CONFIGURATION MIXING WITH RELATIVISTIC SCMF MODELS

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Contents

Outline

- (Relativistic) nuclear energy density functional
  - Adjusting the model parameters
  - Applications: ground-state properties
  - Applications: giant resonances
- collective Hamiltonian model based on the self-consistent RMF
  - Applications: ^{240}Pu isotope
  - Applications: Pt isotopes
- Summary and outlook

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Relativistic energy density functional

Energy density functional consists of the mean-field and the pairing contribution

$$\mathcal{E} = \mathcal{E}_{RMF}[j_\mu, \rho_s] + \mathcal{E}_{pp}(\kappa, \kappa^*)$$

Elementary building blocks

$$(\bar{\psi}\mathcal{O}_\tau\Gamma\psi) \quad \mathcal{O}_\tau \in \{1, \tau_i\} \quad \Gamma \in \{1, \gamma_\mu, \gamma_5, \gamma_5\gamma_\mu, \sigma_{\mu\nu}\}$$

Isoscalar-scalar density

$$\rho_s(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r})\psi_k(\mathbf{r})$$

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Isoscalar-vector current

$$j_\mu(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r})\gamma_\mu\psi_k(\mathbf{r})$$

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Isovector-scalar density

$$\vec{\rho}_s(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r})\vec{\tau}\psi_k(\mathbf{r})$$

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Kinetic energy term

$$\mathcal{E}_{kin} = \sum_i v_i^2 \int \bar{\psi}_i(\boldsymbol{r}) \left(-\gamma \boldsymbol{\nabla} + m\right) \psi_i(\boldsymbol{r})$$

Relativistic energy density functional

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$$\mathcal{E} = \mathcal{E}_{RMF}[j_\mu, \rho_s] + \mathcal{E}_{pp}(\kappa, \kappa^*)$$

Second order terms

$$\mathcal{E}_{2nd} = \frac{1}{2}\int \left[\alpha_v(\rho_v)\rho_v^2 + \alpha_s(\rho_v)\rho_s^2 + \alpha_{tv}(\rho_v)\rho_{tv}^2\right] d\mathbf{r}$$

Relativistic energy density functional

Energy density functional consists of the mean-field and the pairing contribution

$$\mathcal{E} = \mathcal{E}_{RMF}[j_\mu, \rho_s] + \mathcal{E}_{pp}(\kappa, \kappa^*)$$

Derivative terms

$$\mathcal{E}_{der} = \frac{1}{2}\int \delta_s \rho_s \triangle \rho_s d\mathbf{r}$$

Relativistic energy density functional

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$$\mathcal{E} = \mathcal{E}_{RMF}[j_\mu, \rho_s] + \mathcal{E}_{pp}(\kappa, \kappa^*)$$

Coulomb interaction

$$E_{coul} = \frac{e}{2} \int j^p_\mu A^\mu d\mathbf{r}$$

Relativistic energy density functional

Energy density functional consists of the mean-field and the pairing contribution

$$\mathcal{E} = \mathcal{E}_{RMF}[j_\mu, \rho_s] + \mathcal{E}_{pp}(\kappa, \kappa^*)$$

Pairing interaction: finite range separable pairing

$$V(\boldsymbol{r}_1, \boldsymbol{r}_2, \boldsymbol{r}'_1, \boldsymbol{r}'_2) = G\delta(\boldsymbol{R} - \boldsymbol{R}')P(\boldsymbol{r})P(\boldsymbol{r}')\frac{1}{2}(1 - P^\sigma)$$

$$\boldsymbol{R} = \frac{1}{2}(\boldsymbol{r}_1 + \boldsymbol{r}_2), \quad \boldsymbol{r} = \boldsymbol{r}_1 - \boldsymbol{r}_2, \quad P(\boldsymbol{r}) = \frac{1}{4\pi a^2}e^{-\frac{r^2}{4a^2}}$$

Parameters a and G are adjusted to reproduce the pairing gap in the symmetric nuclear matter calculated using the Gogny force.

Relativistic energy density functional

Couplings are density-dependent

$$\alpha_i(\rho_v) = a_i + (b_i + c_i x)\, e^{-d_i x}, \quad x = \rho/\rho_{sat}, \quad i \equiv s,\ v,\ tv$$

Model parameters

$$a_s,\ b_s,\ c_s,\ d_s,\ a_v,\ b_v,\ d_v,\ b_{tv},\ d_{tv},\ \delta_s$$

Adjusted to empirical ground-state properties of finite nuclei.

Empirical ground-state properties of finite nuclei can only determine a small set of parameters.

