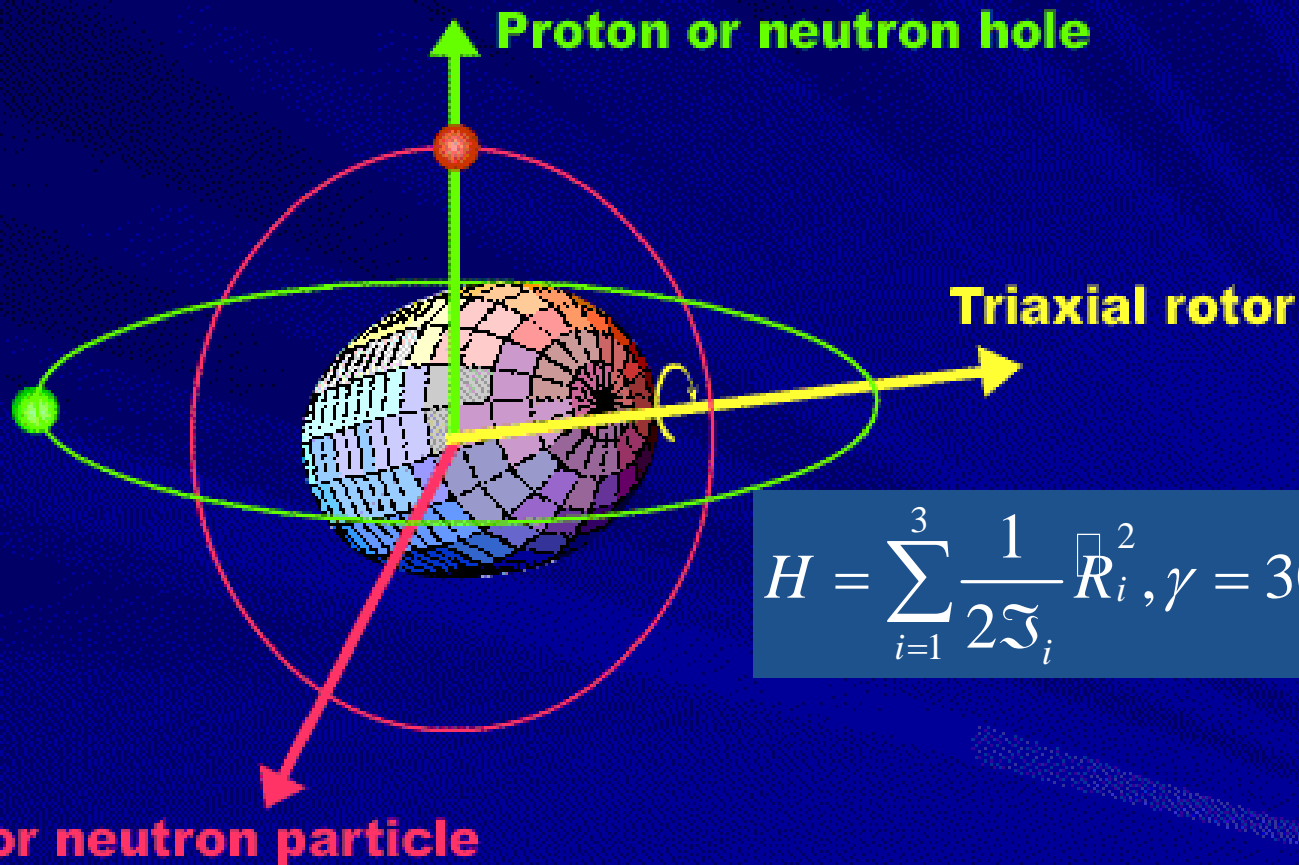
Evidence for "Magnetic Rotation" in Nuclei: Lifetimes of States  
in the M1 bands of  $^{198,199}\text{Pb}$





# Generalizing to triaxial deformed case

$$H_{sp} = -\frac{1}{2}C \left\{ \cos \gamma \left\{ j_3^2 - \frac{j(j+1)}{3} \right\} + \frac{\sin \gamma}{2\sqrt{3}} (j_+^2 + j_-^2) \right\}$$



$$H = \sum_{i=1}^3 \frac{1}{2\mathfrak{I}_i} R_i^2, \gamma = 30^\circ$$

$$H_{sp} = \frac{1}{2}C \left\{ \cos \gamma \left\{ j_3^2 - \frac{j(j+1)}{3} \right\} + \frac{\sin \gamma}{2\sqrt{3}} (j_+^2 + j_-^2) \right\}$$

# Chiral symmetry in atomic nuclei



Nuclear Physics A 617 (1997) 131-147

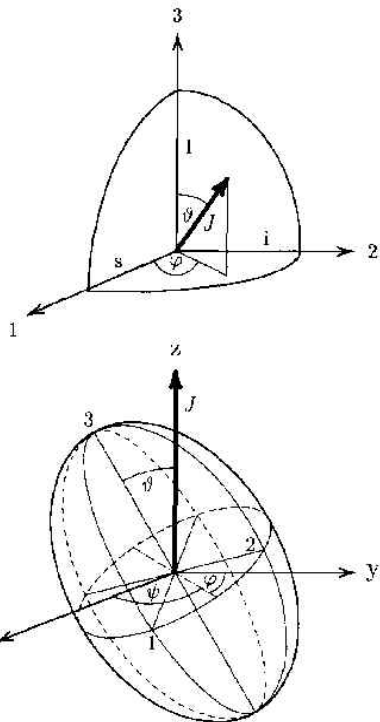
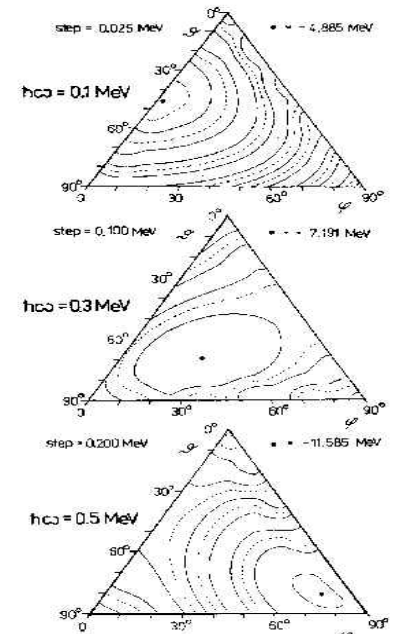
NUCLE  
PHYSIC

## Tilted rotation of triaxial nuclei

S. Frauendorf, Jie Meng<sup>1</sup>

Institut für Kern- und Hadronenphysik, Forschungszentrum Rossendorf e.V.,  
PF 510119, 01314 Dresden, Germany

S. Frauendorf, J. Meng / Nuclear Physics A 617 (1997) 131-147



### Abstract

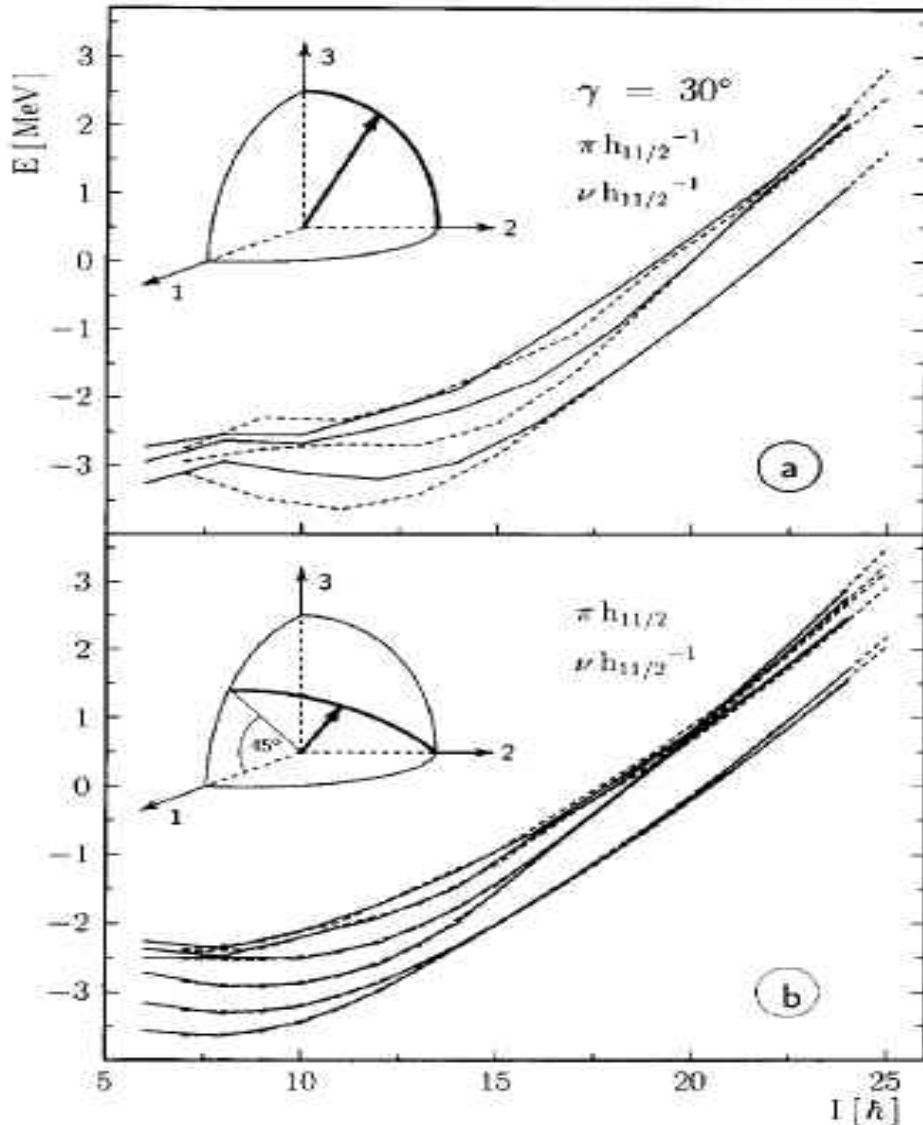
The Tilted Axis Cranking theory is applied to the model of two particles coupled to a triaxial rotor. Comparing with the exact quantal solutions, the interpretation and quality of the mean field approximation is studied. Conditions are discussed when the axis of rotation lies inside or outside the principal planes of the triaxial density distribution. The planar solutions represent  $\Delta I = 1$  bands, whereas the aplanar solutions represent pairs of identical  $\Delta I = 1$  bands with the same parity. The two bands differ by the chirality of the principal axes with respect to the angular momentum vector. The transition from planar to chiral solutions is evident in both the quantal and the mean field calculations. Its physical origin is discussed. © 1997 Elsevier Science B.V.

PACS: ...

Keywords: Tilted axis cranking; Triaxiality; Chirality

# Chiral doublets bands: experimental signal

S. Frauendorf, J. Meng / Nuclear Physics A 617 (1997) 131-147

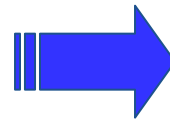
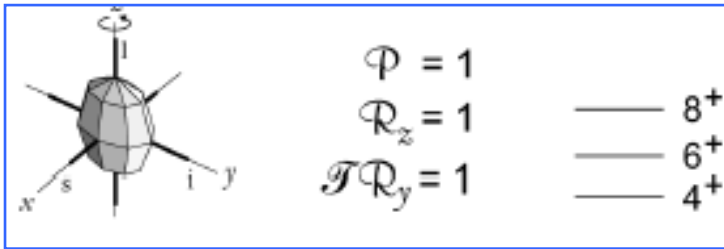


**hh+triaxial rotor=MR**

**ph+triaxial rotor=Chiral doublets**

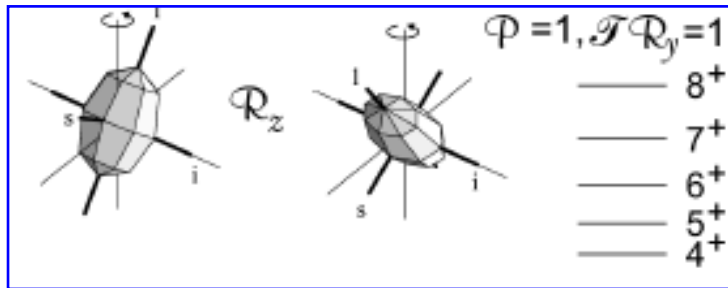
# Chiral doublets bands, Why ?

