

# Collective quadrupole excitations within a selfconsistent approach

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## Microscopic models

phenomenological potentials (eg Nilsson),  
 mean field theory ([Skyrme int.](#), Gogny, RMF)  
 pairing interactions: [constant  \$G\$](#) ,  $\delta$  force.

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↓  
 ATDHFB ([cranking](#)), GCM  
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## Hamiltonian in the collective space

quadrupole variables:  $\beta, \gamma$  + Euler angles

"pairing" variables:  $\Delta_{p,n}, \phi_{p,n}$

Example:  $^{128}\text{Xe}$



## 1. ATDHFB

$$R = \begin{pmatrix} \rho & \kappa \\ -\kappa^* & 1 - \rho^* \end{pmatrix}, \quad \rho_{\mu\nu} = u_\mu^2 \delta_{\mu\nu}, \quad \kappa_{\mu\nu} = s_{\bar{\mu}} u_\mu v_\mu \delta_{\bar{\mu}\nu}$$

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$R(q_k(t))$ , collective variables  $q_k$ ;

Collective (classical) energy

$$E_c = E_{\text{kin}} + V = \frac{1}{2} \sum_{kj} \dot{q}_j \dot{q}_k B_{kj} + V(q_k)$$

Quantum hamiltonian

$$H = -\frac{\hbar^2}{2} \frac{1}{\sqrt{\det B}} \sum_{ij} \frac{\partial}{\partial q_i} \sqrt{\det B} (B^{-1})_{ij} \frac{\partial}{\partial q_j} + V(q_k)$$



## Mass parameters

$$B_{kj} = \frac{\hbar^2}{2} \sum_{\mu\nu} \frac{f_{j,\mu\nu} f_{k,\mu\nu}^* + f_{j,\mu\nu}^* f_{k,\mu\nu}}{(E_\mu + E_\nu)}$$

$$\begin{pmatrix} 0 & f_k \\ -f_k^* & 0 \end{pmatrix} = \left( \frac{\partial R_0}{\partial q_k} \right)_{qp}$$



## 2. Collective variables

Quadrupole part:  $q_\mu = \langle Q_{2\mu} \rangle \rightarrow q_0, q_2, \Omega$  — Euler angles

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$$q_0 = \left\langle \sum_{n=1}^A (3z_n^2 - r_n^2) \right\rangle$$

$$q_2 = \left\langle \sum_{n=1}^A (x_n^2 - y_n^2) \right\rangle$$

$$\beta \cos \gamma = q_0 \sqrt{\pi/5} / A \langle r^2 \rangle$$

$$\beta \sin \gamma = q_2 \sqrt{3\pi/5} / A \langle r^2 \rangle$$

$$\min \langle | \mathcal{H} - \sum_i \alpha_i Q_i | \rangle \quad \text{with} \quad \langle | Q_i | \rangle = q_i \rightarrow R(q_i)$$



Pairing part:  $\Delta, \phi$

$v_i \rightarrow v_i(\Delta, \lambda)$  as in BCS,  $\lambda$  from  $\langle N \rangle = N_0$

$$R(\Delta, \phi) = \begin{pmatrix} \rho & e^{-2i\phi} \kappa \\ -e^{2i\phi} \kappa^* & \rho \end{pmatrix}$$

Other possibility

$$\tilde{\Delta} \sim \left\langle \sum_{k>0} e^{-2i\phi} s_k c_k^+ c_k^+ + e^{2i\phi} s_k^* c_k c_k^- \right\rangle \sim \sum_{k>0} u_k v_k$$



### 3. Density matrix derivatives

$$f_{\alpha,\mu\nu} = \frac{s_\nu}{E_\mu + E_\nu} Q_{\mu\bar{\nu}} (u_\mu v_\nu + v_\mu u_\nu)$$

$$f_{\Delta,\mu\nu} = s_\nu \delta_{\mu\bar{\nu}} \frac{1}{2E_\mu^2} ((\epsilon_\mu - \lambda) + \Delta \partial_\Delta \lambda)$$

$$f_{\phi,\mu\nu} = i s_\nu \delta_{\mu\bar{\nu}} \frac{\Delta}{E_\mu}$$

$$\frac{\partial q}{\partial \alpha} = \sum_{\gamma\mu} \frac{|Q_{\bar{\mu}\gamma}|^2 (u_\gamma v_\mu + v_\gamma u_\mu)^2}{E_\mu + E_\gamma}$$

$$\frac{\partial q}{\partial \Delta} = \Delta \sum_{\mu>0} \frac{Q_{\mu\mu} ((\epsilon_\mu - \lambda) + \Delta \partial_\Delta \lambda)}{E_\mu^3}$$