Nuclear many-body correlations

Implicitly included in the EDF

- short-range $\rightarrow$ hard repulsive core of the NN-interaction
- long-range $\rightarrow$ mediated by nuclear resonance modes (giant resonances)
- the corresponding corrections vary smoothly with the number of nucleons $\rightarrow$ absorbed in the model parameters

- heavy deformed systems present best examples of mean-field nuclei
- high density of states reduces the shell effects

Adjusting the model parameters

Empirical mass formula

The calculated masses of finite nuclei are primarily sensitive to three leading terms in the empirical mass formula

$$E_B = a_v A + a_s A^{2/3} + a_4 \frac{(N-Z)^2}{4A} + \cdots$$

Fitting strategy

- generate families of effective interactions that are characterized by different values of a_v, a_s and the symmetry energy $S_2(0.12\text{fm}^{-3})$
- determine which parametrization minimizes the deviation from empirical binding energies of a large set of deformed nuclei

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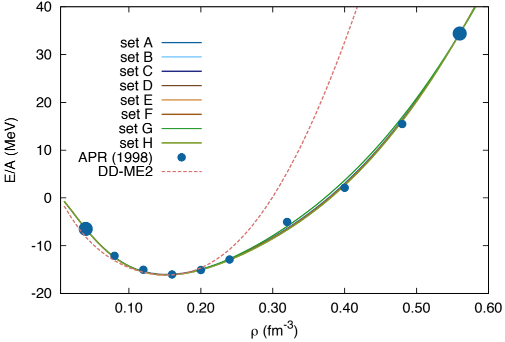
Adjusting the model parameters

Two points from the microscopic EoS curve of Akmal, Pandharipande and Ravenhall are kept fixed.

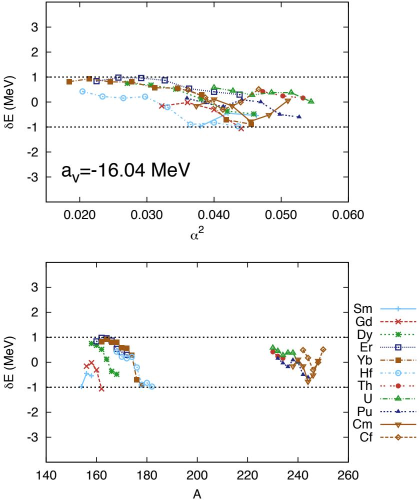


$$\rho_{sat} = 0.152 \text{ fm}^{-3}$$
$$m_D = 0.58m$$
$$a_4 = 33 \text{ MeV}$$

$$K_{nm} = 230 \text{ MeV}$$

$$a_v = -16.04 \text{ MeV}, \ldots, -16.14 \text{ MeV}$$

Adjusting the model parameters



Rare-earth region

Sm (Z=62), Gd (Z=64), Dy (Z=66), Er (Z=68), Yb (Z=70), Hf (Z=72)

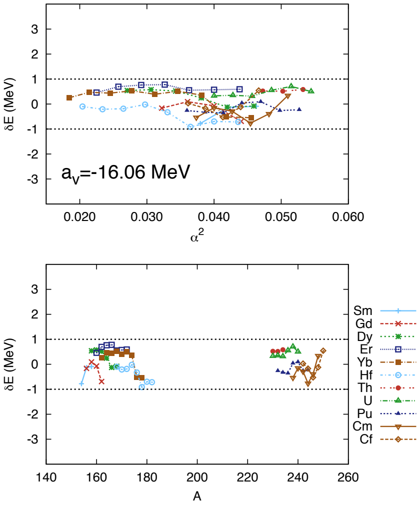
Actinides

Th (Z=90), U (Z=92), Pu (Z=94), Cm (Z=96), Cf (Z=98)

Total

64 isotopes used in the fit

Adjusting the model parameters



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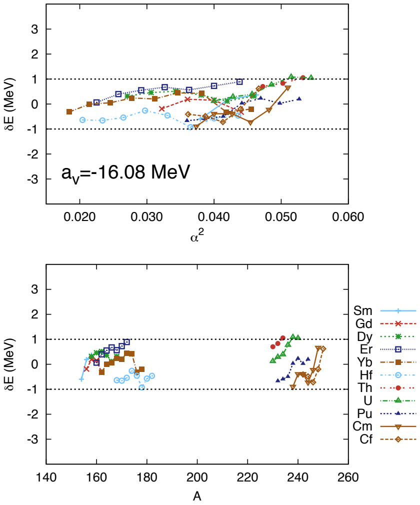

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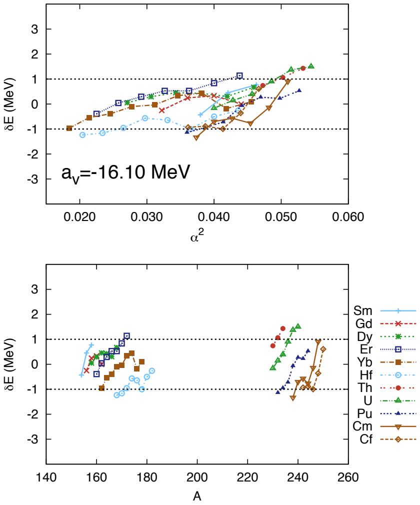

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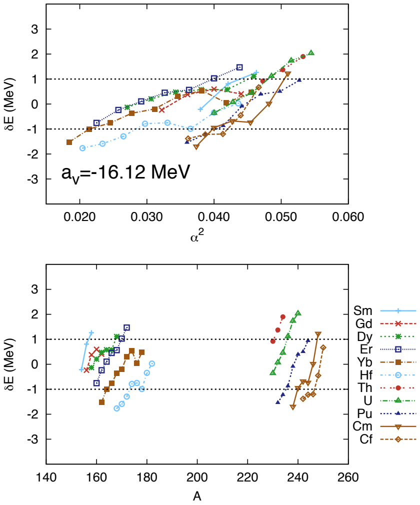