## 1. J lies along a principal axis

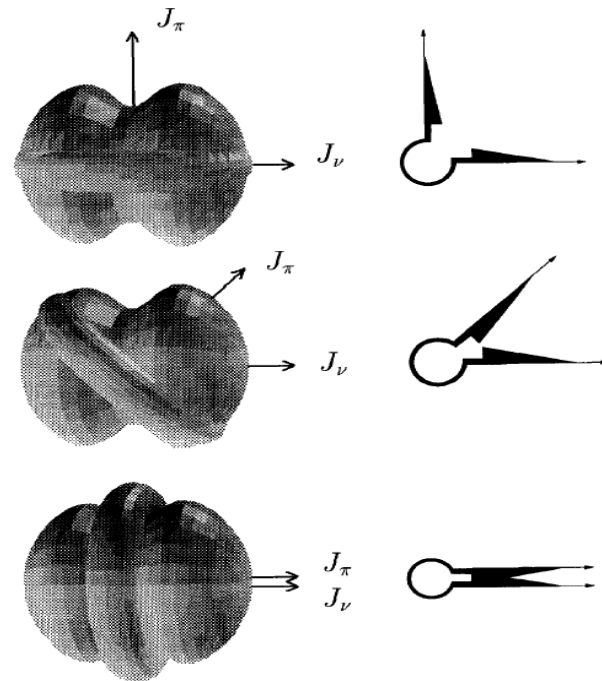


$\Delta I = 2$  rotational bands

## 2. J lies in a principal plane

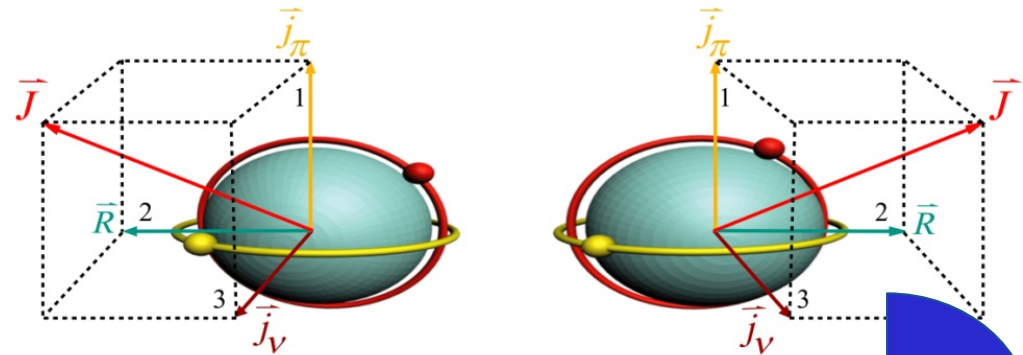
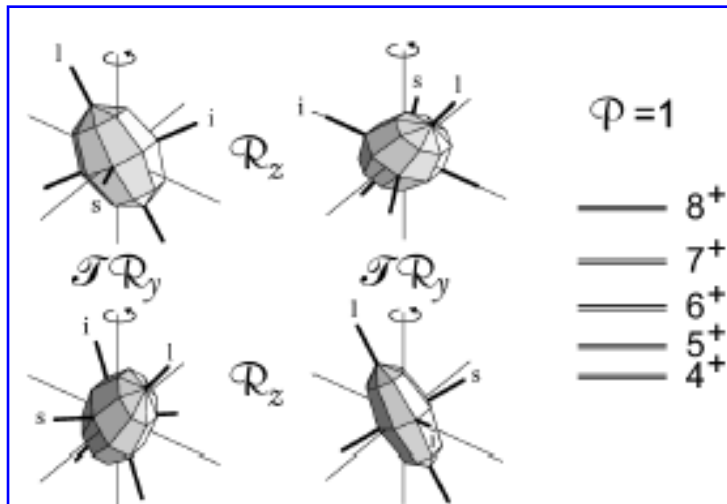


$\Delta I = 1$  magnetic  
 rotational bands  
**Shears bands**



# Chiral doublets bands, Why ?

## 3. J does not lie in any of principal plane



Chiral bands

Experimental signature:  
 Degenerate pairs of  $\Delta I=1$  bands

$$\Xi = T P_y(\pi)$$

# Spontaneous Symmetry Breaking of Chiral symmetry

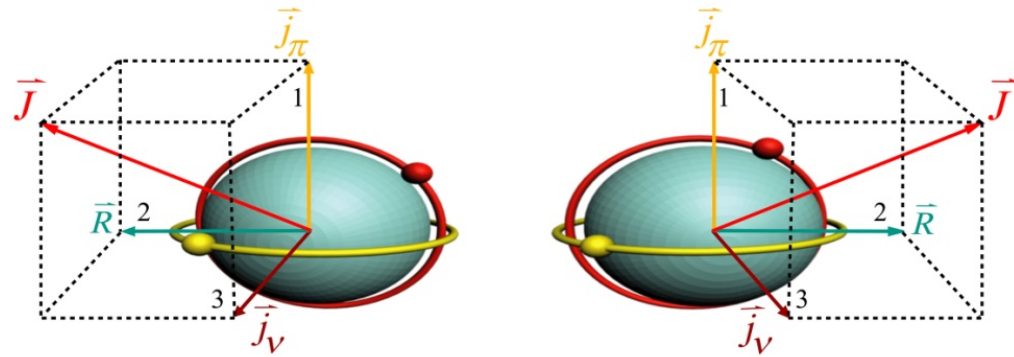
$$[\chi, H] = 0, \chi = TR_y(\pi)$$

$$H |R\rangle = \varepsilon_R |R\rangle, H |L\rangle = \varepsilon_L |L\rangle$$

$$|R\rangle = \chi |L\rangle, |L\rangle = \chi |R\rangle$$

$$\varepsilon_R = \varepsilon_L$$

Ground State (vacuum)  $\Rightarrow$  Chiral



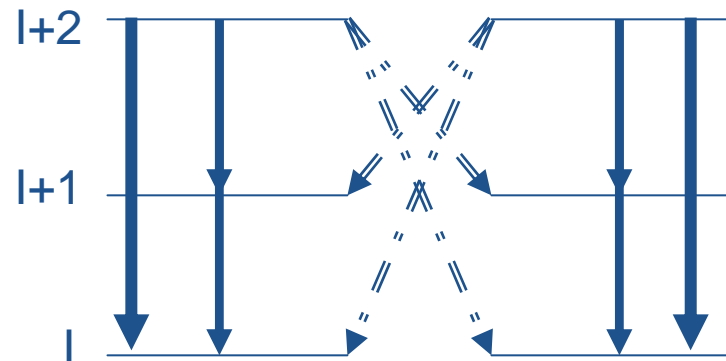
$$|IM +\rangle = \frac{1}{\sqrt{2}} (|R\rangle + |L\rangle)$$

$$|IM -\rangle = \frac{i}{\sqrt{2}} (|R\rangle - |L\rangle)$$

$$H |IM \pm\rangle = \varepsilon_{\pm}^{IM} |IM \pm\rangle$$

$$\chi |IM \pm\rangle = |IM \pm\rangle$$

$$\varepsilon_{+}^{IM} = \varepsilon_{-}^{IM}$$



Chiral doublet  $\Rightarrow$  Restoration



# First Observations of the Chiral doublets bands

VOLUME 86, NUMBER 6

PHYSICAL REVIEW LETTERS

5 FEBRUARY 2001

## Chiral Doublet Structures in Odd-Odd $N = 75$ Isotones: Chiral Vibrations

K. Starosta,<sup>1,\*</sup> T. Koike,<sup>1</sup> C.J. Chiara,<sup>1</sup> D.B. Fossan,<sup>1</sup> D.R. LaFosse,<sup>1</sup> A.A. Hecht,<sup>2</sup> C.W. Beausang,<sup>2</sup> M.A. Caprio,<sup>2</sup>

New sideband partners of the yrast bands built on the  $\pi h_{11/2} \nu h_{11/2}$  configuration were identified in  $^{55}\text{Cs}$ ,  $^{57}\text{La}$ , and  $^{61}\text{Pm}$   $N = 75$  isotones of  $^{134}\text{Pr}$ . These bands form with  $^{134}\text{Pr}$  unique doublet-band systematics suggesting a common basis. Aplanar solutions of 3D tilted axis cranking calculations for triaxial shapes define left- and right-handed chiral systems out of the three angular momenta provided by the valence particles and the core rotation, which leads to spontaneous chiral symmetry breaking and the doublet bands. Small energy differences between the doublet bands suggest collective chiral vibrations.

Static chirality exists in molecules composed of four different atoms and is common for biological and pharmaceutical molecules with important consequences. In particle physics, chirality is a dynamic property which distinguishes for massless fermions the parallel or antiparallel orientation of spin and momentum. Intriguing nuclear structure effects associated with spontaneous chiral symmetry breaking have been predicted recently from angular momentum coupling considerations in odd-odd nuclei having triaxial shapes [1]. These predictions relate to configurations where

\*On leave from Institute of Experimental Physics, Warsaw University, Hoza 69, 00-681 Warsaw, Poland.

†Also at Clark University, Worcester, MA 10610.

‡On leave from Faculty of Physics, Sofia University, BG-1164 Sofia, Bulgaria.

[1] S. Frauendorf and J. Meng, Nucl. Phys. **A617**, 131 (1997).

[2] A. Bohr and B. Mottelsson, *Nuclear Structure* (Benjamin, New York, 1975), Vol. 2.

[3] C.M. Petrache *et al.*, Nucl. Phys. **A597**, 106 (1996); C.M. Petrache *et al.*, Z. Phys. A **344**, 227 (1992).

# Observations in Odd-A Nucleus

VOLUME 91, NUMBER 13

PHYSICAL REVIEW LETTERS

week ending  
26 SEPTEMBER 2003

## A Composite Chiral Pair of Rotational Bands in the Odd-A Nucleus $^{135}\text{Nd}$

S. Zhu,<sup>1</sup> U. Garg,<sup>1</sup> B. K. Nayak,<sup>1</sup> S. S. Ghugre,<sup>2</sup> N. S. Pattabiraman,<sup>2</sup> D. B. Fossan,<sup>3</sup> T. Koike,<sup>3</sup> K. Starosta,<sup>3</sup> C. Vaman,<sup>3</sup>  
R. V. F. Janssens,<sup>4</sup> R. S. Chakrawarthy,<sup>5</sup> M. Whitehead,<sup>5</sup> A. O. Macchiavelli,<sup>6</sup> and S. Frauendorf<sup>1</sup>

<sup>1</sup>*Physics Department, University of Notre Dame, Notre Dame, Indiana 46556, USA*

<sup>2</sup>*IUCDAEF-Calcutta Center, Calcutta 700 094, India*

<sup>3</sup>*Department of Physics and Astronomy, State University of New York, Stony Brook, New York 11794, USA*

<sup>4</sup>*Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA*

<sup>5</sup>*Schuster Laboratory, University of Manchester, Manchester M13 9PL, United Kingdom*

<sup>6</sup>*Nuclear Physics Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

(Received 20 February 2003; revised manuscript received 14 May 2003; published 25 September 2003)

# Observations in A~100 Nucleus

VOLUME 92, NUMBER 3

PHYSICAL REVIEW LETTERS

week ending  
23 JANUARY 2004

## Chiral Degeneracy in Triaxial $^{104}\text{Rh}$

C. Vaman, D. B. Fossan, T. Koike, and K. Starosta\*

*Department of Physics and Astronomy, SUNY at Stony Brook, Stony Brook, New York 11794-3800, USA*

I. Y. Lee and A. O. Macchiavelli

*Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

(Received 11 June 2003; published 22 January 2004)

Chiral doublet bands based on the  $\pi g_{9/2} \otimes \nu h_{11/2}$  configuration that achieve degeneracy at spin  $I = 17$  in the odd-odd triaxial  $^{104}\text{Rh}$  nucleus have been observed. Experimental verification of the interpretation has been tested against specific fingerprints of chirality in the intrinsic system.

# Self-consistent rotating mean field chiral solutions

VOLUME 84, NUMBER 25

PHYSICAL REVIEW LETTERS

19 JUNE 2000

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## Chirality of Nuclear Rotation

V. I. Dimitrov,\* S. Frauendorf, and F. Dönau

*Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556*

*and Institute for Nuclear and Hadronic Physics, Research Center Rossendorf, PB 51 01 19, 01314 Dresden, Germany*

(Received 18 January 2000)

It is shown that the rotating mean field of triaxial nuclei can break the chiral symmetry. Two nearly degenerate  $\Delta I = 1$  rotational bands originate from the left-handed and right-handed solutions.