## Matrix of mass parameters

$$\begin{pmatrix} (B_{\text{rot}}) & 0 & 0 \\ 0 & (B_{\text{vib}}) & 0 \\ 0 & 0 & (B_{\text{rotpair}}) \end{pmatrix}$$

$$B_{\text{vib}} = \begin{pmatrix} B_{\beta\beta} & B_{\beta\gamma} & B_{\beta\Delta_p} & B_{\beta\Delta_n} \\ B_{\gamma\beta} & B_{\gamma\gamma} & B_{\gamma\Delta_p} & B_{\gamma\Delta_n} \\ B_{\Delta_p\beta} & B_{\Delta_p\gamma} & B_{\Delta_p\Delta_p} & B_{\Delta_p\Delta_n} \\ B_{\Delta_n\beta} & B_{\Delta_n\gamma} & B_{\Delta_n\Delta_p} & B_{\Delta_n\Delta_n} \end{pmatrix}$$

$$B_{\text{rotpair}} = \begin{pmatrix} B_{\phi_p\phi_p} & 0 \\ 0 & B_{\phi_n\phi_n} \end{pmatrix}$$





## 4. Basis functions

$$F_j(\beta, \gamma, \Omega) f_k^{(p)}(\Delta_p) f_m^{(n)}(\Delta_n)$$

Eigenfunctions of the hamiltonian without deformation-pairing coupling

$$\begin{pmatrix} (B_{\text{def}}(q, \Delta)) & 0 \\ 0 & (B_{\text{pair}}(q, \Delta)) \end{pmatrix}$$

Matrix elements of some operators (eg  $1/J_k$ )

$$\begin{aligned} \langle \Psi_j | P | \Psi_k \rangle &= \int dq_{\text{def}} F_j F_k \int P(q_{\text{def}}, \Delta_n, \Delta_p) |f_0^p|^2 |f_0^n|^2 g_{\text{pair}}^p g_{\text{pair}}^n d\Delta_p d\Delta_n = \\ &= \int dq_{\text{def}} F_j F_k P_{\text{av}}(q_{\text{def}}) \end{aligned}$$

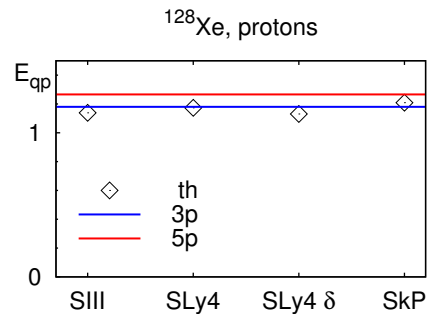
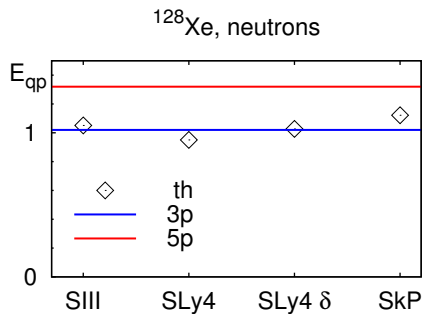


Example:  $^{128}\text{Xe}$

Skyrme interactions: SIII, SLy4, SkP

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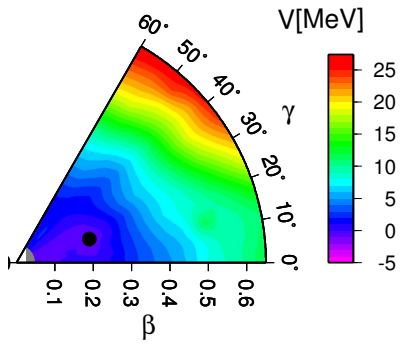
Pairing strength



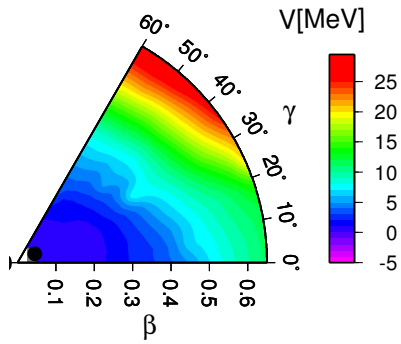
$$E_{\text{qp}} = \langle E_{\text{qp},\text{min}} \rangle$$