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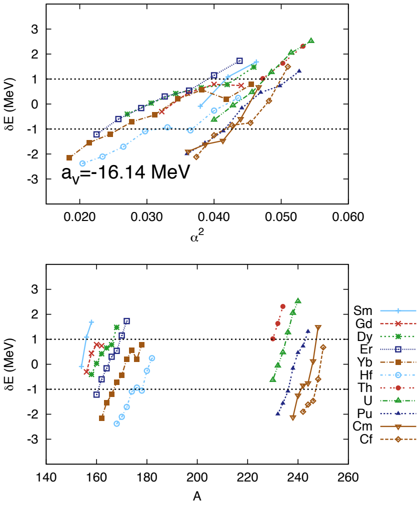

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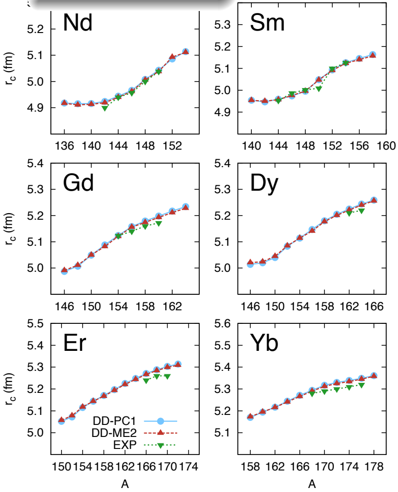
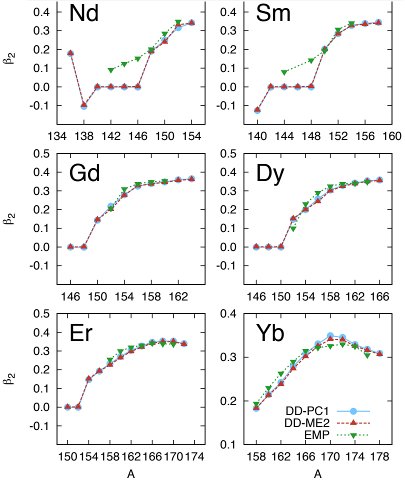
Ground-state properties

Charge radii

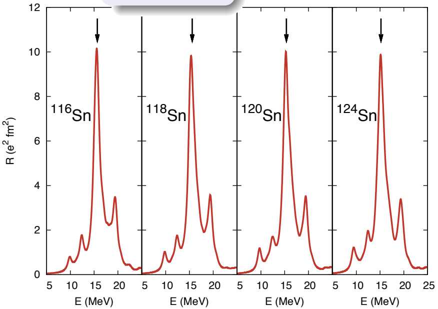


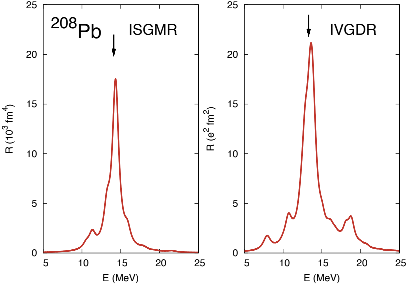
Quadrupole deformations



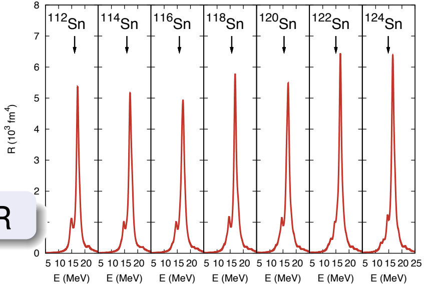
Excitation energies of collective modes

ISGMR





IVGDR



Implementation of the collective Hamiltonian model based on the SCRMF

Collective Hamiltonian

$$\mathcal{H}_{coll} = \mathcal{T}_{rot} + \mathcal{T}_{vib} + \mathcal{V}_{coll}$$

Rotational energy

$$\mathcal{T}_{rot} = \frac{1}{2} \sum_{k=1}^{3} \frac{\hat{J}_k^2}{\mathcal{I}_k}$$

The moments of inertia are calculated by using the Inglis-Belyaev formula.

Implementation of the collective Hamiltonian model based on the SCRMF

Collective Hamiltonian

$$\mathcal{H}_{coll} = \mathcal{T}_{rot} + \mathcal{T}_{vib} + \mathcal{V}_{coll}$$

Vibrational energy

$$\mathcal{T}_{vib} = -\frac{\hbar^2}{2\beta^4\sqrt{wr}}\left[\partial_\beta\sqrt{\frac{r}{w}}\beta^4 B_{\gamma\gamma}\partial_\beta - \partial_\beta\sqrt{\frac{r}{w}}\beta^3 B_{\beta\gamma}\partial_\gamma\right]$$
$$-\frac{\hbar^2}{\sin 3\gamma\sqrt{wr}}\left[-\frac{1}{\beta^2}\partial_\gamma\sqrt{\frac{r}{w}}\sin 3\gamma B_{\beta\gamma}\partial_\beta + \frac{1}{\beta}\partial_\gamma\sqrt{\frac{r}{w}}\sin 3\gamma B_{\beta\beta}\partial_\gamma\right]$$

The mass parameters are calculated in the cranking approximation .