PACS numbers: 21.10.Re, 11.30.Rd, 23.20.Lv, 27.60.+j



# Critical frequency in Chiral rotations

VOLUME 93, NUMBER 5

PHYSICAL REVIEW LETTERS

week ending  
30 JULY 2004

## Critical Frequency in Nuclear Chiral Rotation

P. Olbratowski,<sup>1,2</sup> J. Dobaczewski,<sup>1,2</sup> J. Dudek,<sup>2</sup> and W. Plóciennik<sup>3</sup>

<sup>1</sup>*Institute of Theoretical Physics, Warsaw University, Hoża 69, PL-00681 Warsaw, Poland*

<sup>2</sup>*Institut de Recherches Subatomiques, CNRS-IN2P3/Université Louis Pasteur, F-67037 Strasbourg Cedex 2, France*

<sup>3</sup>*The Andrzej Soltan Institute for Nuclear Studies, PL-05400 Świerk, Poland*

(Received 12 March 2004; published 29 July 2004)

Self-consistent solutions for the so-called planar and chiral rotational bands in  $^{132}\text{La}$  are obtained for the first time within the Skyrme-Hartree-Fock cranking approach. It is suggested that the chiral rotation cannot exist below a certain critical frequency which under the approximations used is estimated as  $\hbar\omega_{\text{crit}} \approx 0.5\text{--}0.6$  MeV. However, the exact values of  $\hbar\omega_{\text{crit}}$  may vary, to an extent, depending on the microscopic model used, in particular, through the pairing correlations and/or calculated equilibrium deformations. The existence of the critical frequency is explained in terms of a simple classical model of two gyroscopes coupled to a triaxial rigid body.

DOI: 10.1103/PhysRevLett.93.052501

PACS numbers: 21.10.Re, 21.30.Fe, 21.60.Jz, 27.60.+j

# Selection rule of EM transitions in Chiral rotations

VOLUME 93, NUMBER 17

PHYSICAL REVIEW LETTERS

week ending  
22 OCTOBER 2004

## Chiral Bands, Dynamical Spontaneous Symmetry Breaking, and the Selection Rule for Electromagnetic Transitions in the Chiral Geometry

T. Koike,<sup>1</sup> K. Starosta,<sup>1,2</sup> and I. Hamamoto<sup>3,4</sup>

<sup>1</sup>*Department of Physics and Astronomy, SUNY at Stony Brook, New York 11794, USA*

<sup>2</sup>*Department of Physics and Astronomy and National Superconducting Cyclotron Laboratory, MSU, East Lansing, Michigan 48824, USA*

<sup>3</sup>*Division of Mathematical Physics, LTH, University of Lund, Sweden*

<sup>4</sup>*The Niels Bohr Institute, Blegdamsvej 17, Copenhagen Ø, DK 2100, Denmark*

(Received 16 March 2004; revised manuscript received 29 July 2004; published 22 October 2004)

A model for a special configuration in triaxial odd-odd nuclei is constructed which exhibits degenerate chiral bands with a sizable rotation, a manifestation of dynamical spontaneous symmetry breaking. A quantum number obtained from the invariance of the model Hamiltonian, which characterizes observable states, is given and selection rules for electromagnetic transition probabilities in chiral bands is derived in terms of this quantum number. The degeneracy of the lowest two bands is indeed obtained in the numerical diagonalization of the Hamiltonian at an intermediate spin range, over which electromagnetic transitions follow exactly the selection rule expected for the chiral geometry.

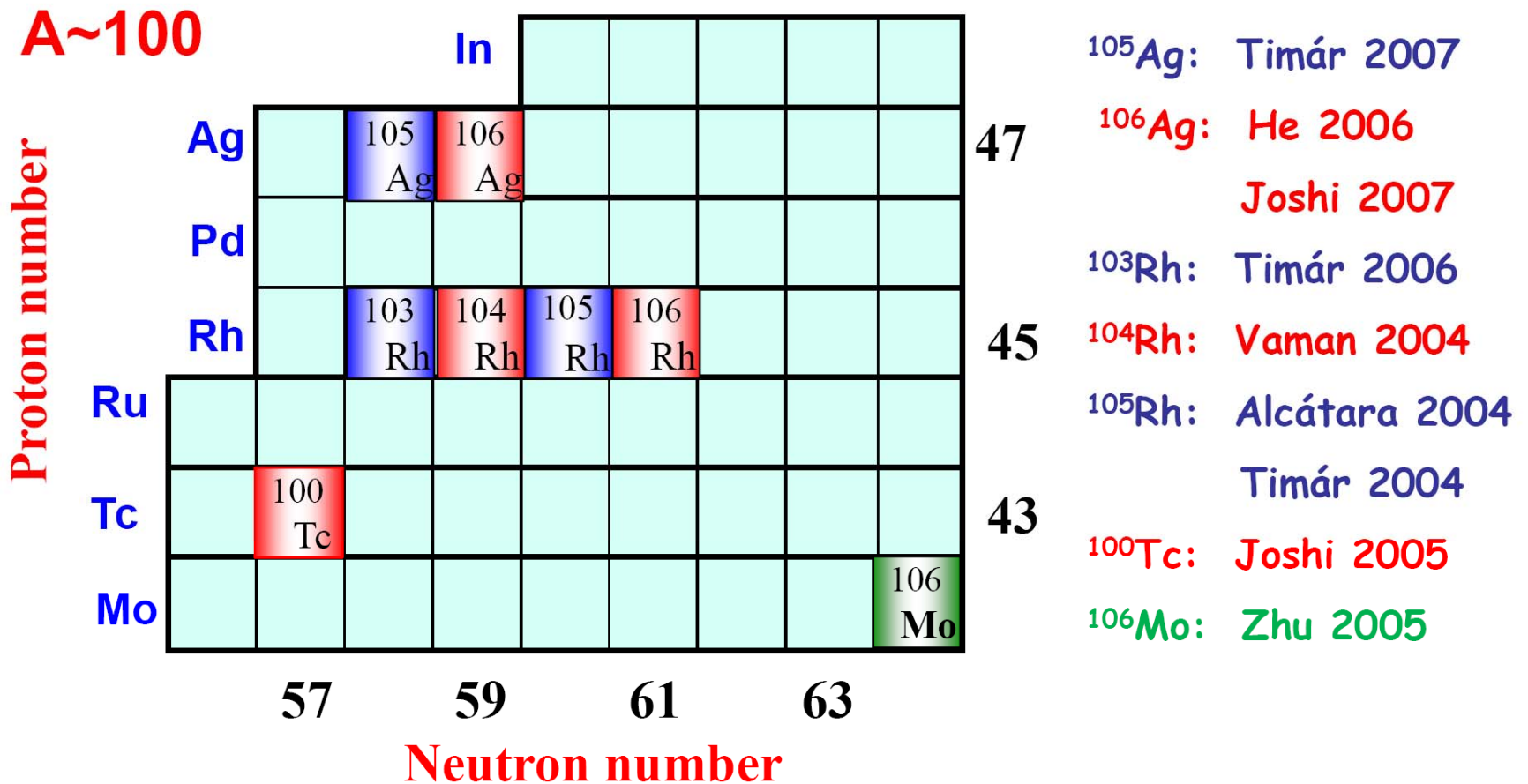
DOI: 10.1103/PhysRevLett.93.172502

PACS numbers: 23.20.Lv, 21.10.Re, 21.60.Ev



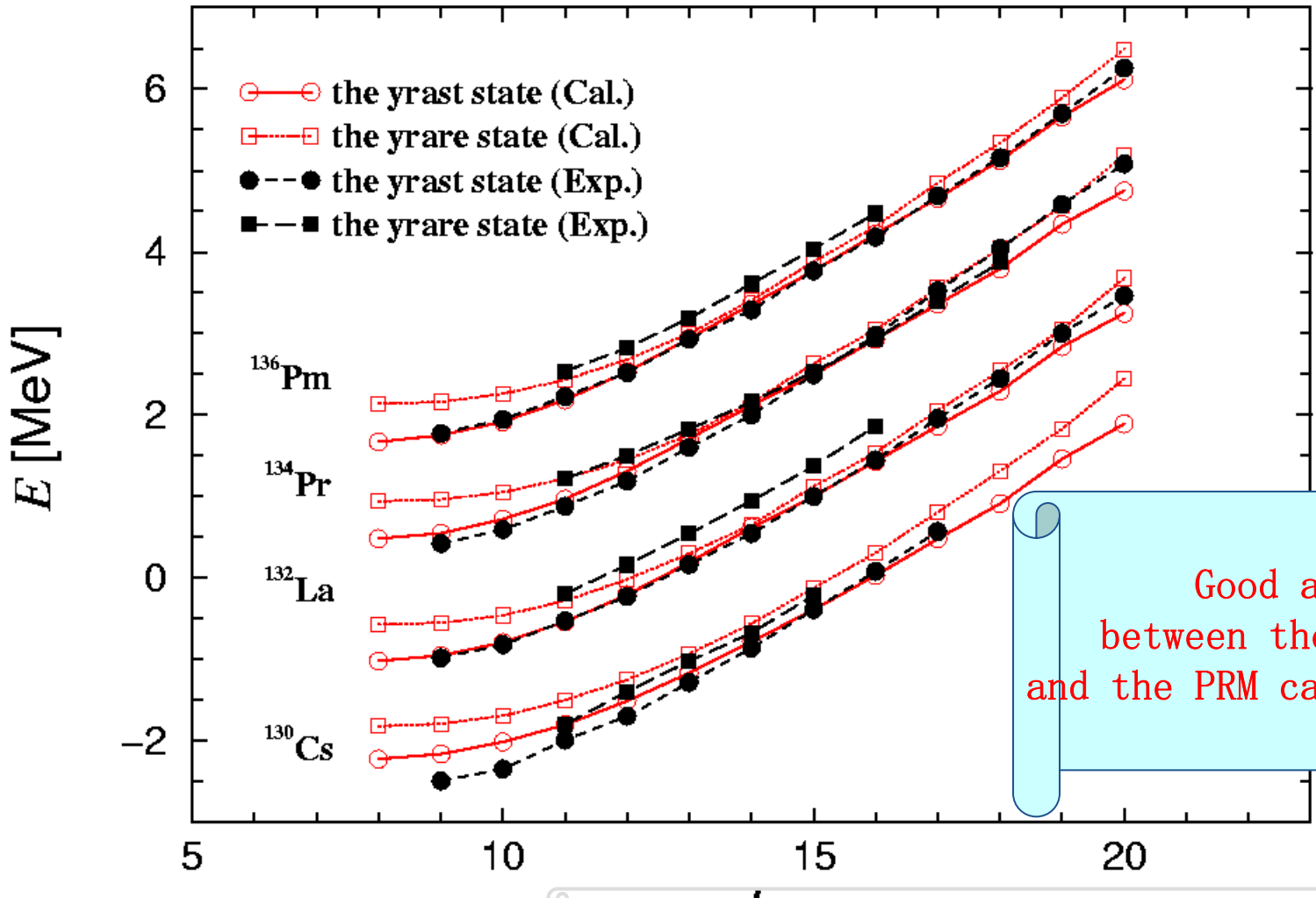


# Candidate Chiral bands in $A \sim 100$ region





# Near energy degeneracy: easy !



Good agreement between the experimental and the PRM calculated results

# Near energy degeneracy – chirality?

PRL **96**, 052501 (2006)

PHYSICAL REVIEW LETTERS

week ending  
10 FEBRUARY 2006

## Transition Probabilities in $^{134}\text{Pr}$ : A Test for Chirality in Nuclear Systems

D. Tonev,<sup>1,2</sup> G. de Angelis,<sup>1</sup> P. Petkov,<sup>2</sup> A. Dewald,<sup>3</sup> S. Brant,<sup>4</sup> S. Frauendorf,<sup>5</sup> D.L. Balabanski,<sup>2,6</sup> P. Pejovic,<sup>3</sup>  
D. Bazzacco,<sup>7</sup> P. Bednarczyk,<sup>8</sup> F. Camera,<sup>9</sup> A. Fitzler,<sup>3</sup> A. Gadea,<sup>1</sup> S. Lenzi,<sup>7</sup> S. Lunardi,<sup>7</sup> N. Marginean,<sup>1</sup> O. Möller,<sup>3</sup>  
D. R. Napoli,<sup>1</sup> A. Paleni,<sup>9</sup> C. M. Petrache,<sup>6</sup> G. Prete,<sup>1</sup> K. O. Zell,<sup>3</sup> Y. H. Zhang,<sup>10</sup> Jing-ye Zhang,<sup>11</sup>  
Q. Zhong,<sup>12</sup> and D. Curien<sup>8</sup>

<sup>1</sup>Laboratori Nazionali di Legnaro, INFN, I-35020 Legnaro, Italy

<sup>2</sup>Institute for Nuclear Research and Nuclear Energy, BAS, 1784 Sofia, Bulgaria

<sup>3</sup>Institut für Kernphysik der Universität zu Köln, D-50937 Köln, Germany

<sup>4</sup>Department of Physics, Faculty of Science, University of Zagreb, 10000 Zagreb, Croatia

<sup>5</sup>Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556, USA

<sup>6</sup>Dipartimento di Fisica, Università di Camerino and INFN Perugia, I-62032 Camerino, Italy

<sup>7</sup>Dipartimento di Fisica, Università and INFN Sezione di Padova, I-35131 Padova, Italy

<sup>8</sup>Institut de Recherches Subatomiques, Boîte Postale 28 F-67037, Strasbourg, France

<sup>9</sup>Dipartimento di Fisica, Università di Milano, I-20133 Milano, Italy

<sup>10</sup>Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 73000, Peoples Republic of China

<sup>11</sup>Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996, USA

<sup>12</sup>Department of Nuclear Physics, China Institute of Atomic Energy, Beijing 102413, Peoples Republic of China

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Excited states in  $^{134}\text{Pr}$  were populated in the fusion-evaporation reaction  $^{119}\text{Sn}(^{19}\text{F}, 4n)^{134}\text{Pr}$ . Recoil distance Doppler-shift and Doppler-shift attenuation measurements using the Euroball spectrometer, in conjunction with the inner Bismuth Germanate ball and the Cologne plunger, were performed at beam energies of 87 MeV and 83 MeV, respectively. Reduced transition probabilities in  $^{134}\text{Pr}$  are compared to the predictions of the two quasiparticle + triaxial rotor and interacting boson fermion-fermion models. The experimental results do not support the presence of static chirality in  $^{134}\text{Pr}$  underlying the importance of shape fluctuations. Only within a dynamical context the presence of intrinsic chirality in  $^{134}\text{Pr}$  can be supported.