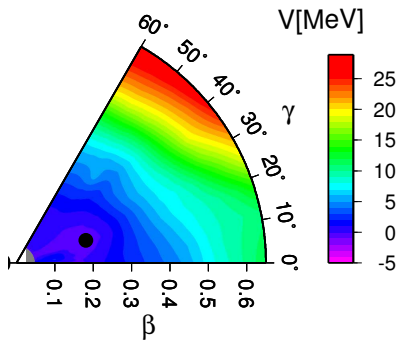
$^{128}\text{Xe}$ , SIII



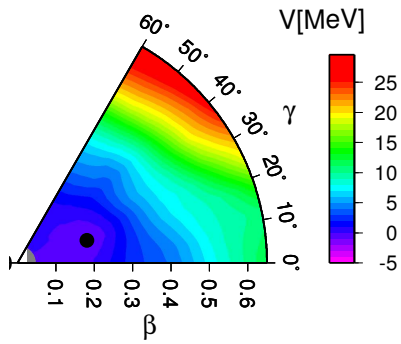
$^{128}\text{Xe}$ , SkP

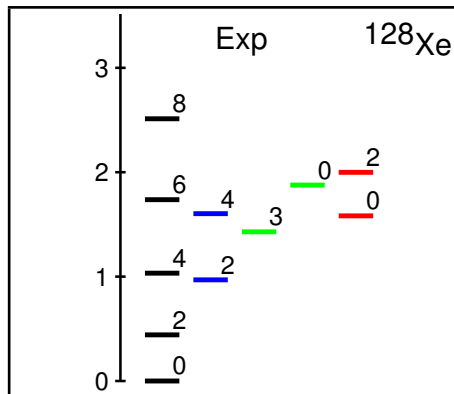
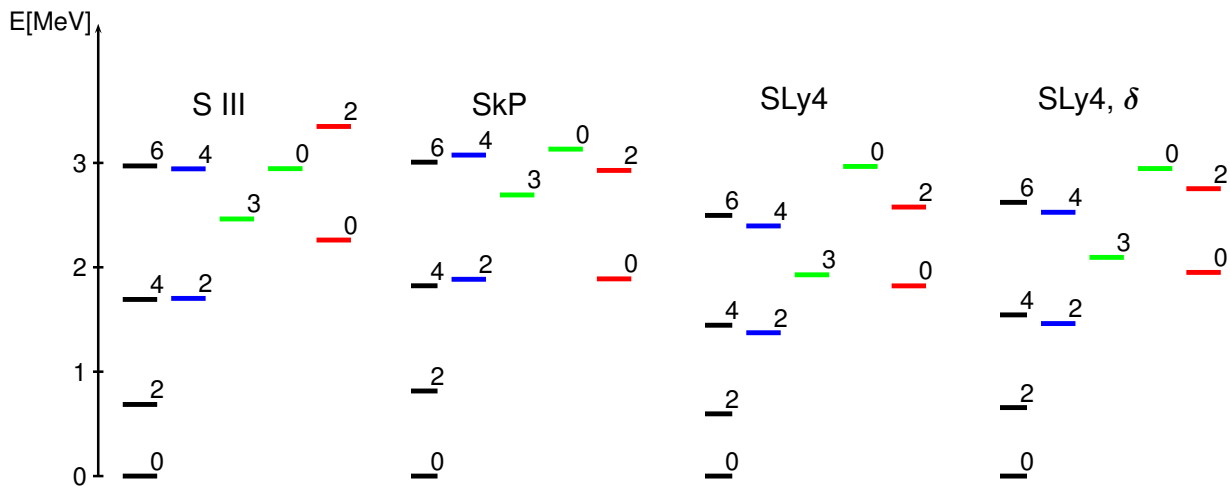