Implementation of the collective Hamiltonian model based on the SCRMF

Collective Hamiltonian

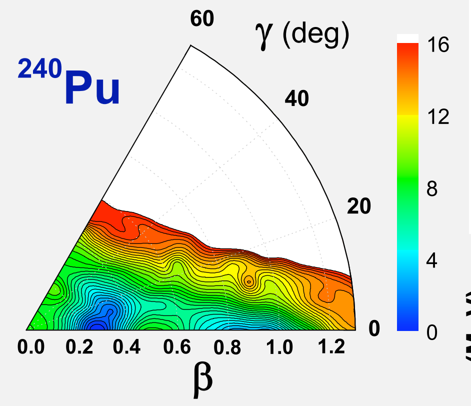
$$\mathcal{H}_{coll} = \mathcal{T}_{rot} + \mathcal{T}_{vib} + \mathcal{V}_{coll}$$

Collective potential

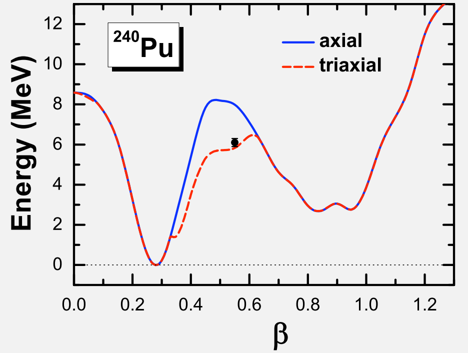
$$\mathcal{V}_{coll}(\beta,\gamma) = E_{tot}(\beta,\gamma) - \Delta V_{vib}(\beta,\gamma) - \Delta V_{rot}(\beta,\gamma)$$

Corresponds to the mean-field potential energy surface with the zero point energy subtracted .

Applications: ^{240}Pu isotope

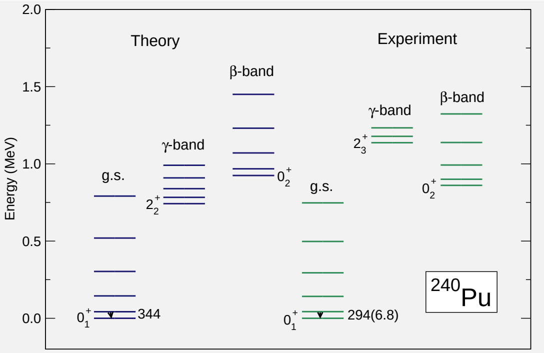


Triaxial effects lower the barrier.



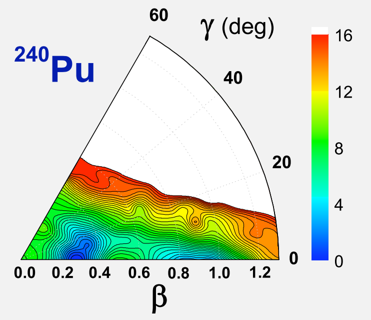
ND and SD minima are separated by the barrier.

Applications: ^{240}Pu isotope



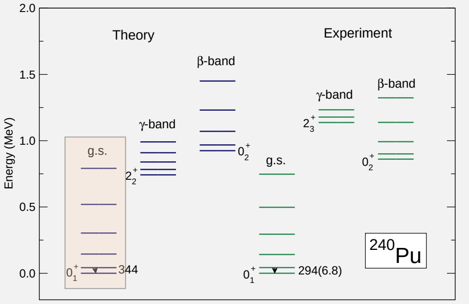
$$E^{th}_{4^+_1} / E^{th}_{2^+_1} = 3.33$$

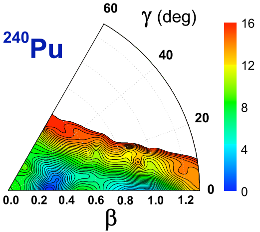
$$E^{exp}_{4^+_1} / E^{exp}_{2^+_1} = 3.31$$



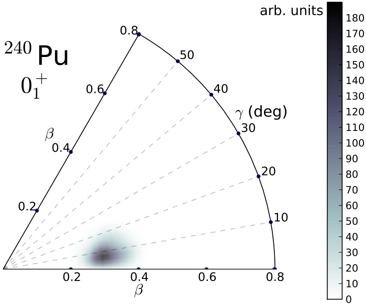
The moments of inertia are renormalized by factor ≈ 1.3 to compensate the difference between IB and TV moments of inertia.