**Risk of Misinterpretation of Nearly Degenerate Pair Bands as Chiral Partners in Nuclei**C. M. Petrache,<sup>1</sup> G. B. Hagemann,<sup>2</sup> I. Hamamoto,<sup>3,2</sup> and K. Starosta<sup>4</sup><sup>1</sup>*Dipartimento di Fisica, Università di Camerino and INFN, Sezione di Perugia, I-62032, Camerino, Italy*<sup>2</sup>*Niels Bohr Institute, Blegdamsvej 17, DK-2100 Copenhagen, Denmark*<sup>3</sup>*Department of Mathematical Physics, Lund Institute of Technology at the University of Lund, S-22362 Lund, Sweden*<sup>4</sup>*Department of Physics and Astronomy and National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing, Michigan 48824, USA*

(Received 19 September 2005; published 21 March 2006)

The experimental information on the observed nearly degenerate bands in the  $N = 75$  isotones, in particular  $^{134}\text{Pr}$  and  $^{136}\text{Pm}$ , which are often considered as the best candidates for chiral bands, is critically analyzed. Most properties of the bands, in particular, the recently measured branching ratios and lifetimes, are in clear disagreement with the interpretation of the two bands as chiral bands. For  $I = 14\text{--}18$  in  $^{134}\text{Pr}$ , where the observed energies are almost degenerate, we have obtained a value of 2.0(4) for the ratio of the transition quadrupole moments of the two bands, which implies a considerable difference in the nuclear shape associated with the two bands. The insufficiency of the near-degeneracy criterion to trace nuclear chirality is emphasized.

DOI: [10.1103/PhysRevLett.96.112502](https://doi.org/10.1103/PhysRevLett.96.112502)

PACS numbers: 21.10.Re, 21.60.Ev, 23.20.Lv, 27.60.+j

**End of the story ?**



# 手征原子核: another best example! ?

PRL **97**, 172501 (2006)

PHYSICAL REVIEW LETTERS

week ending  
27 OCTOBER 2006

## $^{128}\text{Cs}$ as the Best Example Revealing Chiral Symmetry Breaking

E. Grodner,<sup>1</sup> J. Srebrny,<sup>1,2</sup> A. A. Pasternak,<sup>1,2,3</sup> I. Zalewska,<sup>1</sup> T. Morek,<sup>1</sup> Ch. Droste,<sup>1</sup> J. Mierzejewski,<sup>2</sup> M. Kowalczyk,<sup>1,2</sup>  
J. Kownacki,<sup>2</sup> M. Kisieliński,<sup>2,4</sup> S. G. Rohoziński,<sup>5</sup> T. Koike,<sup>6</sup> K. Starosta,<sup>7</sup> A. Kordyasz,<sup>2</sup> P. J. Napiorkowski,<sup>2</sup>  
M. Wolińska-Cichocka,<sup>2</sup> E. Ruchowska,<sup>4</sup> W. Plóciennik,<sup>4,\*</sup> and J. Perkowski<sup>8</sup>

<sup>1</sup>*Institute of Experimental Physics, Warsaw University, ul. Hoża 69, PL-00681, Warsaw, Poland*

<sup>2</sup>*Heavy Ion Laboratory, Warsaw University, ul. Pasteura 5A, 02-093 Warsaw, Poland*

<sup>3</sup>*A. F. Ioffe Physical Technical Institute, 194021 St. Petersburg, Russia*

<sup>4</sup>*The A. Sołtan Institute for Nuclear Studies, 05-400, Świerk, Poland*

<sup>5</sup>*Institute of Theoretical Physics, Warsaw University, ul. Hoża 69, PL-00681, Warsaw, Poland*

<sup>6</sup>*Graduate School of Science, Tohoku University, 980-8578, Japan*

<sup>7</sup>*National Superconducting Cyclotron Laboratory and Department of Physics and Astronomy, Michigan State University,  
164 S. Shaw Lane, East Lansing, Michigan 48825-1321, USA*

<sup>8</sup>*Division of Nuclear Physics, University of Łódź, 90-236 Łódź, Poland*

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The results of the Doppler-shift attenuation method lifetime measurements in partner bands of  $^{128}\text{Cs}$  and  $^{132}\text{La}$  are presented. Experimental reduced transition probabilities in  $^{128}\text{Cs}$  are compared with theoretical calculations done in the frame of the core-quasiparticle coupling model. The electromagnetic properties, energy and spin of levels belonging to the partner bands show that  $^{128}\text{Cs}$  is the best known example revealing the chiral symmetry breaking phenomenon.

DOI: [10.1103/PhysRevLett.97.172501](https://doi.org/10.1103/PhysRevLett.97.172501)

PACS numbers: 21.10.Re, 21.10.Tg, 23.20.-g, 27.60.+j

# 手征原子核: another best example! ?

PRL **99**, 172501 (2007)

PHYSICAL REVIEW LETTERS

week ending  
26 OCTOBER 2007

## From Chiral Vibration to Static Chirality in $^{135}\text{Nd}$

S. Mukhopadhyay,<sup>1,2</sup> D. Almehed,<sup>1</sup> U. Garg,<sup>1</sup> S. Frauendorf,<sup>1</sup> T. Li,<sup>1</sup> P. V. Madhusudhana Rao,<sup>1</sup> X. Wang,<sup>1,3</sup> S. S. Ghugre,<sup>2</sup> M. P. Carpenter,<sup>3</sup> S. Gros,<sup>3</sup> A. Hecht,<sup>3,4</sup> R. V.F. Janssens,<sup>3</sup> F.G. Kondev,<sup>5</sup> T. Lauritsen,<sup>3</sup> D. Seweryniak,<sup>3</sup> and S. Zhu<sup>3</sup>

<sup>1</sup>Physics Department, University of Notre Dame, Notre Dame, Indiana 46556, USA

<sup>2</sup>UGC-DAE Consortium for Scientific Research, Kolkata Centre, Kolkata 700098, India

<sup>3</sup>Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

<sup>4</sup>Department of Chemistry and Biochemistry, University of Maryland, College Park, Maryland 20742, USA

<sup>5</sup>Nuclear Engineering Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

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Electromagnetic transition probabilities have been measured for the intraband and interband transitions in the two sequences in the nucleus  $^{135}\text{Nd}$  that were previously identified as a composite chiral pair of rotational bands. The chiral character of the bands is affirmed and it is shown that their behavior is associated with a transition from a vibrational into a static chiral regime.

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PACS numbers: 21.10.Tg, 11.30.Rd, 21.60.Ev, 27.60.+j

Chirality is a well-known phenomenon in chemistry and biology as a geometric property of many molecules, in particular of complex biomolecules. Particle physics is another domain of chirality, where it describes a kinematical feature of massless particles. Both in chemistry and particle physics, space inversion changes left-handed into right-handed systems. Nuclei have been considered as achiral, because their shapes are, generally, too simple. However, it was pointed out some time ago that a triaxial nucleus becomes chiral if it rotates about an axis that lies outside the three planes spanned by the principal axes of its

excitation energies of the two partner bands never approach each other. In  $^{134}\text{Pr}$ , the two bands come very close around spin 14; however, the intraband  $B(E2)$  values were found to differ by about a factor of 2 [13,14]. The latter observation was interpreted as suggesting that the change in orientation of the angular momentum vector must be accompanied by a change in shape [13], or even that the results rendered the chiral interpretation itself doubtful [14].

The study of magnetic rotation [1] suggests that chirality is better developed if the nonplanar geometry is generated

# Observations of the Chiral doublets bands in China

CHIN.PHYS.LETT.

Vol. 19, No. 12 (2002) 1779

## Search for the Chiral Band in the $N = 71$ Odd-Odd Nucleus $^{126}\text{Cs}$ \*

LI Xian-Feng(李险峰)<sup>1</sup>, MA Ying-Jun(马英君)<sup>1</sup>, LIU Yun-Zuo(刘运祚)<sup>1,2</sup>, LU Jing-Bin(陆景彬)<sup>1</sup>, ZHAO Guang-Yi(赵广义)<sup>1</sup>, YIN Li-Chang(尹利长)<sup>1</sup>, MENG Rui(孟锐)<sup>1</sup>, ZHANG Zhen-Long(张振龙)<sup>1</sup>, WEN Li-Jun(文立军)<sup>1</sup>, ZHOU Xiao-Hong(周小红)<sup>2</sup>, GUO Ying-Xiang(郭应祥)<sup>2</sup>, LEI Xiang-Guo(雷相国)<sup>2</sup>, LIU Zhong(刘忠)<sup>2</sup>, HE Jian-Jun(何建军)<sup>2</sup>, ZHENG Yong(郑勇)<sup>2</sup>

<sup>1</sup>Department of Physics, Jilin University, Changchun 130023

PHYSICAL REVIEW C **74**, 017302 (2006)

## Candidate chiral doublet bands in the odd-odd nucleus $^{126}\text{Cs}$

Shouyu Wang,<sup>1</sup> Yunzuo Liu,<sup>1,2</sup> T. Komatsubara,<sup>3</sup> Yingjun Ma,<sup>1</sup> and Yuhu Zhang<sup>2</sup>

<sup>1</sup>Department of Physics, Jilin University, Changchun 130021, People's Republic of China

<sup>2</sup>Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, People's Republic of China

<sup>3</sup>Institute of Physics, Tandem Accelerator Center, University of Tsukuba, Ibaraki 305, Japan

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The candidate chiral doublet bands recently observed in  $^{126}\text{Cs}$  have been extended to higher spins, several new

Question: Chiral bands ?

Reproduction of Spectra & transition in 2qp coupled with triaxial rotor ! Examine the orientation of the AM

# Nuclear Chirality: Based on the geometry for one particle and one hole coupled to a triaxial rotor with $\gamma=30^\circ$

1. nearly degenerate doublet bands

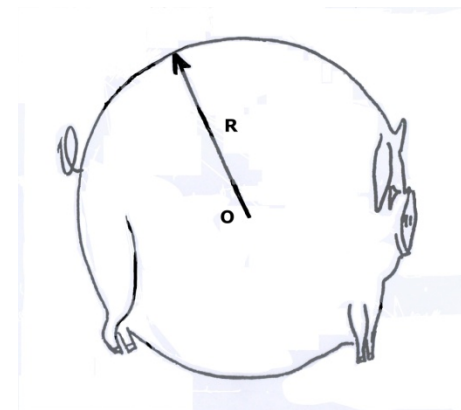
2.  $S(I)$  independent of spin

3. identical spin alignments

4. identical  $B(M1)$ ,  $B(E2)$  values

5. staggering of  $B(M1)/B(E2)$  ratios

6. interband  $B(E2)=0$  at high spin



## chiral bands

It is necessary to have a model with many-particles plus many-holes coupled with triaxial rotor!