$^{128}\text{Xe}$ , SLy4

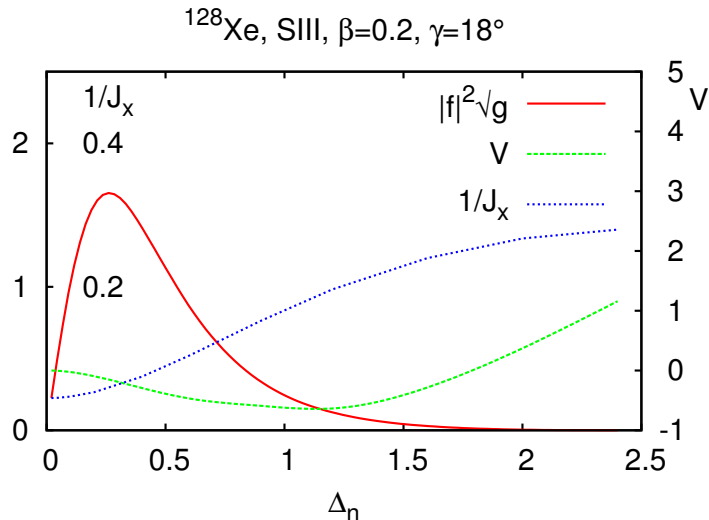


$^{128}\text{Xe}$ , SLy4,  $\delta$  force





## Ground state functions of the collective pairing hamiltonian

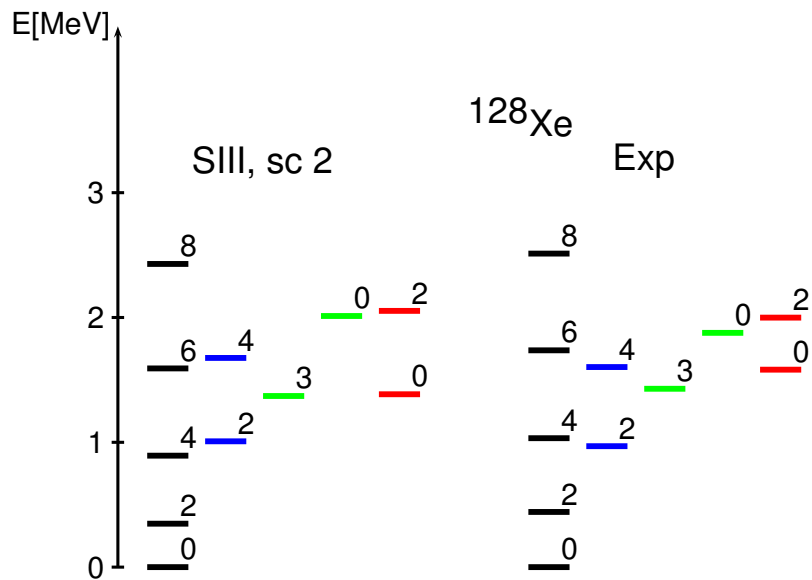


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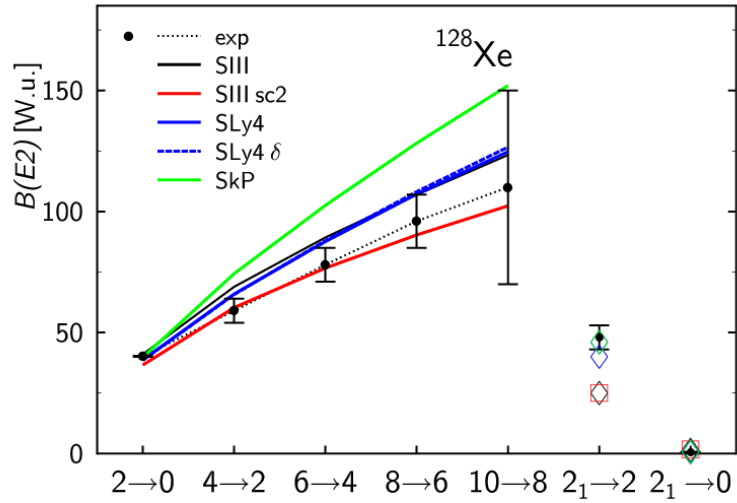
Average scaling factor for  $1/J_k$ ,  $B_{\mu\nu}^{-1} \sim 2$

Estimates for  $^{178}\text{Hf}$ ,  $^{238}\text{U}$ :  $\sim 1.15$





## E2 transitions



## 5. Conclusions

Possible large effect of pairing dynamical degrees of freedom on mass parameters

To do

Larger scale calculations, question of adiabaticity of 'pairing' variables,  $\delta$  interaction, ...