Applications: ^{240}Pu isotope



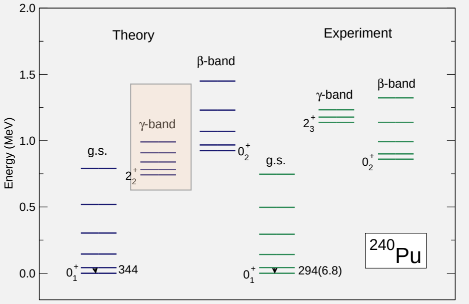


Probability distribution

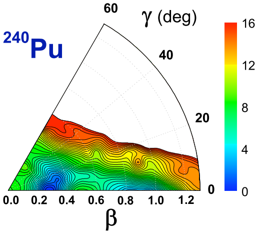


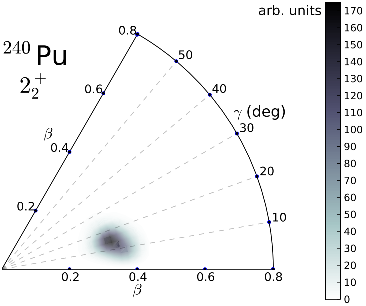
g.s. band head

Applications: ^{240}Pu isotope



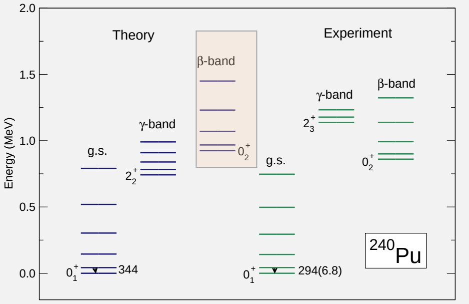
Probability distribution

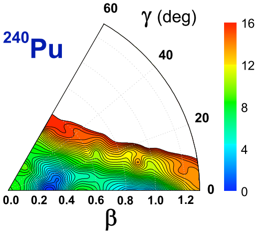




γ band head

Applications: ^{240}Pu isotope





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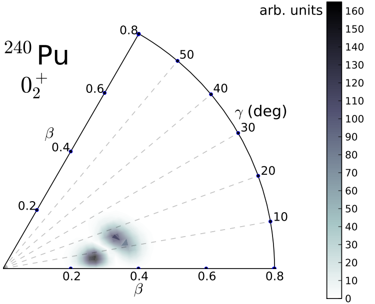
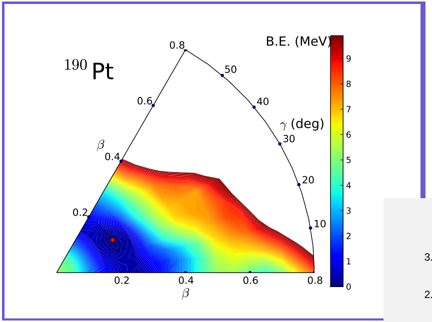
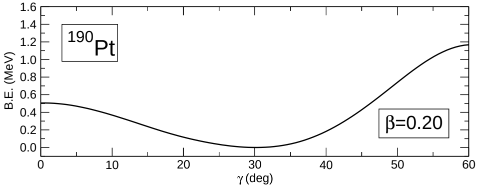
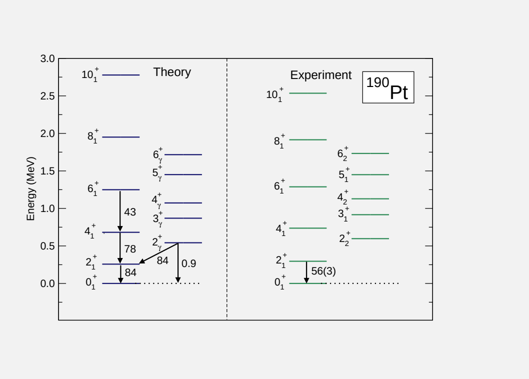


β band head

Applications: Pt isotopes



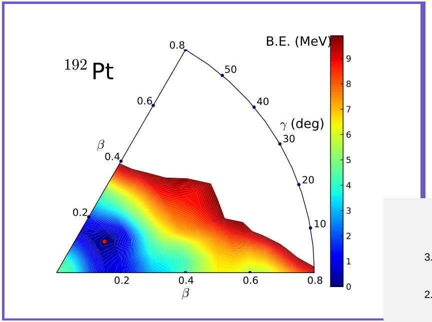


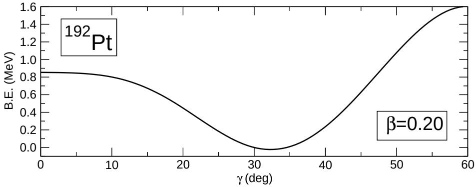


$$E^{th}_{4_1^+} / E^{th}_{2_1^+} = 2.39$$

$$E^{exp}_{4_1^+} / E^{exp}_{2_1^+} = 2.49$$

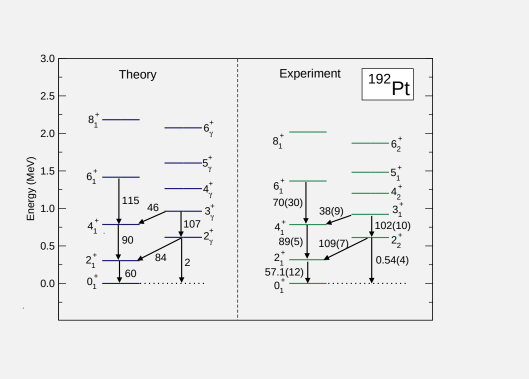
Applications: Pt isotopes



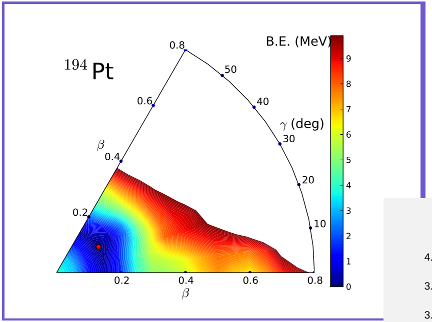


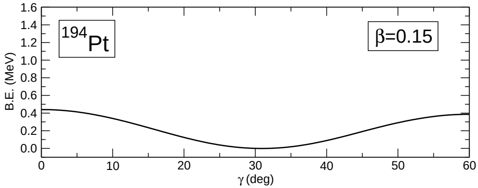
$$E^{th}_{4_1^+} / E^{th}_{2_1^+} = 2.58$$