# Theoretical Progress

- **3D Tilted Axis Cranking**

- S.Frauendorf and J.Meng, Nucl. Phys. A420, 173 (1984)

- **Relativistic Mean Field**

- Madokoro, Meng, Matsuzaki, Yamaji, Phys. Rev. C 62 (2000) 061301

- Peng et al., Phys. Rev. C 77, 024309 (2008)

- **self- consistent rotating mean field solutions**

- Dimitrov, Frauendorf, and F. Doenau, Phys. Rev. Lett. 84, 5732 (2000)

- **Skyrme-Hatree-Fock Cranking**

- P. Olbratowski et al., Phys. Rev. Lett. 93 052501 (2004)

- **TAC+RPA ( $^{135}\text{Nd}$ )** Phys. Rev. Lett. 99, 172501 (2007)

- **Particle-hole plus triaxial rotor model**

- S.Frauendorf and J.Meng, Nucl. Phys. A420, 173 (1984)

- T. Koike, K. Starosta, and I. Hamamoto, Phys. Rev. Lett. 93, 172502 (2004)

- K. Starosta et al., Phys. Rev. C65 044328 (2004)

- S.Q.Zhang et al., Phys. Rev. C75, 044307 (2007)

- B. Qi et al., Phys. Lett. B675 175 (2009)

- E. A. Lawrie et al., Phys. Rev. C 78 021305 (2008)

- **Pair Truncated Shell Model / Quadrupole Coupling Model**

- K. Higashiyama, N. Yoshinaga, and K. Tanabe, Phys. Rev. C 72, 024315 (2005)

- N. Yoshinaga and K. Higashiyama, Eur. Phys. J A 30,343 (2006)

- **Interacting Boson Fermion Fermion Model (IBFFM)**

- S. Brant, D. Vretenar, and A. Ventura, Phys. Rev. C 69, 017304 (2004)

- **Interacting Vector Boson Model (IVBM)**

- H. G. Ganev, A. I. Georgieva, S. Brant, and A. Ventura, Phys. Rev. C 79, 044322 (2009)

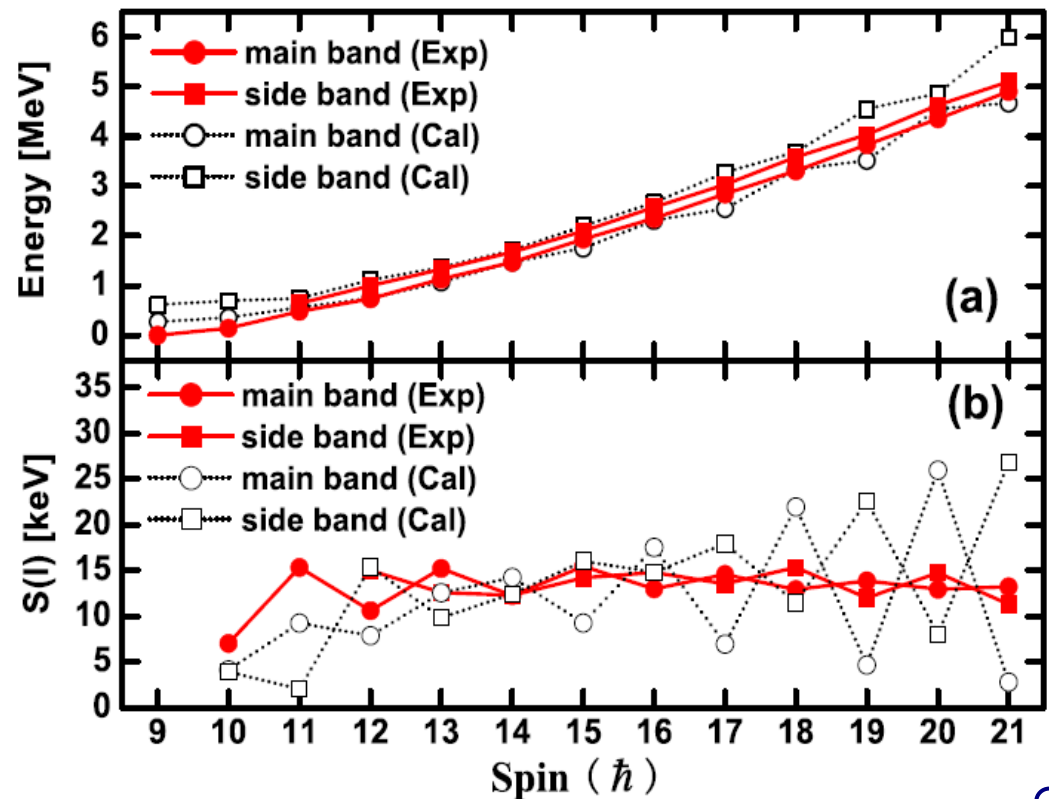
# Simulating n-particles-n-holes by pairing

configuration:

one proton in  $h_{11/2}$

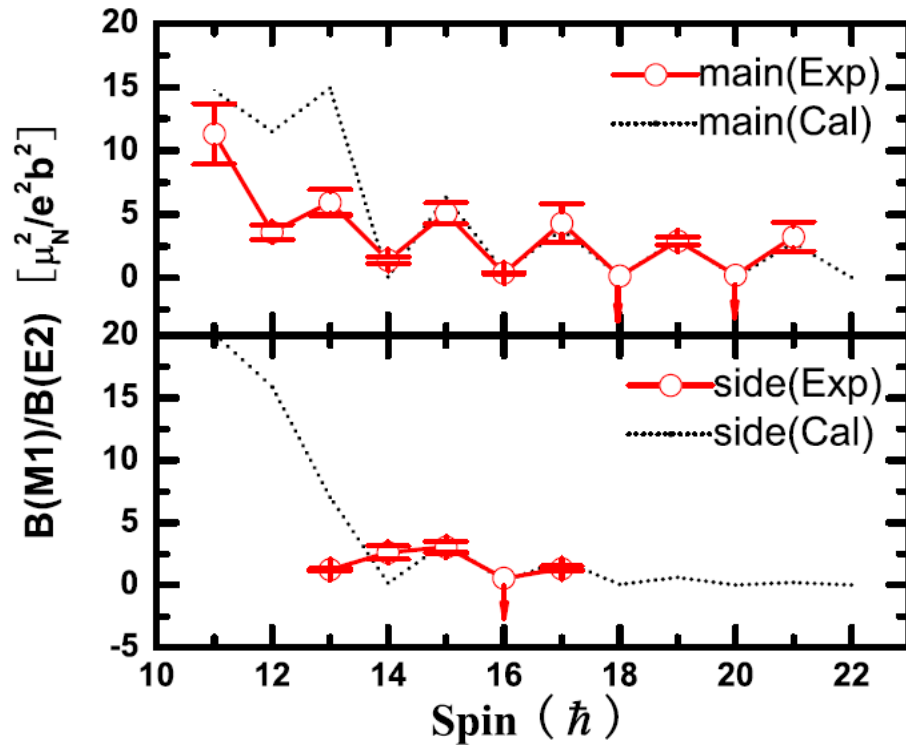
one quasi-neutron in  $h_{11/2}$  : simulate n-neutron holes

Energy spectra  
of  $^{126}\text{Cs}$



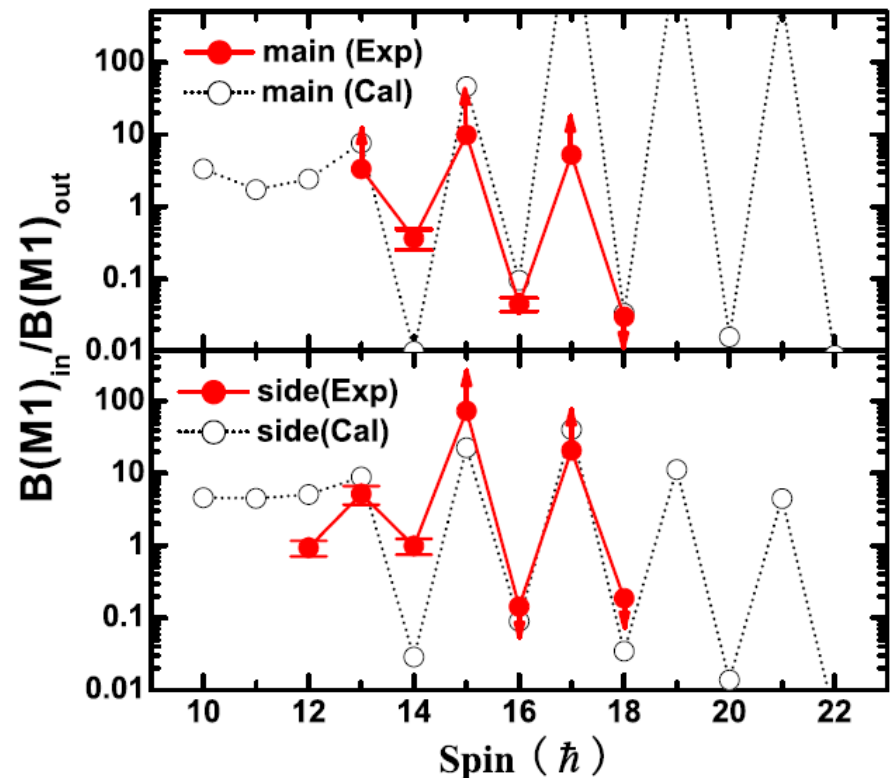


# Simulating n-particles-n-holes by pairing



S.Y. Wang, S.Q. Zhang, B. Qi, J. Meng,  
Phys. Rev. C 75, 024309 (2007)

Electromagnetic  
properties of  $^{126}\text{Cs}$



# $^{135}\text{Nd}$ as an example

PRL **99**, 172501 (2007)

PHYSICAL REVIEW LETTERS

week ending  
26 OCTOBER 2007

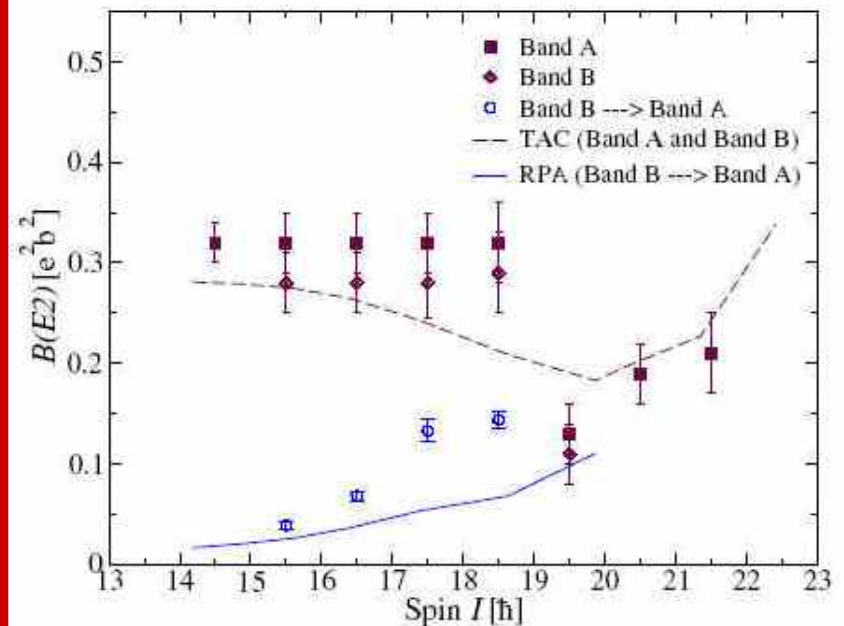
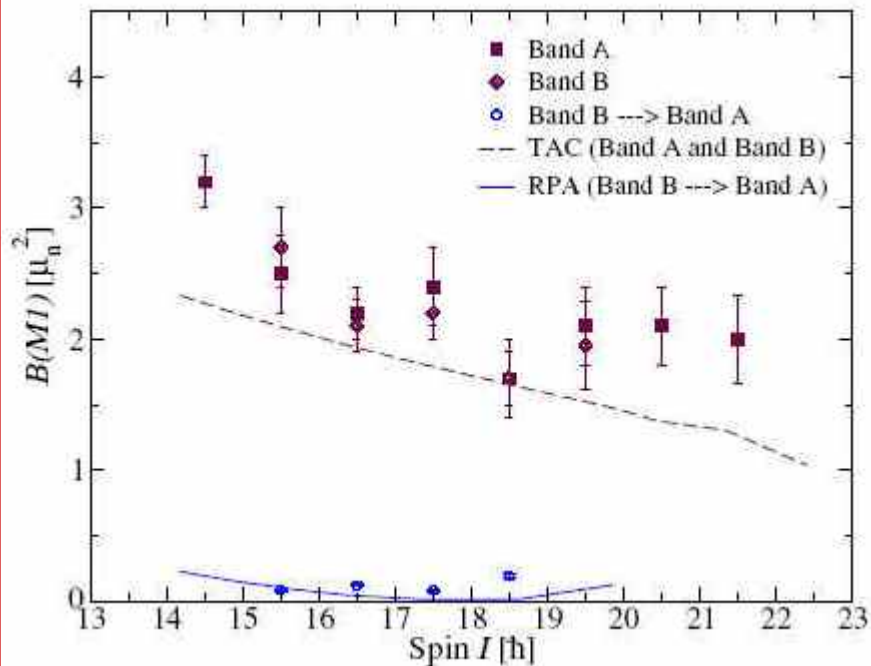
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<sup>3</sup>Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA



# Odd-A Chiral Nucleus $^{135}\text{Nd}$

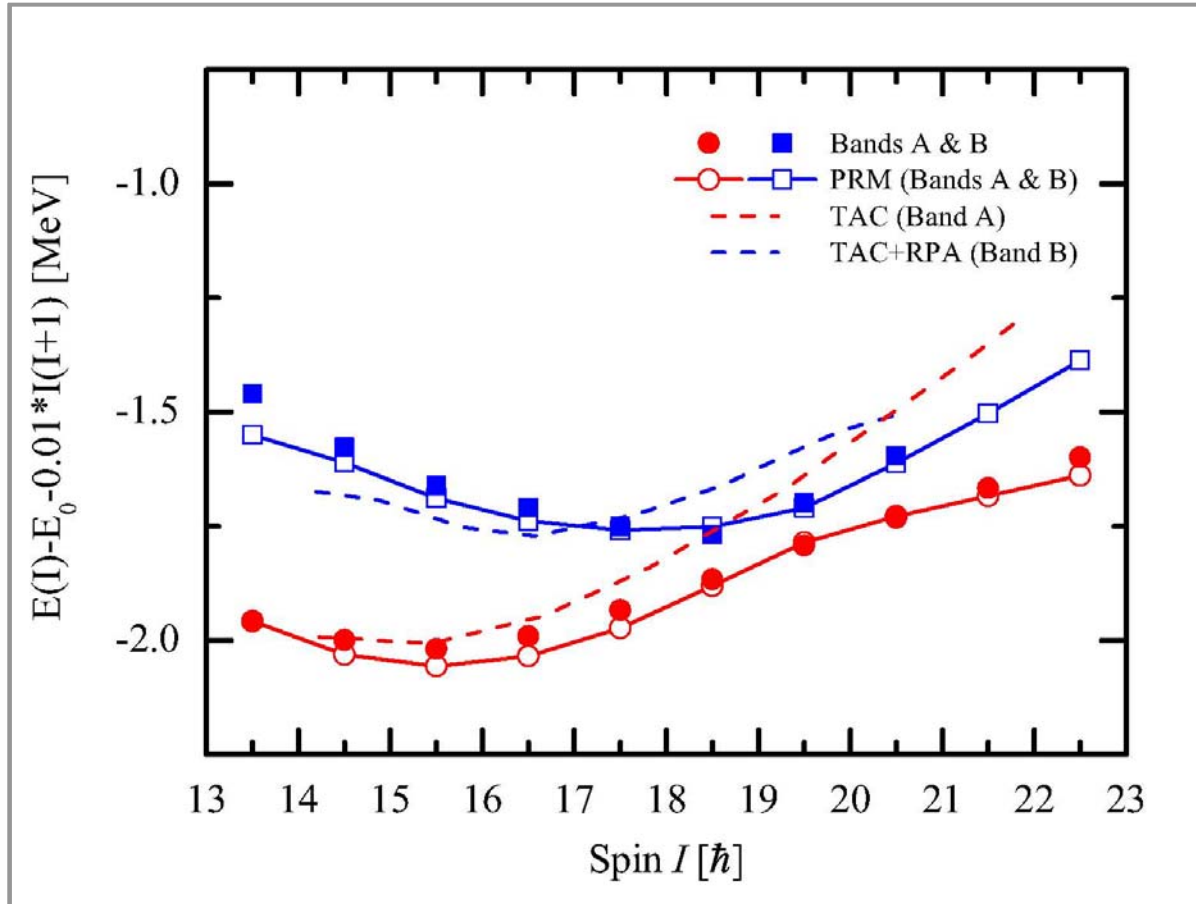
Many-particle-many-hole & triaxial rotor

## Numerical details

- Config:  $\pi h_{11/2}^2 \otimes \nu h_{11/2}^{-1}$
- deformation  $\beta = 0.235, \gamma = 22.4^\circ$  with RMF
- MOI  $29\hbar^2 / \text{MeV}$  justified by data
- intrinsic  $Q$   $Q_0 = (3/\sqrt{5\pi})R_0^2 Z\beta = 4.0eb$
- g-factor  $g_R = Z/A, g_p = 1.21, g_n = -0.21$

# Odd-A Chiral Nucleus $^{135}\text{Nd}$

## 奇-A 核 $^{135}\text{Nd}$ 中候选者征带的能谱

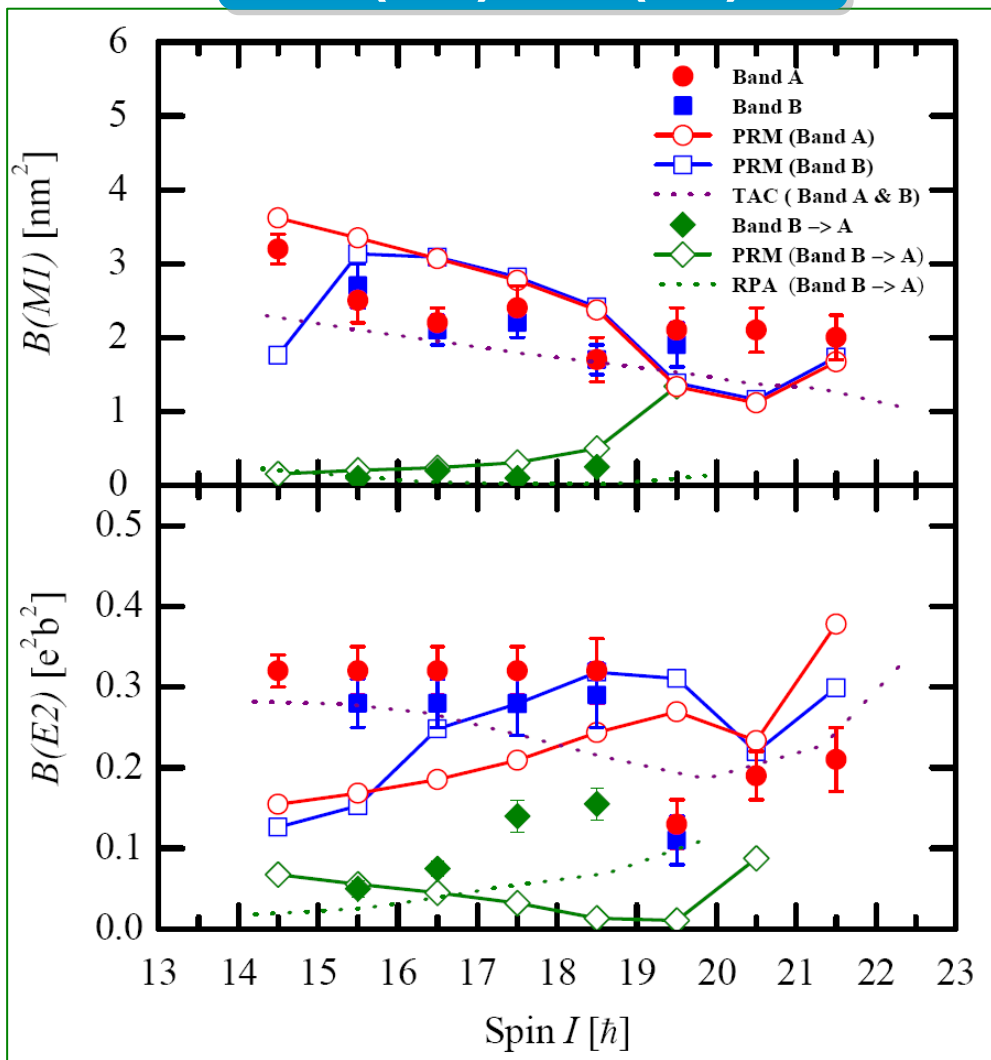


Data: Zhu *PRL* 2003  
TAC: Mukhopahyay *PRL* 2007

Chirality in odd-A nucleus  $^{135}\text{Nd}$  in particle rotor model,  
B.Qi, S.Q. Zhang, J. Meng, S.Y. Wang, S. Frauendorf, *Phys. Lett. B* 675 (2009) 175-180

# Odd-A Chiral Nucleus $^{135}\text{Nd}$

## B(M1) & B(E2)

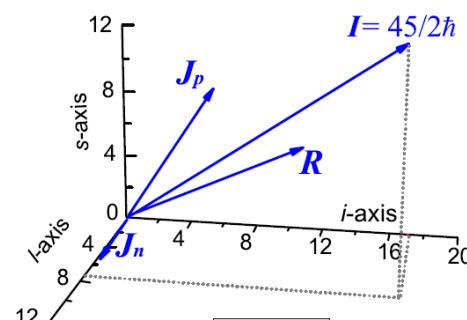
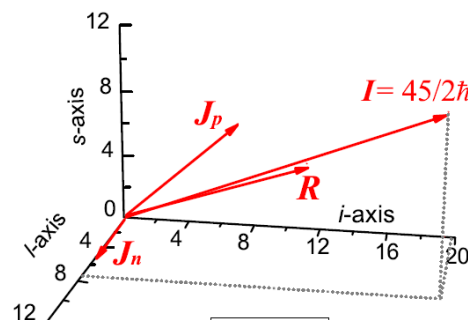
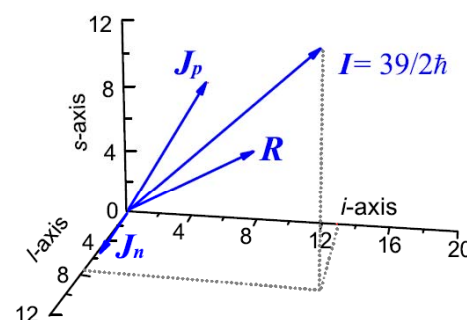
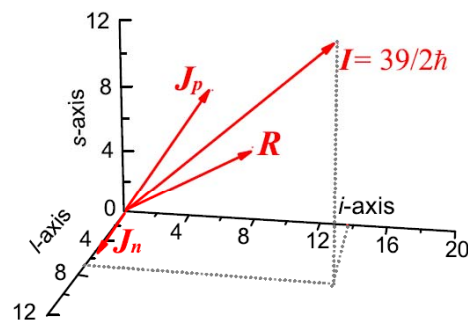
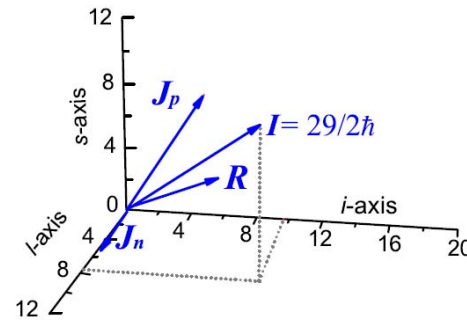
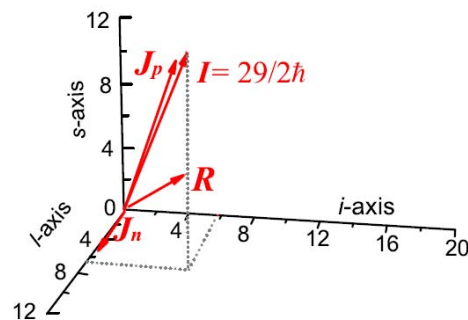