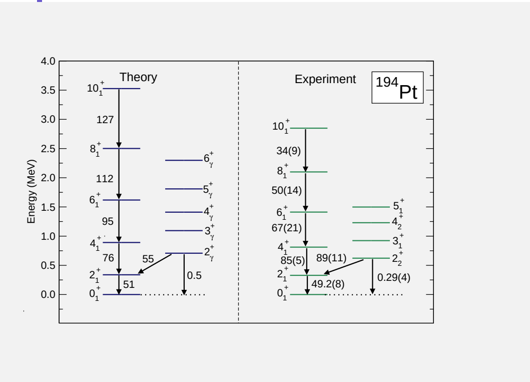
$$E^{exp}_{4_1^+} / E^{exp}_{2_1^+} = 2.48$$



Applications: Pt isotopes



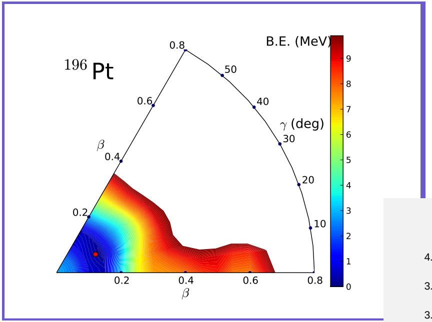


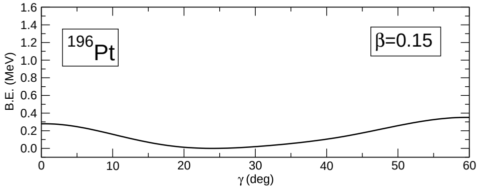


$$E^{th}_{4^+_1} / E^{th}_{2^+_1} = 2.63$$

$$E^{exp}_{4^+_1} / E^{exp}_{2^+_1} = 2.47$$

Applications: Pt isotopes





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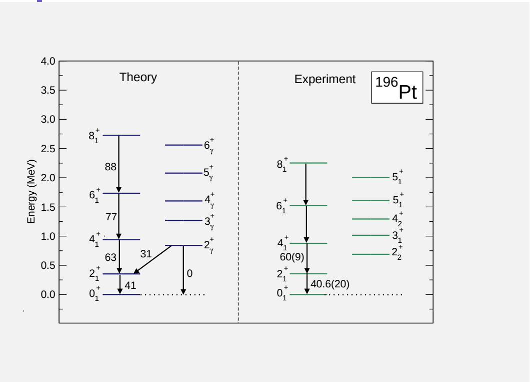
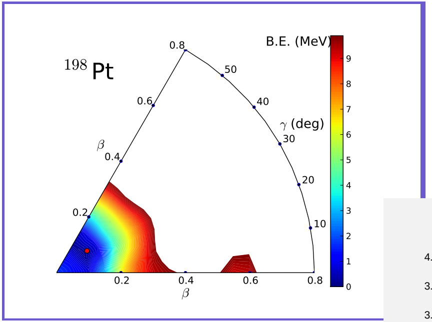
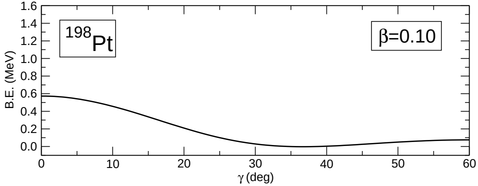
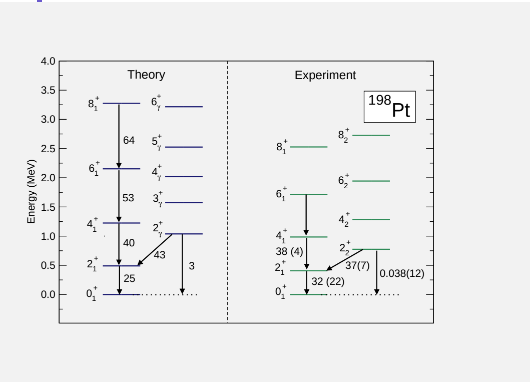


Applications: Pt isotopes



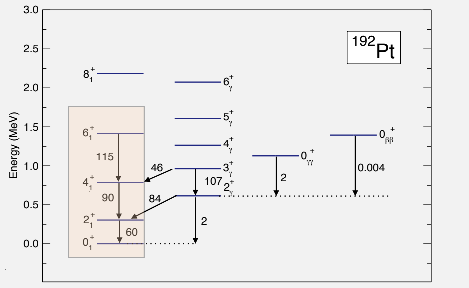


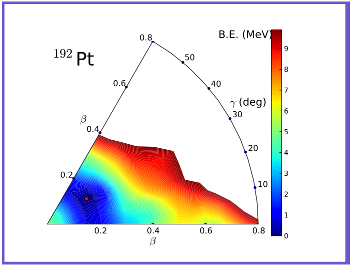


$$E^{th}_{4^+_1} / E^{th}_{2^+_1} = 2.49$$

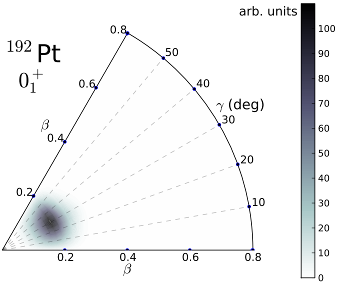
$$E^{exp}_{4^+_1} / E^{exp}_{2^+_1} = 2.42$$