B(M1) & B(E2)的实验  
特征被很好的再现



# Odd-A Chiral Nucleus $^{135}\text{Nd}$

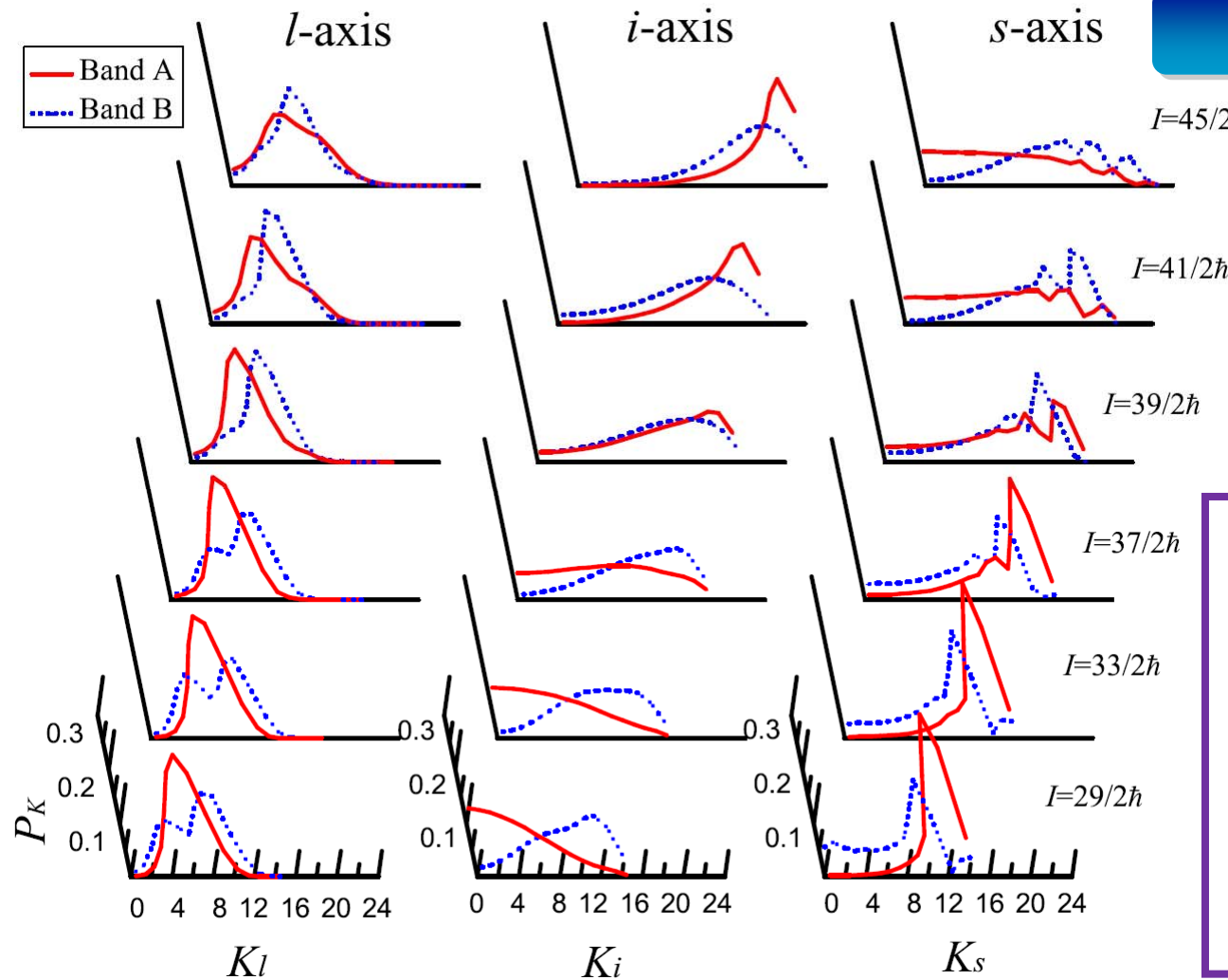
## 角动量期待值



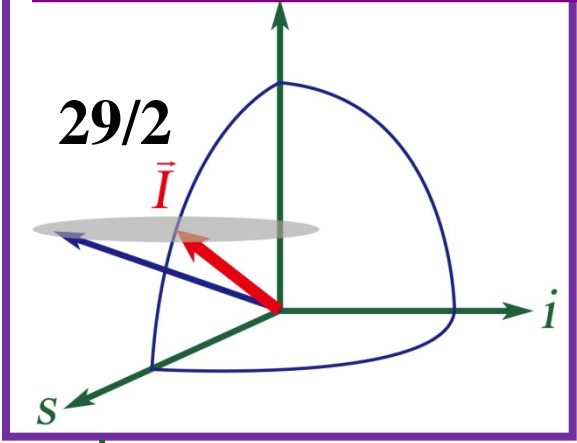
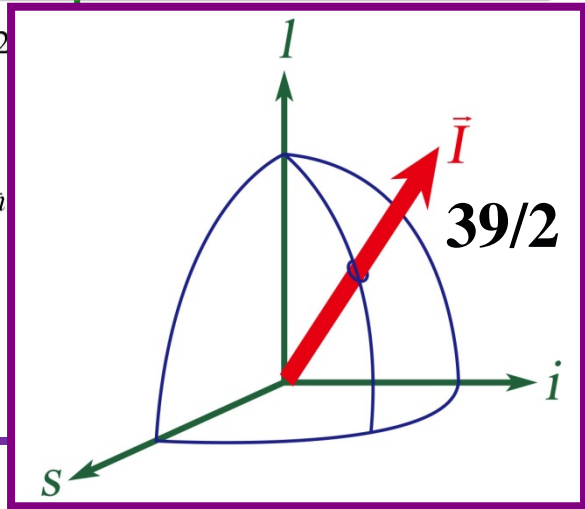
Band A

Band B

# Odd-A Chiral Nucleus $^{135}\text{Nd}$



角动量投影值



动态手性  $\rightarrow$   $I=39/2$ 附近的静态手性  $\rightarrow$  动态手性

# Microscopic deformation calculation for chiral

PHYSICAL REVIEW C 73, 037303 (2006)

## Possible existence of multiple chiral doublets in $^{106}\text{Rh}$

# $M\chi D$

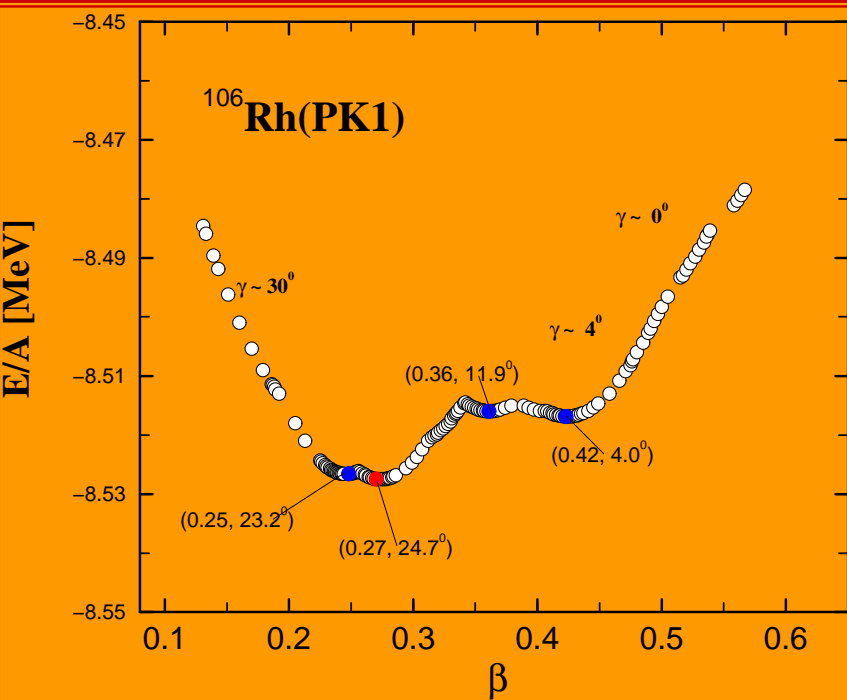
J. Meng,<sup>1,2,3,\*</sup> J. Peng,<sup>1</sup> S. Q. Zhang,<sup>1</sup> and S.-G. Zhou<sup>2,3</sup>

<sup>1</sup>*School of Physics, Peking University, Beijing 100871, China*

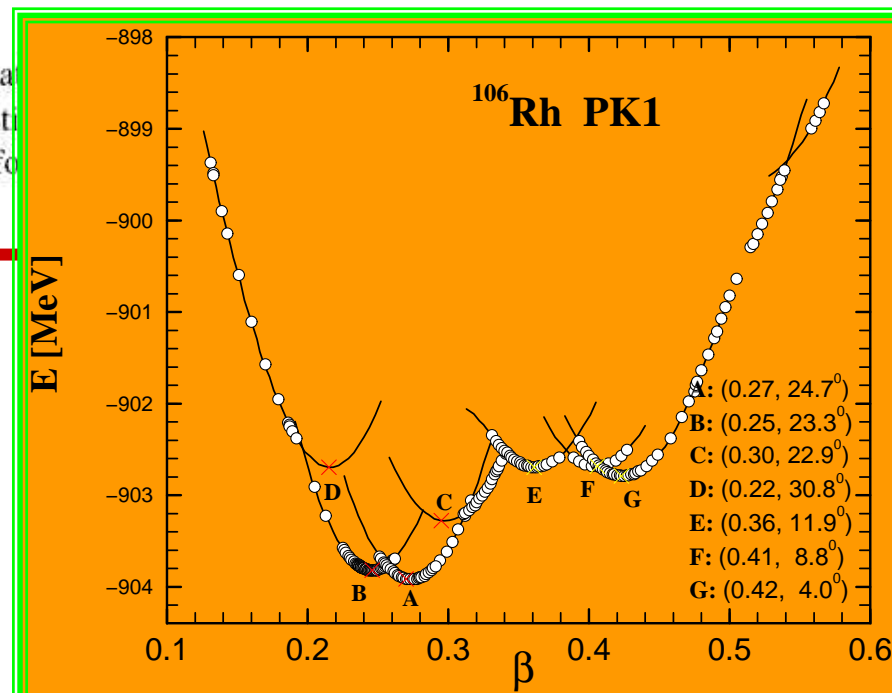
<sup>2</sup>*Institute of Theoretical Physics, Chinese Academy of Science, Beijing 100080, China*

<sup>3</sup>*Center of Theoretical Nuclear Physics, National Laboratory of Heavy Ion Accelerator, Lanzhou 730000, China*

(Received 30 March 2005; published 15 March 2006)



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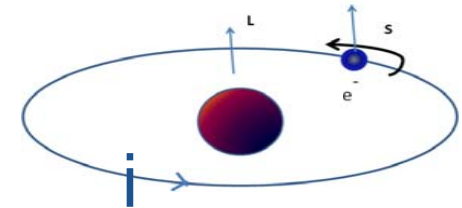
# Relativistic description of magnetic moment

- Nuclear magnetic moment is obtained from the effective electromagnetic current

$$j_i^\mu(x) = \underbrace{Q\bar{\psi}_i\gamma^\mu\psi_i}_{\mathbf{j}_D} + \frac{\lambda_a}{2M}\partial_\nu(\underbrace{\bar{\psi}_i\sigma^{\mu\nu}\psi_i}_{\mathbf{j}_A})$$

- Static magnetic dipole moment is determined by

$$\vec{\mu} = \sum_i \frac{1}{2} \int d^3r [\vec{r} \times \vec{j}_i]$$



- Nuclear magnetic moment can naturally be divided into the **Dirac** and **anomalous** parts,

$$\mu = \begin{cases} \mu_D, & \text{for a relativistic point particle} \\ \mu_A, & \text{from intrinsic structure} \end{cases}$$

# Isoscalar magnetic moment

- The enhanced isoscalar magnetic moment: the over-enhancement for the Dirac magnetic moment

$$\vec{\mu}_D = \frac{1}{2} \int d^3r \vec{r} \times \vec{j}_D = \frac{1}{2} \int d^3r \frac{1}{M^*} \bar{\psi} (\vec{L} + \vec{\Sigma}) \psi$$

With  $M^* \approx 0.6M$

- Time-odd RMF approach: the space-like components of vector meson field leads to non-vanishing current

Isoscalar magnetic moment ( $\mu_N$ )  $\mu_s = [\mu(Z, N) + \mu(Z + 1, N - 1)] / 2$

A	Exp.	Tri.	Axi.	Sph.	Sch.
15	0.22	0.19	0.18	0.32	0.19
17	1.41	1.45	1.48	1.57	1.44
39	0.71	0.67	0.64	0.94	0.64
41	1.92	1.96	1.97	2.21	1.94

The isoscalar magnetic moment can be reproduced quite well

Yao, et al., PRC 74, 024307 (2006).