Applications: ^{192}Pt isotope





Collective probability



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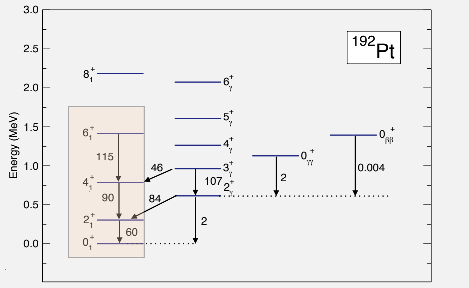
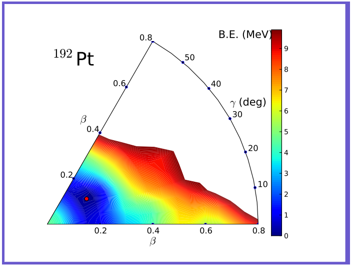
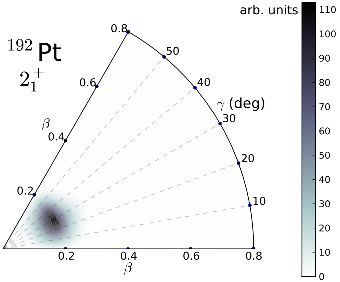




Collective probability



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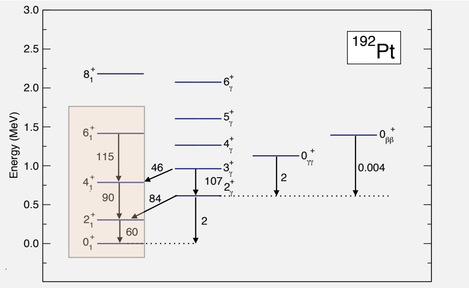

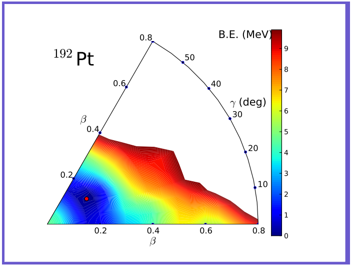

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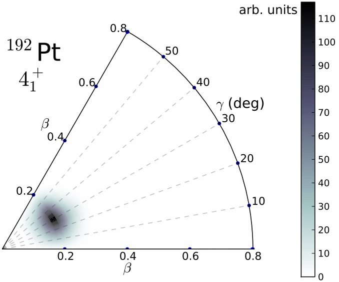

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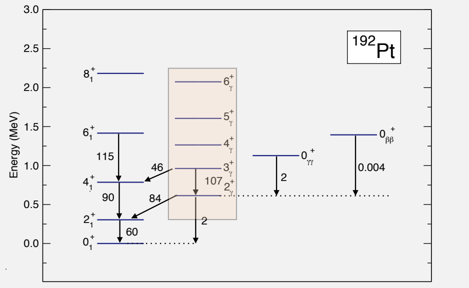
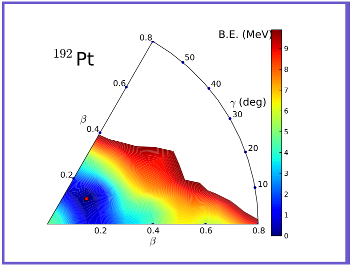
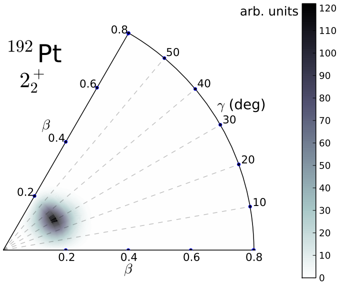




Collective probability



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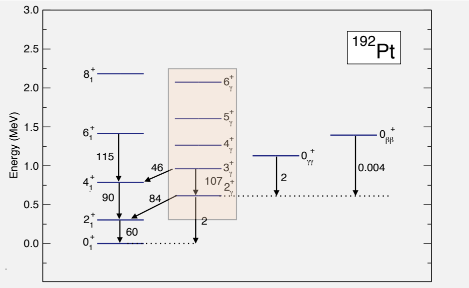

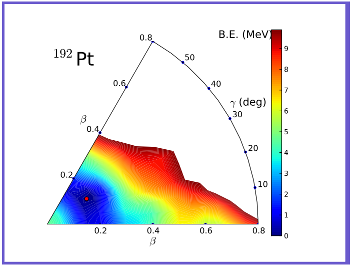

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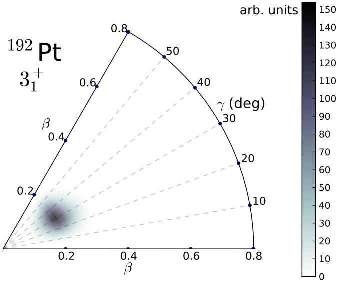