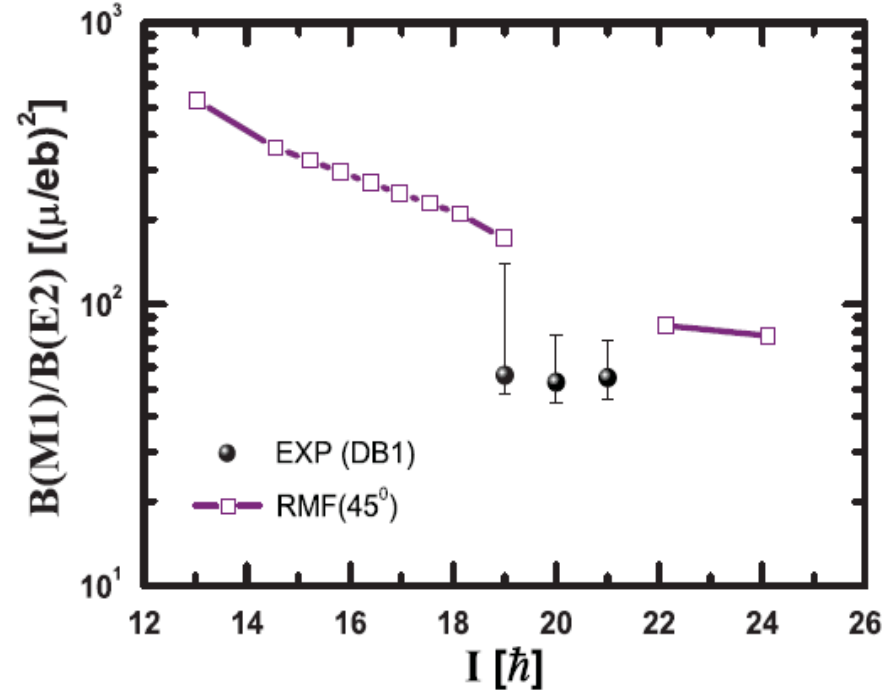
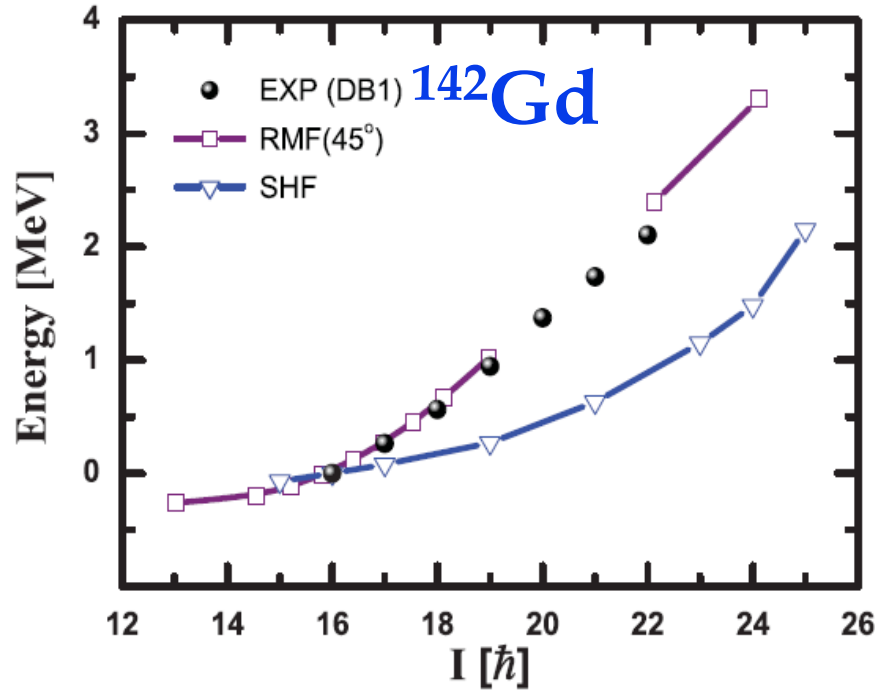


# Magnetic Rotation(MR)

PHYSICAL REVIEW C 78, 024313 (2008)

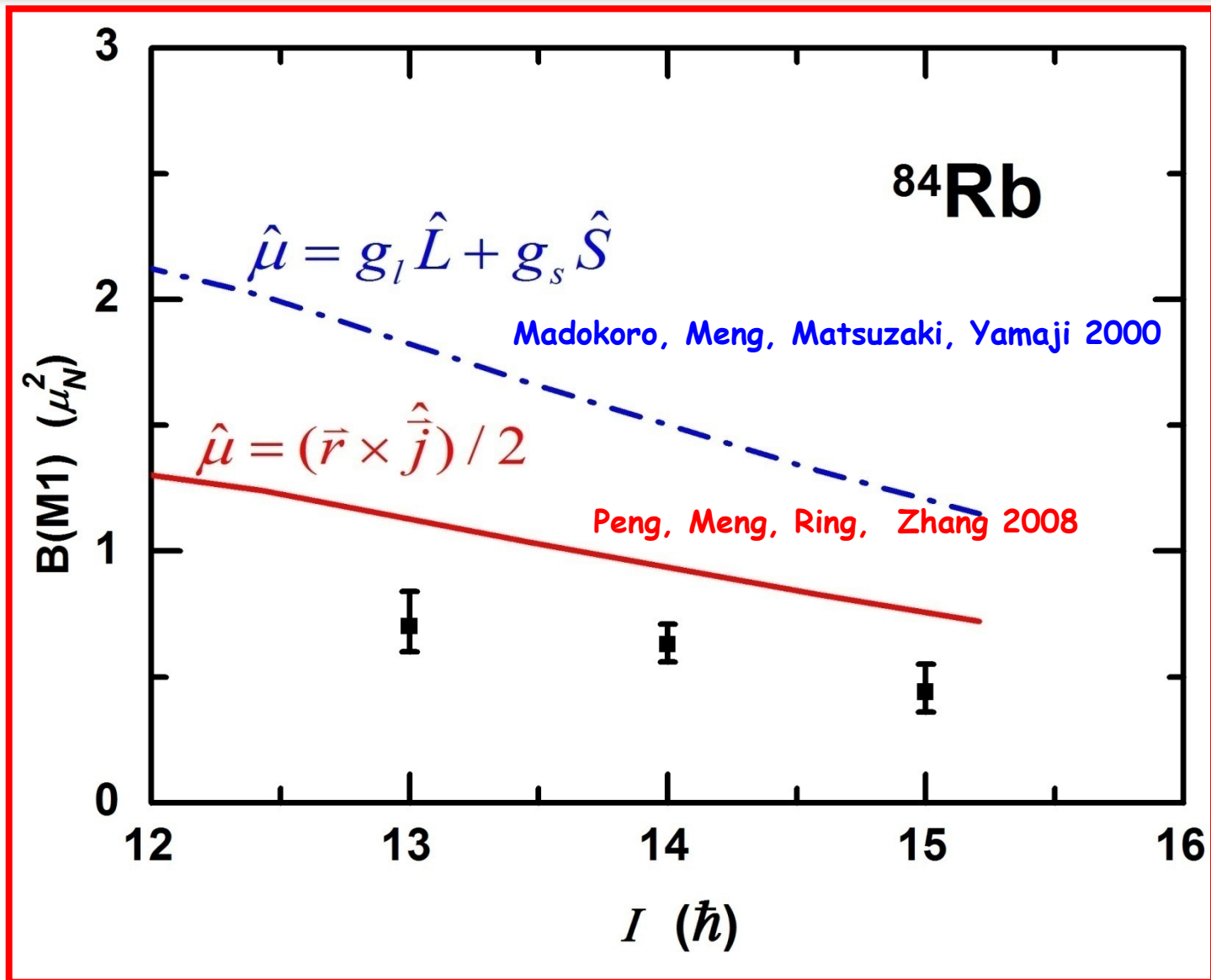
## Covariant density functional theory for magnetic rotation

J. Peng (彭婧),<sup>1,2,3</sup> J. Meng (孟杰),<sup>2,4,5,\*</sup> P. Ring,<sup>3,6,†</sup> and S. Q. Zhang (张双全)<sup>2</sup>



$$\mu = \sum_{i=1}^A \int d^3r \left[ \frac{Mc^2}{\hbar c} Q \psi_i^\dagger(\mathbf{r}) \mathbf{r} \times \boldsymbol{\alpha} \psi_i(\mathbf{r}) + \kappa \psi_i^\dagger(\mathbf{r}) \beta \boldsymbol{\Sigma} \psi_i(\mathbf{r}) \right] \quad B(M1) = \frac{3}{8\pi} (\mu_x \sin \theta_J - \mu_z \cos \theta_J)^2,$$

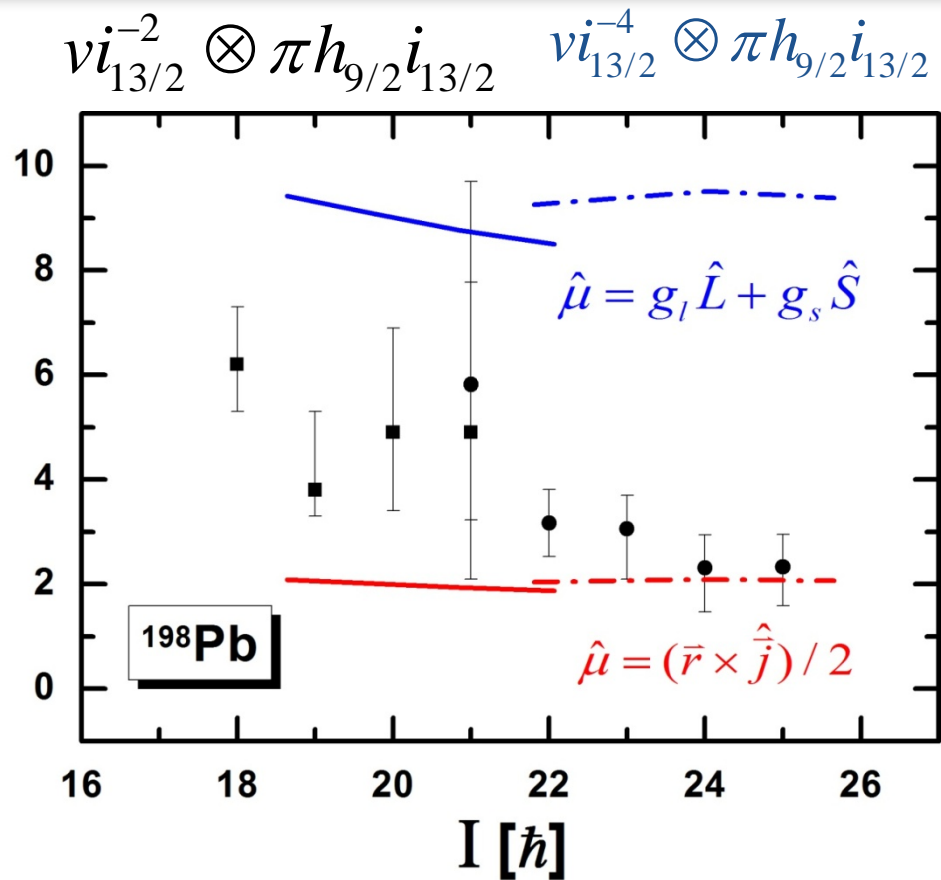
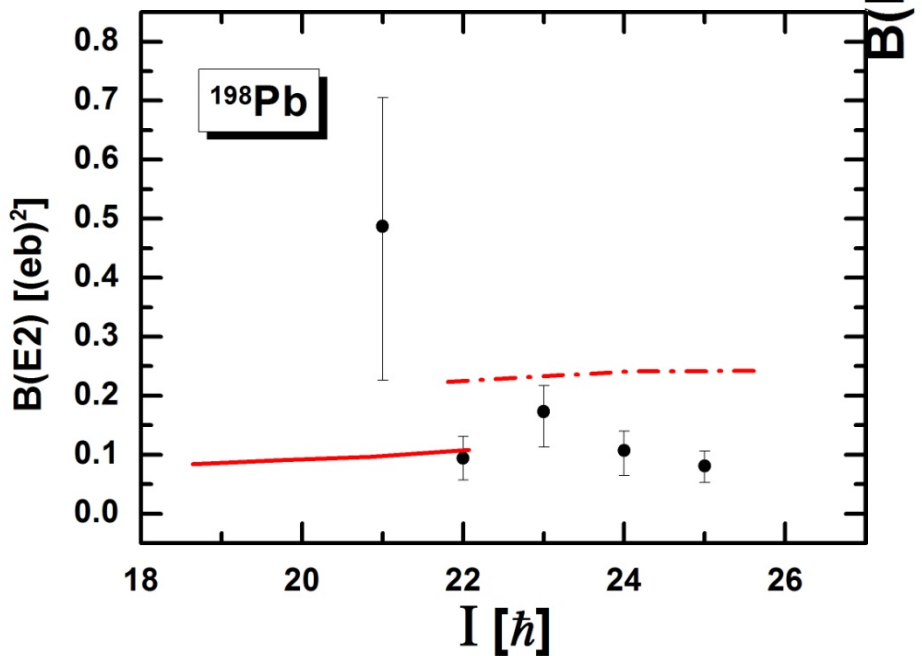
# Relativistic description for B(M1)



# Relativistic description for B(M1) in $^{198}\text{Pb}$

PC-PK1

Nf=12



# Summary & Perspectives

IOP PUBLISHING

JOURNAL OF PHYSICS G: NUCLEAR AND PARTICLE PHYSICS

J. Phys. G: Nucl. Part. Phys. 37 (2010) 064025 (11pp)

doi:10.1088/0954-3899/37/6/064025

## Open problems in understanding the nuclear chirality

Jie Meng<sup>1,2,3</sup> and S Q Zhang<sup>2</sup>

<sup>1</sup> School of Physics and Nuclear Energy Engineering, Beihang University, Beijing 100191, People's Republic of China

<sup>2</sup> State Key Laboratory of Nuclear Physics and Technology, School of Physics, Peking University, Beijing 100871, People's Republic of China

<sup>3</sup> Department of Physics, University of Stellenbosch, Stellenbosch, South Africa

E-mail: [mengj@pku.edu.cn](mailto:mengj@pku.edu.cn)

- Chirality in atomic nuclei is a hot topic
- Lots of theoretical and experimental progress on chirality in atomic nuclei are achieved
- Efforts to understand nuclear chirality more are appreciated



**Thank You!**