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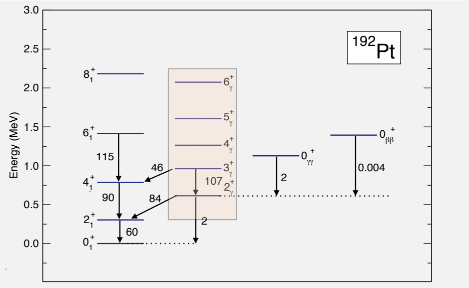
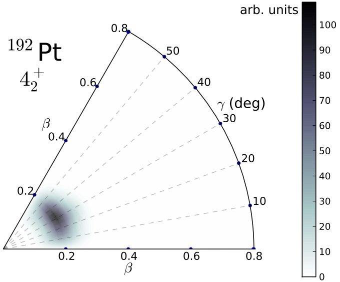
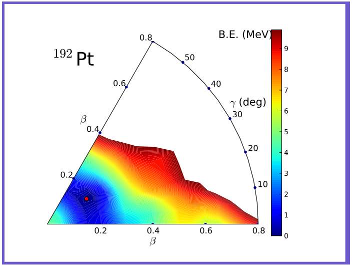


Collective probability





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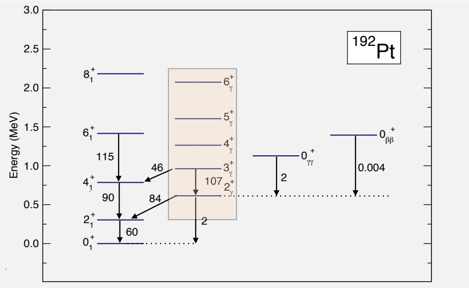

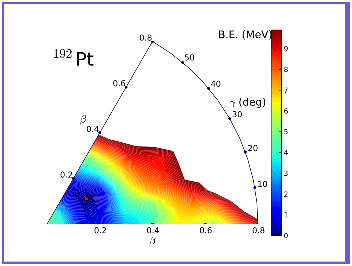

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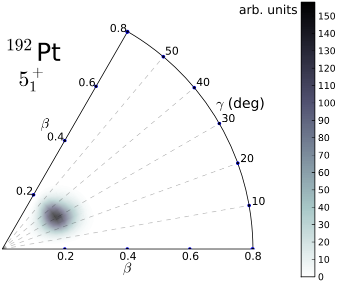

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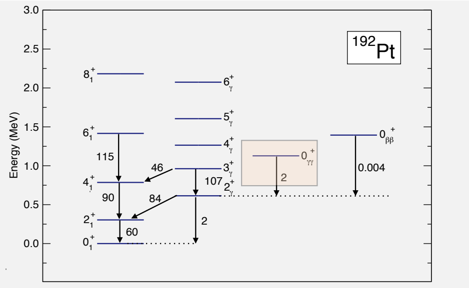

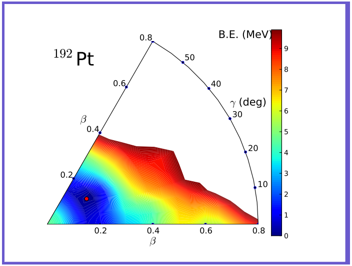

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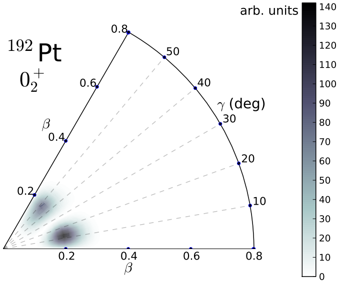

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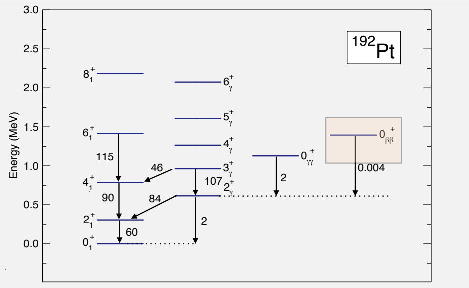

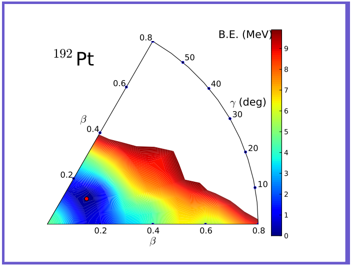

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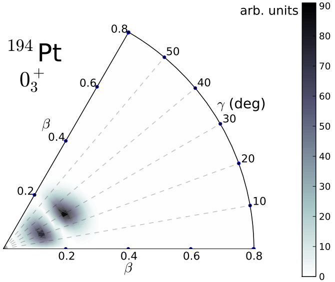

Summary and outlook

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Unified microscopic description of the structure of stable and nuclei far from stability, and reliable extrapolations toward the drip lines.

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When extended to take into account collective correlations, it describes deformations and shape-coexistence phenomena associated with shell evolution.

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Further improvements of the model and more systematic calculations.

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Collaborators

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